Introduction to Spectral Methods
Numerical Solution of Poisson's Equation

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Outline

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When $f = 0$, we have Laplace's equation as a special case.

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- And much more ...

Outline

Poisson's Equation in one dimension

$$\frac{d^2u}{dx^2} = f(x)$$

Subject to $u(a) = g_1$ and $u(b) = g_2$.

It is not always possible to integrate twice to find the solution.

We must translate this continuous problem into a discrete one.

For the space coordinate $x \in [a, b]$, we choose n nodes $\{x_i\}_{i=1}^n$ on which we would like to know the value of $u(x_i) = u_i$.

Input u_1, u_n, $\vec{x} = \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_n \end{bmatrix}$, and $\vec{f} = \begin{bmatrix} f(x_1) \\ f(x_2) \\ \vdots \\ f(x_n) \end{bmatrix}$; output $\vec{u} = \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_n \end{bmatrix}$.

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Is there a way to express $u(x)$ in terms of u_i?

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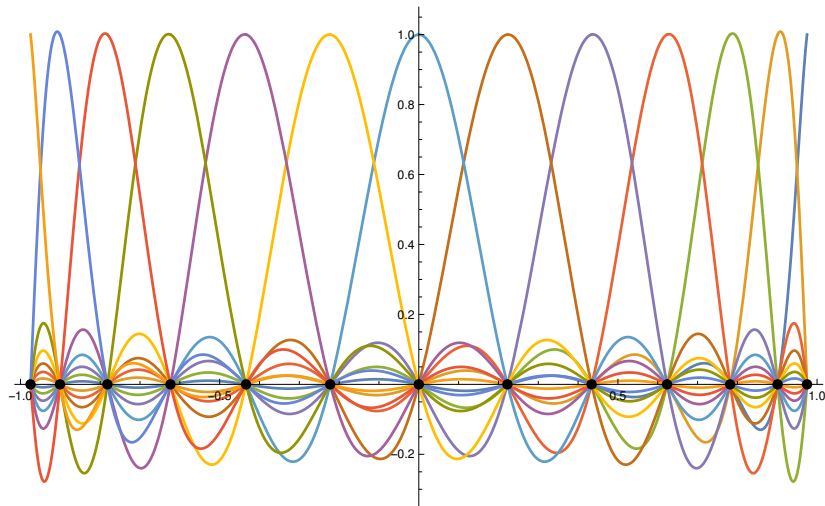
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$$\delta_j(x_i) = \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases}$$

This is how the cardinal functions look like.

Now differentiate the expansion in cardinal functions.

$$u'(x_i) = \sum_j u_j \delta_j'(x_i)$$

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$$\begin{bmatrix} D^2_{11} & D^2_{1j} & D^2_{1n} \\ D^2_{i1} & D^2_{ij} & D^2_{in} \\ D^2_{n1} & D^2_{nj} & D^2_{nn} \end{bmatrix} \begin{bmatrix} u_1 \\ u_j \\ u_n \end{bmatrix} = \begin{bmatrix} D^2_{11} \\ D^2_{i1} \\ D^2_{n1} \end{bmatrix} u_1 + \begin{bmatrix} D^2_{1j} \\ D^2_{ij} \\ D^2_{nj} \end{bmatrix} \begin{bmatrix} u_j \end{bmatrix} + \begin{bmatrix} D^2_{1n} \\ D^2_{in} \\ D^2_{nn} \end{bmatrix} u_n = \begin{bmatrix} f_1 \\ f_i \\ f_n \end{bmatrix}$$

With i and j ranging from 2 to $n-1$.

We have $n-2$ unknowns, but still n equations.

Discard the first and last rows.

$\nabla^2 u = f$ will not hold true at the boundary, but we will ensure the boundary conditions.

$$\left[D_{ij}^2\right]\left[u_j\right] = \left[f_i\right] - u_1\left[D_{i1}^2\right] - u_n\left[D_{in}^2\right]$$
$$\left[u_j\right] = \left[D_{ij}^2\right]^{-1}\left[f_i - u_1 D_{i1}^2 - u_n D_{in}^2\right]$$

With i and j ranging from 2 to $n-1$.

Cardinal functions are obtained from a mother function.

Usually we chose an orthogonal polynomial. It all depends on the domain of your solution.

Chebysheb	$\Psi_n(x) = T_n(x)$	$x \in [-1, 1]$
Legendre	$\Psi_n(x) = P_n(x)$	$x \in [-1, 1]$
Hermite	$\Psi_n(x) = H_n(x)e^{-x^2/2}$	$x \in (-\infty, \infty)$
Laguerre	$\Psi_n(x) = xL_{n-1}^{(1)}(x)e^{-x/2}$	$x \in [0, \infty)$

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Let $\Psi_n(x)$ be a function with n roots on $\{x_i\}_{i=1}^n$, i. e. $\Psi_n(x_i) = 0$. By expanding $\Psi_n(x)$ around its roots we can find a expression for $\delta_j(x)$.

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$$\Psi_n(x) = 0 + \Psi'_n(x_j)(x - x_j) + \frac{\Psi''_n(x_j)}{2}(x - x_j)^2 + O((x - x_j)^3)$$

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$$\delta_j(x) = \frac{\Psi_n(x)}{(x - x_j)\Psi'_n(x_j)} = 1 + \frac{\Psi''_n(x_j)}{2\Psi'_n(x_j)}(x - x_j) + O((x - x_j)^2)$$

Differentiate to obtain the elements of D.

$$D_{ij} = \delta'_j(x_i) = \frac{1}{x_i - x_j} \frac{\Psi'_n(x_i)}{\Psi'_n(x_j)}, \qquad i \neq j$$

The diagonal is a bit tricky.

$$D_{jj} = \delta'_j(x_j) = \frac{\Psi''_n(x_j)}{2\Psi'_n(x_j)}$$

MATLAB code

Solve Poisson's equation

$$\frac{d^2u}{dx^2} = 1000\cos(5\pi x)e^{-x^2}$$

with boundary conditions $u(-1) = 2$, $u(1) = -1$.

MATLAB code
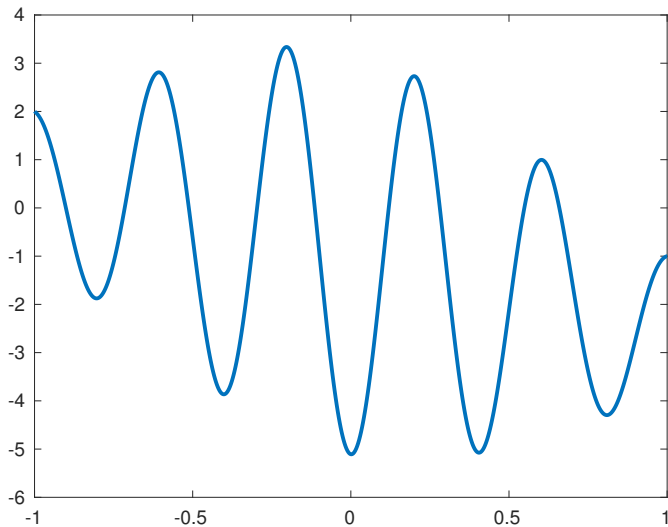
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```
1  [D,x]=chebD(n); D2=D*D;   % Differentiation matrix and nodes
2  f=1000*cos(5*pi*x).*exp(-x.^2); % Source term
3  u=zeros(n,1); u([1,n])=[-1,2]; % Impose boundary conditions
4  % Solve and plot
5  u(2:n-1)=D2(2:n-1,2:n-1)\(f(2:n-1)-D2(2:n-1,[1,n])*u([1,n]));
6  plot(x,u);
```

Here is our solution.

Outline

Poisson's Equation in two dimensions

Our domain of solution will be the rectangle $[a, b] \times [c, d]$

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = f(x, y)$$

Subject to boundary conditions

$$u(x, y) = \begin{cases} g_1(y), & x = a \\ g_2(y), & x = b \\ h_1(x), & y = c \\ h_2(x), & y = d \end{cases}$$

In 2D we use matrices for u and f.

$$u(x_i, y_j) = u_{ij}$$
$$f(x_i, y_j) = f_{ij}$$

$$u(x, y) = \sum_{k,l} u_{kl} \delta_k(x) \delta_l(y)$$

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$$\nabla^2 u = D^2 U + U(D^2)^\mathsf{T}$$

Now solve the system of linear equations.

$$D^2 U + U(D^2)^\mathsf{T} = F$$

Boundary conditions are the tricky part.

$$U = \tilde{U} + U_B = \begin{bmatrix} 0 & 0 & 0 \\ 0 & u_{ij} & 0 \\ 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} u_{11} & u_{1j} & u_{1n} \\ u_{i1} & 0 & u_{in} \\ u_{n1} & u_{nj} & u_{nn} \end{bmatrix}$$

$$D^2 \tilde{U} + \tilde{U}(D^2)^\mathsf{T} = F - D^2 U_B - U_B(D^2)^\mathsf{T} = \tilde{F}$$

Discard first and last rows and columns.

$\nabla^2 u = f$ will not hold true at the boundary, but we will ensure the boundary conditions.

$$(D_{ij}^2)\tilde{U}_{ij} + \tilde{U}_{ij}(D_{ij}^2)^{\mathsf{T}} = \tilde{F}_{ij}$$

With i and j ranging from 2 to $n - 1$.

Now we just have to solve $(n - 2)^2 \times (n - 2)^2$ system of linear equations.

Matrix inversion does not help.

A matrix equation of the form $AX + XB = C$ is known as a Sylvester equation.

There's a built-in MATLAB function that solves them.

```
X=sylvester(A,B,C)
```

This algorithm uses the Schur decomposition to solve a block-triangular matrix very efficiently.

Two dimensional example

Solve Laplace's equation on $(x, y) \in [-1, 1] \times [-1, 1]$

$$\nabla^2 u = 0$$

with boundary conditions

$$u(x, y) = \begin{cases} \sin^4(\pi x), & y = 1 \text{ and } -1 < x < 0, \\ \frac{1}{5} \sin(3\pi y), & x = 1, \\ 0, & \text{otherwise.} \end{cases}$$

MATLAB Code

```matlab
1  [D,x]=chebD(n); D2=D*D; y=x';
2  [xx, yy]=ndgrid(x);
3
4  % Boundary conditions
5  g=[0.2*sin(3*pi*y); 0*y];
6  h=[(x<0).*sin(pi*x).^4, 0*x];
7  uu=zeros(n);
8  uu([1 n],:)=g;
9  uu(:,[1 n])=h;
10
11 % Solve Laplace's equation
12 F=zeros(n);
13 RHS=F-D2(:,[1 n])*g-h*D2(:,[1 n])';
14 uu(2:n-1, 2:n-1)=sylvester(D2(2:n-1, 2:n-1), ...
15 D2(2:n-1, 2:n-1)', RHS(2:n-1, 2:n-1));
16
17 surfl(xx,yy,uu,'light'); colormap(jet(256));
18 shading interp; axis square;
```

Here is our solution.
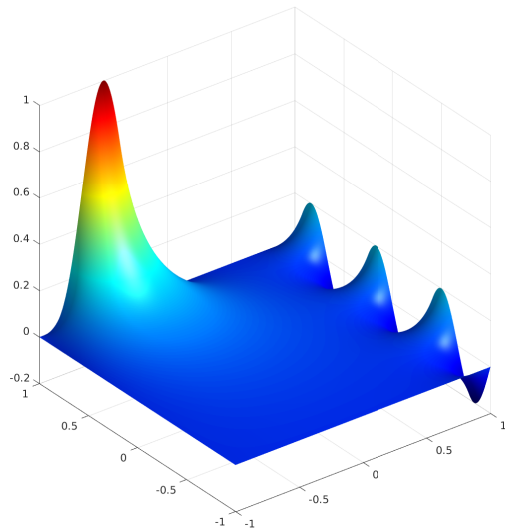
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