

Calculation of derivative properties in XMS-CASPT2

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Shiozaki group meeting
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Calculation of matrix elements for many electron operators



- Need a generic approach for evaluating terms of the form

$$\sum_{\substack{ijkl \\ wxyz}} \sum_J T_{ijkl}^\dagger g_{wxyz} q_{uv} \langle I | a_i a_j a_k^\dagger a_l^\dagger a_w^\dagger a_x^\dagger a_y a_z a_u^\dagger a_v | J \rangle c_J$$

- SMITH3 generates code of evaluating these integrals.
- Jae developed and implemented an algorithm such that the most costly term is

$$\sum_{ijklmn}^{act} \sum_J \langle I | a_i^\dagger a_j^\dagger a_k^\dagger a_l a_m a_n | J \rangle c_J A_{ijklmn}$$

- This is highly efficient, but unfortunately, it is not as well suited to the cases with spin-flipping interactions.

Outline of new approach



- Aim to produce routine which can calculate arbitrary BraKet:

$$\langle M | \hat{A} \hat{B} \hat{C} \dots | N \rangle$$

- Should take advantage of symmetries, and range constraints.
- The program uses the following information to construct an object associated with each operator:
 - Number of indices.
 - Position of creation and annihilation operators.
 - Range of indices.
 - List of symmetry functions and constraints.
 - If the operator couples spin and spatial components.
- A bracket object is then built from a vector of such objects.
- The program then uses this information to build a list of operations (tensor contractions and additions) which need to be performed.

Distinction from current approach



- Cannot calculate CI-derivatives for four component wavefunctions using the current approach.
- New approach should incorporate symmetry more effectively, particularly symmetries associated with spin.
- New approach repeatedly reshuffles the indexes in order to take advantage of constraints on the indices.
- Final expressions are not normal ordered, and a different algorithm is used for evaluating the CI derivatives.
- Using a generic algorithm, not generated code.

Outline of method



Step 1 : Get operator information.

- Split input tensors up into blocks.
- Generate all possible contractions of these blocks.
- Identify which blocks/contracted blocks are equivalent.

Step 2 : Manipulate indexes.

- Rearrange indexes to normal order, applying range constraints.
- Put indexes into alternating order and group common indexes.

Step 3 : Generate computational procedure.

- Merge all contributions from different brackets.
- Generate a task list to obtain contracted A-tensors.

Step 4 : Execute computational procedure.

- Contract A-tensors with gamma matrices.

Step 5 : Loop over spin sectors and states.

- Repeat for all necessary spin sectors.

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Generation of contraction task list



- Most terms are too complicated to evaluate directly, e.g.,

$$\sum_{\substack{ijkl \\ wxyz \\ uv}} \sum_J T_{ijkl}^{\dagger} g_{wxyz} q_{uv} \langle I | a_i a_j a_k^{\dagger} a_l^{\dagger} a_w^{\dagger} a_x^{\dagger} a_y a_z a_u^{\dagger} a_v | J \rangle c_J$$

- Reordering indexes to represent as sum of contractions:

$$\begin{aligned} \rightarrow \sum_{\substack{ijkl \\ wxyz \\ uv}} \gamma_{ijklwxyzuv}^I A_{ijklwxyzuv} &+ \sum_{abcdefgh} \gamma_{abcdefgh}^I A_{abcdefgh} \\ &+ \sum_{abcdef} \gamma_{abcdef}^I A_{abcdef} + \sum_{ab} \gamma_{ab}^I A_{ab} \end{aligned}$$

Generation of contraction task list



- Each A-tensor is defined by the original operators, T_{ijkl}^\dagger , g_{wxyz} , q_{uv} , and a set, X , of contraction indices and factors:

$$A_{abcd} = \sum_x^{x \in X} \sum_{efghmn} \delta_{ef} \delta_{gh} \delta_{mn} T_{ijkl}^\dagger g_{wxyz} q_{uv}$$

$$X = \{(\{(a_1, b_1), (c_1, d_1), (e_1, f_1)\}, s_1), (\{(a_2, b_2), (c_2, d_2), (e_2, f_2)\}, s_2), \dots\}$$

- Each contribution, A^x , to the A-tensor is obtained by performing a recursive sequence of operations:

$$A_{abcd}^x = B(A_{abcdefgh}^x, (g, h)) =$$

$$= B(B(A_{abcdefgh}^x, (g, h)), (e, f)) = B(B(B(A_{abcdefgh}^x, (g, h)), (e, f)), (c, d))$$

Generation of contraction task list



- The sequence of the operations is the task list.
- How the task list is built can significantly influence the computational cost.
- Currently using the following tactics to reduce cost:
 - Standardize order in which contractions are performed (facilitates term reuse).
 - Reorder indexes so that similar indexes are together (facilitates term reuse).
 - Perform all intra-tensor contractions before any inter tensor contractions (reduces maximum dimension of tensor to be contracted).
 - Never store product tensors.

Contraction tree



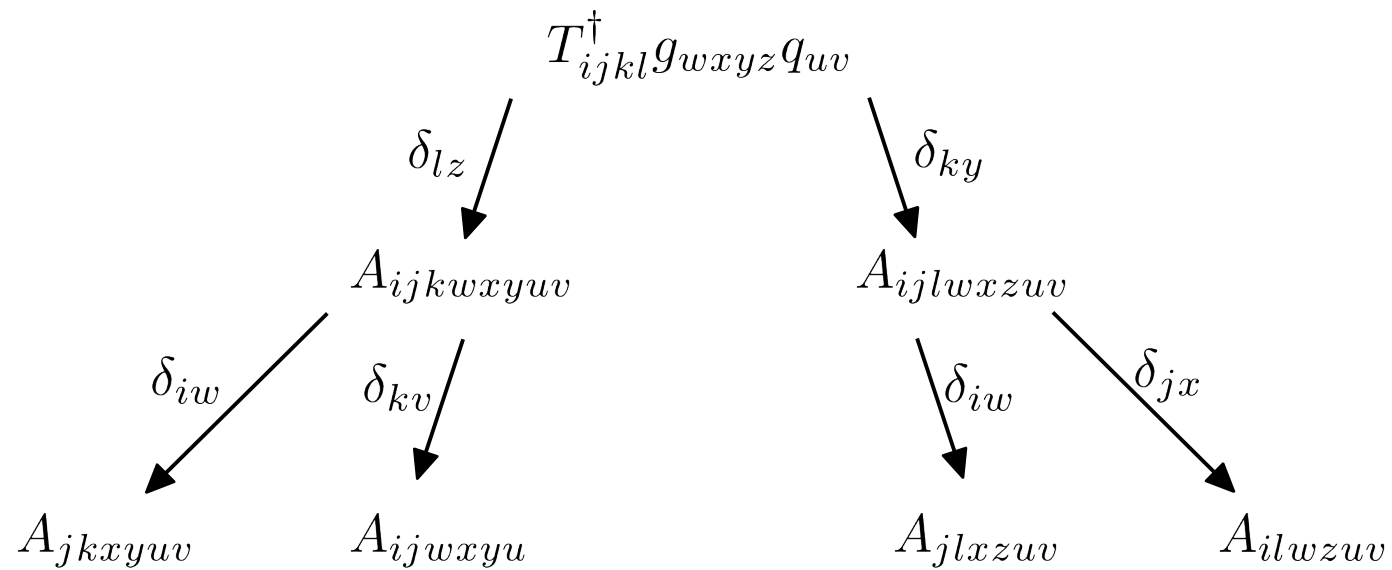
$$T_{ijkl}^{\dagger} g_{wxyz} q_{uv}$$

Contraction tree

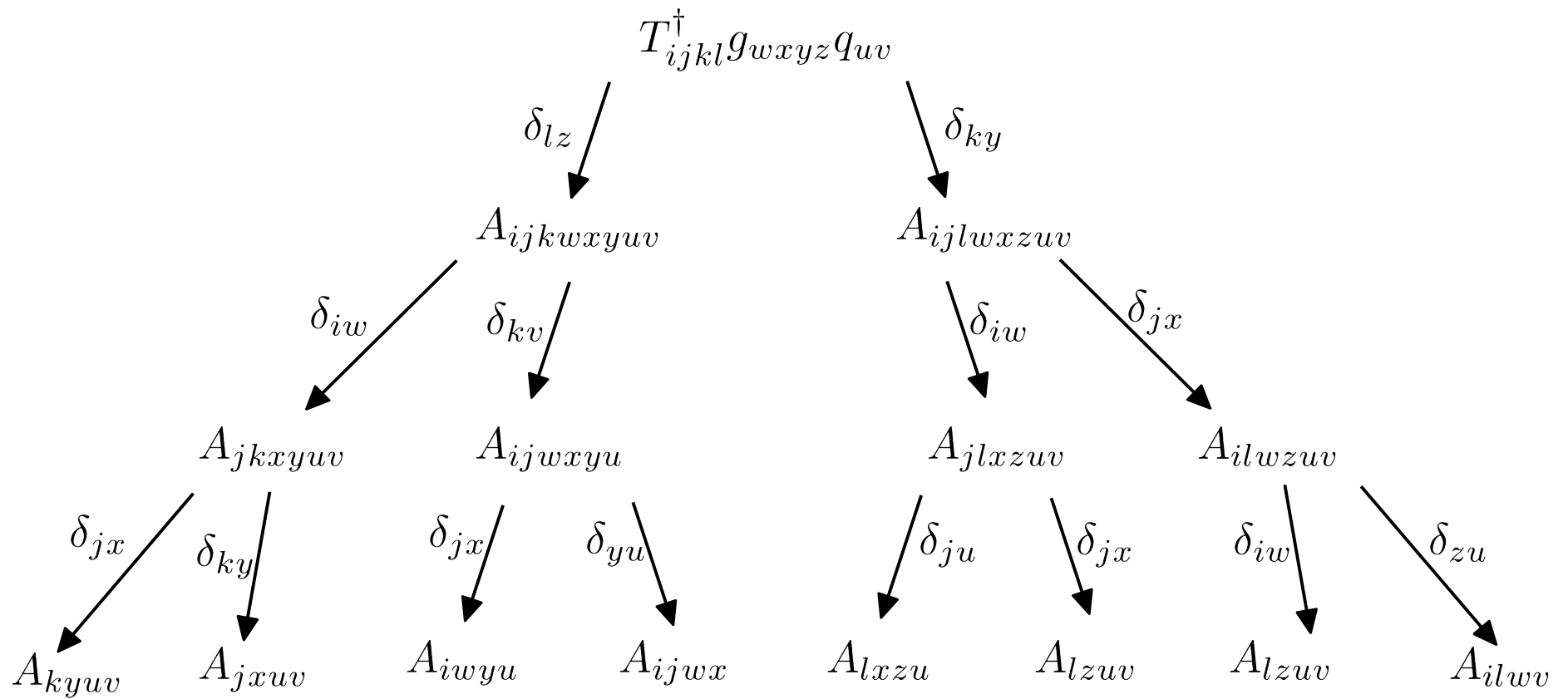


$$\begin{array}{ccc} & T_{ijkl}^{\dagger} g_{wxyz} q_{uv} & \\ \delta_{lz} \swarrow & & \searrow \delta_{ky} \\ A_{ijkwxyuv} & & A_{ijlwxyzuv} \end{array}$$

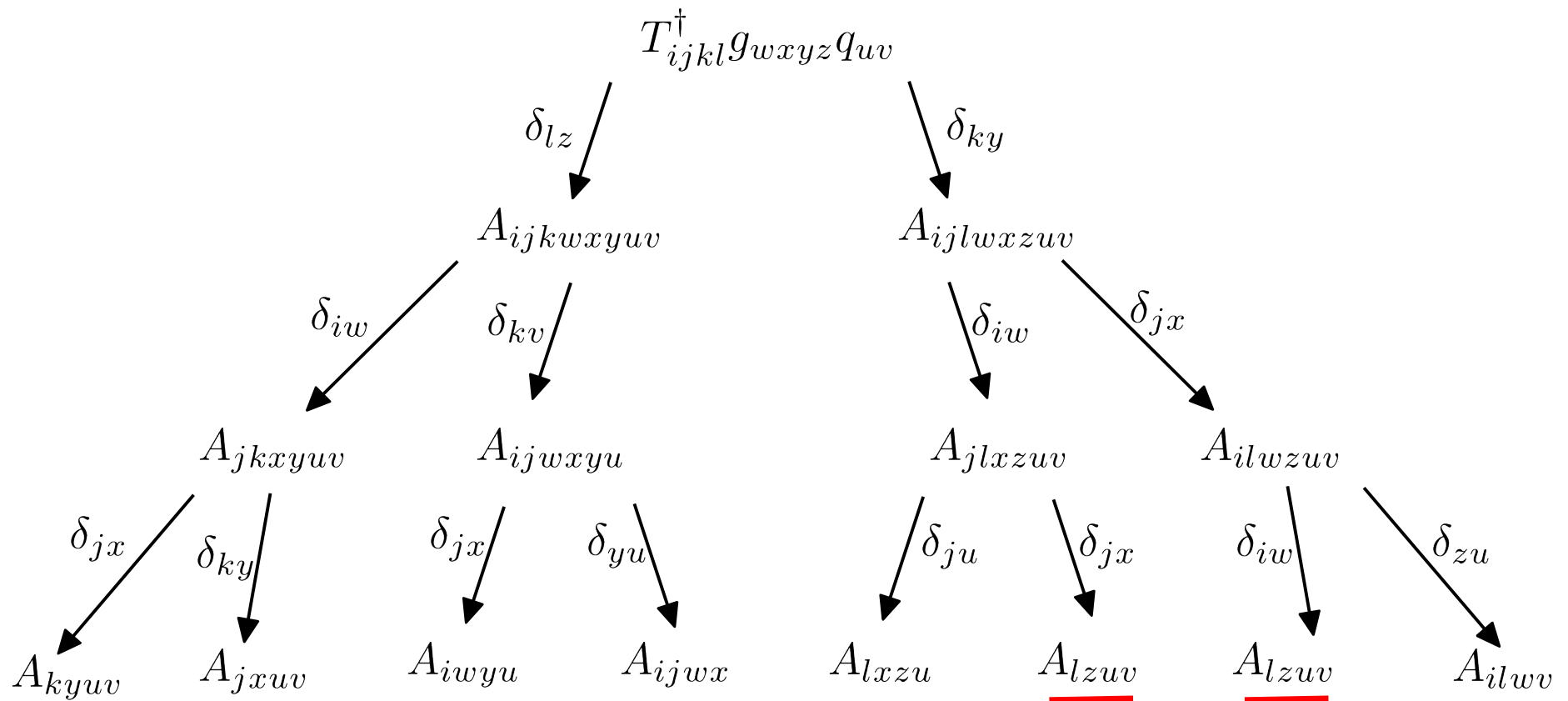
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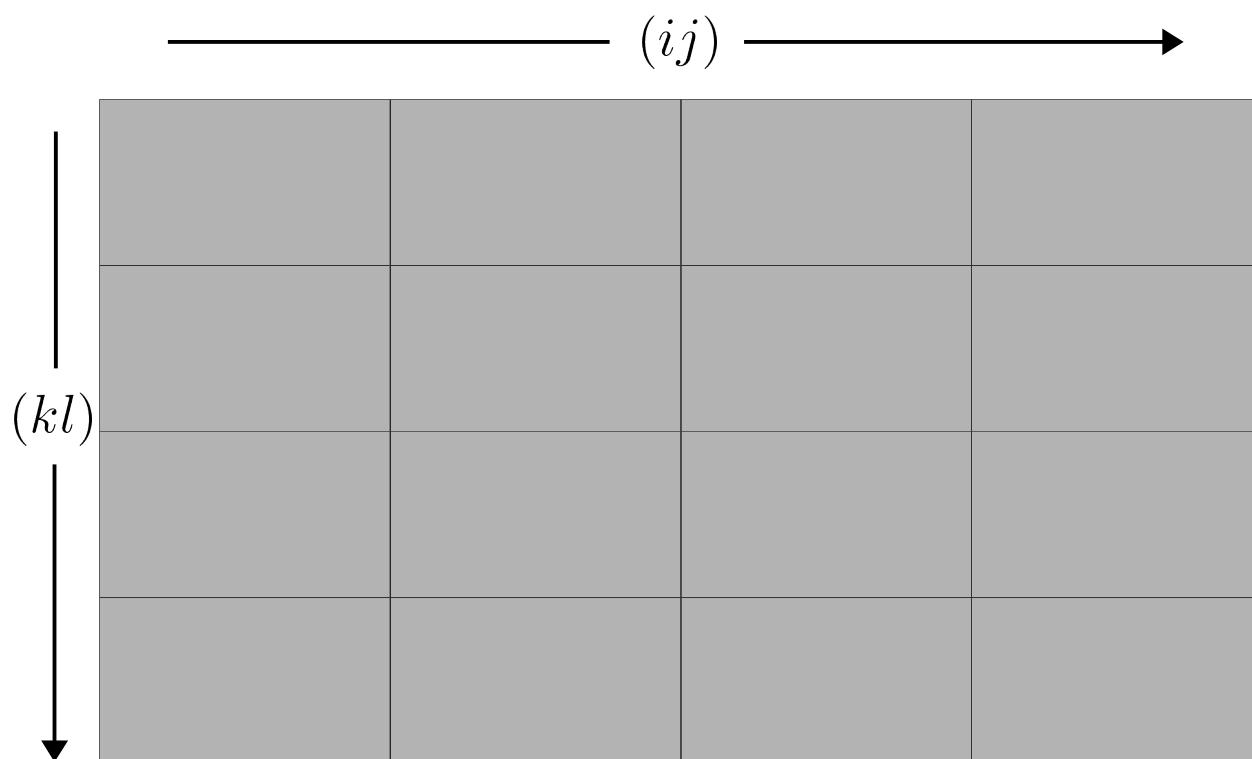
$$X = \{(\{(a_1, b_1), (c_1, d_1), (e_1, f_1)\}, s_1), (\{(a_2, b_2), (c_2, d_2), (e_2, f_2)\}, s_2), \dots\}$$

- Each path from top to bottom corresponds to a member of X .
- Order of contractions is not important.
- Reorder the individual members of X to get the most efficient tree.

Application of block symmetry



Example : CASPT2 perturbation coefficients $T_{(ij),(kl)}$.



Possible transitions

closed \rightarrow *active*

closed \rightarrow *virtual*

active \rightarrow *active*

active \rightarrow *virtual*

Application of block symmetry



Example : CASPT2 perturbation coefficients $T_{(ij),(kl)}$.

	$c \rightarrow a$	$c \rightarrow v$	$a \rightarrow a$	$a \rightarrow v$
$c \rightarrow a$				
$c \rightarrow v$				
$a \rightarrow a$				
$a \rightarrow v$				

Possible transitions

closed \rightarrow active

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Application of symmetry



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Possible transitions

closed \rightarrow active

closed \rightarrow virtual

active \rightarrow active

active \rightarrow virtual

Application of block symmetry



Gaunt (or Breit) term:

	$\alpha \rightarrow \alpha$	$\alpha \rightarrow \beta$	$\beta \rightarrow \alpha$	$\beta \rightarrow \beta$
$\alpha \rightarrow \alpha$				
$\alpha \rightarrow \beta$				
$\beta \rightarrow \alpha$				
$\beta \rightarrow \beta$				

Possible transitions

$$\alpha \rightarrow \alpha$$

$$\alpha \rightarrow \beta$$

$$\beta \rightarrow \alpha$$

$$\beta \rightarrow \beta$$

Application of block symmetry



Example : Gaunt (or Breit) term: $H_{(ij),(kl)}^{Gaunt}$

	$\alpha \rightarrow \alpha$	$\alpha \rightarrow \beta$	$\beta \rightarrow \alpha$	$\beta \rightarrow \beta$
$\alpha \rightarrow \alpha$				
$\alpha \rightarrow \beta$				
$\beta \rightarrow \alpha$				
$\beta \rightarrow \beta$				

Possible transitions

$$\alpha \rightarrow \alpha$$

$$\alpha \rightarrow \beta$$

$$\beta \rightarrow \alpha$$

$$\beta \rightarrow \beta$$

Contraction of tensor blocks



Contract over all possible pairs of indexes, e.g.,

$$\hat{H}\hat{T} \rightarrow \sum_{wxyz} \sum_{ijkl} H_{wxyz} a_w^\dagger a_x^\dagger a_y a_z T_{ijkl} a_i^\dagger a_j^\dagger a_k a_l$$

$$A_{abcdef}^{HT,st} = \sum_{wxyz} \sum_{ijkl} H_{wxyz} T_{ijkl} \delta_{st}$$

$$s := w, x, i, \text{ or } j$$

$$t := y, z, k \text{ or } l$$

Use the block range constraints to rule out possible contractions, e.g.

$$A_{abcdef}^{HT,st} = \begin{cases} A_{abcdef}^{HT,st} & \text{if } rng(s) = rng(t) \\ 0 & \text{otherwise} \end{cases}$$

Contraction list is different for every range block.

Contraction of tensor blocks



- Distinct range blocks may become equivalent following contraction, e.g.

$$A_{\text{wxyzjkl}}^{HT} = \sum_{\text{wxyz}} \sum_{\text{ijkl}} H_{\text{wxyz}} T_{\text{ijkl}} \delta_{\text{zi}} \quad \text{Active}$$

$$A_{\text{wxyzjkl}}^{HT} = \sum_{\text{wxyz}} \sum_{\text{ijkl}} H_{\text{wxyz}} T_{\text{ijkl}} \delta_{\text{zi}} \quad \text{Core}$$

$$A_{\text{wxyzjkl}}^{HT} = \sum_{\text{wxyz}} \sum_{\text{ijkl}} H_{\text{wxyz}} T_{\text{ijkl}} \delta_{\text{zi}} \quad \text{Virtual}$$

- Following the contraction the range block is the same.
- Important to combine contraction lists following this; otherwise can potentially increase the complexity.

Generation of contraction list



Step 1 : First take creation and annihilation operators in order specified by input operator, e.g.,

$$\langle I | a_l a_i^\dagger a_m a_j^\dagger a_n a_k^\dagger | J \rangle$$

Step 2 : Rewrite in normal order:

$$\begin{aligned} = & \langle I | a_i^\dagger a_j^\dagger a_k^\dagger a_l a_m a_n | J \rangle + (s_{m'n'} \delta_{m'n'} + \dots) \langle I | a_i^\dagger a_j^\dagger a_{k'} a_{l'} | J \rangle \\ & + (s_{k'l'm'n'} \delta_{m'n'} \delta_{k'l'} + \dots) \langle I | a_i^\dagger a_j^\dagger | J \rangle + \dots \end{aligned}$$

Step 3 : Remove terms which known to be zero from index ranges.

Step 4 : Repeat process, but reorder to alternating order

$$\begin{aligned} \langle I | a_i^\dagger a_k^\dagger a_j^\dagger a_l a_m a_n | J \rangle & \rightarrow \langle I | a_i^\dagger a_l a_j^\dagger a_m a_k^\dagger a_n | J \rangle \\ & = \gamma_{ijklmn}^{IJ} \end{aligned}$$

Grouping of indexes



- Takes too long in relativistic framework;

$$N_{blocks} \approx N_{range}^{N_{id}} = 6^{10}$$

- Furthermore, applying spin symmetry was difficult.
- Rewrote contraction list generation routines to avoid this problem.
- Gamma generation routine now takes range constraints.
- Possible to immediately rule out majority of ranges.

Grouping of indexes



- Introduced grouping of indexes to reduce number of distinct terms, e.g.,

$$\langle I | a_{\mathbf{i}}^{\dagger} a_{\mathbf{j}} a_{\mathbf{k}}^{\dagger} a_{\mathbf{l}} a_{\mathbf{m}}^{\dagger} a_{\mathbf{n}} a_{\mathbf{o}}^{\dagger} a_{\mathbf{p}} | J \rangle c_J \rightarrow \langle I | a_{\mathbf{i}}^{\dagger} a_{\mathbf{n}} a_{\mathbf{o}}^{\dagger} a_{\mathbf{l}} a_{\mathbf{m}}^{\dagger} a_{\mathbf{j}} a_{\mathbf{k}}^{\dagger} a_{\mathbf{p}} | J \rangle c_J$$

$$\gamma_{\mathbf{ij}} \gamma_{\mathbf{kl}} \gamma_{\mathbf{mn}} \gamma_{\mathbf{op}} \rightarrow \gamma_{\mathbf{in}} \gamma_{\mathbf{ol}} \gamma_{\mathbf{mj}} \gamma_{\mathbf{kp}}$$

- Colours denote different ranges.
- Immediately useful for spin, but can be used in other contexts, e.g., point symmetry.

Contraction of γ_{ijklmn}^I and A_{ijklmn}



Perform contraction in two steps:

$$\sum_{\sigma_3} \sum_{ijklmn} \gamma_{ij,kl,mn}^{IJ\sigma_3} A_{ij,kl,mn}^{\sigma_3} c_J = \sum_{\sigma_1} \sum_{ij} \sum_K \gamma_{ij}^{IK\sigma_1} \tilde{\gamma}_{ij}^{K\sigma_1}$$

- Where

$$\tilde{\gamma}_{ij}^{K\sigma_1} = \sum_{\sigma_2} \sum_{LJ} \sum_{klmn} \gamma_{kl,mn}^{KJ\sigma_2} A_{ij,kl,mn}'^{\sigma_1 \otimes \sigma_2}$$

and

$$\sigma_3 = \sigma_1 \cup \sigma_1 \cup \sigma_1 = \{s_1 s_2 s_3 s_4 s_5 s_6\}$$

- This step should be comparable in speed to the most expensive step of the current algorithm.

Contraction of γ_{ijklmn}^I and A_{ijklmn}



Contract with A-tensor whilst generating the gamma matrices:

$$\begin{aligned}
 & \sum_{ijklmn} \gamma_{ij,kl,mn}^{IJ} A_{ij,kl,mn} c_J \\
 &= \sum_{ij} \sum_K \langle I | a_i^\dagger a_j | K \rangle \sum_{kl} \sum_L \langle K | a_k^\dagger a_l | L \rangle \sum_{mn} \sum_J \langle L | a_m^\dagger a_n | J \rangle A_{ij,kl,mn} c_J \\
 &= \sum_{ij} \sum_K \langle I | a_i^\dagger a_j | K \rangle \sum_{kl} \sum_L \langle K | a_k^\dagger a_l | L \rangle \tilde{A}_{ij,kl}^L \\
 &= \sum_{ij} \sum_K \langle I | a_i^\dagger a_j | K \rangle \tilde{A}_{ij}^K
 \end{aligned}$$

- Minimizes the number of gamma matrices which need to be generated and stored.