

# Parquet Theory for Molecular Systems

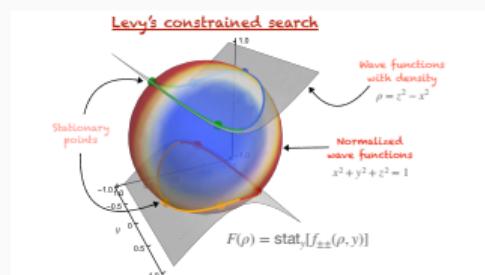
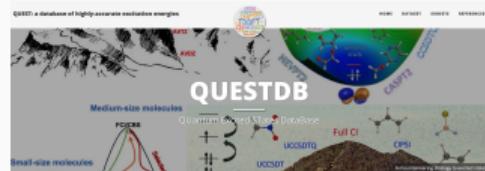
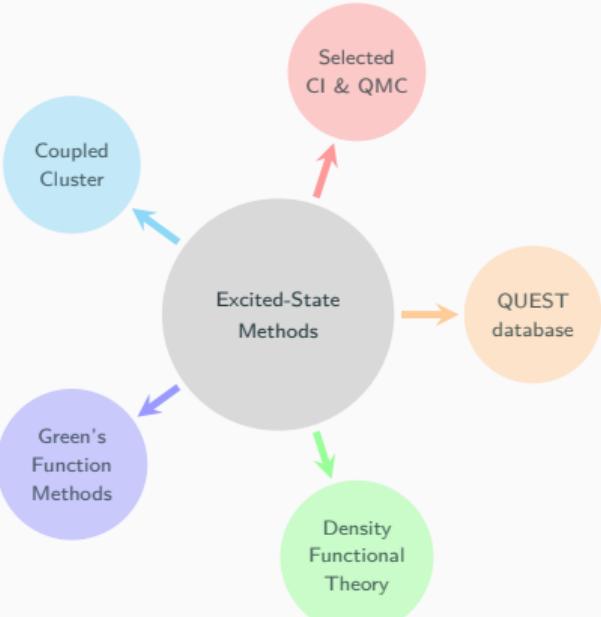
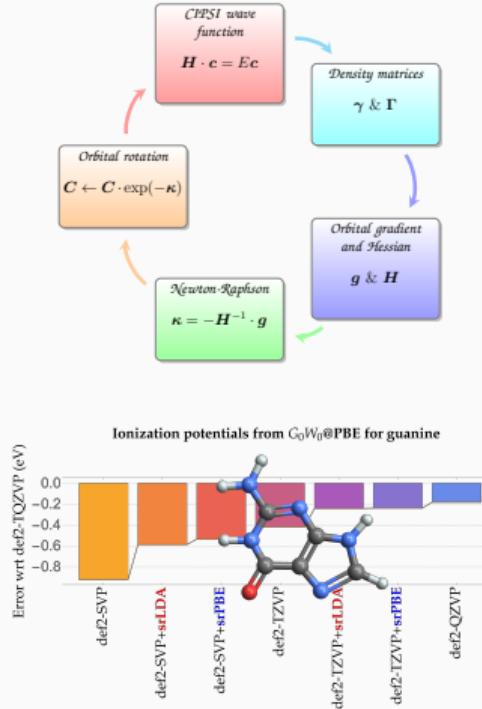
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[https://pfloos.github.io/WEB\\_LOOS](https://pfloos.github.io/WEB_LOOS)

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# General Overview of our Research Group



<https://lcpq.github.io/PTEROSOR>

# Green's Function Methods



Antoine Marie (PhD)

# The Man, The Myth, The Legend...



# Electronic Schrödinger Equation

## Wave Function Theory

$$\hat{H} \Psi(\mathbf{r}_1, \dots, \mathbf{r}_N) = E \Psi(\mathbf{r}_1, \dots, \mathbf{r}_N)$$

Hamiltonian      Energy  
                        ↓  
                        ↑  
                        Wave function

$$\hat{H} = \hat{T} + \hat{W}_{ee} + \hat{V}_{ext} \Rightarrow E = E_T + E_w + E_v$$

kinetic      external potential  
                        ↓  
                        ↑  
                        electron repulsion

## Density Functional Theory

$$N \int \cdots \int \Psi^*(\mathbf{r}, \dots, \mathbf{r}_N) \Psi(\mathbf{r}, \dots, \mathbf{r}_N) d\mathbf{r}_2 \cdots d\mathbf{r}_N = n(\mathbf{r})$$

electron density  
↓

Wave Function Theory (WFT)  $\sim$  Density Functional Theory (DFT)

$$E = E_T + E_W + E_V$$

✗ ✗ ✓

Hohenberg & Kohn, Phys. Rev. 1964 (B864) 136

### Density Matrix Functional Theory

$$N \int \cdots \int \Psi^*(\mathbf{r}, \dots, \mathbf{r}_N) \Psi(\mathbf{r}', \dots, \mathbf{r}_N) d\mathbf{r}_2 \cdots d\mathbf{r}_N = n_1(\mathbf{r}, \mathbf{r}')$$

1st-order reduced density matrix  
↓

Wave Function Theory (WFT)  $\leadsto$  Reduced Density Matrix Functional Theory (RDMF)

$$E = E_T + E_W + E_V$$

✓ ✗ ✓

Gilbert, Phys. Rev. B 12 (1975) 2111

# One-Body Green's Function: The Sweet Spot?

## One-Body Propagator in the Time Domain

$$\text{one-body Green's function} \quad \xrightarrow{\text{time-ordering}} \quad G(\mathbf{r}, \mathbf{r}'; t - t') = -i \left\langle \Psi_0^N \right| \hat{T} \begin{bmatrix} \hat{\psi}(\mathbf{r}t) & \hat{\psi}^\dagger(\mathbf{r}'t') \end{bmatrix} \left| \Psi_0^N \right\rangle$$

Field operators

N-electron ground state

$$G(\mathbf{r}, \mathbf{r}'; t - t') = \begin{cases} -i \langle \Psi_0^N | \hat{\psi}(\mathbf{r}t) \hat{\psi}^\dagger(\mathbf{r}'t') | \Psi_0^N \rangle & \text{for } t > t' \\ +i \langle \Psi_0^N | \hat{\psi}^\dagger(\mathbf{r}'t') \hat{\psi}(\mathbf{r}t) | \Psi_0^N \rangle & \text{for } t' < t \end{cases}$$

- $\langle \Psi_0^N | \hat{\psi}(\mathbf{r}t) \hat{\psi}^\dagger(\mathbf{r}'t') | \Psi_0^N \rangle$  measures the propagation of an **electron** (electron branch)
- $\langle \Psi_0^N | \hat{\psi}^\dagger(\mathbf{r}'t') \hat{\psi}(\mathbf{r}t) | \Psi_0^N \rangle$  measures the propagation of a **hole** (hole branch)

Martin, Reining & Ceperley, "Interacting Electrons"

## One-Body Propagator in the Frequency Domain

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\nu} \frac{\mathcal{I}_{\nu}(\mathbf{r}) \mathcal{I}_{\nu}^*(\mathbf{r}')}{\omega - (E_0^N - E_{\nu}^{N-1}) - i\eta} + \sum_{\nu} \frac{\mathcal{A}_{\nu}(\mathbf{r}) \mathcal{A}_{\nu}^*(\mathbf{r}')}{\omega - (E_{\nu}^{N+1} - E_0^N) + i\eta}$$

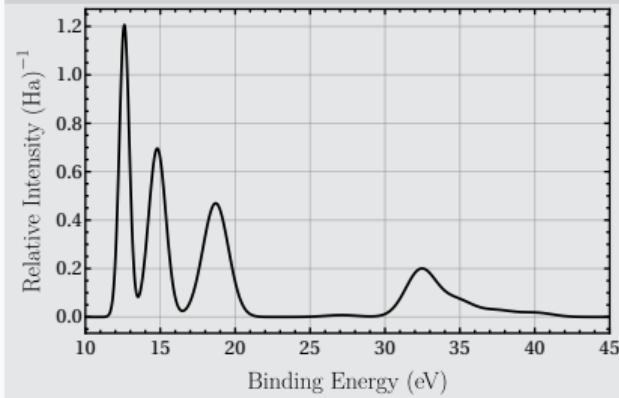
$\xrightarrow{\nu\text{th ionization potential (IP)}}$        $\xrightarrow{\nu\text{th electron affinity (EA)}}$

## Spectral function

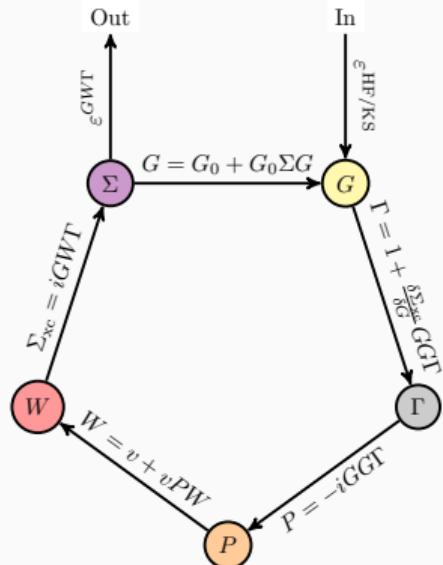
$$A(\omega) = \frac{1}{\pi} \int d\mathbf{r} d\mathbf{r}' |\text{Im } G(\mathbf{r}, \mathbf{r}'; \omega)|$$

Marie & Loos, JCTC 20 (2024) 4751

## Photoemission spectrum of water



# Hedin's Pentagon



Hedin, Phys. Rev. 139 (1965) A796

## Hedin's Equations

$$\underbrace{G(12)}_{\text{Green's function}} = G_0(12) + \int G_0(13)\Sigma(34)G(42)d(34)$$

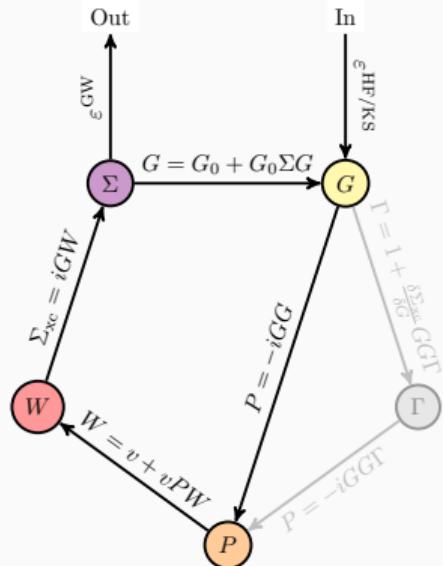
$$\underbrace{\Gamma(123)}_{\text{vertex}} = \delta(12)\delta(13) + \int \frac{\delta\Sigma_{xc}(12)}{\delta G(45)} G(46)G(75)\Gamma(673)d(4567)$$

$$\underbrace{P(12)}_{\text{polarizability}} = -i \int G(13)\Gamma(342)G(41)d(34)$$

$$\underbrace{W(12)}_{\text{screening}} = v(12) + \int v(13)P(34)W(42)d(34)$$

$$\underbrace{\Sigma_{xc}(12)}_{\text{self-energy}} = i \int G(14)W(13)\Gamma(423)d(34)$$

# Hedin's Square



Hedin, Phys. Rev. 139 (1965) A796

## The $GW$ Approximation

$$\underbrace{G(12)}_{\text{Green's function}} = G_0(12) + \int G_0(13)\Sigma(34)\mathcal{G}(42)d(34)$$

$$\underbrace{\Gamma(123)}_{\text{vertex}} = \delta(12)\delta(13)$$

$$\underbrace{P(12)}_{\text{polarizability}} = -iG(12)\mathcal{G}(21)$$

$$\underbrace{W(12)}_{\text{screening}} = v(12) + \int v(13)\mathcal{P}(34)\mathcal{W}(42)d(34)$$

$$\underbrace{\Sigma_{xc}(12)}_{\text{self-energy}} = iG(12)\mathcal{W}(12)$$

# Self-Energy

## Self-Energy as a Function of the Bare Coulomb Operator

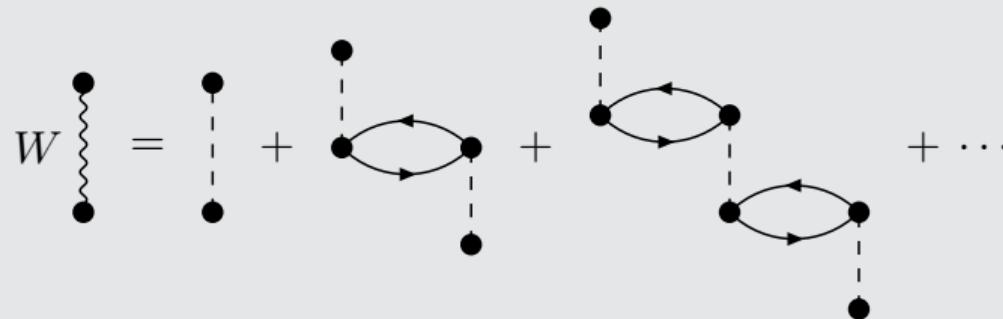
$$\Sigma(11') = \underbrace{-i\bar{v}(12; 1'2') G(2'2)}_{\text{first-order terms}} + \underbrace{\frac{1}{2} \bar{v}(12; 3'2') G(3'3) G(4'2) G(2'4) \bar{v}(34; 1'4')}_{\text{second-order terms}} + \dots$$

## Diagrammatic Representation



# Hedin's Equations

## GW Approximation

$$W \text{ (wavy line)} = \text{ (dashed line)} + \text{ (dashed line with loop)} + \text{ (dashed line with two loops)} + \dots$$


Hedin, Phys. Rev. 139 (1965) A796

## pp T-matrix Approximation

$$T \text{ (square with } T\text{)} = \text{ (dashed line)} + \text{ (dashed line with two arrows)} + \text{ (dashed line with three arrows)} + \dots$$


Marie, Romaniello & Loos, PRB 110 (2024) 115155

# How to Compute $G$ ?

## The Dyson Equation

$$G(11') = G_0(11') + \int d(22') G_0(12) \Sigma(22') G(2'1')$$

mean-field propagator

one-body Green's function

self-energy

$$G^{-1}(11') = G_0^{-1}(11') - \Sigma(11')$$

## Quasi-Particle Equation

$$\left[ H_0 + \Sigma(\omega = \epsilon_p) \right] \psi_p(x) = \epsilon_p \psi_p(x),$$

mean-field Hamiltonian

self-energy

Dyson orbitals

poles of the Green's function

## Two-Body Green's Function

### Two-Body Propagator in the Time Domain

two-body Green's function

$$G_2(12; 1'2') = (-i)^2 \left\langle \Psi_0^N \right| \hat{T} \left[ \hat{\psi}(2) \hat{\psi}^\dagger(2') \hat{\psi}(1) \hat{\psi}^\dagger(1') \right] \left| \Psi_0^N \right\rangle$$

### Propagation of electron-hole pairs ( $t_{1'} > t_1$ and $t_{2'} > t_2$ )

$$G_2^{eh}(12; 1'2') = (-i)^2 \left\langle \Psi_0^N \right| \hat{\psi}^\dagger(1') \hat{\psi}(1) \hat{\psi}^\dagger(2') \hat{\psi}(2) + \hat{\psi}^\dagger(2') \hat{\psi}(2) \hat{\psi}^\dagger(1') \hat{\psi}(1) \left| \Psi_0^N \right\rangle$$

### Propagation of electron-electron and hole-hole pairs ( $t_{1'} > t_{2'}$ and $t_1 > t_2$ )

$$G_2^{ee}(12; 1'2') = (-i)^2 \left\langle \Psi_0^N \right| \hat{\psi}(1) \hat{\psi}(2) \hat{\psi}^\dagger(1') \hat{\psi}^\dagger(2') \left| \Psi_0^N \right\rangle$$

$$G_2^{hh}(12; 1'2') = (-i)^2 \left\langle \Psi_0^N \right| \hat{\psi}^\dagger(1') \hat{\psi}^\dagger(2') \hat{\psi}(1) \hat{\psi}(2) \left| \Psi_0^N \right\rangle$$

# The Electron-Hole Channel

## Electron-Hole Correlation Function

eh correlation function

$$L(12; 1'2') = -G_2(12; 1'2') + G(11')G(22')$$
$$L(\mathbf{r}_1 \mathbf{r}_2; \mathbf{r}_1' \mathbf{r}_2'; \omega) = \sum_{\nu > 0} \frac{L_\nu^N(\mathbf{r}_2 \mathbf{r}_{2'}) R_\nu^N(\mathbf{r}_1 \mathbf{r}_{1'})}{\omega - (E_\nu^N - E_0^N - i\eta)} - \sum_{\nu > 0} \frac{L_\nu^N(\mathbf{r}_2 \mathbf{r}_{2'}) R_\nu^N(\mathbf{r}_1 \mathbf{r}_{1'})}{\omega - (E_0^N - E_\nu^N + i\eta)}$$

$\nu$ th (de)excitation energy

## Electron-Hole Bethe-Salpeter Equation (eh-BSE)

$$L(12; 1'2') = \underbrace{L_0(12; 1'2')}_{G(12')G(21')} + \int d(33'44') L_0(13'; 1'3) \Xi^{eh}(34'; 3'4) \underbrace{L(42; 4'2')}_{eh \text{ kernel}}$$

Strinati, Riv. Nuovo Cimento 11 (1988) 1; Blase et al. JPCL 11 (2020) 7371; Marie et al. JCP 162 (2025) 134105

## Schwinger-Dyson Relationship

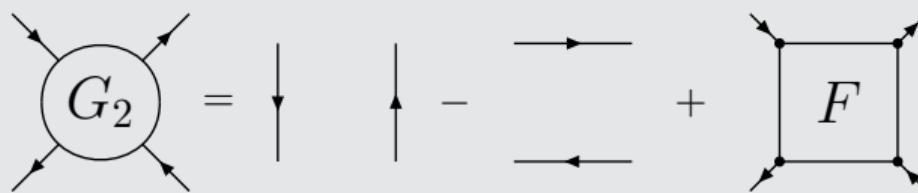
$$G^{-1}(11') = G_0^{-1}(11') - \Sigma(11')$$

$$\Sigma(11') = -i\nu(12; 3'2') G_2(3'2'; 32) G^{-1}(31')$$

## Two-body Vertex

$$G_2(12; 34) = G(13)G(24) - G(14)G(23) - G(11')G(3'3) F(1'2'; 3'4') G(4'4)G(22')$$

Full two-body vertex



Parquet theory aims at computing  $F$ , hence  $G_2$ , through Dyson equations

## Bethe-Salpeter Equations

$$\begin{array}{c} \text{Diagram: } F \\ \text{Diagram: } \Gamma^{\text{eh}} \\ \text{Diagram: } \Gamma^{\text{eh}} \text{ (horizontal loop)} \\ \text{Diagram: } F \end{array} = \begin{array}{c} \text{Diagram: } \Gamma^{\text{eh}} \\ \text{Diagram: } \Gamma^{\text{eh}} \end{array} + \begin{array}{c} \text{Diagram: } \Gamma^{\text{eh}} \text{ (vertical loop)} \\ \text{Diagram: } F \end{array}$$

$$\begin{array}{c} \text{Diagram: } F \\ \text{Diagram: } \Gamma^{\overline{\text{eh}}} \\ \text{Diagram: } \Gamma^{\overline{\text{eh}}} \text{ (vertical loop)} \\ \text{Diagram: } F \end{array} = \begin{array}{c} \text{Diagram: } \Gamma^{\overline{\text{eh}}} \\ \text{Diagram: } \Gamma^{\overline{\text{eh}}} \end{array} + \begin{array}{c} \text{Diagram: } \Gamma^{\overline{\text{eh}}} \\ \text{Diagram: } F \end{array}$$

$$\begin{array}{c} \text{Diagram: } F \\ \text{Diagram: } \Gamma^{\text{pp}} \\ \text{Diagram: } \Gamma^{\text{pp}} \text{ (crossed loop)} \\ \text{Diagram: } F \end{array} = \begin{array}{c} \text{Diagram: } \Gamma^{\text{pp}} \\ \text{Diagram: } \Gamma^{\text{pp}} \end{array} + \begin{array}{c} \text{Diagram: } \Gamma^{\text{pp}} \\ \text{Diagram: } F \end{array}$$

### Parquet Decomposition

$$F(12; 34) = \Lambda(12; 34) + \underbrace{\Phi^{eh}(12; 34) + \Phi^{\overline{eh}}(12; 34) + \Phi^{pp}(12; 34)}_{\text{can be computed with Bethe-Salpeter equations}}$$

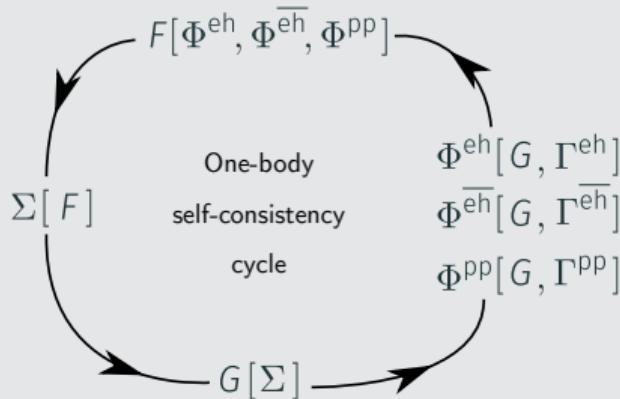
↑  
Irreducible vertex

Proper way to account for different correlation channels in the self-energy without double counting!

De Dominicis & Martin, J. Math. Phys. 5 (1964) 14; ibid 5 (1964) 31

Bickers, “*Self-consistent many-body theory for condensed matter systems*” in Theoretical Methods for Strongly Correlated Electrons (2004) 237

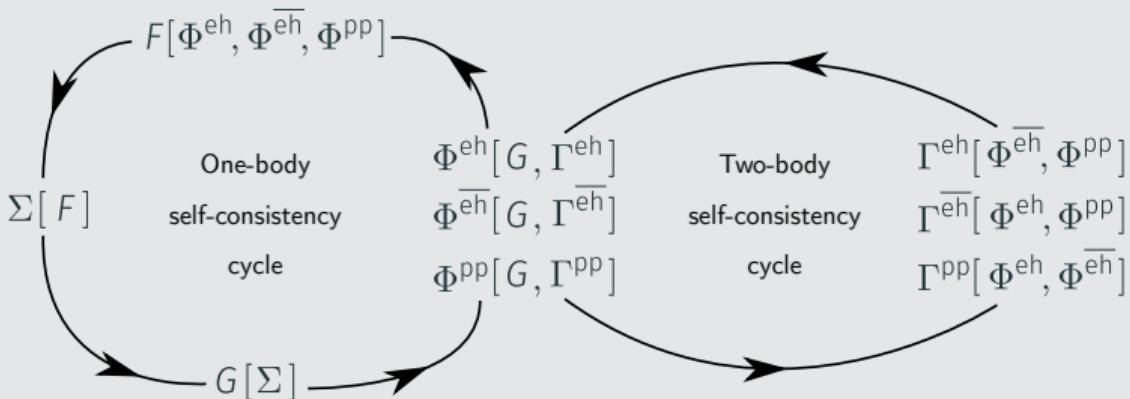
## Self-Consistent Algorithm



## Approximations

- Parquet approximation  
 $\Lambda = -i\bar{v}$
- One-shot approximation
- Static kernel  
approximation for  $\Gamma$

## Self-Consistent Algorithm

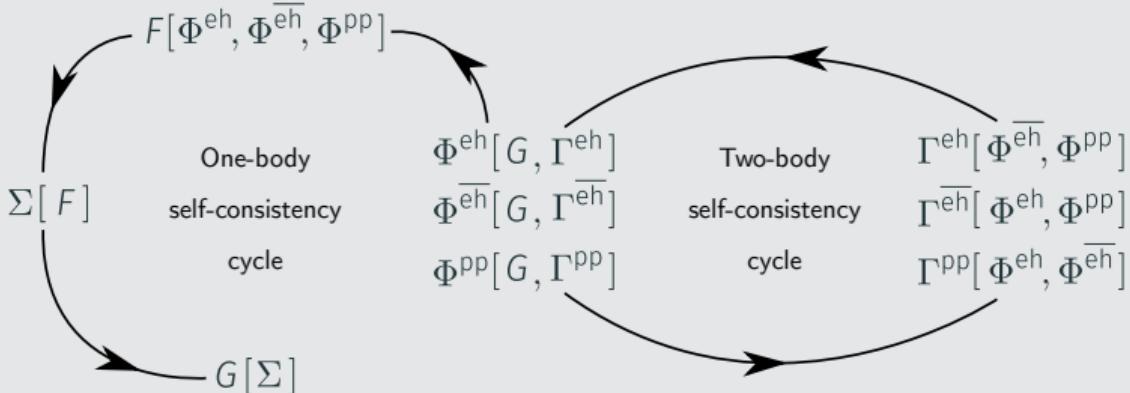


## Approximations

- Parquet approximation  
 $\Lambda = -i\bar{v}$
- One-shot approximation
- Static kernel approximation for  $\Gamma$

# Parquet Algorithm

## Self-Consistent Algorithm



## Approximations

- Parquet approximation  
 $\Lambda = -i\bar{v}$
- One-shot approximation
- Static kernel approximation for  $\Gamma$

## One-shot parquet approximation (osPA)

Full two-body self-consistency, single one-body iteration in the diagonal approximation

## Preliminary Results on Principal IPs

Preliminary statistics on 20 IPs in the aug-cc-pVTZ basis set

| Method | osPA | $G_0 W_0$ | $G_0 T_0$ |
|--------|------|-----------|-----------|
| MAE    | 0.29 | 0.37      | 0.34      |

Marie & Loos, JCP 163 (2025) 194115

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- Enzo Monino
- Roberto Orlando
- Raúl Quintero-Monsebaiz

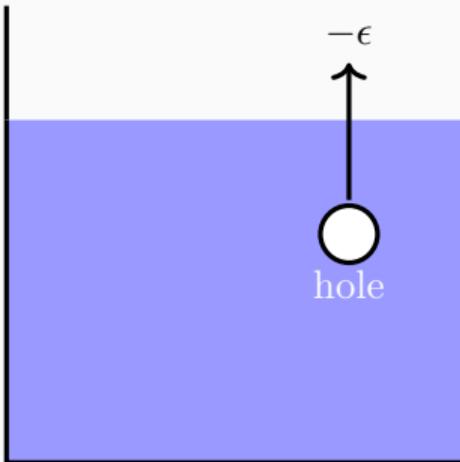


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<https://lcpq.github.io/PTEROSOR>

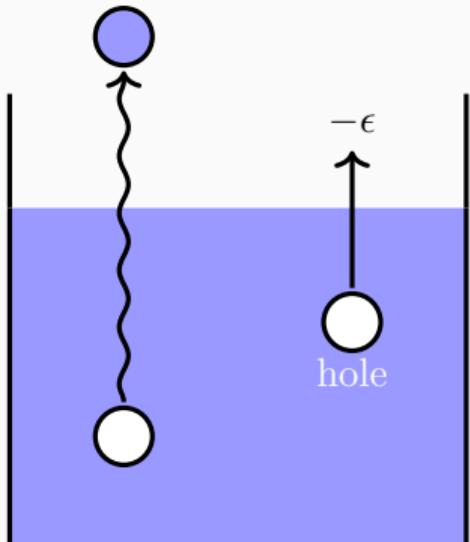
# Quasiparticle Concept



electron removal

- Link to electron-boson Hamiltonian:  
Langreth, PRB 1 (1970) 471  
Hedin, JPCM 11 (1999) R489
- Link to coupled-cluster theory:  
Lange & Berkelbach, JCTC 14 (2018) 4224  
Quintero-Monsebaiz et al. JCP 157 (2022) 231102  
Tolle & Chan, JCP 158 (2023) 124123

# Quasiparticle Concept

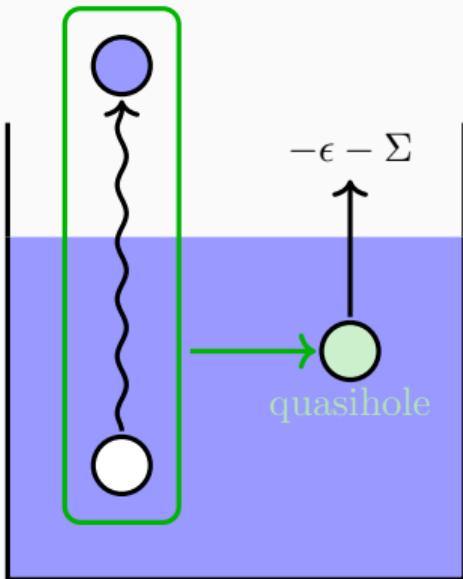


electron removal

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# Quasiparticle Concept

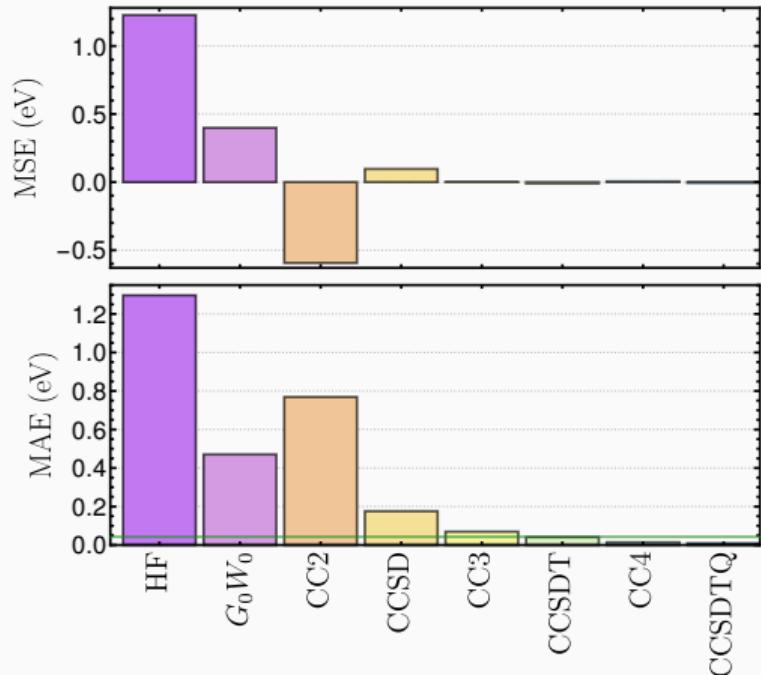
RPA excitation



electron removal

- Link to electron-boson Hamiltonian:  
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Lange & Berkelbach, JCTC 14 (2018) 4224  
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# Inner- and Outer-valence IPs (aug-cc-pVTZ) for 23 small molecules (FCI reference)



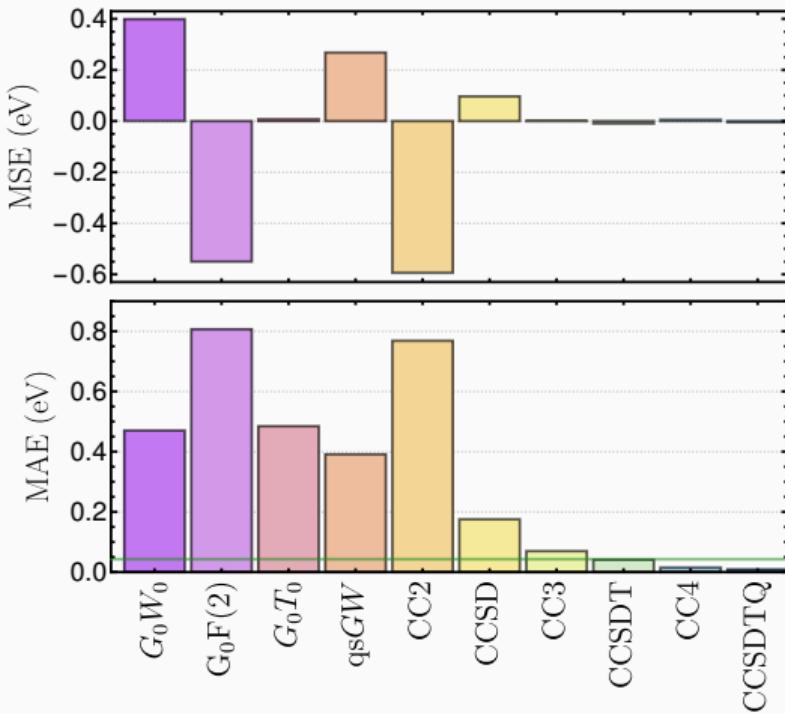
## Computational cost

- HF  $\mathcal{O}(K^4)$
- $G_0W_0$   $\mathcal{O}(K^6) \rightarrow \mathcal{O}(K^4)$
- IP-EOM-CC2  $\mathcal{O}(K^5)$
- IP-EOM-CCSD  $\mathcal{O}(K^6)$
- IP-EOM-CCSDT  $\mathcal{O}(K^8)$

## Some issues:

- Highly starting point dependent!
- Systematic improvable?

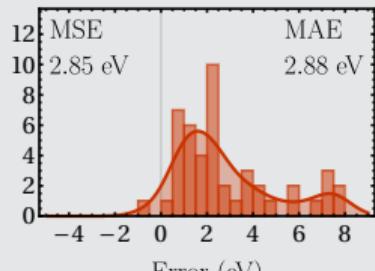
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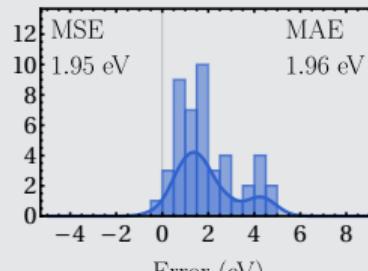
# Singlet and Triplet DIPs (aug-cc-pVTZ) for 23 small molecules (FCI reference)

## Effect of the Quasiparticle Energies

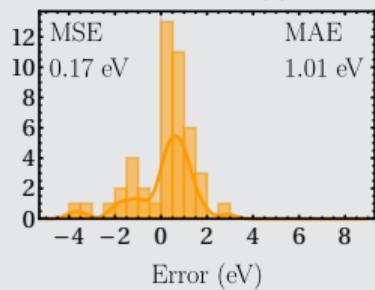
ppRPA@HF



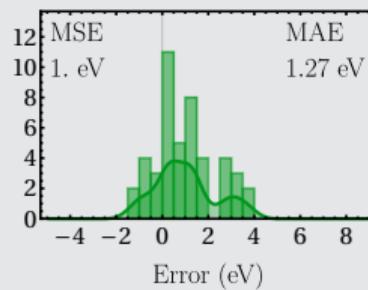
ppRPA@GW



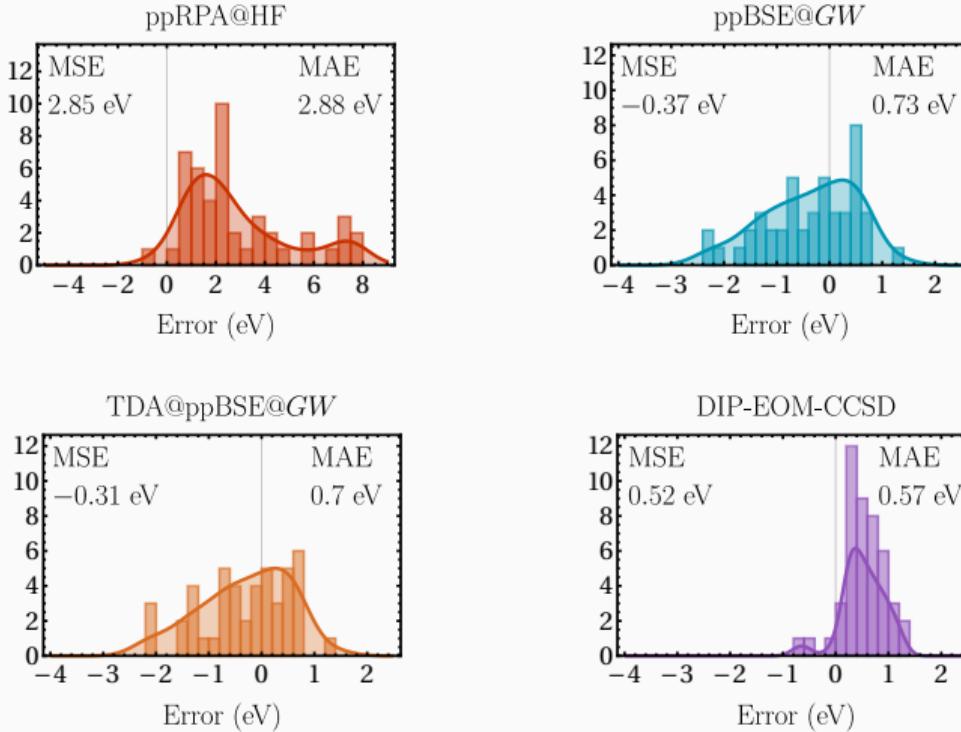
ppRPA@GF(2)



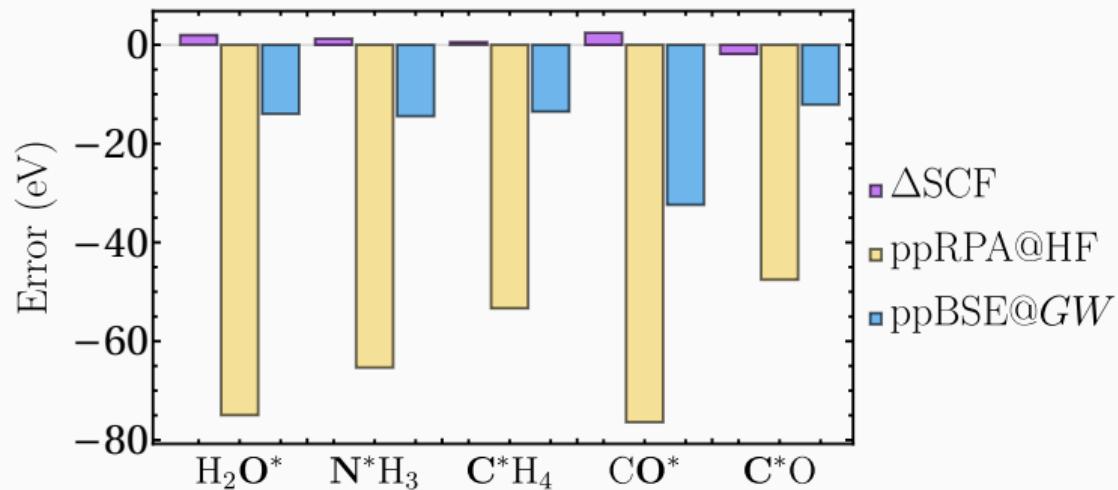
ppRPA@GT



# Singlet and Triplet DIPs (aug-cc-pVTZ) for 23 small molecules (FCI reference)



## (Single-Site) Double Core Holes (aug-cc-pCVTZ & CVS-FCI reference)



Cederbaum et al. JCP 85 (1986) 6513; Marie et al. JCP 162 (2025) 134105

# Parquet Algorithm

