

The latent variable proximal point algorithm for variational problems with inequality constraints

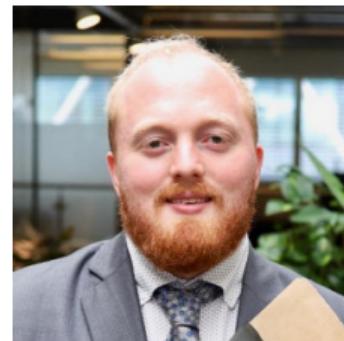
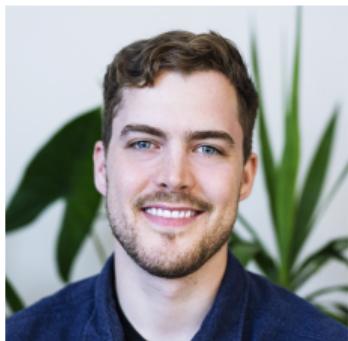
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Jørgen S. Dokken⁴

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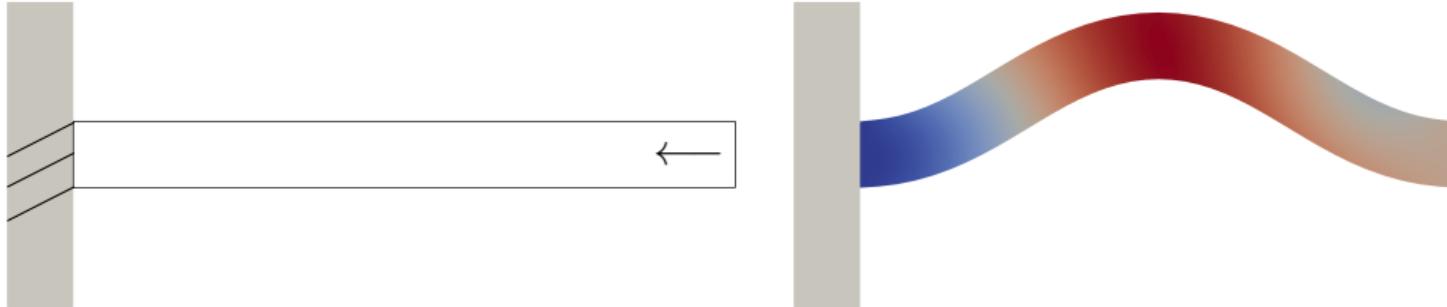
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⁵WIAS

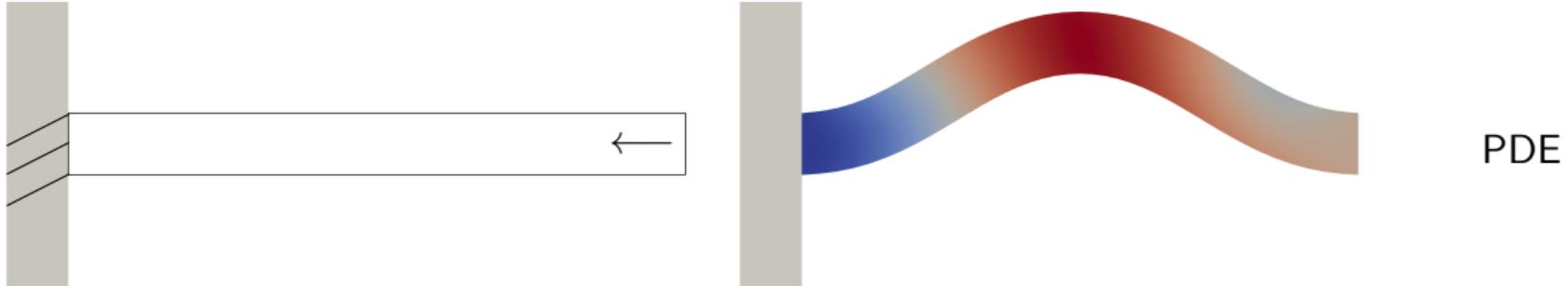
Problems with inequality constraints are formulated as *variational inequalities* (VIs).



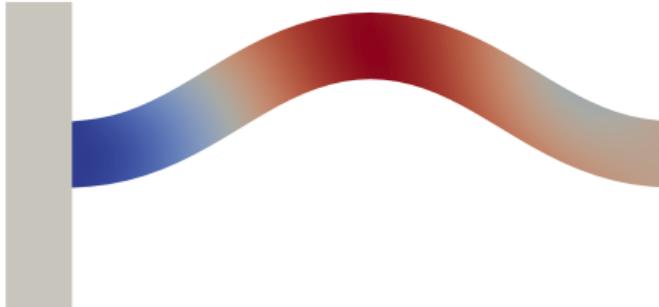
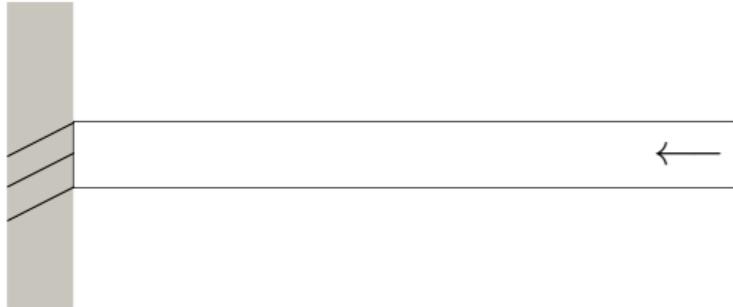
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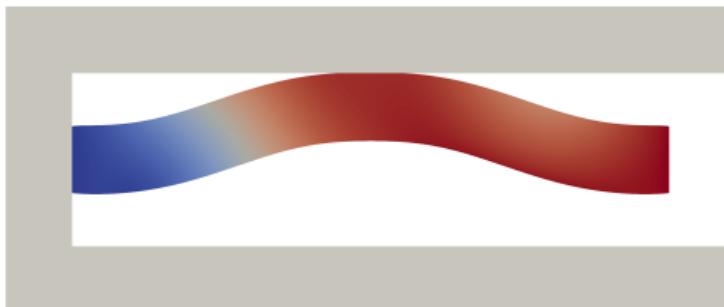
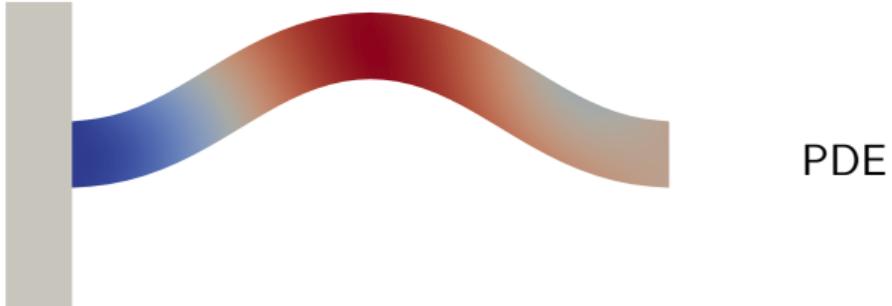
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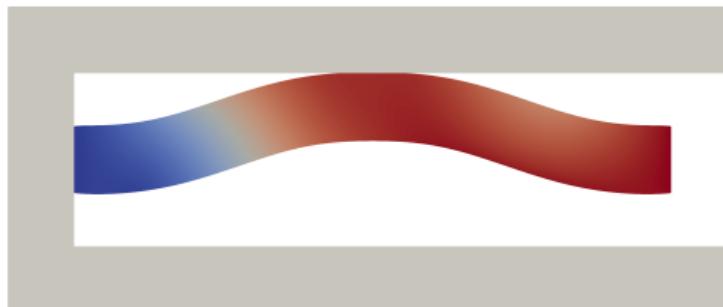
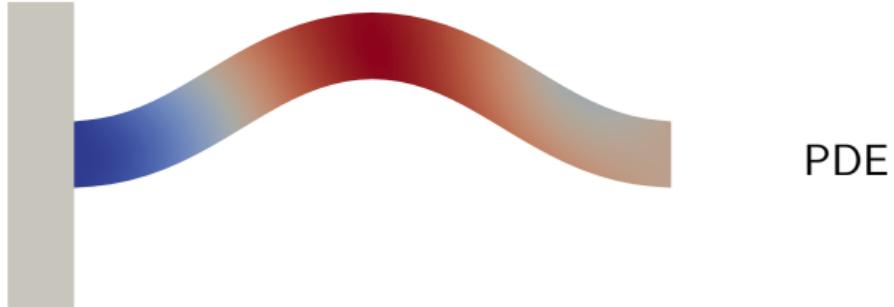
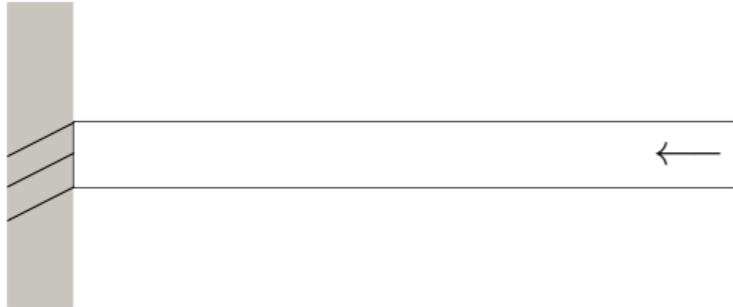
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This talk

A new framework for solving infinite-dimensional variational inequalities.

Let's see an example. Minimising the Dirichlet-type energy

$$J(u) = \frac{1}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx - \int_{\Omega} fu \, dx$$

over the constrained set

$$K = \{v \in H_0^1(\Omega) \mid v \geq \phi \text{ a.e. in } \Omega\} \subsetneq V := H_0^1(\Omega)$$

is known as the obstacle problem.

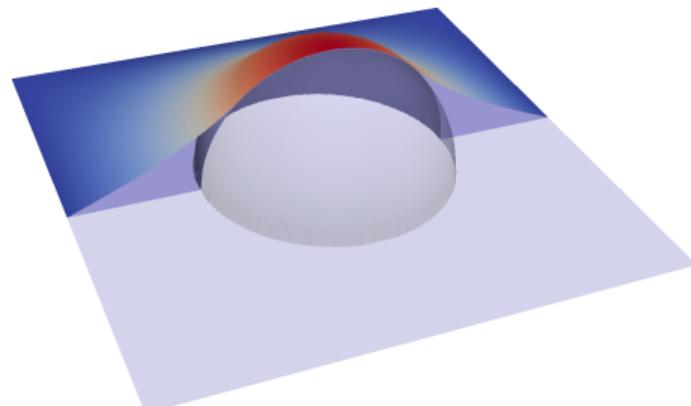
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$$f = 0, \phi = \circlearrowleft$$

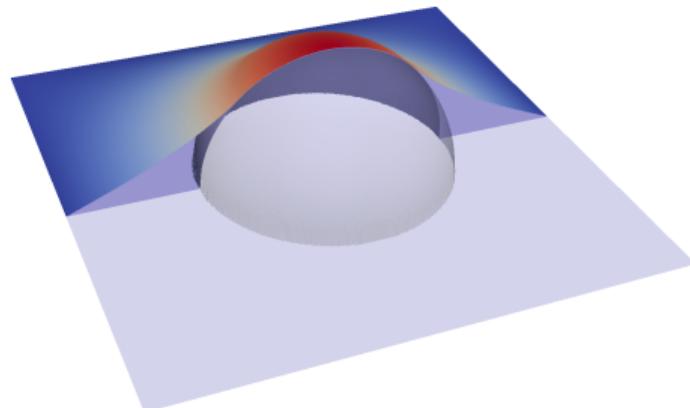
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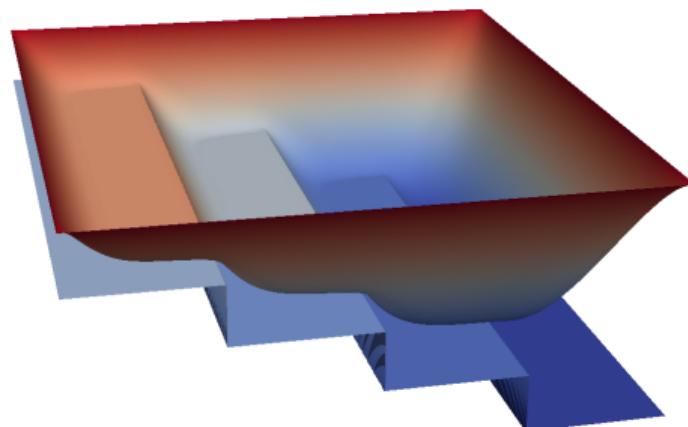
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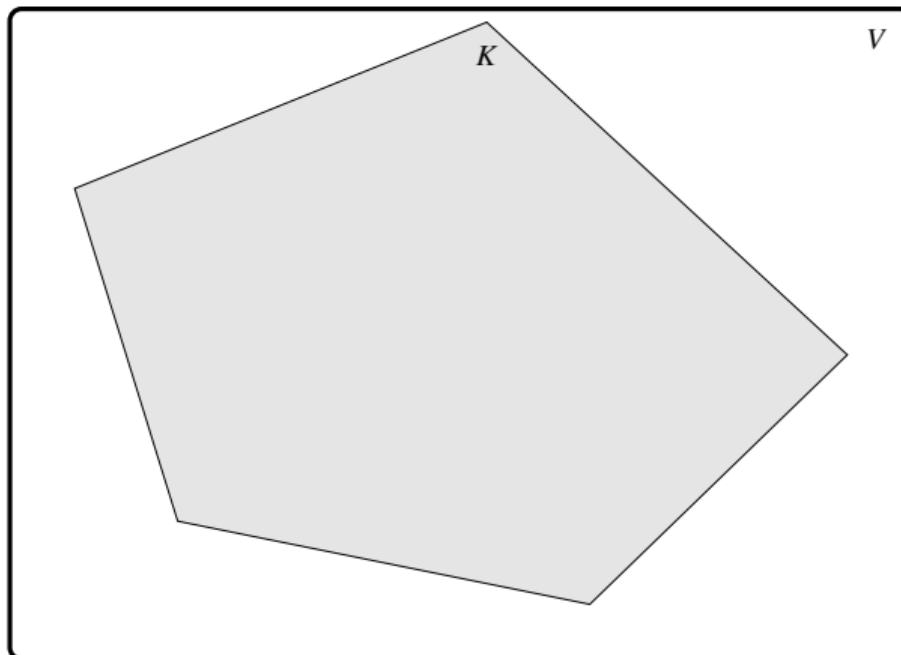
$$f = -10, \phi = \nwarrow$$

The optimality condition for this problem is the variational inequality

$$u \in K : \quad J'(u; v - u) \geq 0 \quad \forall v \in K.$$

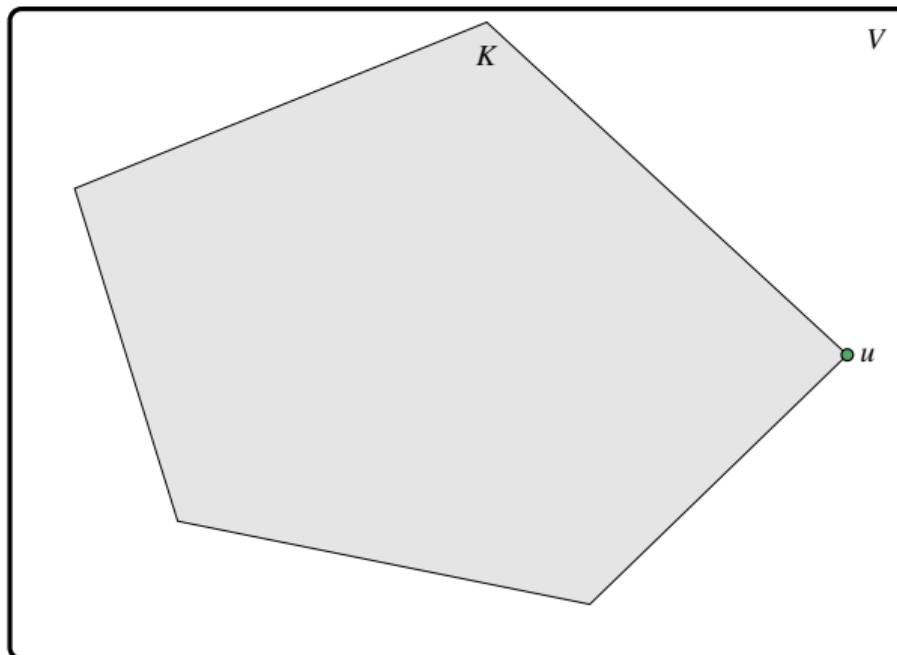
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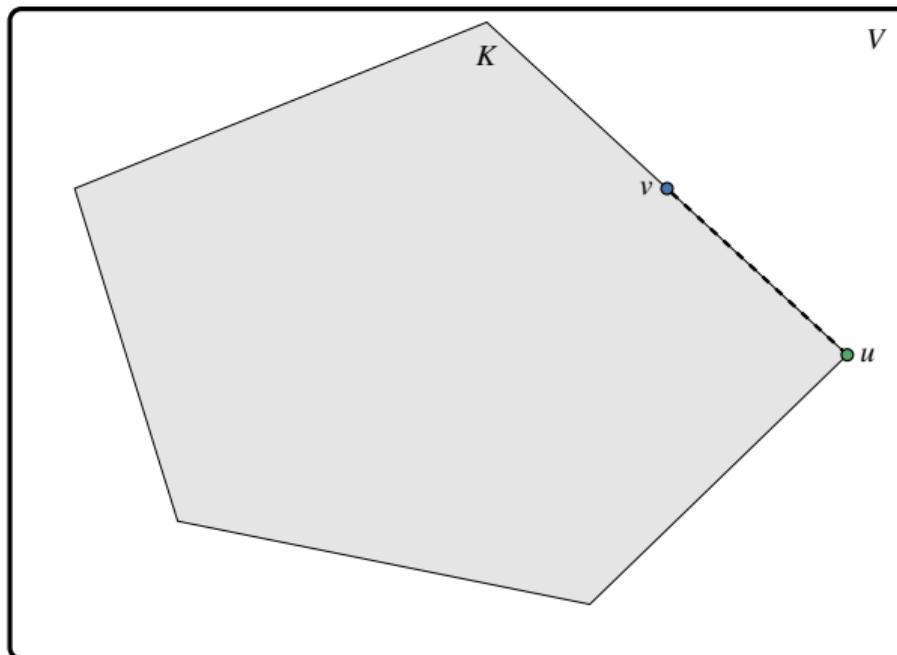
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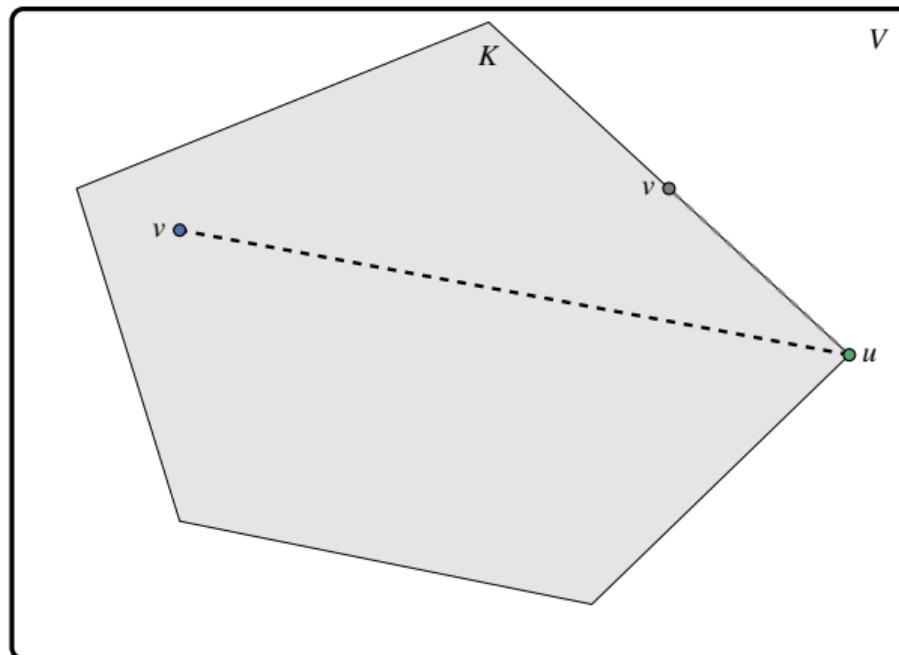
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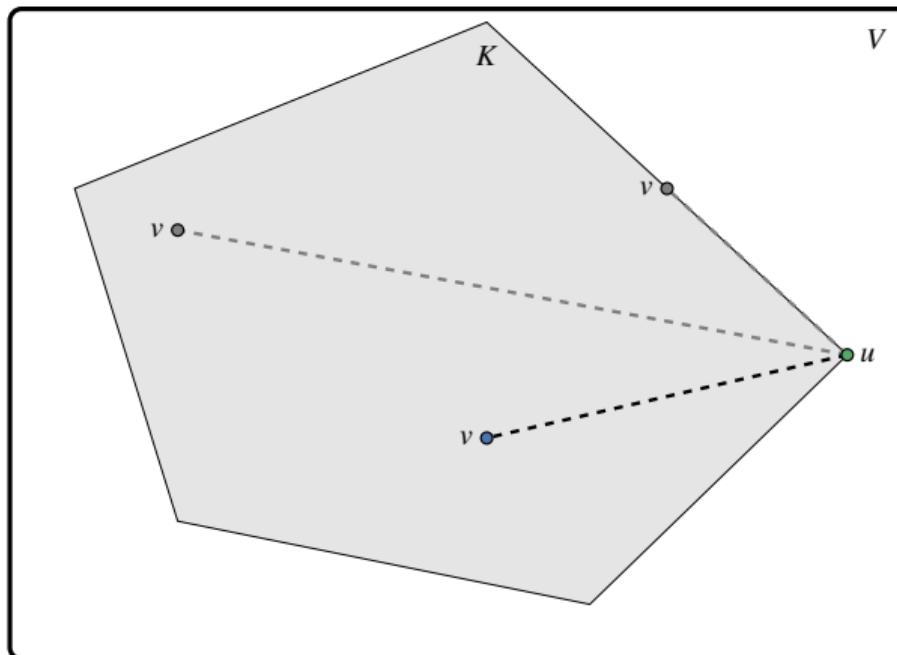
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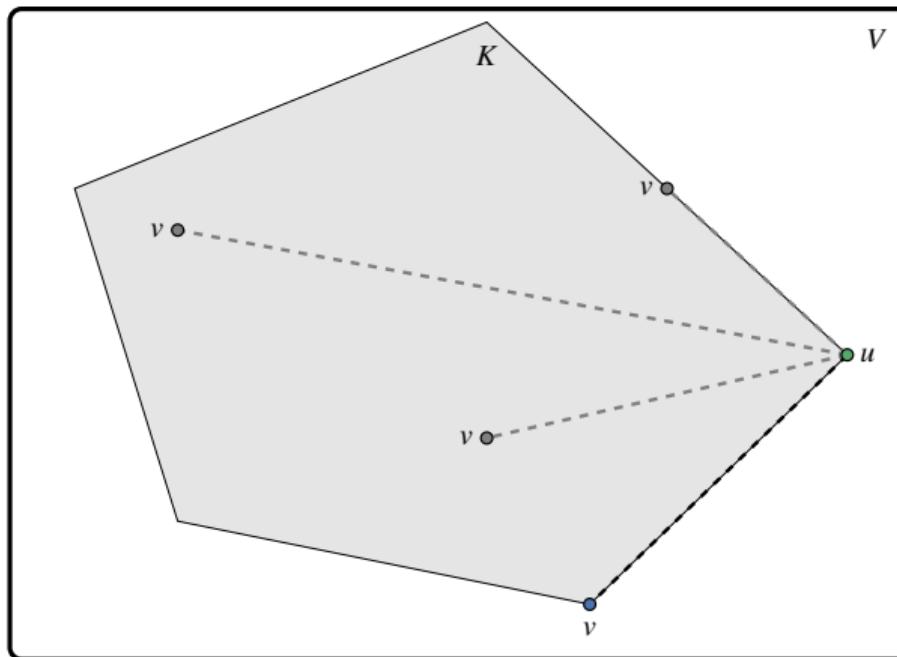
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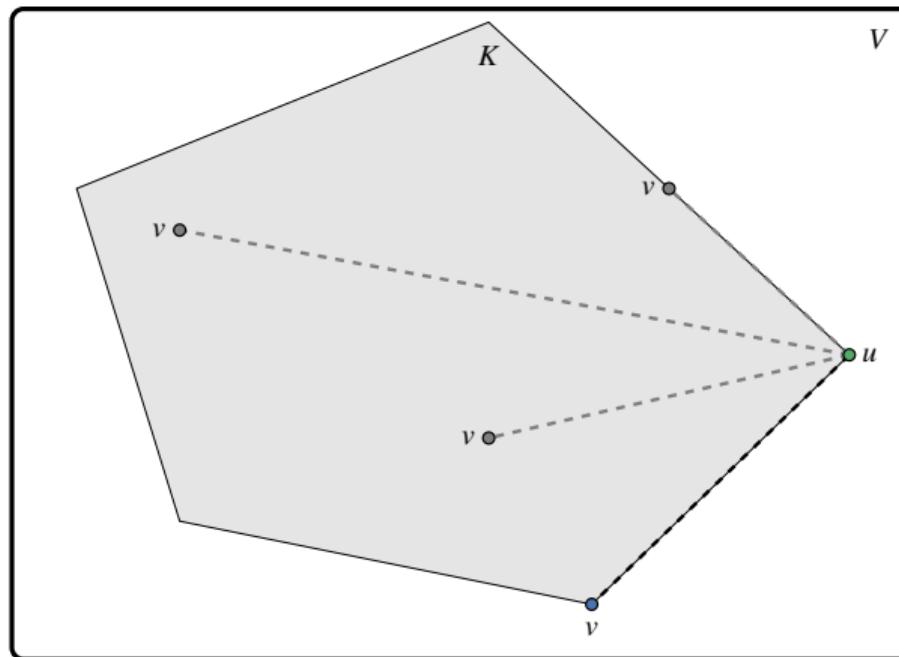
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At a solution u , the function must not decrease *in any feasible direction* $v - u$.

Section 2

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How it works

LVPP breaks down a VI into a sequence of nonlinear PDE solves.

We then use Newton's method to break down nonlinear PDE solves into linear PDE solves.

Subsection 1

Legendre functions

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Our general feasible set

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Here $C(x) \subset \mathbb{R}^m$, $\operatorname{int} C(x) \neq \emptyset$, $C(x)$ convex is the *feasible image* at x .

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$R_x : C_x \rightarrow \mathbb{R} \cup \{\infty\}$ has domain C_x and *singular gradient* ∇R_x on ∂C_x .

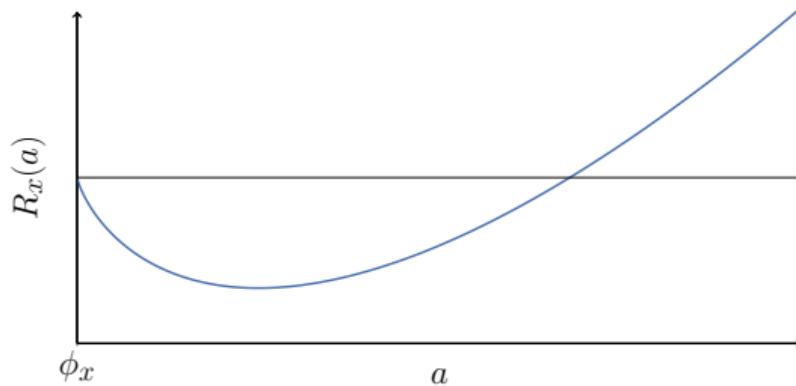
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Example

For the obstacle problem, $C_x = [\phi_x, \infty)$, and we choose a modified Shannon entropy:

$$R_x(a) = (a - \phi_x) \log(a - \phi_x) - (a - \phi_x), \quad \nabla R_x(a) = \log(a - \phi_x).$$



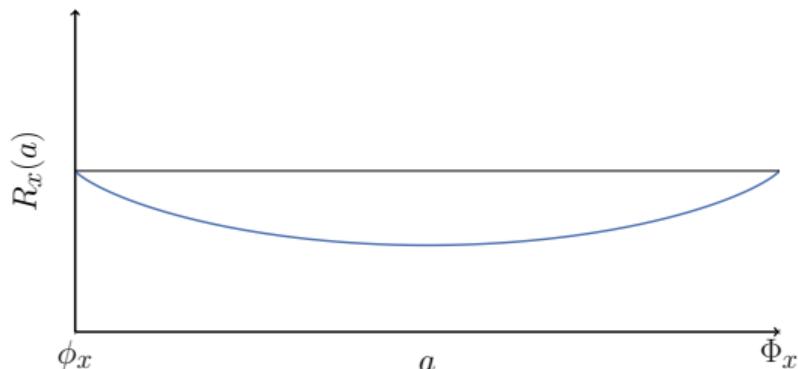
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Example

For the double obstacle problem, $C_x = [\phi_x, \Phi_x]$, and we choose the Fermi–Dirac entropy:

$$R_x(a) = (a - \phi_x) \log(a - \phi_x) + (\Phi_x - a) \log(\Phi_x - a), \quad \nabla R_x(a) = \log(a - \phi_x) + \log(\Phi_x - a).$$



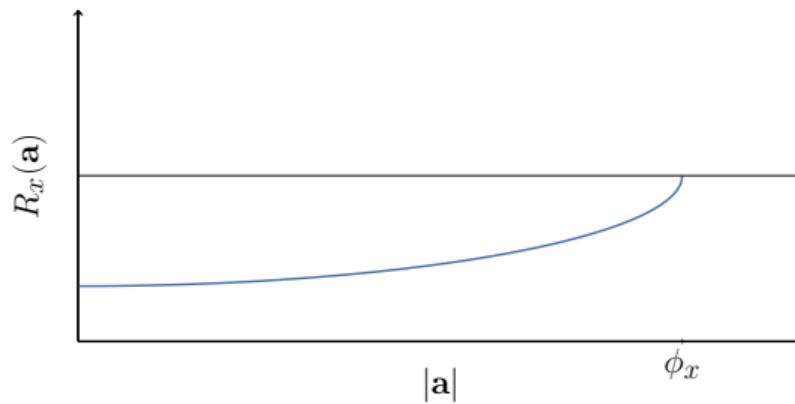
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Example

For gradient constraints, $B = \nabla$, $C_x = \mathcal{B}(0, \phi_x)$, and we choose a modified Hellinger entropy:

$$R_x(\mathbf{a}) = -\sqrt{\phi_x^2 - |\mathbf{a}|^2}, \quad \nabla R_x(\mathbf{a}) = \mathbf{a}/\sqrt{\phi_x^2 - |\mathbf{a}|^2}.$$



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Theorem (Rockafellar (1967))

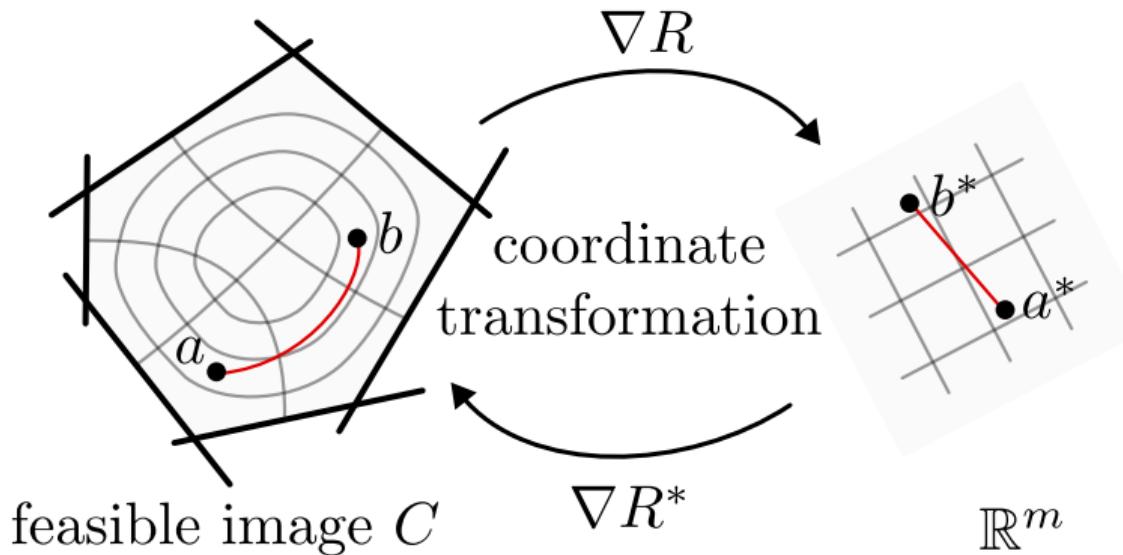
A proper convex function R is a Legendre function if and only if its convex conjugate R^ is also a Legendre function. Moreover,*

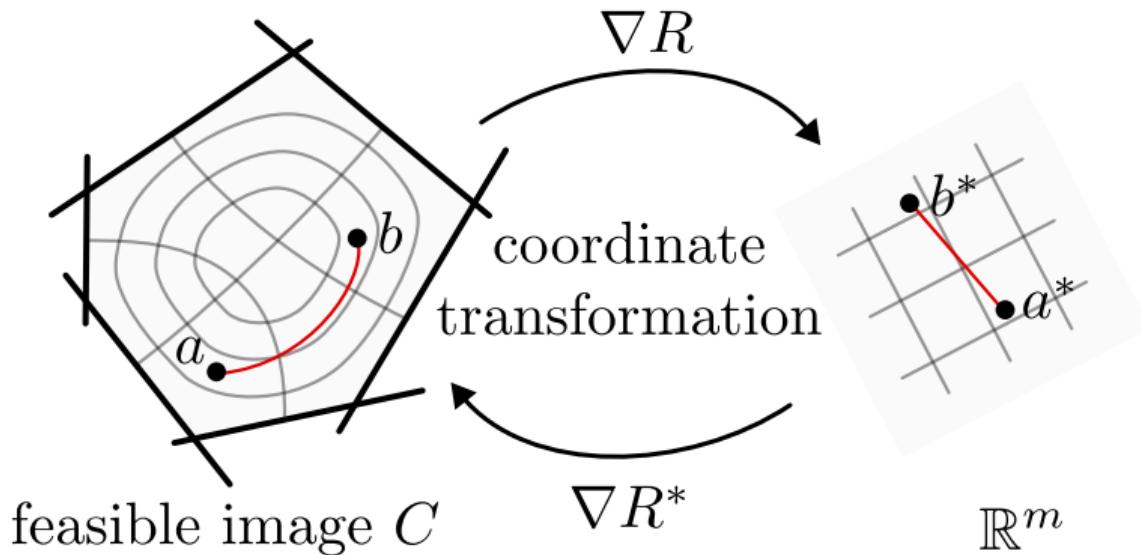
$$\nabla R: \text{int}(\text{dom } R) \rightarrow \text{int}(\text{dom } R^*)$$

is a topological isomorphism with $(\nabla R)^{-1} = \nabla R^$.*



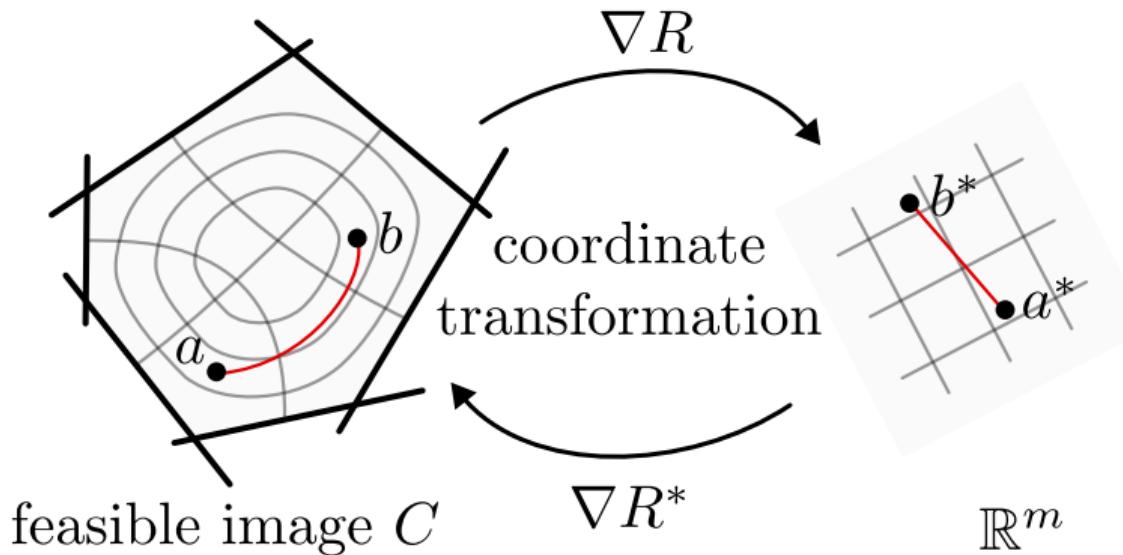
R. Tyrell Rockafellar





Good news

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We can represent any feasible function with a *latent variable* in a \mathbb{R}^m -valued Banach space!

...or more precisely any *strictly* feasible function.

Subsection 2

Proximal point

Proximal point is a fundamental algorithm in nonsmooth, convex optimisation.

To solve

$$u \in \operatorname{argmin}_{v \in K} J(v)$$



Bernard Martinet



Osman Güler

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Amazingly, for convex J , this converges in J arbitrarily quickly:

$$J(u^k) - J(u) \leq \frac{\|u^0 - u\|_V^2}{\sum_{i=1}^k \alpha_i}.$$



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Bad news

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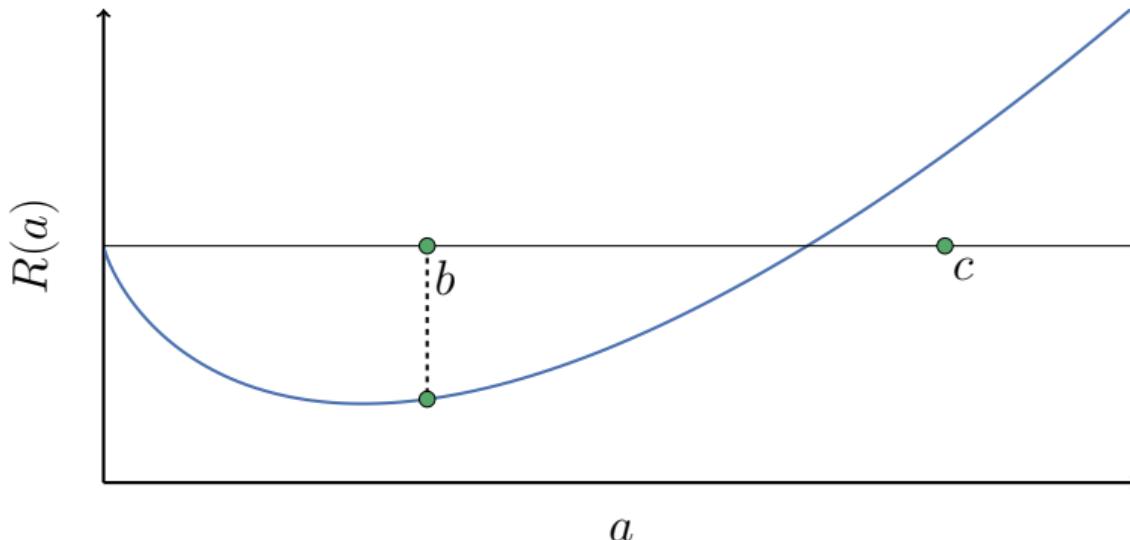
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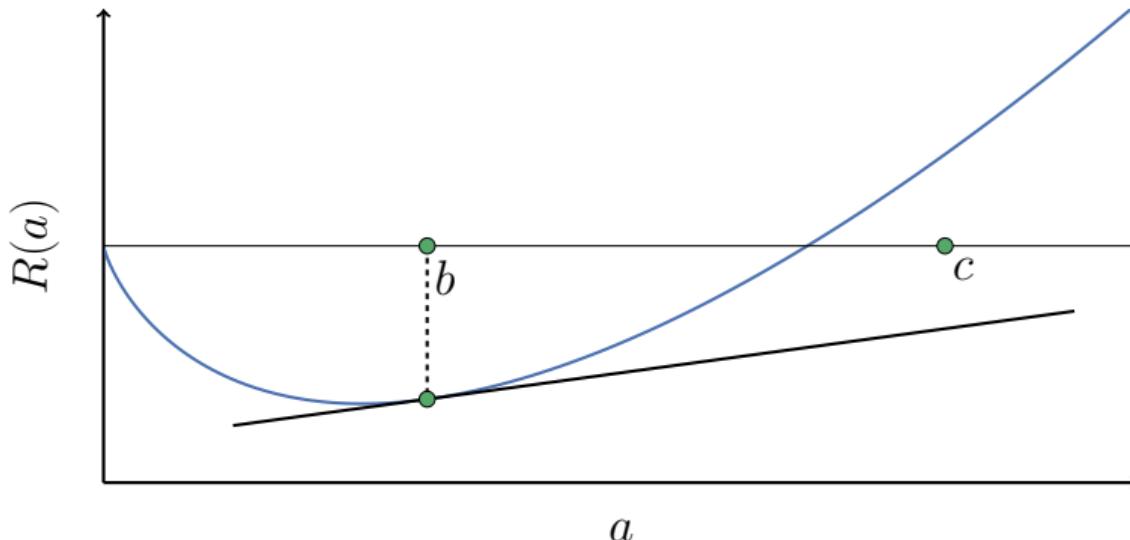
The Legendre function gives a notion of distance where the subproblems *do* simplify.

To define the *Bregman distance* $D_R(c, b)$ between b (base) and c , proceed as follows.



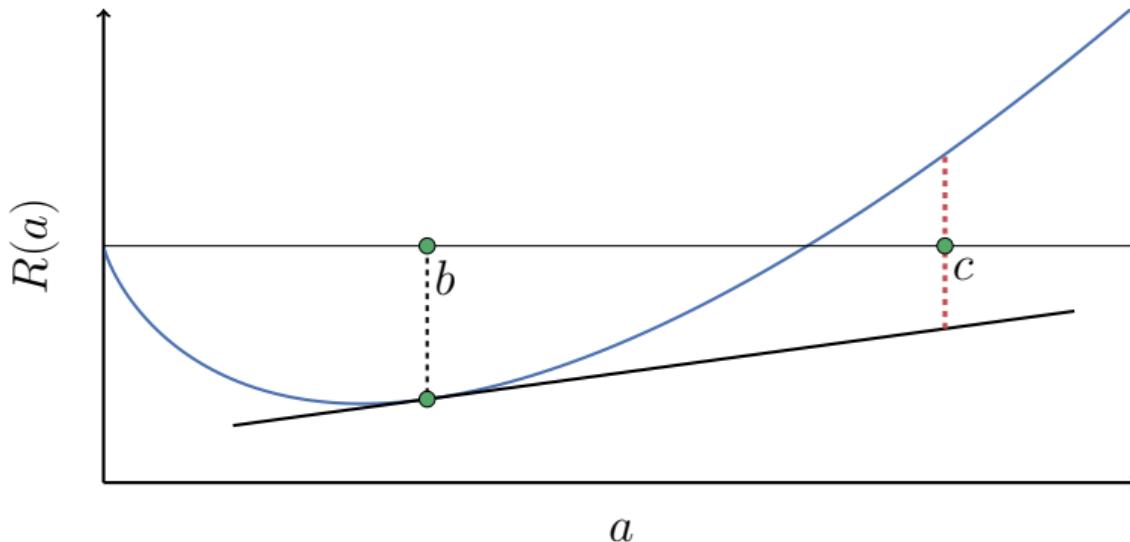
Start with the Legendre function R .

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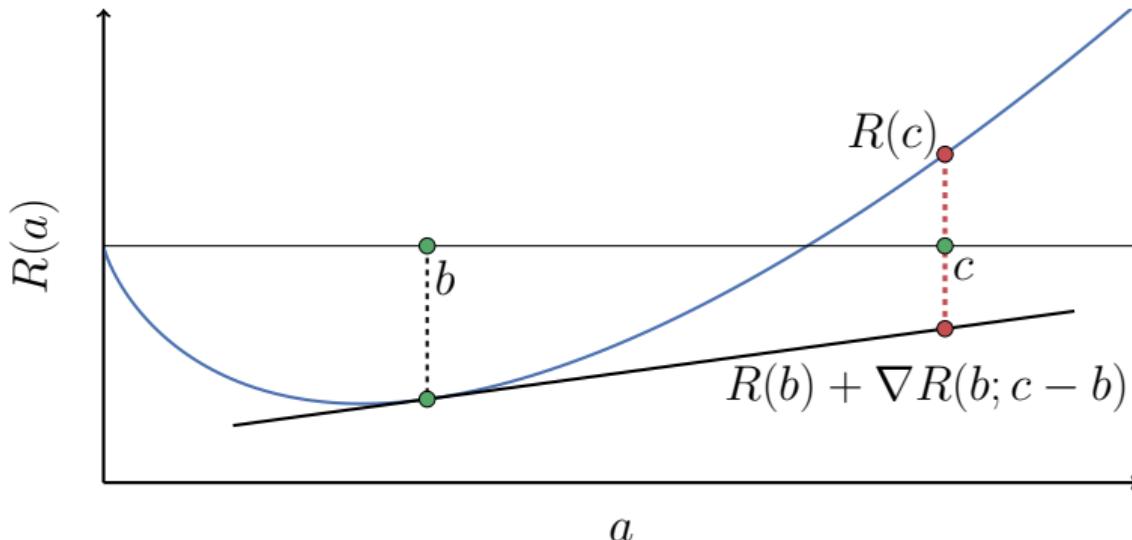
Build the tangent at b .

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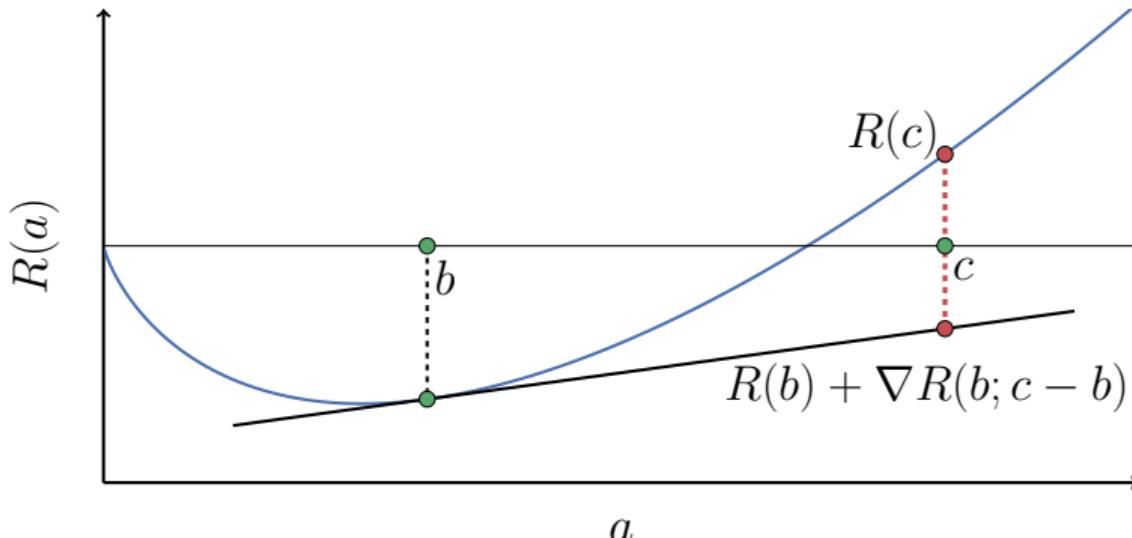
Measure distance between tangent and R at c .

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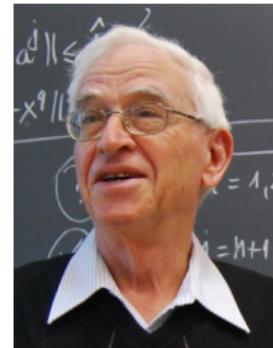


$$D_R(c, b) = R(c) - R(b) - \nabla R(b; c - b).$$

If R is the Shannon entropy, D_R is the Kullback–Leibler divergence.

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Yair Censor



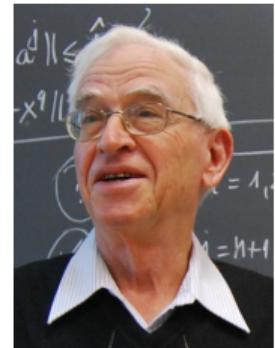
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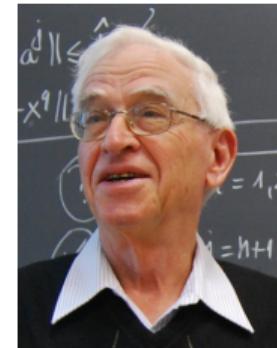
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For many problems, this forces u^k to be strictly feasible, and the subproblem optimality condition *becomes a PDE*:

$$u^k \in K : \alpha^k J'(u^k) + B^* \nabla R(Bu^k) - B^* \nabla R(Bu^{k-1}) = 0.$$



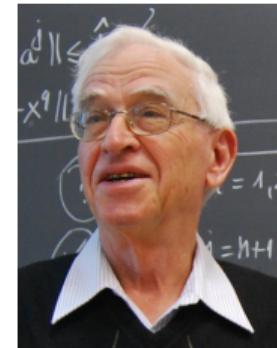
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$$u^k \in K : \alpha^k J'(u^k) + B^* \nabla R(Bu^k) - B^* \nabla R(Bu^{k-1}) = 0.$$



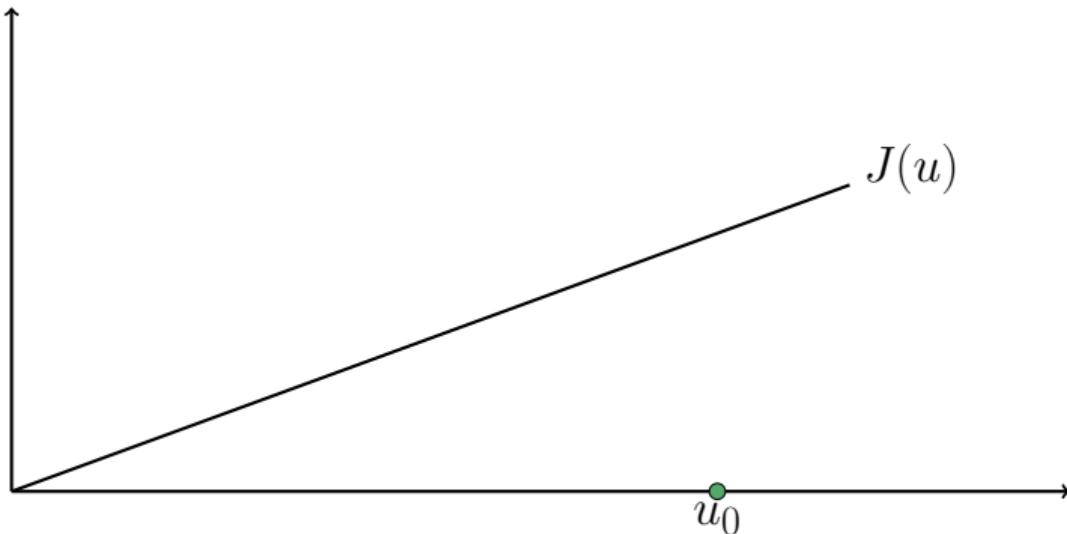
This breaks down the VI into a sequence of nonlinear PDEs!

Stavros Zenios

Consider the toy problem

$$u \in \operatorname{argmin}_{v \in [0, \infty)} J(v) = v$$

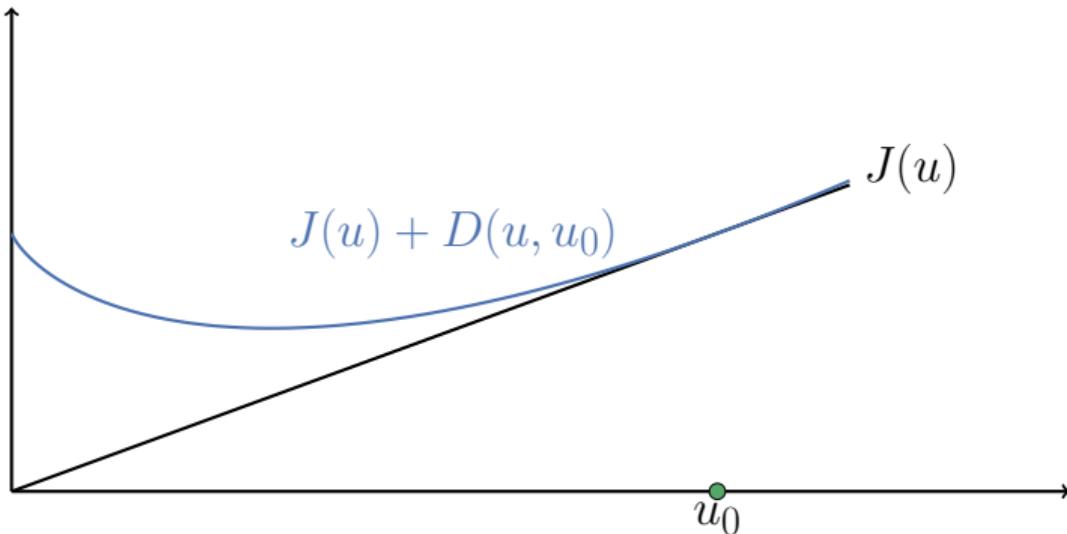
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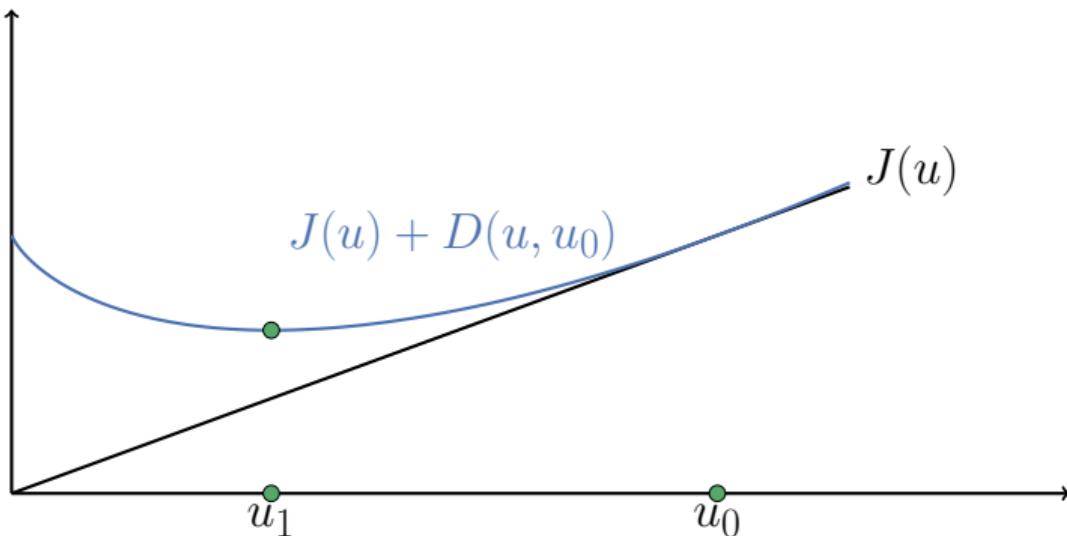
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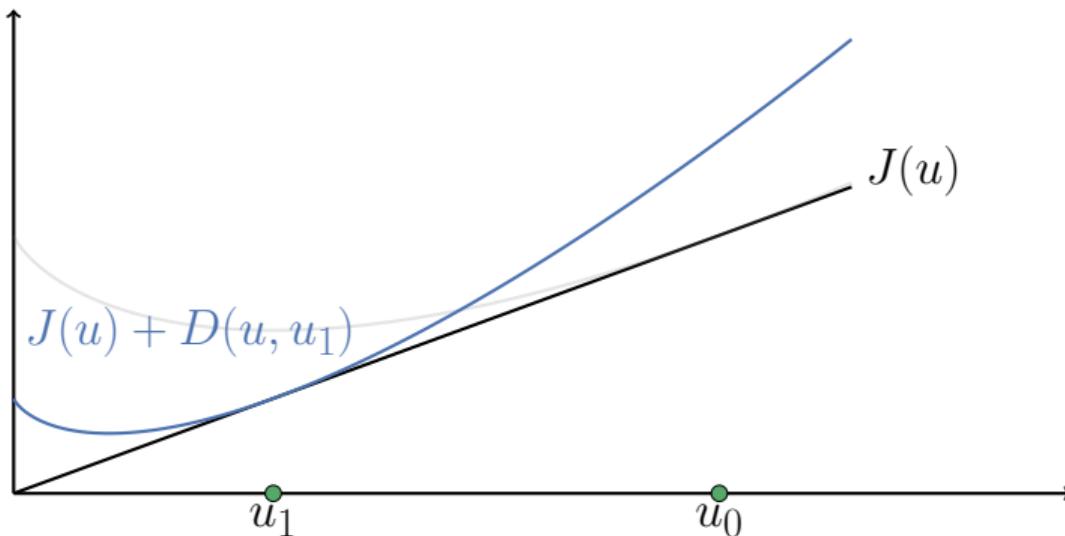
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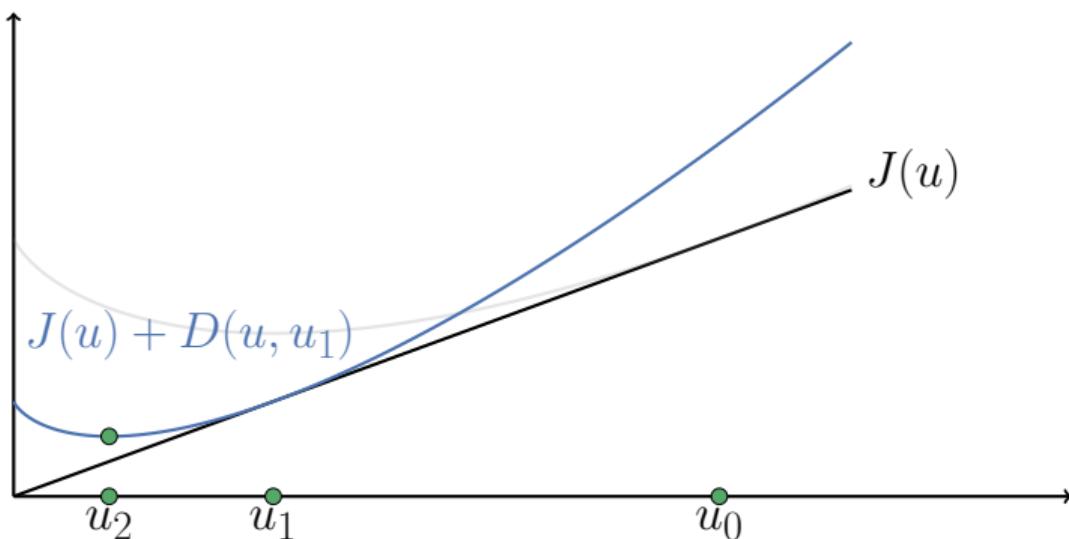
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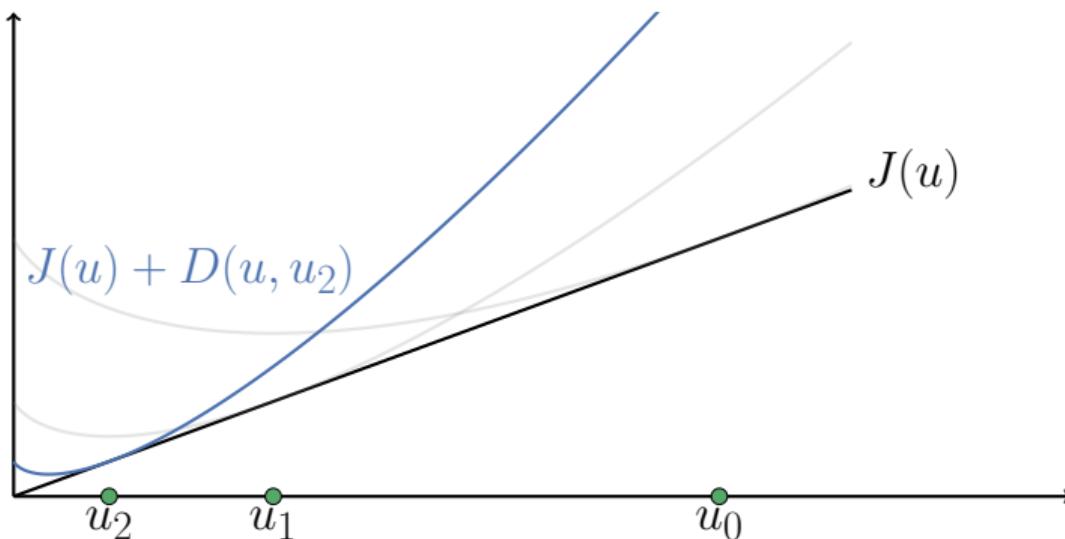
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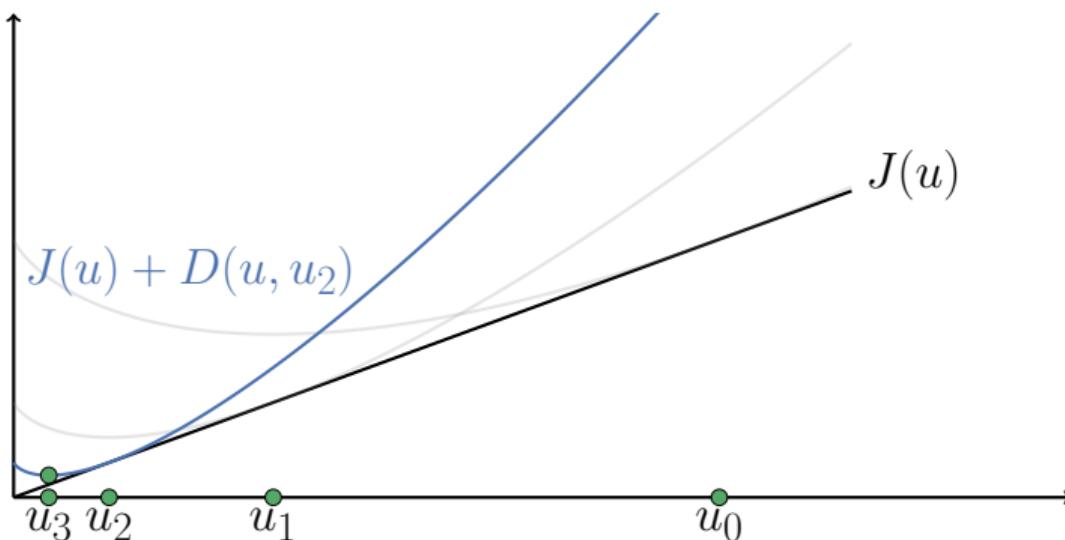
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Subsection 3

Mixed formulation

Good news

The iterates are strictly feasible, so the optimality condition becomes a PDE.

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Idea

Introduce latent variable $\psi \in W$ and enforce $Bu = \nabla R^*(\psi)$!

Latent variable proximal point

For some $\psi^0 \in W$, find $(u^k, \psi^k) \in V \times W$ s. t.

$$\begin{aligned}\alpha_k J'(u^k) + B^* \psi^k &= B^* \psi^{k-1}, \\ Bu^k - \nabla R^*(\psi^k) &= 0.\end{aligned}$$

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Important observation

This only requires discretising $u \in V$, not $u \in K$!

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Important observation

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Two approximations

When discretised, this gives *two* approximations of the observable Bu :

$$Bu_h \neq \nabla R^*(\psi_h).$$

In particular, if $B = I$, $\nabla R^*(\psi_h)$ is strictly feasible (while u_h probably is not).

Subsection 4

Comparison

There are lots of algorithms for obstacle-type VIs. How does LVPP compare?

	feasible?	inf-dim?	mesh-indep?	no param to 0/ ∞ ?
► active set/semismooth Newton	✓	✓	✗	✓
penalty/augmented Lagrangian	✗	✓	✓	✗
monotone multigrid	✓	✗	✓	✓
interior point	✓	✓	✓	✗
latent variable proximal point	✓*	✓	✓	✓

Active set/NCP function + semismooth Newton

Often mesh-dependent if applied directly because Lagrange multipliers are not smooth enough.

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interior point	✓	✓	✓	✗
latent variable proximal point	✓*	✓	✓	✓

Penalty/augmented Lagrangian methods

Always approximate from outside the feasible set; penalty $\rightarrow \infty$ even with AL in ∞ -dim.

There are lots of algorithms for obstacle-type VIs. How does LVPP compare?

	feasible?	inf-dim?	mesh-indep?	no param to 0/ ∞ ?
active set/semismooth Newton	✓	✓	✗	✓
penalty/augmented Lagrangian	✗	✓	✓	✗
► monotone multigrid	✓	✗	✓	✓
interior point	✓	✓	✓	✗
latent variable proximal point	✓*	✓	✓	✓

Monotone multigrid

Inherently discrete; requires hierarchy; specialised components required for each problem.

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monotone multigrid	✓	✗	✓	✓
► interior point	✓	✓	✓	✗
latent variable proximal point	✓*	✓	✓	✓

Interior point

Often involves mixed subproblems; line search very finicky as $h \rightarrow 0$, barrier $\rightarrow 0$.

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interior point	✓	✓	✓	✗
► latent variable proximal point	✓*	✓	✓	✓

Latent variable proximal point

Applies to very wide class of problems; requires mixed subproblems; $B \neq I$ not feasible.

Section 3

Bound constraints

We consider again the obstacle problem:

$$u \in \operatorname{argmin}_{v \in K} J(v) = \frac{1}{2} \int_{\Omega} \nabla v \cdot \nabla v \, dx - \int_{\Omega} fv \, dx,$$

for feasible set

$$K = \{v \in H_0^1(\Omega) \mid v \geq \phi \text{ a.e. in } \Omega\} \subsetneq V := H_0^1(\Omega).$$

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The mixed LVPP formulation becomes: for $\psi^0 = 0$, find $(u^k, \psi^k) \in H_0^1(\Omega) \times L^\infty(\Omega)$
s. t.

$$\begin{aligned}\alpha_k(\nabla u^k, \nabla v) + (\psi^k, v) &= \alpha_k(f, v) + (\psi^{k-1}, v), \\ (u^k, w) - (\exp(\psi^k) + \phi, w) &= 0,\end{aligned}$$

for all $(v, w) \in H_0^1(\Omega) \times L^\infty(\Omega)$.

We can discretise the LVPP iterations with a Galerkin scheme:

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Open question

What is the 'best' inf-sup stable finite element pair?

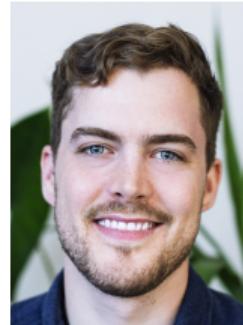
Method	Degree $p = 1$			Degree $p = 2$		
	h	$h/2$	$h/4$	h	$h/2$	$h/4$
Proximal Galerkin	15	13	12	15	16	12
Active Set (PETSc RSLS)	11	16	25			
Trust-Region (GALAHAD)	6	12	19	Not bound preserving		
Interior Point (IPOPT)	9	9	8			
IPOPT without Hessian	90	260	500			

Total number of linear system solves for popular solvers using various mesh sizes h .

For the proof that the Bregman proximal point iterations are PDEs, not VIs, see



B. Keith and T. M. Surowiec. “Proximal Galerkin: a structure-preserving finite element method for pointwise bound constraints”. In: *Foundations of Computational Mathematics* (2024). DOI: [10.1007/s10208-024-09681-8](https://doi.org/10.1007/s10208-024-09681-8).



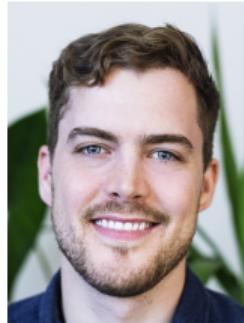
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Brendan Keith



Thomas Surowiec

For the proof that the convergence of the outer loop of proximal Galerkin is mesh-independent, see

-  **B. Keith, R. Masri, and M. Zeinhofer.** *A priori error analysis of the proximal Galerkin method.* arXiv:2507.13516. 2025.



Rami Masri



Marius Zeinhofer

You can easily apply other discretisations, too.

Mesh size h	2^{-1}	2^{-2}	2^{-3}	2^{-4}	2^{-5}	2^{-6}
Finite Difference	10	15	13	15	16	16

Degree p	8	16	24	32	40	48
Spectral Method	16	17	16	16	16	15

Total number of linear system solves for other discretisations.

Subsection 2

Fracture

Variational fracture is an important model with a non-convex energy.

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We consider the anti-plane shear test and only solve for vertical component of displacement:

$$(u, c) \in \underset{(v, d) \in K}{\operatorname{argmin}} J(v, d) = \frac{G}{2} \int_{\Omega} (\epsilon + (1 - \epsilon)(1 - d)^2) |\nabla v|^2 \, dx + \frac{G_c}{2} \int_{\Omega} \ell |\nabla d|^2 + \ell^{-1} d^2 \, dx,$$

for feasible set

$$K = \left\{ (u, c) \in H_D^1(\Omega) \times H^1(\Omega) \mid 0 \leq c_{\text{prev}} \leq c \leq 1 \text{ a.e. in } \Omega \right\}.$$

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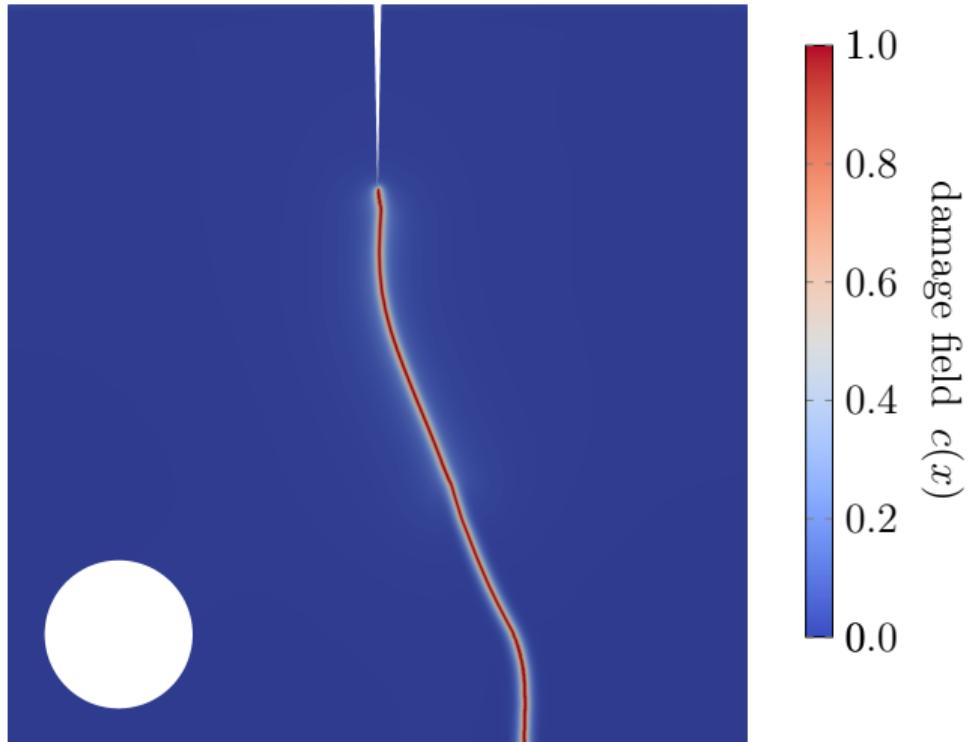
$$K = \left\{ (u, c) \in H_D^1(\Omega) \times H^1(\Omega) \mid 0 \leq c_{\text{prev}} \leq c \leq 1 \text{ a.e. in } \Omega \right\}.$$

To put this in our general abstraction, we take

$$B = (0, \text{id}), \quad \Omega_d = \Omega, \quad C(x) = \mathbb{R} \times [c_{\text{prev}}(x), 1].$$

$$u = -L$$

$$u = +L$$



The final damage field for the fracture problem considered.

We introduce a latent variable for the bound constraint on c .

The LVPP system becomes: find $(u^k, c^k, \psi^k) \in H_D^1(\Omega) \times H^1(\Omega) \times L^\infty(\Omega)$ s. t.

$$\alpha_k G((\epsilon + (1 - \epsilon)(1 - c^k)^2) \nabla u^k, \nabla v) = 0,$$

$$-\alpha_k G((1 - \epsilon)(1 - c^k)|\nabla u^k|^2, d) + \alpha_k G_c(\ell(\nabla c^k, \nabla d) + \ell^{-1}(c^k, d)) + (\psi^k, d) = (\psi^{k-1}, d),$$

$$(c^k, w) - \left(\frac{c_{\text{prev}} + \exp(\psi^k)}{\exp(\psi^k) + 1}, w \right) = 0,$$

for all $(v, d, w) \in H_0^1(\Omega) \times H^1(\Omega) \times L^\infty(\Omega)$.

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for all $(v, d, w) \in H_0^1(\Omega) \times H^1(\Omega) \times L^\infty(\Omega)$.

CG₁-CG₁-CG₁ discretisation

Each loading step takes an average of 2.85 proximal iterations (but with large variance).

Each proximal iteration takes an average of 5.44 Newton iterations.

Subsection 3

A quasi-variational inequality

Quasi-variational inequalities (QVIs) are even harder.

$$\text{VI: } u \in K \subsetneq V : F(u; v - u) \geq 0 \quad \forall v \in K$$



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The obstacle explicitly appears in the LVPP equations, so it can depend on other variables.

In a thermoforming QVI, a heated membrane is pushed upwards into a mold with force f .

The mold deforms with temperature:

$$K(T) = \{v \in H_0^1(\Omega) \mid v \leq \Phi := \Phi_0 + \xi T\}.$$

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The thermoforming QVI is to find $(u, T) \in K(\textcolor{red}{T}) \times H^1(\Omega)$ s. t.

$$\begin{aligned} (\nabla T, \nabla q) + \beta(T, q) &= (g(\Phi_0 + \xi T - u), q), \\ (\nabla u, \nabla(v - u)) &\geq (f, v - u), \end{aligned}$$

for all $(v, q) \in K(\textcolor{red}{T}) \times H^1(\Omega)$.

We again use the Shannon entropy (with signs switched).

The LVPP iterations become: find $(u^k, \psi^k, T^k) \in H_0^1(\Omega) \times L^\infty(\Omega) \times H^1(\Omega)$ s. t.

$$(\nabla T^k, \nabla q) + \beta(T^k, q) = (g(\exp(-\psi^k)), q),$$

$$\alpha_k(\nabla u^k, \nabla v) + (\psi^k, v) = \alpha_k(f, v) + (\psi^{k-1}, v),$$

$$(u^k, w) + (\exp(-\psi^k), w) = (\Phi_0 + \xi T^k, w).$$

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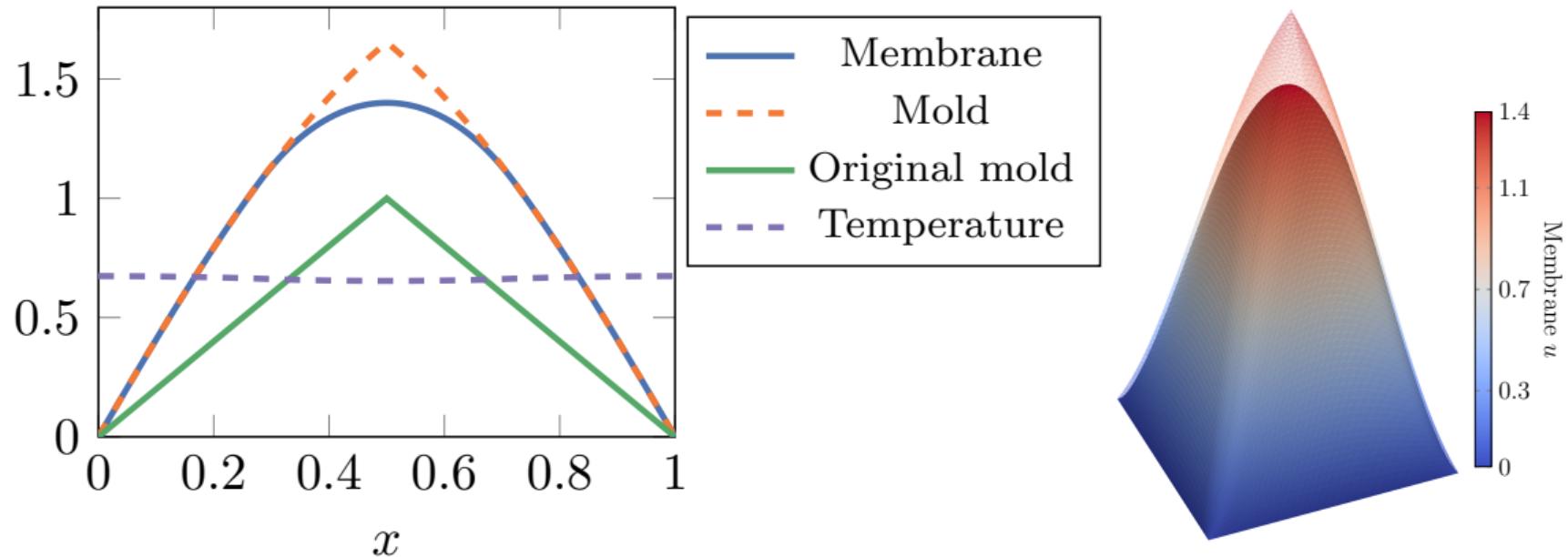
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for all $(v, w, q) \in H_0^1(\Omega) \times L^\infty(\Omega) \times H^1(\Omega)$.

Good news

Mesh-independent convergence and small iteration counts (~ 20 linear solves total).



Section 4

Gradient constraints

We consider again the Dirichlet energy:

$$u \in \operatorname{argmin}_{v \in K} J(v) = \frac{1}{2} \int_{\Omega} \nabla v \cdot \nabla v \, dx - \int_{\Omega} fv \, dx,$$

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but now impose *gradient* constraints:

$$K = \left\{ v \in H_0^1(\Omega) \mid |\nabla v| \leq \Phi \text{ a.e. in } \Omega \right\}.$$

Here $B = \nabla$ and $C(x)$ is the Euclidean ball of radius $\Phi(x)$.

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Here $B = \nabla$ and $C(x)$ is the Euclidean ball of radius $\Phi(x)$.

Numerical methods are scarce.

We use the modified Hellinger entropy $R(\mathbf{a}) = -\sqrt{\Phi^2 - |\mathbf{a}|^2}$.

The LVPP iteration becomes: find $(u^k, \Psi^k) \in H_0^1(\Omega) \times L^\infty(\Omega, \mathbb{R}^n)$ ($n = \text{spatial dimension}$) s. t.

$$\alpha_k(\nabla u^k, \nabla v) + (\Psi^k, \nabla v) = \alpha_k(f, v) + (\Psi^{k-1}, \nabla v),$$

$$(\nabla u^k, W) - \left(\frac{\Phi \Psi^k}{\sqrt{1 + |\Psi^k|^2}}, W \right) = 0$$

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for all $(v, W) \in H_0^1(\Omega) \times L^\infty(\Omega, \mathbb{R}^n)$.

Good news

Again, we observe (but do not yet prove) robust mesh-independent convergence.

LVPP extends to *intersections of constraints*. Take again the Dirichlet energy

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but now impose *both obstacle and gradient* constraints:

$$K = \left\{ v \in H_0^1(\Omega) \mid v \geq \phi \text{ and } |\nabla v| \leq \Phi \text{ a.e. in } \Omega \right\}.$$

LVPP extends to *intersections of constraints*. Take again the Dirichlet energy

$$u \in \operatorname{argmin}_{v \in K} J(v) = \frac{1}{2} \int_{\Omega} \nabla v \cdot \nabla v \, dx - \int_{\Omega} fv \, dx,$$

but now impose *both obstacle and gradient* constraints:

$$K = \{v \in H_0^1(\Omega) \mid v \geq \phi \text{ and } |\nabla v| \leq \Phi \text{ a.e. in } \Omega\}.$$

Now $B = (\operatorname{id}, \nabla)^{\top}$ and $C(x) = [\phi(x), \infty) \times \mathcal{B}(0, \Phi(x))$, \mathcal{B} the Euclidean ball.

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Legendre functions for intersection

Legendre functions for intersections of sets are additive:

$$R(a, \mathbf{a}) = (a - \phi) \log(a - \phi) - (a - \phi) - \sqrt{\Phi^2 - |\mathbf{a}|^2}.$$

The induced isomorphism has two components:

$$\nabla R^*((a^*, \mathbf{a}^*)) = \begin{pmatrix} \phi + \exp a^* \\ \Phi \mathbf{a}^* \\ \sqrt{1 + |\mathbf{a}^*|^2} \end{pmatrix}.$$

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The LVPP iteration becomes: find $(u^k, \psi^k, \Psi^k) \in H^1(\Omega) \times L^\infty(\Omega) \times L^\infty(\Omega, \mathbb{R}^n)$ s. t.

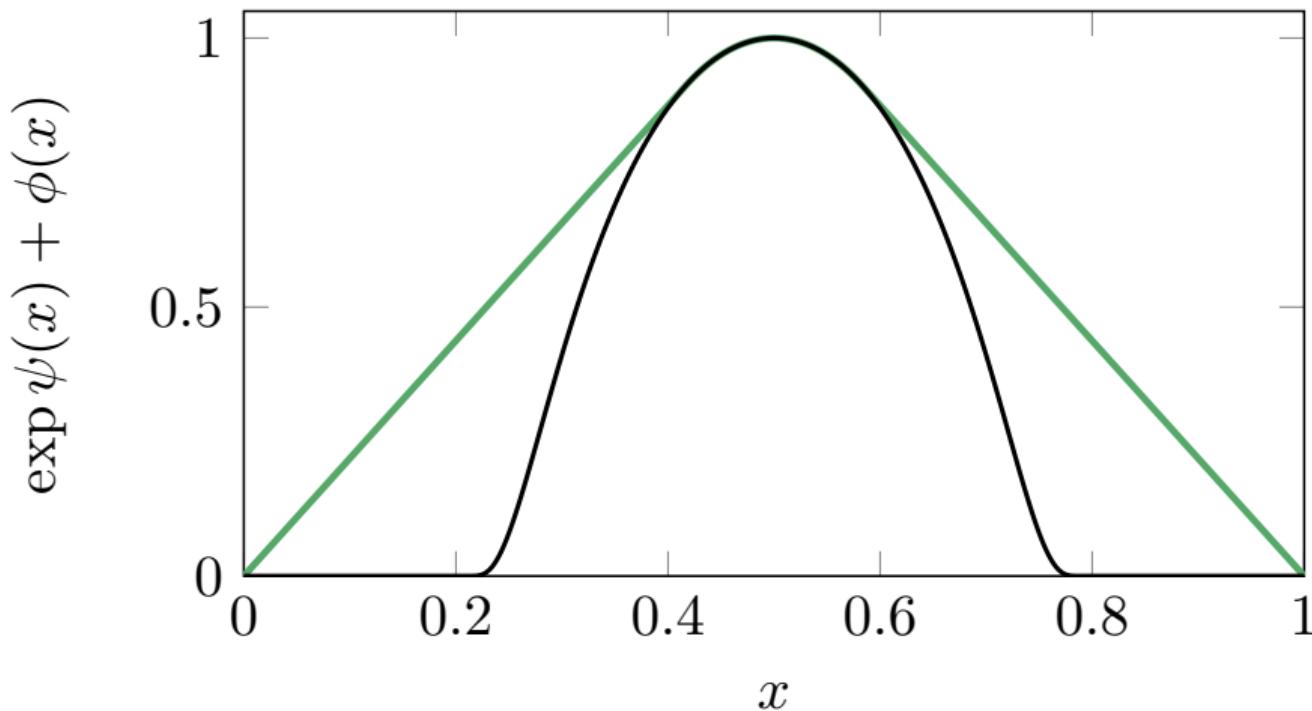
$$\alpha_k (\nabla u^k, \nabla v) + (\psi^k, v) + (\Psi^k, \nabla v) = (\psi^{k-1}, v) + (\Psi^{k-1}, \nabla v)$$

$$(u^k, w) - (\exp \psi^k, w) - (\phi, w) = 0$$

$$(\nabla u^k, W) - \left(\frac{\Phi \Psi^k}{\sqrt{1 + |\Psi^k|^2}}, W \right) = 0$$

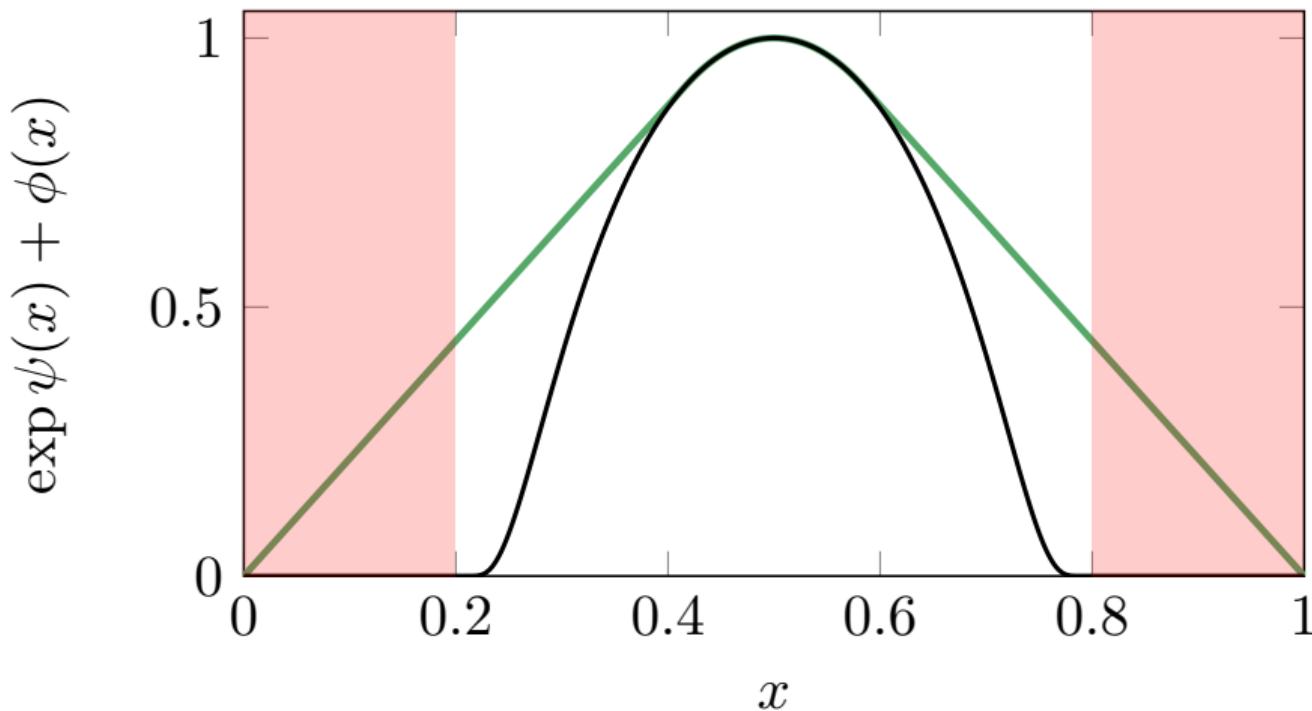
for all $(v, w, W) \in H^1(\Omega) \times L^\infty(\Omega) \times L^\infty(\Omega, \mathbb{R}^n)$.

We set $f = 0$ and vary the gradient constraints.



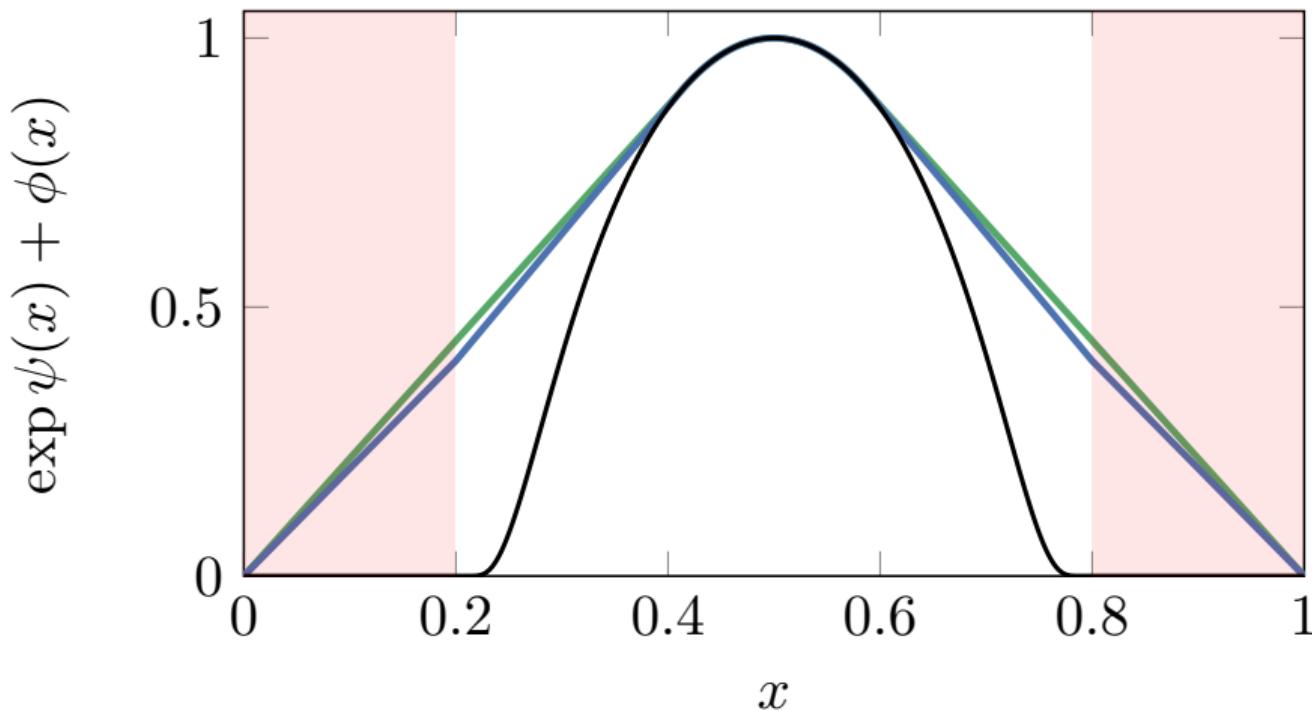
No gradient constraints.

We set $f = 0$ and vary the gradient constraints.



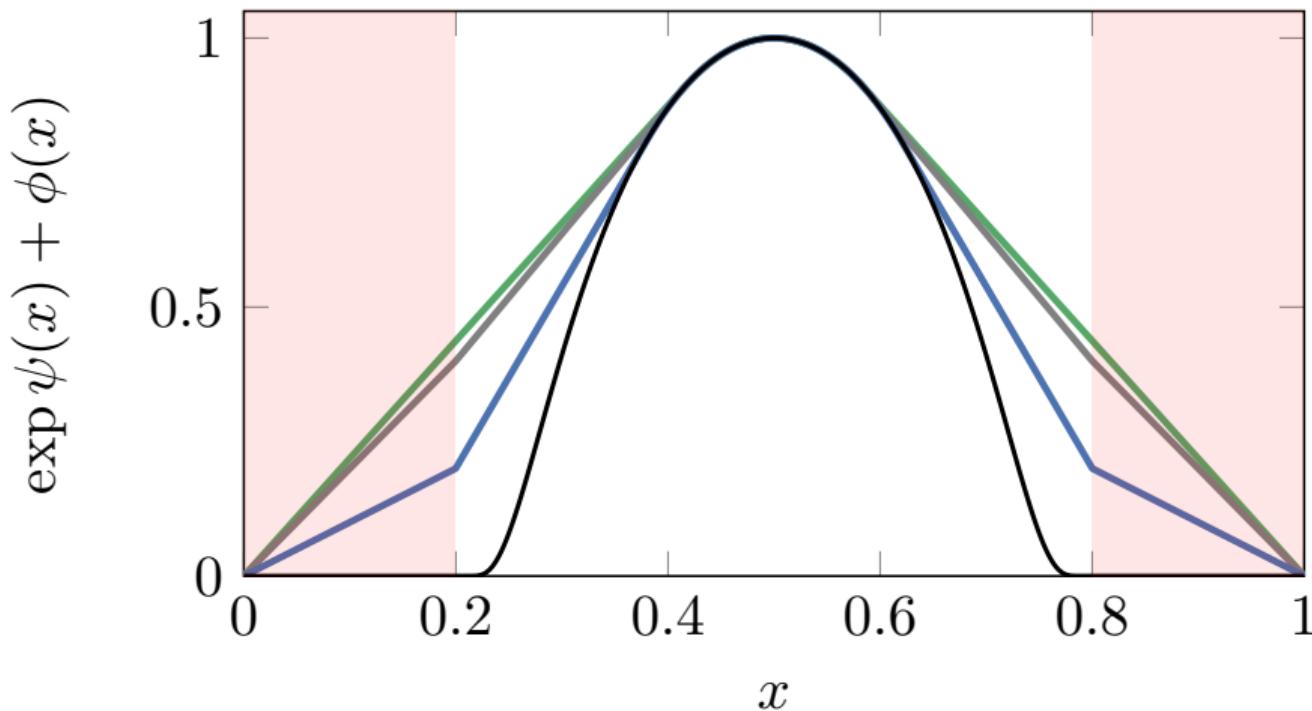
Apply gradient constraints on $[0, 0.2] \cup [0.8, 1]$.

We set $f = 0$ and vary the gradient constraints.



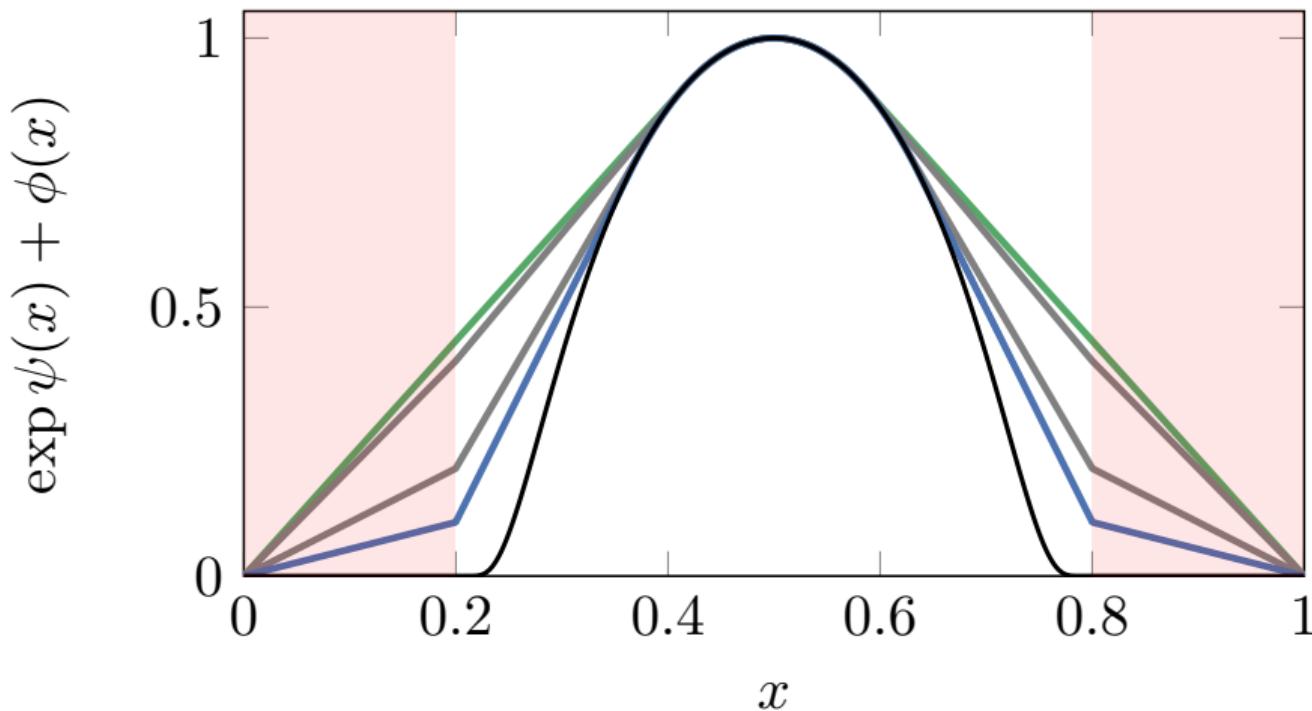
Light gradient constraints.

We set $f = 0$ and vary the gradient constraints.



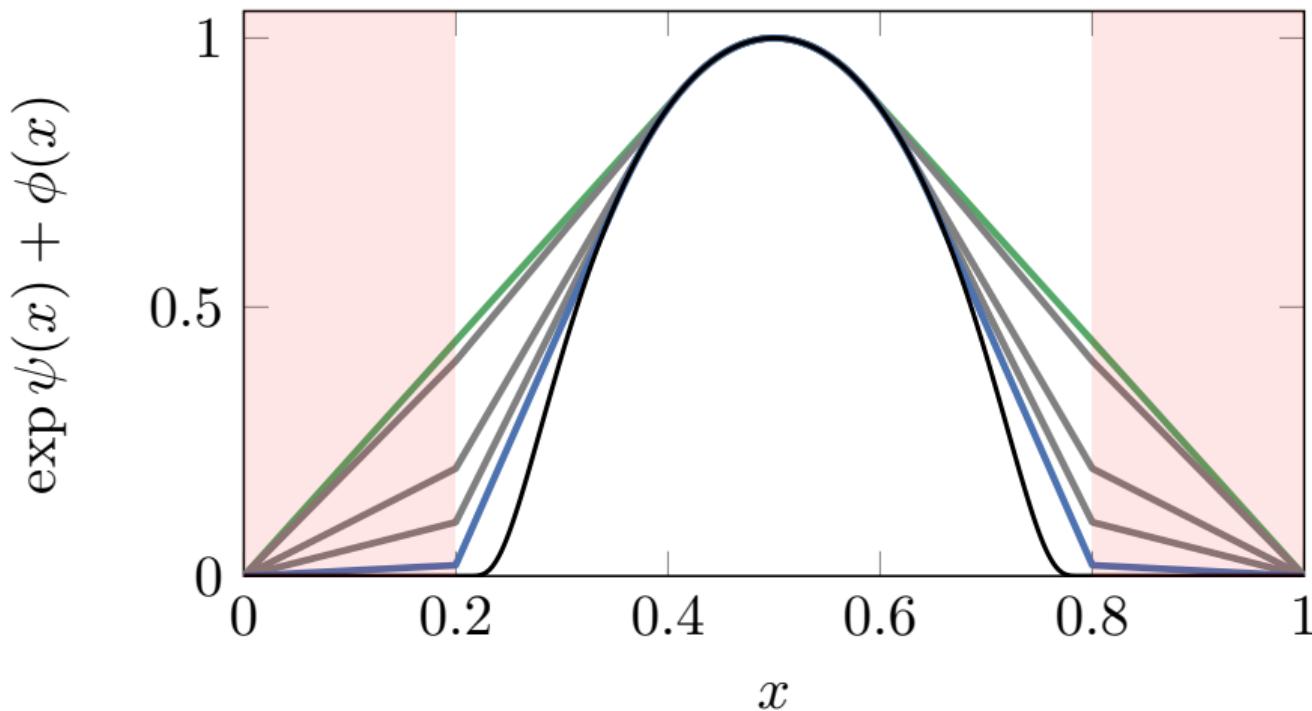
Medium gradient constraints.

We set $f = 0$ and vary the gradient constraints.



Heavy gradient constraints.

We set $f = 0$ and vary the gradient constraints.



Extreme gradient constraints.

Section 5

Eigenvalue constraints

Eigenvalue constraints are very important, but very difficult to enforce numerically.

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The Landau–de Gennes model of nematic liquid crystals minimises

$$J(Q) = \frac{1}{2} \int_{\Omega} \nabla Q : \nabla Q \, dx + \frac{1}{2} \int_{\Omega} A \operatorname{tr}(Q^2) \, dx + \frac{1}{4} \int_{\Omega} C(\operatorname{tr}(Q^2))^2 \, dx$$

for a symmetric traceless matrix field Q .

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for a symmetric traceless matrix field Q .

To be physical, Q must satisfy eigenvalue constraints (n = spatial dimension)

$$\lambda_i(Q) \in [-1/n, (n-1)/n], \quad i = 1, \dots, n,$$

but this is usually ignored as too difficult.

Fix $n = 2$ for simplicity. We employ as Legendre function

$$R(A) = \text{tr} \left((A + I/2) \log(A + I/2) + (I/2 - A) \log(I/2 - A) \right),$$

with $\nabla R^*(A^*) = \text{tanhm}(A^*/2)/2$,

where \log and tanhm are the matrix logarithm and hyperbolic tangent functions.

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where \log and \tanhm are the matrix logarithm and hyperbolic tangent functions.

The LVPP iteration becomes: find $(Q^k, \psi^k) \in H_D^1(\Omega, \mathbb{R}_{\text{sym},\text{tr}}^{2 \times 2}) \times L^\infty(\Omega, \mathbb{R}_{\text{sym},\text{tr}}^{2 \times 2})$ s. t.

$$\alpha_k J'(Q; V) + (\psi^k, V) = (\psi^{k-1}, V)$$

$$(Q, w) - \left(\frac{1}{2} \tanhm(\psi/2), w \right) = 0$$

for all $(V, w) \in H_0^1(\Omega, \mathbb{R}_{\text{sym},\text{tr}}^{2 \times 2}) \times L^\infty(\Omega, \mathbb{R}_{\text{sym},\text{tr}}^{2 \times 2})$.

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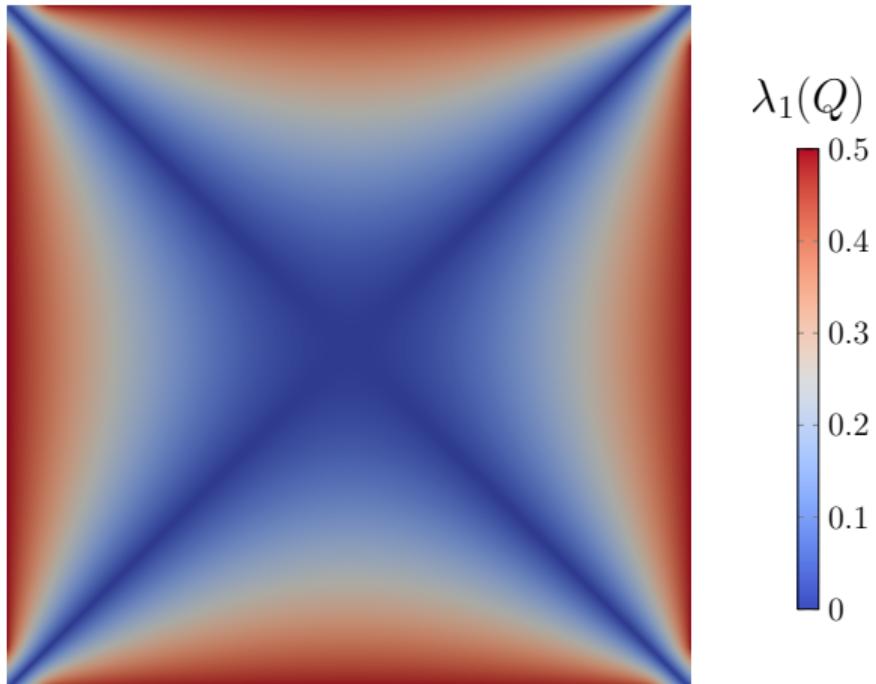
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All of this extends straightforwardly to $n > 2$.

Good news

Mesh-independent convergence, ~ 6 proximal steps, ~ 11 Newton iterations.



The larger eigenvalue $\lambda_1(Q)$. Both eigenvalues satisfy the inequality constraints.

Section 6

Conclusions

Conclusion

Latent variable proximal point is a powerful framework for problems with inequality constraints.



J. S. Dokken et al. "The latent variable proximal point algorithm for variational problems with inequality constraints". In: *Computer Methods in Applied Mechanics and Engineering* 445 (2025), p. 118181. DOI: 10.1016/j.cma.2025.118181.

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Good news

Many open questions remain! Proofs, discretisations, solvers, nonconvex constraints,

Section 7

Monge–Ampère

For uniformly positive $\rho \in C(\overline{\Omega})$ and $g \in C^3(\overline{\Omega})$, find the unique $u \in K$ such that

$$\det(\nabla^2 u) = \rho \quad \text{in } \Omega, \quad u = g \quad \text{on } \partial\Omega.$$

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where

$$K = \{u \in H^2(\Omega) \cap H_g^1(\Omega) \mid \nabla^2 u \succeq 0 \text{ a.e. in } \Omega\}$$

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This set has Legendre function

$$R(A) = \operatorname{tr}(A \ln A - A), \quad \text{with } \nabla R^*(A^*) = \exp A^*.$$

So, introduce a latent variable

$$\psi = \ln \nabla^2 u \quad \iff \quad \exp \psi = \nabla^2 u.$$

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Using the identity $\det(\exp \psi) = \exp(\operatorname{tr} \psi)$, the main PDE becomes an algebraic relation:

$$\operatorname{tr} \psi = \ln \rho.$$

This leads to the variational formulation: find $u \in H^2(\Omega) \cap H_g^1(\Omega)$ and $\psi \in L^\infty(\Omega; \mathbb{R}_{\text{sym}}^{n \times n})$ s. t.

$$\begin{aligned} (\nabla^2 u, w) - (\exp \psi, w) &= 0, \\ (\operatorname{tr} \psi, v) &= (\ln \rho, v), \end{aligned}$$

for all $v \in H^2(\Omega) \cap H_0^1(\Omega)$ and $w \in L^\infty(\Omega; \mathbb{R}_{\text{sym}}^{n \times n})$.

Most codes cannot discretise $u \in H^2$, so introduce $T = \nabla u$:

Find $u \in H_g^1(\Omega)$, $T \in H^1(\Omega, \mathbb{R}^n)$, and $\psi \in L^\infty(\Omega; \mathbb{R}_{\text{sym}}^{n \times n})$ s. t.

$$\begin{aligned}(T, S) - (\nabla u, S) &= 0, \\ (\nabla T, w) - (\exp \psi, w) &= 0, \\ (\operatorname{tr} \psi, v) &= (\ln \rho, v),\end{aligned}$$

for all $v \in H_0^1(\Omega)$, $S \in H^1(\Omega, \mathbb{R}^n)$, and $w \in L^\infty(\Omega; \mathbb{R}_{\text{sym}}^{n \times n})$.

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Numerical experiments with CG _{p} -[CG _{$p+1$}] ^{n} -[CG _{p}]_{sym} ^{$n \times n$} show excellent convergence, for $p \geq 2$.

Good news

Robust Newton convergence, error of $\sim 10^{-13}$ for $p = 14$.