

Semester 2 Notes

Connor Stocks

February 7, 2025

Contents

1	Analysis II	3
1.1	Propositional Logic	3
1.2	Continuous Functions	3
2	waves and fields	4
2.1	Periodic Motion	4
3	Modern Physics	9
3.1	Black Body Radiation	9
3.2	The Photoelectric Effect	10
4	Intro to Quantum Physics	12
4.1	A Brief History	12
4.2	Rutherford And Compton Scattering	13
4.3	Milikan's Oil Drop	14
5	Mathematical Methods II	16
5.1	Non Cartesian Co-Ordinates	16
6	Derivations	17
6.1	Analysis II	17
6.2	Waves and Fields	17
6.3	Modern Physics	17
6.4	Intro to Quantum Physics	18
6.5	Maths Methods II	19

1 Analysis II

1.1 Propositional Logic

27/1/25

- $\neg(p \cup q) = (\neg p) \cap (\neg q)$
- $\neg(p \cap q) = (\neg p) \cup (\neg q)$
- $\neg(\exists p) = \forall p$
- $p \Rightarrow q = (p \cap q) \cup \neg q$
- so $\neg(p \Rightarrow q) = (\neg p \cup \neg q) \cap \neg \neg q = \neg p \cup q$

1.2 Continuous Functions

31/1/25

Continuity

- For $f : D \rightarrow \mathbb{R}$, f is continuous at a point $a \in D$ if:
 $\forall \epsilon > 0, \exists \delta > 0 : \forall x \in D, |x - a| < \delta \Rightarrow |f(x) - f(a)| < \epsilon$
- f is uniformly continuous if continuous $\forall a$, or
 $\forall \epsilon > 0, \exists \delta > 0 : \forall x, y \in D, |x - y| < \delta \Rightarrow |f(x) - f(y)| < \epsilon$
- In general, continuity \nRightarrow uniform continuity
- A function $f : D \rightarrow \mathbb{R}$ is continuous at $a \in \mathbb{R}$ iff:
 $\forall (x_n)_{n \in \mathbb{N}}, x_n \in D, \lim_{n \rightarrow \infty} x_n = a, \lim_{n \rightarrow \infty} f(x_n) = f(a)$

3/2/25

Theorem 1 (Intermediate Value Theorem) for $f : D \rightarrow C \subseteq \mathbb{R}$
then $u \in C \Rightarrow \exists x \in D : u = f(x)$

2 waves and fields

2.1 Periodic Motion

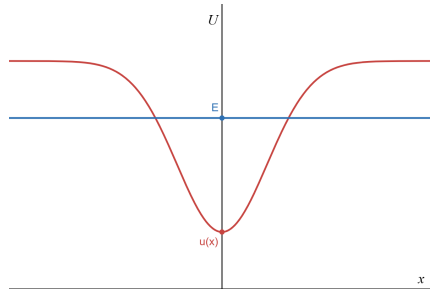
Periodic Motion In 1D

28/1/25

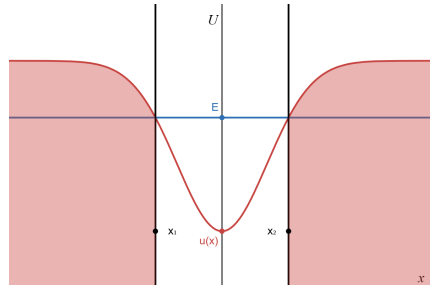
- Periodic motion is one in which the position repeats values regularly
- It can be characterised like so:

$$r(t + nT) = r(t), \forall t \in \mathbb{R}, n \in \mathbb{N}, T = \text{period}$$

- If we consider a situation in a conservative field, with a potential $u(x)$, and a system with a total energy E :



- We can see that the system would be bound within the region $U < E$:



- This naturally leads to periodic motion, as at $x = x_1, v = 0$
- At this point, the energy is not minimised, and so the system tends towards $x = 0$, where $v \neq 0$, causing periodic motion.
- One way to achieve this is through simple harmonic motion (SHM)
- SHM is any system where $\ddot{x} = -x$
- We can see here that $u(x) = -\int f dx = \int \frac{\ddot{x}}{m} dx$, or $f = -\frac{du}{dx} \Rightarrow m\ddot{x} = -\frac{du}{dx}$
- We can then add another term to take energy out of the system:
- $m\ddot{x} = -\frac{du}{dx} - \gamma\dot{x}$ in general.

30/1/25

-

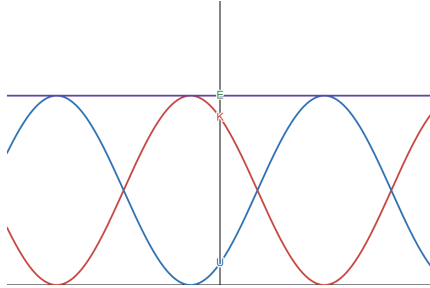
-

- Gravity is a conservative force so $u = u(\vec{r})$, and angular momentum is conserved
- Evidently, from $\vec{L} = \vec{r} \times \vec{p}$, the position \vec{r} is perpendicular to the constant angular momentum \vec{L} , and thus the system is confined to a single plane.
- Let v_ϕ = the component of v perpendicular to \vec{r} , and v_r = the component of v parallel to \vec{r}
- Then $L = m\vec{r} \times (r\dot{\phi})$
- Since $K = \frac{1}{2}m\dot{\vec{r}}^2 = \frac{1}{2}m(\dot{\phi}^2 + \dot{r}^2) = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2)$
- $\dot{\phi} = \frac{L}{mr^2} \Rightarrow K = \frac{1}{2}m(\dot{r}^2 + r^2(\frac{L}{mr^2})^2)$
- So $E = K + u = \frac{1}{2}m(\dot{r}^2 + (L/mr)^2 + u(\vec{r}))$

Simple Harmonic Motion

31/1/25

- Simple harmonic motion is any motion of the form $x(t) = A\cos(\omega t + \phi)$
- From this, we can see that $\ddot{x}(t) = -A\omega^2 \cos(\omega t + \phi) = -\omega^2 x$
- This differential equation, $\ddot{x} \propto x$, characterises simple harmonic oscillation(SHO)
- Then, $f = m\ddot{x} = -m\omega^2 x$, is only dependant on position
- Then, the potential energy: $U(x) = -\int f dx = m\omega^2 \int x dx = \frac{1}{2}m\omega^2 x^2$
- Then, $E = K + U = \frac{1}{2}m\dot{x}^2 + \frac{1}{2}m\omega^2 x^2 = \frac{1}{2}mA^2\omega^2$
- E, K, and U will look like this over time, with $\overline{U} = \overline{K}$ over one period:



Simple Pendulums and Small oscillators

- For a simple pendulum, $|\tau| = I\ddot{\theta} = -f_{\perp} \cdot r = -f_g \sin(\theta) \cdot l$
- Then, using the small angle approximation $\sin(\theta) = \theta$:

$$I\ddot{\theta} = -f_g l \theta$$
- Another example of simple harmonic motion is the torsion pendulum
- This is a mass on a string that can twist around its axis:

$$\tau = -c\theta \Rightarrow I\ddot{\theta} = -c\theta$$

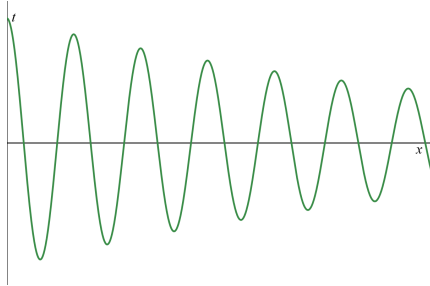
Damped Harmonic Motion

6/2/25

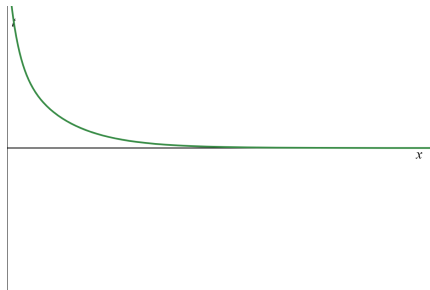
- We expect the amplitude A of a damped harmonic oscillator to decrease over time.
- For DHM, The 2nd law of motion looks like this:
- $m\ddot{x} = -kx - b\dot{x}$
- Setting $b=0$, we recover the SHM characteristic equation
- For $b^2 - 4km < 0$, This is called underdamping

- Underdamping will have the following form of motion:

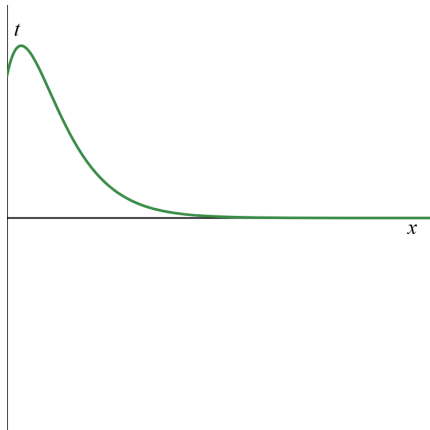
$$x(t) = Ae^{-\frac{b}{2m}t} \cos(\omega t + \varphi) \quad (2.1)$$



- Then considering the quantity $\frac{2m}{b}$, we see $|x(t + \frac{2m}{b})| = e^{-1}|x(t)|$
- If $b^2 > 4mk$, we get overdamped motion:
- $x(t) = A_1 e^{-\lambda_1 t} + A_2 e^{-\lambda_2 t}$



- and if $b^2 = 4mk$, critically damped Motion:
- $x(t) = (A_1 + A_2 t)e^{-\frac{b}{2m}t}$



Forced Harmonic Oscillation & Resonance

7/2/25

- A forced harmonic oscillator (FHO) is once subject to an external driving force f_d
- $\ddot{x} + \frac{b}{m}\dot{x} + \frac{k}{m}x = \frac{1}{m}f_d$
- Considering $f_d = \cos(\omega_d t + \varphi_d)$:

$$\omega_r = \omega_0 \sqrt{\frac{k}{m} - \frac{b^2}{2m^2}} \quad (2.2)$$

- Where A is maximised at $\omega_d = \omega_r$

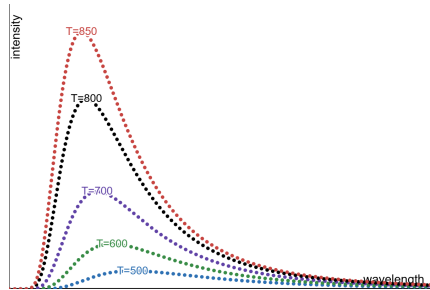
3 Modern Physics

3.1 Black Body Radiation

31/1/25

Observations

- Black bodies satisfy 3 criteria:
 - absorbs all incident radiation
 - emits as much radiation as possible at any wavelength
 - emits radiation isotropically
- By studying the spectra of black bodies at different temperatures:



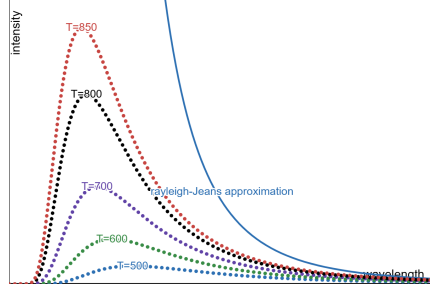
- We can observe 2 main laws:
 - Wein's displacement law: $\lambda_{max} = \frac{b}{T}$
 - Stephans law: $P(T) = \sigma AT^4$

Rayleigh-Jeans Law

- In June 1900, John Rayleigh discovered that the spectral radiance $\propto \lambda^{-4}$
- In 1905, James Jean derived the full Rayleigh-Jeans law including constants
- The Rayleigh-Jeans law is a classical approximation of the radiance of a black body
- Rayleigh-Jeans law can be expressed as so:

$$B(\lambda, T) = \frac{2ck_bT}{\lambda^4} \quad (3.1)$$

- This produced the following prediction for spectral radiance at T=800:



- This presents many problems, mainly that it doesn't match for $\lambda < \lambda_p$, and $I \rightarrow \infty$ as $\lambda \rightarrow 0$

Planck's Law

- Rayleigh-Jeans law, as seen in the derivation, can be expressed as:
- $R(\lambda) \propto \rho(f) \cdot E_{avg}$
- Where ρ and E are the number of waves and average energy of each wave
- In the Rayleigh-Jeans law, $E_{avg} = kT$
- This was due to the equipartition theorem
- Max Planck's great insight was to abandon the equipartition theorem, instead writing:
- $E = n\epsilon \Rightarrow \overline{E} = \frac{hf}{e^{\frac{hf}{kT}} - 1}$
- And so:

$$R(\lambda) = \frac{2hc}{\lambda^5} \cdot \frac{1}{e^{\frac{hc}{\lambda kT}} - 1} \quad (3.2)$$

3.2 The Photoelectric Effect

6/2/25

- In 1887, Heinrich Hertz discovered that light rays incident on a metal plate produces an emf across the material
- This is due to electrons being liberated by the light
- This works with metals due to the delocalised electron bonds
- The electrons would be liberated with a range of kinetic energies up to a maximum
- There was also an observed threshold frequency, below which a ray can't liberate electrons

classical Predictions

- $V_s \propto I$
- $K_{max} \propto A$
- $\frac{dN}{dt} \propto f$

Einstein's Explanation

- In 1905, Einstein published his explanation of the photoelectric effect
- He explained that the light interacts in a 1:1 manner, with $E = hf$
- This is insightful as Max Planck had only quantised matter oscillation
- Einstein quantised light
- This explains the intensity profile, as the particle light must interact with the electrons 1:1
- It also explained the threshold frequency phenomenon, as light with insufficient energy just reflects off of the metal
- $E_\gamma = hf = \varphi + \frac{1}{2}m_e v_{max}^2$ or $hf = \varphi + K_{max}$

4 Intro to Quantum Physics

4.1 A Brief History

Ultraviolet Catastrophe

29/1/25

- In 1860, Gustav Kirchhoff discovered what is now called black body radiation
- This is where an object with internal energy radiates electromagnetic waves correlated with its thermal properties.
- Classical predictions, using namely the Rayleigh-Jeans law, predicted an infinite amount of low-frequency light waves:
- In 1900, Max Planck produced a solution to this problem which correctly predicted the spectral radiance of black body radiation
- This theory worked by quantising light, which he expected to disappear in the final solution
- He still did not believe that light was actually quantised
- He earned the Nobel Prize in 1918

Photoelectric Effect

- In 1887, Heinrich Hertz discovered the photoelectric effect
- This is when you expose a metal plate to high intensity light and observe a voltage over the plate
- The contemporary theories predicted that the energy transferred should be proportional to the frequency
- Instead, there is an observed threshold frequency, below which no energy is transferred
- In 1905, Albert Einstein explained this by employing Max Planck's quantum light
- He won the Nobel Prize in 1921

Atomic Models

- In 1911, Ernst Rutherford discovered positively charged atomic nuclei
- The presence of the electron in the atom, then, predicted that atoms were unstable, decaying in $10^{-11}s$ as the electron would fall into the nuclei
- In 1913, Niels Bohr predicted that angular momentum was quantised, restricting electrons to specific "orbits" and stopping atomic decay

- Neils Bohr could not come up with a mechanism for angular momentum quantisation, however
- He won the Nobel Prize in 1913

DeBroglie's Electron Strings

- In 1924, Louis DeBroglie suggested wave particle duality
- This suggests that sufficient restrictions upon an electrons position force it to exhibit wave like behaviour
- When employed on an electron in orbit, the electron waves superpose, producing a standing wave
- Standing waves can only exist with an integer multiplier on the fundamental frequency
- This multiplier corresponds to the energy the electron can have, quantising orbital angular momentum
- He won the Nobel Prize in 1929

4.2 Rutherford And Compton Scattering

30/1/25

Rutherford Scattering

- In 1897, Thompson devised the Plum Pudding model.
- In 1906 in Manchester, Ernst Rutherford devised the α scattering experiment
- He shot α particles at a thin sheet of gold and observe the distribution of scattering angles
- The Plum Pudding model predicts the following distribution:
- It predicted little to no scattering due to the diffusivity of the positive charge
- Whereas the distribution Rutherford found was the following:
- In response to this, rutherford devised the nuclear model
- This explained that the nucleus repels the α particles due to the like charges, sometimes up to a full π^c

Compton Scattering

- In 1923, Arthur Compton discovered Compton scattering
- This is where x-rays incident on an atom reflect at a different wavelength
- The classical model, with $p_\gamma = 0$, predicted no shift:
- But in reality, there was a shift:
- This showed that light does carry momentum, and therefore supported the photon model
- Using conservation of momentum and conservation of energy, you can arrive at the following equation:

$$\lambda - \lambda' = \frac{h}{m_0 \cdot c} [1 - \cos(\theta)] \quad (4.1)$$

4.3 Milikan's Oil Drop

5/2/25

- In 1897, Robert Milikan devised an experiment to measure the charge of the electron:
- The 2 plates are connected to a power source, producing an emf
- A nozzle containing oil is small to be sprayed
- There is a camera or microscope to observe the drops
- The drops become charged due to friction with the nozzle
- Oil drops that move up are large charge or low mass, and vice versa
- There is also the bouyant and viscous forces
- $\vec{f}_b = \frac{4}{3}\pi gr^3 \rho_{air} \hat{e}_z$ since the oil is displacing the air
- from Stoke's Law, $\vec{f}_v = -6\pi\eta_{air} \cdot rv$
- $\vec{f}_v = -6\pi\eta_{air}$
- Since it is hard to measure the properties of the drops explicitly, we can instead formulate expressions for them:

$$r = \sqrt{\frac{6\eta_{air}v_{off}}{2(\rho_{oil} - \rho_{air})g}} \quad (4.2)$$

$$q = \frac{g\sqrt{2}(v_{on} + v_{off})d}{V_0} \cdot \sqrt{\frac{\eta_{air}^3 v_{off}}{(\rho_{oil} - \rho_{air})g}} \quad (4.3)$$

- You then repeat this multiple times and record the charge of each drop 6/2/25
- There will be a random distribution of charges, however the smallest difference in charge will be $e = 1.60 \cdot 10^{-19}$

5 Mathematical Methods II

5.1 Non Cartesian Co-Ordinates

Polar Co-Ordinates

31/1/25

- for $x\hat{i} + y\hat{j}$, the equivalent polar basis vectors are:
 $\hat{r} = (\cos(\theta), \sin(\theta))$, $\hat{\theta} = (-\sin(\theta), \cos(\theta))$ where $\tan(\theta) = \frac{y}{x}$
- Then, $\cos(\theta)\hat{r} - \sin(\theta)\hat{\theta} = \hat{i}$ and $\sin(\theta)\hat{r} + \cos(\theta)\hat{\theta} = \hat{j}$

3D Polar Co-Ordinates

- There are 2 main 3D Polar Coordinate Bases:
- Spherical $(\hat{r}, \hat{\theta}, \hat{\phi})$
- And Cylindrical $(\hat{\rho}, \hat{\theta}, \hat{z})$

6 Derivations

6.1 Analysis II

6.2 Waves and Fields

eqn 2.1 Underdamping Motion Form

- $m\ddot{x} = -kx - b\dot{x}$
 - $\Rightarrow \ddot{x} + \frac{b}{m}\dot{x} + \frac{k}{m}x = 0$
 - This has a characteristic equation $\lambda^2 + \frac{b}{m}\lambda + \frac{k}{m}$
 - With solutions $\lambda = -\frac{b}{2m} \pm \frac{1}{2}\sqrt{\left(\frac{b}{m}\right)^2 - 4\frac{k}{m}}$
 - For case $b^2 < 4mk$, $\lambda = -\frac{b}{2m} \pm i\sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$
- $$\Rightarrow x(t) = \text{Re}[Ae^{-\frac{b}{2m}t}e^{\pm i\sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}}]$$
- $\Rightarrow \boxed{x(t) = Ae^{-\frac{b}{2m}t} \cos(\omega t + \varphi)}$

eqn 2.2 Resonance Frequency

DE

- For $\ddot{x} + \frac{b}{m}\dot{x} + \frac{k}{m}x = \frac{1}{m}f_d$, $f_d = d \cos(\omega_d t + \varphi)$
- assume $x(t) = \mathcal{A}e^{i\omega_d t}$, $f_d = \mathcal{D}e^{i\omega}$
- Into DE:
- $-\omega_d^2 \mathcal{A}e^{i\omega_d t} + \frac{b}{m}i\omega_d \mathcal{A}e^{i\omega_d t} + \frac{k}{m}\mathcal{A}e^{i\omega_d t} = \frac{1}{m}C e^{i\omega_d t}$
- $\Rightarrow A(\omega_d^2 - \frac{b}{m}i\omega_d) - \frac{k}{m} = -\frac{1}{m}C$
-

6.3 Modern Physics

eqn 3.1 Rayleigh-Jeans Law

- Firstly, we begin by noting that $R(\lambda) = \rho(f) \cdot \frac{c}{4\pi}$ where $\rho(f)$ is the energy density inside a Jeans cube
- Then, $\rho(f) = N(f) \cdot E_{avg}$
- $N(f) = \frac{8\pi f^2}{c^3}$, and $E_{avg} = kT$
- Then $\rho(f) = \frac{8\pi f^2 kT}{c^3}$
- $= \frac{8\pi kT}{\lambda^2}$

- $\Rightarrow \boxed{\frac{2ckT}{\lambda^4}}$

eqn 3.2 Planck's Law

6.4 Intro to Quantum Physics

eqn 4.1 Compton Wavelength

Milikan's Experiment

For plates with p.d. V_0 , separation d , and droplet of charge q and radius r moving at velocity v :

- $\vec{f}_e = q \cdot \vec{E} = q \cdot \frac{V_0}{d} \cdot \hat{e}_z$
- $\vec{f}_g = mg \cdot \hat{e}_z$ and $m = V \cdot \rho_{oil} = \frac{4}{3}\pi r^3 \rho_{oil} \Rightarrow \vec{f}_g = \frac{4}{3}\pi r^3 \rho_{oil} g \hat{e}_z$
- oil displacing air $\Rightarrow \vec{f}_b = w_{displaced} = v_{air} \rho_{air}$
- $v_{air} = v_{oil} \Rightarrow \vec{f}_b = \frac{4}{3}\pi r^3 \rho_{air} g \hat{e}_z$
- from Stoke's Law, $\vec{f}_v = -6\pi\eta_{air} \cdot rv$

eqn 4.2 Milikan's Oil Drop Radius

If we set $V_0 = 0$ and consider the terminal velocity of the drop v_{off}

- $\vec{f}_g + \vec{f}_b + \vec{f}_v = 0$
- Note, $\vec{f}_g + \vec{f}_b + \vec{f}_v = 0 \Rightarrow f_v = f_g - f_b$
- $\frac{4}{3}\pi \rho_{oil} g r^3 \hat{e}_z - \frac{4}{3}\pi \rho_{air} g r^3 \hat{e}_z - 6\pi \eta_{air} v_{off} \hat{e}_z = 0$
- $\frac{4}{3}\pi (\rho_{oil} - \rho_{air} g) r^3 = 6\eta_{air} v_{off} r$
- $r = \sqrt{\frac{6\eta_{air} v_{off}}{2(\rho_{oil} - \rho_{air})g}}$

eqn 4.3 Milikan's Oil Drop Charge

apply voltage V_0 such that drop moves upwards to a terminal velocity v_{on}

- $\vec{f}_e + \vec{f}_v' + \vec{f}_b + \vec{f}_e = 0$
- $f_g \hat{e}_z + f_v' \hat{e}_z - f_b \hat{e}_z - f_e \hat{e}_z = 0$
- $f_e = f_v' + (f_g - f_b) = f_v' + f_v$, where f_v = viscous force from off case
- $\frac{qV_0}{d} = 6\pi\eta_{air} r v_{on} + 6\pi\eta_{air} v_{off} = 6\pi\eta_{air} r (v_{on} + v_{off})$

- Leading to:

- $$q = \frac{g\sqrt{2}(V_{on} + V_{off}d)}{V_0} \cdot \sqrt{\frac{\eta_{air}^3 V_{off}}{(\rho_{oil} - \rho_{air})g}}$$

6.5 Maths Methods II