Einstein's equations from a Newtonian perspective

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1 Introduction

- * If space metric and chronometry are influenced by matter, then all "rates of change" are determined by matter.
 - * Rate of change of momentum reinterpreted as inertial force (Mach)
- * "Motion" is not clear any longer. Topologically we can consider a bulk of matter as static and the metric among its elements as changing, or vice versa.
- * The problem is that only the relations of the geometrical objects with respect to one another is observable; but such relations are not all arbitrary, there are constraints among them. In Newtonian mechanics we use a reference frame a rigid body to find a compact description that gets rid of the constraints [1–5]. In general relativity there is no "rigidity" and we have more choice.
 - * Do the equations for the metric have constitutive parts?
- * If mass is twisted 3-form, then (analogously to displacement current) gravitational potential must be related to a twisted 2-form [6;7, ch. V, pp. 192-195]. Then is there an equivalent "magnetic field strength" that relates to momentum?

$$\mathrm{d} m=0, \qquad \partial_t \rho + \mathrm{d} j=0, \tag{1}$$

$$dF = \rho, \qquad \partial_t F + dH = j$$
 (2)

* Extrinsic curvature can be understood (kinematically) as the timederivative of 3-metric (no Einstein eqns needed). Setting lapse $\alpha = 1$ and shift $\beta = 0$:

$$\partial_t \mathbf{h} = -2\mathbf{K} \cdot \mathbf{h} \tag{3}$$

$$\partial_t \mathbf{K} = \mathbf{R} + \mathbf{K} \operatorname{tr} \mathbf{K} + 4\pi (\operatorname{tr} \mathbf{T} - E - 2\mathbf{T})$$
(4)

$$\operatorname{tr} \mathbf{R} + (\operatorname{tr} \mathbf{K})^2 - \mathbf{K} : \mathbf{K} = 16\pi E \tag{5}$$

$$\nabla \cdot \mathbf{K} - \nabla \operatorname{tr} \mathbf{K} = 8\pi \mathbf{p} \tag{6}$$

the second should come from

$$\partial_t(\mathbf{K} \cdot \mathbf{h}) = \mathbf{R} \cdot \mathbf{h} + \mathbf{h} \cdot \mathbf{K} \operatorname{tr} \mathbf{K} - 2\mathbf{h} \cdot \mathbf{K} \cdot \mathbf{K} + 4\pi [\mathbf{h}(\operatorname{tr} \mathbf{T} - E) - 2\mathbf{T} \cdot \mathbf{h})$$
 (7)

$$\partial_t \boldsymbol{p} = \frac{1}{8\pi} (\nabla \cdot \boldsymbol{R} - \partial_t \nabla \operatorname{tr} \boldsymbol{K} + \operatorname{tr} \boldsymbol{K} \nabla \cdot \boldsymbol{K}) + \frac{1}{2} \nabla \operatorname{tr} \boldsymbol{T} - \frac{1}{2} \nabla E - \nabla \cdot \boldsymbol{T}$$
(8)
$$\partial_t \boldsymbol{p} - \boldsymbol{p} \operatorname{tr} \boldsymbol{K} = \frac{1}{8\pi} [\nabla \cdot \boldsymbol{R} - \partial_t \nabla \operatorname{tr} \boldsymbol{K} + \frac{1}{2} \nabla (\operatorname{tr} \boldsymbol{K})^2] + \frac{1}{2} \nabla \operatorname{tr} \boldsymbol{T} - \frac{1}{2} \nabla E - \nabla \cdot \boldsymbol{T}$$
(9)

$$\partial_t \rho + \nabla \cdot (\rho v) = 0, \tag{10a}$$

$$\partial_t(\rho v) + \nabla \cdot (\rho v \otimes v) - \nabla \cdot \mathbf{T} - \rho \mathbf{b} = 0, \tag{10b}$$

$$\mathbf{T}^{\mathsf{T}} - \mathbf{T} = 0,\tag{10c}$$

$$\partial_t(\rho u) + \nabla \cdot (\rho v u) - \mathbf{T} : \nabla v + \nabla \cdot \mathbf{q} + \rho Q = 0, \tag{10d}$$

$$\begin{array}{l} \partial_t \left(\rho u + \frac{1}{2} \rho v^2 \right) + \nabla \cdot \left(\rho v u + \frac{1}{2} \rho v v^2 \right) - \\ \nabla \cdot \left(\boldsymbol{T} \cdot \boldsymbol{v} \right) - \rho \boldsymbol{b} \cdot \boldsymbol{v} - \nabla \cdot \boldsymbol{q} - \rho \boldsymbol{Q} = 0, \end{array} \tag{10e}$$

$$\partial_t(\rho s) + \nabla \cdot (\rho v s) - \nabla \cdot (q/\Theta) - \rho Q/\Theta \geqslant 0.$$
 (10f)

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