

19 #1

Since $\Delta u = 0$ in the weak sense, u is weakly harmonic.

Recall that a weakly harmonic function is also harmonic.

By the Mean Value Property, for each $x \in \mathbb{R}^n$,

$$u(x) = \frac{1}{|B(x,1)|} \int_{B(x,1)} u(y) dy$$

and hence

$$|u(x)| \leq \frac{1}{|B(x,1)|} \int_{B(x,1)} |u(y)| dy \leq \frac{C}{|B(x,1)|}.$$

As $|B(x,1)|$ does not depend on x , it follows that u

is bounded on \mathbb{R}^n and hence by Liouville's Theorem,

u is constant.

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Eqn #2 We have

$$f(x) = (\Delta - aI) \int K_a(x-y) f(y) dy$$

$$f(x) = (\Delta - bI) \int K_b(x-y) f(y) dy$$

Now that

$$(\Delta - aI)(\Delta - bI) \int (c_1 K_a(x-y) + c_2 K_b(x-y)) f(y) dy$$

$$= c_1 (\Delta - bI) f(x) + c_2 (\Delta - aI) f(x)$$

$$= (c_1 + c_2) \Delta f(x) - c_1 b f(x) - c_2 a f(x) \quad (a)$$

Taking $c_1 = (a-b)^{-1}$ and $c_2 = (b-a)^{-1}$, we have that $(a) = f(x)$.

Thus $(a-b)^{-1} K_a + (b-a)^{-1} K_b$ is a fundamental solution of $(\Delta - aI)(\Delta - bI)$.

We have $\Delta^2 - \Delta = \Delta(\Delta - I)$. So the fundamental solution is $-K_0 + K_1$, where K_0 is the fundamental solution for Δ and K_1 is the fundamental sol. for $\Delta - I$.

When $n=3$, $K_0 = \frac{1}{4\pi|x|}$.

$K_1 \neq \frac{e^{ix}}{4\pi|x|} \rightarrow$ Fundamental solution of Helmholtz eqn. in \mathbb{R}^3 .

Bessel Potential

Formally we solve

$$\Delta u - u = \delta_0 \rightarrow -\Delta u + u = -\delta_0$$

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Take the Fourier transform of both sides to get

$$(1+|y|^2) \hat{u} = -\hat{\delta}_0$$

$$\rightarrow u = \left(\frac{-\hat{\delta}_0}{(1+|y|^2)} \right)^\vee = \frac{1}{(2\pi)^{3/2}} (-\delta_0 + B) = -\frac{1}{(2\pi)^{3/2}} B(x)$$

where $B = \frac{1}{1+|y|^2}$. Then as $\frac{1}{1+|y|^2} = \int_0^\infty e^{-t} e^{-t|y|^2} dt$,

$$B = \left(\frac{1}{1+|y|^2} \right)^\vee = \frac{1}{(2\pi)^{3/2}} \int_0^\infty e^{-t} \int_{\mathbb{R}^3} e^{ix \cdot y - t|y|^2} dy dt.$$

$$\rightarrow u = -\frac{1}{(2\pi)^3} \int_0^\infty e^{-t} \int_{\mathbb{R}^3} e^{ix \cdot y - t|y|^2} dy dt = K_1$$

Pr 99 #3:

Claim: $\lambda = \min_{\substack{w \\ \frac{\partial w}{\partial \nu} = 0 \text{ on } \partial U}} \frac{\langle Lw, w \rangle}{\langle w, w \rangle}$ is the smallest eigenvalue of L .

Pr: We first show $\min_w \frac{\langle Lw, w \rangle}{\langle w, w \rangle}$ is indeed an eigenvalue. Let u be the function with $\frac{\partial u}{\partial \nu} = 0$ on ∂U s.t. $\frac{\langle Lw, w \rangle}{\langle w, w \rangle}$ attains its minimum. Then for all v with $\frac{\partial v}{\partial \nu} = 0$ on ∂U

$$0 = \frac{d}{d\varepsilon} \frac{\langle Lu + \varepsilon Lv, u + \varepsilon v \rangle}{\langle u + \varepsilon v, u + \varepsilon v \rangle} \Big|_{\varepsilon=0}$$

Then let $F(w) := \frac{\langle Lw, w \rangle}{\langle w, w \rangle}$. Then for all v with $\frac{\partial v}{\partial \nu} = 0$ on ∂U ,

$$\begin{aligned} 0 &= \lim_{\varepsilon \rightarrow 0} \frac{F(u + \varepsilon v) - F(u)}{\varepsilon} = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \left[\frac{\langle Lu + \varepsilon Lv, u + \varepsilon v \rangle}{\langle u + \varepsilon v, u + \varepsilon v \rangle} - \frac{\langle Lu, u \rangle}{\langle u, u \rangle} \right] \\ &= \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \left[\frac{\langle Lu, u \rangle + \varepsilon \langle Lu, v \rangle + \varepsilon \langle Lv, u \rangle + \varepsilon^2 \langle Lv, v \rangle}{\langle u, u \rangle + 2\varepsilon \langle u, v \rangle + \varepsilon^2 \langle v, v \rangle} - \frac{\langle Lu, u \rangle}{\langle u, u \rangle} \right] \end{aligned}$$

$$\rightarrow 0 = \langle Lu, v \rangle \langle u, u \rangle + \langle Lv, u \rangle \langle u, u \rangle - 2 \langle Lu, u \rangle \langle u, v \rangle.$$

Since L is self-adjoint, we have

$$0 = 2 \langle Lu, v \rangle \langle u, u \rangle - 2 \langle Lu, u \rangle \langle u, v \rangle$$

and hence

$$\int (Lu) v \, dx = \int \frac{\langle Lu, u \rangle}{\langle u, u \rangle} u v \, dx \quad \forall v.$$

Thus $Lu = \frac{\langle Lu, u \rangle}{\langle u, u \rangle} u$. If μ was another eigenvalue of L , then

$$\lambda \leq \frac{\langle Lw, w \rangle}{\langle w, w \rangle} = \frac{\mu \langle w, w \rangle}{\langle w, w \rangle} = \mu. \quad \#$$

Self-adjointness of L is proven as follows:

$$\int (Lu) v \, dx = - \int \sum_{i,j} \partial_{x_j} (a^{ij}(x) \partial_{x_i} u) v \, dx = - \sum_{i,j} \int \partial_{x_j} (a^{ij}(x) \partial_{x_i} u) v \, dx$$

$$\stackrel{\text{since } \frac{\partial u}{\partial \nu} = 0}{=} - \sum_{i,j} \int a^{ij} \partial_{x_i} u \partial_{x_j} v \, dx \stackrel{\text{since } \frac{\partial v}{\partial \nu} = 0}{=} - \sum_{i,j} \int \partial_{x_j} (a^{ij}(x) \partial_{x_i} v) u \, dx = \int (Lv) u \, dx$$

F99#3 cont:

Suppose all eigenvalues of L are ≥ 0 . Then $\lambda \geq 0$. Thus by the claim,

$$\min_{w: \frac{\partial w}{\partial \nu} = 0 \text{ on } \partial U} \frac{\langle Lw, w \rangle}{\langle w, w \rangle} \geq 0.$$

Fix a u with $\frac{\partial u}{\partial \nu} = 0$ and consider $u+a$ with a to be chosen later. We have

$$\begin{aligned} \langle L(u+a), u+a \rangle &= \langle Lu + c(x)a, u+a \rangle \\ &= \langle Lu, u \rangle + a \int Lu \, dx + a \int c(x)u \, dx + a^2 \int c(x) \, dx. \quad (*) \end{aligned}$$

Since $\int c(x) \, dx \leq 0$ and $c(x) \neq 0$ for some $x \in U$, $c(x)$ smooth, choose a sufficiently large, we can make $(*) < 0$. Therefore

$$0 \leq \min_{w: \frac{\partial w}{\partial \nu} = 0 \text{ on } \partial U} \frac{\langle Lw, w \rangle}{\langle w, w \rangle} \leq \frac{\langle L(u+a), u+a \rangle}{\langle u+a, u+a \rangle} < 0,$$

a contradiction. Therefore \exists a negative eigenvalue. $\#$

Eqn #4:

We solve

$$\begin{cases} u_t + a(x)u_x = 0 \\ u(x, 0) = f(x) \end{cases}$$

by method of characteristics. We have

$$F(p, q, z, x, t) = q + a(x)p.$$

$$\dot{t}(s) = 1 \quad t(0) = 0$$

$$\dot{x}(s) = a(x) \quad x(0) = x_0$$

$$\dot{z}(s) = 0 \quad z(0) = f(x_0).$$

Let $a(x) := x^2 + 1$. Then

$$\begin{aligned} \dot{x} &= x^2 + 1 & \tan^{-1}(x) &= s + \tan^{-1}(x_0) \\ x(0) &= x_0 & \rightarrow \end{aligned}$$

Since $t = s$, the ~~path~~ characteristic curves are given by

$$\tan^{-1}x = t + \tan^{-1}x_0$$

Since $\tan^{-1}x_0 \in (-\frac{\pi}{2}, \frac{\pi}{2})$ for any choice of x_0 , the characteristic curves do not propagate beyond time $t = \pi$. Thus the values of $u(x, t)$ for $t > \pi$ do not depend on the values $u(x, 0)$.
Then the solution of the Cauchy problem is not unique.

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Ex 99 #5 :

We have

$$\hat{u}(m, n, t) = \int_{[0, 2\pi]^2} u(x, y, t) e^{-i(xm + ny)} dx dy.$$

Then

$$\begin{aligned} \hat{u}_t(m, n, t) &= (-\varepsilon)(-m^2 - n^2) \hat{u}(m, n, t) \\ &\quad - (m^4 + 2m^2 n^2 + n^4) \hat{u}(m, n, t). \end{aligned}$$

$$\hat{u}_t(m, n, t) = \varepsilon(m^2 + n^2) \hat{u}(m, n, t) - (m^2 + n^2)^2 \hat{u}(m, n, t).$$

$$\hat{u}_t(m, n, t) = (m^2 + n^2) \left[\frac{\varepsilon - (m^2 + n^2)}{\varepsilon - (m^2 + n^2)} - 1 \right] \hat{u}(m, n, t)$$

$$\hat{u}(m, n, t) = e^{[(m^2 + n^2)/(\varepsilon - (m^2 + n^2))]t} \hat{u}(m, n, 0)$$

Therefore

$$u(x, y, t) = \sum_{m, n \in \mathbb{Z}} e^{(m^2 + n^2)(\varepsilon - m^2 - n^2)t} \hat{u}(m, n, 0) e^{i(mx + ny)}.$$

i) A solution is e^{t+ix} . Indeed,

$$\partial_t(e^{t+ix}) = e^{t+ix}$$

$$\Delta u = \partial_{xx}(e^{t+ix}) = e^{t+ix}(-1)$$

$$\Delta^2 u = e^{t+ix}.$$

Therefore

$$-2\Delta u - \Delta^2 u = 2e^{t+ix} - e^{t+ix} = e^{t+ix}.$$

F99 #5 cont;

ii) We are

$$\epsilon - m^2 - n^2 < 0 \quad \forall m, n.$$

$$\epsilon < m^2 + n^2.$$

Therefore take $\epsilon_0 = 1$. #

F99 #6:

a) We have

$$\partial_t p = -sp'$$

$$\partial_x u = u'$$

↓ deriv. w.r.t y

$$\partial_t u = -su'$$

$$\partial_x(pu) = p_x u + pu_x = p'u + pu'$$

Thus

$$-sp' + u' = 0$$

$$\rightarrow u' = sp'$$

$$p' = \frac{1}{s}u'$$

$$-su' + p'u + pu' = u''$$

~~pe~~

$$u = sp + C_1$$

$$\rightarrow -su' + \frac{1}{s}u'u + (\frac{u}{s} - C_1)u' = u''$$

$$-su' + \frac{2}{s}u'u - C_1 u' = u''$$

Therefore

$$u'' = -(s + C_1)u' + \frac{1}{s}(u^2)'$$

Thus

$$u' = -(s + C_1)u + \frac{1}{s}u^2 + C_2 \quad (*)$$

b) Now (*) is of the form $u' + Au^2 + Bu = C$ with $A = -1/s$, $B = s + C_1$, $C = C_2$. One way to solve (*) is to observe (*) is separable, so,

$$\int \frac{1}{\frac{1}{s}u^2 - (s + C_1)u + C_2} du = \int dc$$

$$\int \frac{s}{u^2 - s(s + C_1)u + sC_2} du = z + \tilde{C}_1$$

but the easiest way is to observe that if we guess the solution is of the form

$$u(y) = u_0 + u_1 \tanh(\alpha y + y_0),$$

F99#6 cont:

$$u' = \alpha u, \operatorname{sech}^2(\alpha y + y_0).$$

Thus

$$u' + Au^2 + Bu = C$$

$$\alpha u, \operatorname{sech}^2(\alpha y + y_0) + \downarrow u_0^2 + 2Au_0 u, \tanh(\alpha y + y_0) + Au_1^2 \tanh(\alpha y + y_0) + Bu_0 + Bu, \tanh(\alpha y + y_0) = C.$$

Our ansatz is a solution if we choose u_0, u_1, α s.t.

$$\begin{aligned} \alpha u_1 &= Au_1^2 & u_1 &= \alpha/A \\ 2Au_0 u_1 + Bu_1 &= 0 & \rightarrow u_0 &= -B/2A \\ Au_1^2 + Au_0^2 + Bu_0 &= C & \alpha &= (AC + B^2/2)^{1/2}. \end{aligned}$$

Thus with the ~~choice~~ free choices of c_1, c_2 , (*) has solutions of the form $u(y) = u_0 + u_1 \tanh(\alpha y + y_0)$.

F99 #7.

i. We want to find ϕ s.t. $\langle Lu, v \rangle_\phi = \langle u, Lv \rangle_\phi$ for u, v satisfying $u(0) = u(1) = 0$ and $v(0) = v(1) = 0$.

We have

$$\begin{aligned} \langle Lu, v \rangle_\phi &= \int_0^1 (u'' + 2u')v\phi \, dx = \int_0^1 u''v\phi + 2u'v\phi \, dx \\ &= -\int_0^1 u'(v\phi)' \, dx + 2\int_0^1 u(v\phi)' \, dx \\ &= \int_0^1 u(v\phi)'' \, dx - 2\int_0^1 u(v\phi)' \, dx \\ &= \int_0^1 u(v''\phi + 2v'\phi' + v\phi'') - 2u(v'\phi + v\phi') \, dx \\ &= \int_0^1 u(v''\phi + 2v'\phi' + v\phi'' - 2v'\phi - 2v\phi') \, dx \\ &= \int_0^1 u(v''\phi + 2v'(\phi' - \phi) + v(\phi'' - 2\phi')) \, dx \\ &= \int_0^1 u\phi(v'' + 2v'(\frac{\phi'}{\phi} - 1) + v(\frac{\phi'' - 2\phi'}{\phi})) \, dx \end{aligned}$$

Thus we would have $\langle Lu, v \rangle_\phi = \langle u, Lv \rangle_\phi$ if

$$\begin{aligned} \frac{\phi'}{\phi} - 1 &= 1 & \phi' &= 2\phi \\ \frac{\phi'' - 2\phi'}{\phi} &= 0 & \phi'' &= 2\phi' \end{aligned}$$

Taking $\phi = e^{2x}$ gives such a ϕ .

ii. To show $L + aI$ is invertible, we will show that $\ker(L + aI)$ is trivial. Let $u \in \ker(L + aI)$. Then

$$\begin{aligned} 0 &= \langle (L + aI)u, u \rangle = \int_0^1 (Lu + au)u\phi \, dx = \int_0^1 (u'' + 2u' + au)u\phi \, dx \\ &= -\int_0^1 u'(u\phi)' \, dx + 2\int_0^1 u'u\phi + au^2\phi \, dx \\ &= -\int_0^1 u'(ue^{2x})' \, dx + 2\int_0^1 u(u e^{2x})' + au^2e^{2x} \, dx \end{aligned}$$

F99#7

cont.

$$\begin{aligned}
 0 &= \langle (L + aI)u, u \rangle = \int_0^1 (Lu + au)u \phi dx = \int_0^1 (u'' + 2u' + au)u \phi dx \\
 &= \int_0^1 (u'' + 2u' + au)ue^{2x} dx = \int_0^1 u(ue^{2x})'' - 2u(ue^{2x})' + au^2e^{2x} dx \\
 &= \int_0^1 -u'(ue^{2x})' + 2u'ue^{2x} + au^2e^{2x} dx \\
 &= \int_0^1 -u'(u'e^{2x} + 2e^{2x}u) + 2u'ue^{2x} + au^2e^{2x} dx \\
 &= \int_0^1 -(u')^2e^{2x} + au^2e^{2x} dx \\
 &= \int_0^1 e^{2x} [- (u')^2 + au^2] dx
 \end{aligned}$$

Since $a < 0$, if $u \neq 0$, \int is < 0 which would lead to a contradiction. Therefore we must have $u \equiv 0$. Thus if $a < 0$, $L + aI$ is invertible.

iii: Take $a = -1$, $u = e^{-x}$.

F99 #8:

We look for a radial symmetric u .

a) We want to solve
$$\begin{cases} -\Delta u = 1 & \text{on } B_{R(t)}(0) \\ u = 0 & \text{on } |x| = R(t) \end{cases}$$

Since we want to find radial u , we solve

$$-(u_{rr} + \frac{1}{r}u_r) = 1$$

$$ru_{rr} + u_r = -r$$

$$(ru_r)' = -r$$

$$ru_r = -\frac{r^2}{2} + C_1$$

$$u_r = -\frac{r}{2} + \frac{C_1}{r}$$

$$u = -\frac{1}{4}r^2 + C_1 \log r + C_2$$

Since we want u to be defined at the origin, $C_1 = 0$. Since also $u = 0$ when $r = R(t)$, we have

$$u = -\frac{1}{4}r^2 + \frac{1}{4}R(t)^2 = -\frac{1}{4}|x|^2 + \frac{1}{4}R(t)^2.$$

b) Using $\frac{dR}{dt} = -u_r|_{r=R}$, we have

$$\frac{dR}{dt} = -u_r|_{r=R} = \frac{1}{2}R.$$

Therefore as $R(0) = R_0$,

$$R(t) = e^{\frac{1}{2}t} R_0.$$