

# Chapter 1

# Chapter 2

# Chapter 3

- Stellar parallax:

$$d = \frac{1 \text{ AU}}{\tan p} \approx \frac{1}{p} \text{ AU} \quad (1)$$

- parsec: distance when parallax angle is 1 arcsecond.
- Hipparcos (ESA): accuracies approaching  $0.001''$
- Gaia: accuracies approaching 4 microarcsec

- apparent magnitude:

- difference of 5 magnitudes = factor of 100 difference in brightness  $\Rightarrow 100^{1/5} = 2.512$
- Solar constant:  $1.4 \text{ kW m}^{-2}$
- 

$$\frac{F_2}{F_1} = 100^{(m_1 - m_2)/5} \Rightarrow m_1 - m_2 = -2.5 \log_{10} \left( \frac{F_1}{F_2} \right) \quad (2)$$

- absolute magnitude:

- apparent magnitude at 10 kpc
- distance modulus

$$100^{(m-M)/5} = \frac{F_{10}}{F} = \left( \frac{d}{10 \text{ pc}} \right)^2 \Rightarrow d = 10^{(m-M+5)/5} \text{ pc}$$

$$\Rightarrow m - M = 5 \log_{10}(d) - 5 = 5 \log_{10} \left( \frac{d}{10 \text{ pc}} \right)$$

- Double slit experiment:

- spacing  $d$  between slits, distance  $L$  to screen. At angle  $\theta$ , extra path length of  $d \sin \theta$  for far slit  $\Rightarrow$  bright if  $d \sin \theta = n\lambda$ , dark if  $d \sin \theta = (n - \frac{1}{2})\lambda$ .

- Poynting vector:

$$\mathbf{S} = \frac{1}{\mu_0} \mathbf{E} \times \mathbf{B} \text{ (SI)} = \frac{c}{4\pi} \mathbf{E} \times \mathbf{B} \text{ (cgs)} \quad (3)$$

- radiation pressure:

$$F_{\text{rad}} = \frac{\langle S \rangle A}{c} \cos \theta \quad (\text{absorption})$$

$$F_{\text{rad}} = \frac{2\langle S \rangle A}{c} \cos^2 \theta \quad (\text{reflection})$$

- Blackbodies:

- Wien's displacement law:

$$\lambda_{\max} T = 0.002897755 \text{ m K} \quad (4)$$

- Stefan-Boltzmann equation:

$$L = A\sigma T^4, \quad \sigma = 5.67 \times 10^{-8} \text{ W m}^{-2}\text{K}^{-4}. \quad (5)$$

- Rayleigh-Jeans law: works for  $\lambda$  large. Take oven with distance  $L$  between walls, then can have standing waves with wavelength  $\lambda = 2L/1, 2L/2, 2L/3, 2L/4, 2L/5, \dots$ , each wavelength with energy  $kT$ . Then (DERIVATION — LEARN)

$$B_\lambda(T) \approx \frac{2ckT}{\lambda^4} \quad (6)$$

- Wien's approximation: works for  $\lambda$  small. (DERIVATION — LEARN)

$$B_\lambda(T) \approx a\lambda^{-5}e^{-b/\lambda T} \quad (7)$$

- Planck:

$$B_\lambda(T) = \frac{a/\lambda^5}{e^{b/kT} - 1} \quad (8)$$

Assume energy must be quantized:  $nh\nu$ . (DERIVATION — LEARN)

$$B_\lambda(T) = \frac{2hc^2/\lambda^5}{e^{hc/\lambda kT} - 1} \quad (9)$$

$\lambda = c/\nu$ , so  $d\lambda = c/\nu^2 d\nu$ :

$$B_\nu(T) = \frac{2h\nu^3/c^2}{e^{h\nu/kT} - 1} \quad (10)$$

- Color index

$$U - B = M_U - M_B$$

$$B - V = M_B - M_V$$

- Bolometric correction:

$$BC = m_{\text{bol}} - V = M_{\text{bol}} - M_V \quad (11)$$

- Sensitivity function  $S(\lambda)$

$$U = -2.5 \log_{10} \left( \int_0^\infty F_\lambda S_U d\lambda \right) + C_U \quad (12)$$

$C_U$  chosen such that Vega has magnitude zero in each filter.

$$U - B = -2.5 \log_{10} \left( \frac{\int F_\lambda S_U d\lambda}{\int F_\lambda S_B d\lambda} \right) + C_{U-B}, \quad C_{U-B} \equiv C_U - C_B \quad (13)$$

- \*\*\* Color index does not depend on distance, so it is a measure solely of the temperature of a model blackbody star

**Chapter 4**

**Chapter 5**

**Chapter 6**

**Chapter 7**

**Chapter 8**

**Chapter 9**

**Chapter 10**