

FOO #1: Let u_1, u_2 be 2 classical solutions of the given equation. Let $w := u_1 - u_2$. Then

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$$\Delta w + a(x)w = 0, \quad x \in \Omega$$

$$\omega|_S = 0 \quad x \in S$$

Therefore

$$\int_{\partial \Omega} |\nabla w|^2 dx = - \int_{\Omega} w \Delta w dx = \int_{\Omega} a(x) w^2 dx$$

$$\leq \|a\|_{L^\infty} \int_{D^*} \omega^2 dx \leq \|a\|_{L^\infty} C(D) \int_D |\nabla \omega|^2 dx$$

↪ Poincaré ineq.

$$\leq \frac{1}{2} \int_0^1 |w|^2 dx$$

if we ~~choose~~ assume $\|a\|_{L^2} C(D) \leq \frac{1}{2}$. Then if $|a(x)|$ was small enough, then $\int_D |aw|^2 dx = 0 \rightarrow \nabla w = 0 \rightarrow w = 0$.

6. Let $H = H_0' \setminus \{0\} = \{u \in H' \setminus \{0\} : u = 0 \text{ on } S\}$ and

$$B[u, v] = \int_D \nabla u \cdot \nabla v - a(x)uv \, dx, \quad u, v \in H_0^1(D).$$

Note that $H_0^1(D)$ is a Hilbert space and $B(u, v) : H_0^1(D) \times H_0^1(D) \rightarrow \mathbb{R}$.
We will use Lax-Milgram to prove the result.

① We have

$$|B[u, v]| \leq \int_{\Omega} |\nabla u| |\nabla v| \, dx + \int_{\Omega} |a(x)| |u| |v| \, dx$$

$$\leq \| \nabla u \|_{L^2} \| \nabla v \|_{L^2} + \| a \|_{L^\infty} \| u \|_{L^2} \| v \|_{L^2}$$

$$\leq \alpha \|u\|_{H^1} \|v\|_{H^1}$$

for some constant $\alpha > 0$ which only depends on $\|g\|_{\infty}$

⑥ We want to show there is a constant $\beta > 0$ s.t. $\beta \|u\|_{H^1}^2 \leq B(u, u)$
 $\forall u \in H_0^1$. We have

F' is 1. var.

$$B(u, u) = \int_{\Omega} |\nabla u|^2 - a(x) u^2 dx = \|\nabla u\|_{L^2}^2 - \int_{\Omega} a(x) u^2 dx.$$

$$\geq \frac{1}{3} \|\nabla u\|_{L^2}^2 + \frac{2}{3} \|\nabla u\|_{L^2}^2 - \|a\|_{L^\infty} \int_{\Omega} u^2 dx.$$

$$= \frac{1}{3} \|\nabla u\|_{L^2}^2 + \frac{2}{3} \|\nabla u\|_{L^2}^2 - \|a\|_{L^\infty} \|u\|_{L^2}^2.$$

$$\geq \frac{1}{3} \|\nabla u\|_{L^2}^2 + \frac{2}{3} C \|\nabla u\|_{L^2}^2 - \|a\|_{L^\infty} \|u\|_{L^2}^2$$

$$= \frac{1}{3} \|\nabla u\|_{L^2}^2 + \left(\frac{2}{3} C - \|a\|_{L^\infty}\right) \|u\|_{L^2}^2.$$

when from Poincaré,
 $\exists C$ depending only on Ω s.t.
 $C \|u\|_{L^2} \leq \|\nabla u\|_{L^2}$.

~~Choose~~ If $|a(x)|$ is sufficiently small, then \exists constant $\beta > 0$ independent of u s.t. $B(u, u) \geq \beta \|u\|_{H^1}^2$.

Therefore since $v \mapsto \int_{\Omega} f v dx$ is a bounded linear functional on $H_0^1(\Omega)$, by Lax-Milgram, if $|a(x)|$ is sufficiently small, \exists a unique $u \in H_0^1(\Omega)$ s.t.

$$\int_{\Omega} \nabla u \cdot \nabla v - a(x) u v dx = - \int_{\Omega} f v dx$$

$\forall v \in H_0^1(\Omega)$. Since

$$\begin{aligned} \int_{\Omega} \nabla u \cdot \nabla v dx &= - \int_{\Omega} v \Delta u dx + \int_{\partial \Omega} \frac{\partial u}{\partial \nu} v d\sigma \\ &= - \int_{\Omega} v \Delta u dx, \end{aligned}$$

it follows that

$$\int_{\Omega} (\Delta u + a(x) u) v dx = \int_{\Omega} f v dx, \quad \forall v \in H_0^1(\Omega).$$

This shows the existence of a solution in $H^1(\Omega)$ assuming $f \in L^2$.

→ assume $\lim_{x \rightarrow \infty} u_1(x, t) = \lim_{x \rightarrow \infty} u_2(x, t)$
for each t .

FOO #2: Let u_1, u_2 be 2 classical bounded solutions of the given equation. Let $w := u_1 - u_2$. Then

$$w_t - \Delta w + w(u_1 + u_2) = 0, \quad x \in \mathbb{R}^N, \quad 0 < t < T$$

$$w(x, 0) = 0.$$

Let

$$E(t) := \frac{1}{2} \int_{\mathbb{R}^N} w(x, t)^2 dx.$$

Then

$$E'(t) = \int_{\mathbb{R}^N} w w_t dx = \int_{\mathbb{R}^N} w (\Delta w - w(u_1 + u_2)) dx$$

$$= \int_{\mathbb{R}^N} w \Delta w dx - \int_{\mathbb{R}^N} w^2 (u_1 + u_2) dx.$$

Since $w(0) = 0$
~~by the divergence theorem~~

$$= - \int_{\mathbb{R}^N} |\nabla w|^2 dx - \int_{\mathbb{R}^N} w^2 (u_1 + u_2) dx$$

$$\leq - \|u_1 + u_2\|_{L^\infty} \int_{\mathbb{R}^N} w^2 dx$$

$$\leq - 2 \|u_1 + u_2\|_{L^\infty} E(t).$$

Thus by Gronwall's Inequality,

$$E(t) \leq E(0) e^{-\int_0^t 2 \|u_1 + u_2\|_{L^\infty} ds} = 0$$

Therefore $w(x, t) = 0 \rightarrow u_1 = u_2$.

FOO #4

Let v_1, v_2, v_3 be smooth compactly supported functions. Then the minimizer u_1, u_2, u_3 satisfies

$$0 = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (F(u + \varepsilon v) - F(u)).$$

$$= \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \left(\int_0^3 \sum_{j,k=1}^3 (u_{j,x_k} + \varepsilon v_{j,x_k})^2 + \alpha \left(\sum_{j=1}^3 (u_j + \varepsilon v_j)^2 - 1 \right)^2 - \int_0^3 \sum_{j,k=1}^3 (u_{j,x_k})^2 + \alpha \left(\sum_{j=1}^3 u_j^2 - 1 \right)^2 \right)$$

Thus

$$\begin{aligned} 0 &= \int_0^3 \sum_{j,k=1}^3 2 u_{j,x_k} v_{j,x_k} + \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \left[\alpha \int_0^3 \left(\left(\sum_{j=1}^3 u_j^2 - 1 \right) + 2 \sum_{j=1}^3 u_j v_j \varepsilon + \sum_{j=1}^3 \varepsilon^2 v_j^2 \right)^2 - \alpha \int_0^3 \left(\sum_{j=1}^3 u_j^2 - 1 \right)^2 \right] \\ &= \int_0^3 \sum_{j,k=1}^3 2 u_{j,x_k} v_{j,x_k} + 4 \left(\sum_{j=1}^3 u_j^2 - 1 \right) \left(\sum_{j=1}^3 u_j v_j \right) dx. \\ &= \int_0^3 - \sum_{j,k=1}^3 2 u_{j,x_k} v_{j,x_k} + 4 \left(\sum_{j=1}^3 u_j^2 - 1 \right) \left(\sum_{j=1}^3 u_j v_j \right) dx. \\ &= \int_0^3 \sum_{j=1}^3 \left[-2 \sum_{k=1}^3 (u_{j,x_k})_{x_k} + 4 \left(\sum_{\ell=1}^3 u_{\ell}^2 - 1 \right) u_j \right] v_j. \end{aligned}$$

Since the v_j are arbitrary, we have

$$\begin{aligned} -\Delta u_j + 4(|\vec{u}|^2 - 1) u_j &= 0 \text{ on } D, j=1,2,3 \\ u_j &= \varphi_j \text{ on } \partial D. \end{aligned}$$

$$\begin{aligned} -\Delta \vec{u} + 4(|\vec{u}|^2 - 1) \vec{u} &= 0 \text{ on } D. \\ \vec{u} &= \vec{\varphi} \text{ on } \partial D. \end{aligned}$$

F00
E99 #6:

We will use Laplace Transform. Since $u(x, 0) = 0$,

$$\begin{aligned}\int_0^{\infty} u_t(x, t) e^{-st} dt &= - \int_0^{\infty} u(x, t) (-s e^{-st}) dt \\ &= s \int_0^{\infty} u(x, t) e^{-st} dt.\end{aligned}$$

Fix on $x \geq 0$. Then as

$$u_t - u_{xx} + au = 0$$

we have

$$\begin{aligned}s \mathcal{L}[u] - \mathcal{L}[u]_{xx} + a \mathcal{L}[u] &= 0 \\ \rightarrow \mathcal{L}[u]_{xx} - (s+a) \mathcal{L}[u] &= 0.\end{aligned}$$

~~then as we want~~
that

We will find a bounded solution to the problem, so we will want $\mathcal{L}[u]$ to be bounded. We have

$$\mathcal{L}[u](x) = A e^{-\sqrt{s+a} x}$$

Since ~~that~~ $u(0, t) = g(t)$, $\mathcal{L}[u(0, t)] = \mathcal{L}[g]$.

Therefore $A = \mathcal{L}[g]$ and hence

$$\mathcal{L}[u] = \mathcal{L}[g] e^{-\sqrt{s+a} x}.$$

$$\rightarrow u(x, t) = g * \mathcal{L}^{-1}[e^{-\sqrt{s+a} x}]$$