Critical Rate Determination in a Dynamical Glaciation Model

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Governing Equations

Our glaciation model is presented in [1], given as:

$$C\dot{T} = \frac{S(t)}{4}(1-\widehat{\alpha}(T)) - \frac{S_0}{4}(1-\alpha_0) - a(T-T_0) + b\log\left[\frac{P}{P_0}\right],\tag{1}$$

$$\dot{P} = V - W(T) \exp\left[k(T - T_0)\right] \left(\frac{P}{P_0}\right)^{\beta}.$$
 (2)

where:

$$\widehat{\alpha}(T) = \alpha_c - (\alpha_c - \alpha_w) \left(\frac{1}{2} + \frac{1}{\pi} \arctan\left[\frac{T - T_i}{\gamma} \right] \right), \quad (3)$$

$$W(T) = W_w \left(\frac{1}{2} + \frac{1}{\pi} \arctan\left[\frac{T - T_i}{\gamma} \right] \right), \quad (4)$$

$$S(t) = S_u - \frac{1}{2}(S_u - S_l)(1 + \tanh[t]).$$
 (5)

Non-dimensionalization

We follow [2] and first non-dimensionalize time:

$$au = \epsilon t$$
 where $\epsilon = \frac{V}{P_0}$. (6)

We then introduce two non-dimensionalized variables for the temperature (fast) and carbon-silicate (slow) equations:

$$x = \frac{T}{T_0}$$
 and $y = \log \left| \frac{P}{P_0} \right|$. (7)

Non-autonomous Dimensionless System

Using this non-dimensionalization, we arrive at the following system:

$$\delta \frac{\mathrm{d}x}{\mathrm{d}\tau} = \frac{\lambda(\tau)}{4} \left(1 - \alpha(x) \right) - \frac{1}{4} \left(1 - \alpha_0 \right) - a(x - 1) + by \qquad (8)$$

$$\frac{\mathrm{d}y}{\mathrm{d}\tau} = \left(1 - w(x) e^{k(x - 1) + \beta y} \right) e^{-y}. \qquad (9)$$

where:

$$\alpha(x) = \alpha_c - (\alpha_c - \alpha_w) \left(\frac{1}{2} + \frac{1}{\pi} \arctan \left[\frac{xT_0 - T_i}{\gamma} \right] \right)$$
(10)
$$w(x) = w_w \left(\frac{1}{2} + \frac{1}{\pi} \arctan \left[\frac{xT_0 - T_i}{\gamma} \right] \right)$$
(11)
$$\lambda(\tau) = s_u - \frac{1}{2} (s_u - s_l) (1 + \tanh \left[\frac{\tau}{\epsilon} \right]).$$
(12)

Dimensionless Parameters

Using this non-dimensionalization, we arrive at the following dimensionless parameters:

Parameter	Expression	Value
а	AT_0/S_0	\sim 0.46
Ь	B/S_0	$\sim 5.9 imes 10^{-3}$
W_W	W_w/V	~ 1.114
k	KT_0	~ 29
s_u	S_u/S_0	\sim 0.941
Sį	S_I/S_0	~ 0.931

Rate Constant

We think of $0 < \delta \ll 1$ as quantifying the ratio of the x and y time scales and defined as:

$$\delta = \frac{CVT_0}{P_0S_0} \sim 3 \times 10^{-4}.$$
 (13)

Also, it is the rate, which we define as $\hat{v}=1/\epsilon$, that achieves some critical value, \hat{v}_c , that is of interest in this work.

Rate Constant

Specifically, we note that the largest rate, at $\tau=0$, is defined for the sigmoid by:

$$\frac{\widehat{v}(s_u - s_l)}{2} = v, \tag{14}$$

from which we can derive as a first estimate:

$$v = \frac{P_0(s_u - s_l)}{2V\Delta t} \sim 7 \times 10^{-4}.$$
 (15)

Dimensionless System

Lastly, we can note that our $\lambda(\tau)$ can be viewed as a solution of a Bernoulli Equation and then define:

$$\delta s = \frac{1}{2} (s_u - s_l) \left(1 + \tanh \left[\widehat{v} \tau \right] \right), \tag{16}$$

which allows us to finally deal with the equations:

$$\delta \frac{\mathrm{d}x}{\mathrm{d}\tau} = \frac{\lambda(\tau)}{4} \left(1 - \alpha(x) \right) - \frac{1}{4} \left(1 - \alpha_0 \right) - a(x - 1) + by \tag{17}$$

$$\frac{\mathrm{d}y}{\mathrm{d}\tau} = \left(1 - w(x)e^{k(x-1) + \beta y}\right)e^{-y},\tag{18}$$

$$\frac{\mathrm{d}\lambda}{\mathrm{d}\tau} = 4\nu \left(\frac{\delta s}{(s_u - s_l)} - \frac{(\delta s)^2}{(s_u - s_l)^2} \right). \tag{19}$$

General Dynamical System

We follow [3] and define general dynamical system:

$$\delta \frac{\mathrm{d}x}{\mathrm{d}\tau} = f(x, y, \lambda(\tau), \delta)$$

$$\frac{\mathrm{d}y}{\mathrm{d}\tau} = g(x, y, \lambda(\tau), \delta),$$
(20)

$$\frac{\mathrm{d}y}{\mathrm{d}\tau} = g(x, y, \lambda(\tau), \delta),\tag{21}$$

where:

x is the fast variable

y is the slow variable

f and g are smooth functions

 δ is a small parameter, $0 < \delta \ll 1$ that defines the ratio of time scales values between x and y

 $\lambda_{min} \leq \lambda(\tau) \leq \lambda_{max}$ is a bounded function that evolves in time on the slow time scale, defined by the time scale: $\tau = \epsilon t$.

General Dynamical System

Now, we can think of the singular limit of $\delta=0$ as defining the one-dimensional critical manifold $S(\lambda)$, on which $\mathrm{d}y/\mathrm{d}\tau=g(x,y,\lambda,0)$ is constrained. We can therefore define the critical manifold by:

$$S(\lambda) : f(x, y, \lambda, 0) = 0.$$
 (22)

In our work, this will look like the function y = h(x).

General Dynamical System

By (a2) of [3], our first task is to find the one steady state which is asymptotically state and varies continuously with λ . This is found by calculating the intersection of the nullclines, $\mathrm{d}y/\mathrm{d}\tau=0$ and $\mathrm{d}x/\mathrm{d}\tau=0$, and determining which equilibrium point is stable.

General Dynamical System

Then, by (a1) of [3], we will determine the single fold $F(\lambda)$ that is tangent to the fast x-direction defined by:

$$\left. \frac{\partial f}{\partial x} \right|_{S} = 0. \tag{23}$$

and:

$$\left. \frac{\partial^2 f}{\partial x^2} \right|_S \neq 0. \tag{24}$$

This will give us values x_F , the one closest to the stable equilibrium we will choose, as is a condition for this theorem. We can then plug this back into our equation for $S(\lambda)$ to get y_F .

General Dynamical System

Lastly, we can determine the critical rate, v_c as:

$$v_c \approx \inf \left\{ v > 0 : m = \left(g \frac{\partial f}{\partial y} + v \frac{\partial f}{\partial \lambda} \frac{\mathrm{d}\lambda}{\mathrm{d}\tau} \right) \Big|_{F = (x_E, v_E)} = 0 \right\}.$$
 (25)

General Dynamical System

At this point, we are really dealing with two functions:

$$f(x_F, y_F, \lambda(\tau), 0)$$
 and $g(x_F, y_F, \lambda(\tau), 0)$, (26)

so we can think of this problem as a bifurcation problem in just the (τ, m) -plane: at which point v_c is the function defined in Eq. (26) satisfied, i.e. for the single real double root of m to come into existence.

Critical Manifold

We first find the critical manifold using (22) to find:

$$S(\lambda) : y = -\frac{1}{b} \left[\frac{\lambda(\tau)}{4} (1 - \alpha(x)) - \frac{1 - \alpha_0}{4} - a(x - 1) \right].$$
 (27)

Equilibria

We next calculate the equilibria by finding the intersection of the nullclines defined by:

x-nullcline:
$$y = -\frac{1}{b} \left[\frac{\lambda(\tau)}{4} (1 - \alpha(x)) - \frac{1 - \alpha_0}{4} - a(x - 1) \right]$$
(28)
$$\text{y-nullcline:} \quad y = \frac{1}{a} \left[k(1 - x) - \log \left[w(x) \right] \right].$$

y-nullcline:
$$y = \frac{1}{\beta} [k(1-x) - \log [w(x)]].$$
 (29)

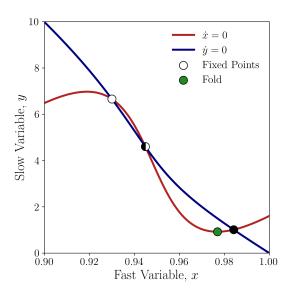
We next compute the fold following (23):

$$F(\lambda): x_f = \frac{T_i}{T_0} + \left[\frac{\gamma(\alpha_c - \alpha_w)}{4a\pi T_0} \lambda(\tau) - \left(\frac{\gamma}{T_0}\right)^2 \right]^{1/2}, \quad (30)$$

satisfying (24). Then, we find the corresponding y_F to be:

$$y_F = -\frac{1}{b} \left[\frac{\lambda(\tau)}{4} (1 - \alpha(x_F)) - \frac{1 - \alpha_0}{4} - a(x_F - 1) \right].$$
 (31)

Critical Rate Determination: Application Equilibria



Numerical Estimation of Critical Rate

We now have a minimization problem of the form:

$$v_c = \min_{v>0} \{v : m(v) = 0\}.$$
 (32)

To calculate this, we implementing a custom minimization scheme in Python that is given by the following pseudo-code:

Initialize some v > 0 sufficiently large

At some tolEp, reduce v per N many steps, N sufficiently large Per step, calculate m(v):

If m(v) < tol0, then break

Otherwise, run until this condition is satisfied.

Numerical Estimation of Critical Rate

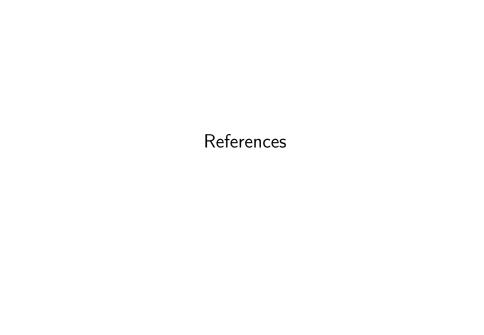
From this routine, we estimate:

$$v_c \sim 8.578 \times 10^{-4},$$
 (33)

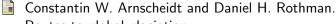
at which rate we have:

$$(x_F, y_F, \lambda_F) = (0.9766, 1.368, 0.9356).$$
 (34)

That's within an order of magnitude of the estimated $\sim 7 \times 10^{-4}...$ look's like we're getting somewhere.



References I



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