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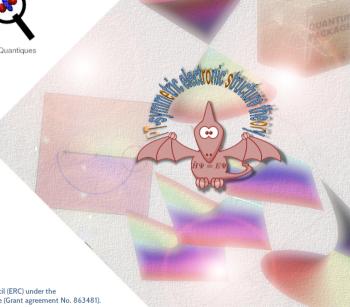
European Research Council Laboratoire de Chimie et Physique Quantiques

Cumulant Green's function methods for molecules

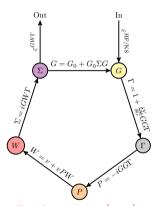
<u>Pierre-François (Titou) Loos,</u> Antoine Marie & Abdallah Ammar

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Laboratoire de Chimie et Physique Quantiques, IRSAMC, UPS/CNRS, Toulouse https://lcpq.github.io/PTEROSOR



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Hedin, Phys Rev 139 (1965) A796

Hedin's equations

$$\underline{\mathsf{G}(12)} = \mathsf{G}_0(12) + \int \mathsf{G}_0(13) \Sigma(34) \underline{\mathsf{G}}(42) d(34)$$

Green's function

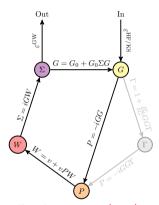
$$\underbrace{\Gamma(123)}_{\text{vertex}} = \delta(12)\delta(13) + \int \frac{\delta\Sigma(12)}{\delta G(45)} \frac{G}{G}(46)G(75)\Gamma(673)d(4567)$$

$$\underbrace{P(12)}_{} = -i \int G(13)\Gamma(342)G(41)d(34)$$

polarizability

$$\underbrace{\mathbf{W}(12)}_{\text{screening}} = \mathbf{v}(12) + \int \mathbf{v}(13)\mathbf{P}(34)\mathbf{W}(42)d(34)$$

$$\underbrace{\Sigma(12)}_{\text{self-energy}} = i \int \mathbf{G}(14)\mathbf{W}(13)\Gamma(423)d(34)$$



Hedin, Phys Rev 139 (1965) A796

The GW approximation

$$\begin{split} \underbrace{\frac{G(12)}{\text{Green's function}}}_{\text{Green's function}} &= G_0(12) + \int G_0(13) \Sigma(34) \frac{G(42) d(34)}{G(42) d(34)} \\ \underbrace{\frac{\Gamma(123)}{\text{vertex}}}_{\text{vertex}} &= \delta(12) \delta(13) \\ \underbrace{\frac{P(12)}{\text{polarizability}}}_{\text{polarizability}} &= -i \frac{G(12) G(21)}{G(21)} \\ \underbrace{\frac{W(12)}{\text{screening}}}_{\text{screening}} &= v(12) + \int v(13) \frac{P(34) W(42) d(34)}{G(42) d(34)} \\ \underbrace{\frac{\Sigma(12)}{\text{self-energy}}}_{\text{self-energy}} &= i \frac{G(12) W(12)}{G(12)} \end{split}$$

- It's cheap and fairly accurate for IPs but it relies on error cancelation
- GW is not suitable for strongly correlated systems Ammar et al. JCP 160 (2024) 114101
- It's complicated to go beyond GW Lewis & Berkelbach JCTC 15 (2019) 2925; Mejuto-Zaera & Vlcek 106 (2022) 165129
- ► It's (extremely) bad for satellite transitions Marie & Loos JCTC 20 (2024) 4751

Cumulant ansatz in the time domain

One-body Green's function
$$G(t) = G_0(t) e^{\frac{Cumulant}{C(t)}}$$
Reference Green's function

- Cumulant is a cheap beyond-GW scheme rooted to electron-boson models Langreth PRB 1 (1970) 471; Hedin JPCM 11 (1999) 489
- It's been very successful for materials to describe satellite features Aryasetiawan et al. PRL 77 (1996) 2268; Guzzo et al. PRL 107 (2011) 166401
- How does it perform for molecules?
 Vlcek et al. PRM 2 (2018) 030801; McClain et al. PRB 93 (2016) 235139

Landau form of the cumulant

$$C_{pp}(t) = \int d\omega \frac{\beta_p(\omega + \epsilon_p^{\rm HF})}{\omega^2} \left(\begin{array}{c} e^{-i\omega t} \\ \end{array} + \begin{array}{c} \text{Quasiparticle shift} \\ \text{i}\omega t \\ \end{array} \right)$$
Renormalization

$$eta_p(\omega) = -rac{1}{\pi} \operatorname{Im} \left[\sum_{pp}^{\mathsf{c}}(\omega) \right]$$

Correlation part of the self-energy

Main outcomes

- ► GW+C sometimes improves upon GW but far from being systematic
- Cumulant estimates satellite energies without solving dynamical equations
- Beyond-GW schemes describing satellites accurately would be useful!

$$\begin{split} \left[\boldsymbol{\epsilon} + \boldsymbol{\Sigma}^{\text{GW}} \left(\boldsymbol{\omega} = \boldsymbol{\epsilon}_{p}^{\text{GW}}\right)\right] \psi_{p}^{\text{GW}} &= \boldsymbol{\epsilon}_{p}^{\text{GW}} \psi_{p}^{\text{GW}} \\ \boldsymbol{\Sigma}^{\text{GW}} (\boldsymbol{\omega}) &= \boldsymbol{V}^{\text{2h1p}} \cdot \left(\boldsymbol{\omega} \mathbf{1} - \boldsymbol{C}^{\text{2h1p}}\right)^{-1} \cdot (\boldsymbol{V}^{\text{2h1p}})^{\dagger} \\ &+ \boldsymbol{V}^{\text{2p1h}} \cdot \left(\boldsymbol{\omega} \mathbf{1} - \boldsymbol{C}^{\text{2p1h}}\right)^{-1} \cdot (\boldsymbol{V}^{\text{2p1h}})^{\dagger} \end{split} \qquad \begin{split} \underbrace{\begin{array}{c} \operatorname{downfolding} \\ \operatorname{upfolding} \end{array}}_{\text{upfolding}} & \begin{cases} \boldsymbol{H} \boldsymbol{\Psi}_{p,s} &= \boldsymbol{\epsilon}_{p,s}^{\text{GW}} \boldsymbol{\Psi}_{p,s} \\ \boldsymbol{H} &= \begin{pmatrix} \boldsymbol{V}^{\text{2h1p}} & \boldsymbol{V}^{\text{2p1h}} \\ (\boldsymbol{V}^{\text{2h1p}})^{\dagger} & \boldsymbol{C}^{\text{2h1p}} & \boldsymbol{0} \\ (\boldsymbol{V}^{\text{2p1h}})^{\dagger} & \boldsymbol{0} & \boldsymbol{C}^{\text{2p1h}} \end{pmatrix} \end{split}$$

$$\mathbf{1h} \, \& \, \mathbf{1p} \, \mathbf{conf}. & \begin{cases} \boldsymbol{\epsilon} & \boldsymbol{V}^{\text{2h1p}} \, \boldsymbol{V}^{\text{2p1h}} \\ \boldsymbol{V}^{\text{2h1p}} \, \boldsymbol{C}^{\text{2h1p}} & \boldsymbol{0} \\ \boldsymbol{V}^{\text{2h1p}} \, \boldsymbol{C}^{\text{2h1p}} & \boldsymbol{0} \\ \boldsymbol{V}^{\text{2p1h}} & \boldsymbol{0} & \boldsymbol{C}^{\text{2p1h}} \end{pmatrix} \end{array} \right] \text{ internal space } \boldsymbol{P}$$

$$\mathbf{2p1h} \, \mathbf{conf}. & \begin{cases} \boldsymbol{V}^{\text{2h1p}} \, \boldsymbol{V}^{\text{2p1h}} & \boldsymbol{0} \\ \boldsymbol{V}^{\text{2p1h}} & \boldsymbol{0} & \boldsymbol{C}^{\text{2p1h}} \end{pmatrix} \end{array} \right.$$

Backhouse et al. JCTC 16 (2020) 1090; Bintrim & Berkelbach, JCP 154 (2021) 041101; Monino & Loos JCP 156 (2022) 231101