Peter Maginot Jean Ragusa and Jim Morel

Texas A&M University- Department of Nuclear Engineering

M&C 2015 April 20, 2015



Goal



### Outline

- Goal/Motivation
- BCSZ Method
- Initial Results

Computational Results

#### Goal

- Long term goal: Accurate methods for multi-dimensional radiative transfer
- Near term goal: Non-negative bilinear DFEM for neutron transport
  - Bilinear DFEM required on quads to maintain thick diffusion limit
  - Radiative transfer cells will almost certainly be optically thick
- History: Extension of non-negative linear (1, x, y) scheme developed on rectangles

# Solution representation, $\widetilde{\psi}_{BCSZ}(s,t)$ :

Solution representation,  $\psi_{BCSZ}(s,t)$ :

$$\widetilde{\psi}_{BCSZ}(s,t) = \begin{cases} \widehat{\psi}_{BCSZ}(s,t) & \widehat{\psi}_{BCSZ}(s,t) > 0\\ 0 & \text{otherwise} \end{cases}$$
 (1)

Bilinear function,  $\widehat{\psi}_{BCSZ}$ , to search for,

$$\widehat{\psi}_{BCSZ}(s,t) = \sum_{i=0}^{3} \psi_{i,BCSZ} B_i(s,t), \qquad (2)$$

# Moment Equation Edge Integration

DFEM moment equations will need to integrate quantities like this on cell edges:

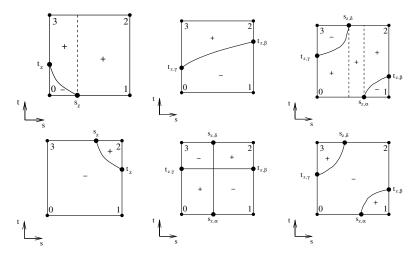
$$(\vec{\Omega}_d \cdot \vec{n}_\alpha) \int_{\alpha} B_i(s,-1) \widetilde{\psi}_{BCSZ} ds$$

- Check vertices for negativity
- ② By definition of  $\widetilde{\psi}_{BCSZ}$ , integrate  $\widehat{\psi}_{BCSZ}$  only over portion of the interval where  $\widehat{\psi}_{BCSZ} \geq 0$
- **3** If  $\psi_I < 0$  or  $\psi_R < 0$ :

$$s_{z} = \frac{\psi_{L} + \psi_{R}}{\psi_{L} - \psi_{R}}$$

# Cell Interior Integration

Must integrate terms like,  $B_i \psi_{BCSZ} |J|$ , over areas like these:



Along every integration area curve:

$$\widehat{\psi}_{BCSZ}(s,t) = 0$$

Transform interpolatory  $\widehat{\psi}_{BCSZ}$  to moment based f(s,t):

$$f(s,t) = f_c + sf_s + tf_t + stf_{st}$$
(3)

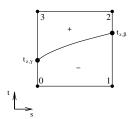
$$\begin{bmatrix} 1 & -1 & -1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \end{bmatrix} \begin{bmatrix} f_c \\ f_s \\ f_t \\ f_{\pm} \end{bmatrix} = \vec{\psi}_{BCSZ}$$
 (4)

Along each integration area defined by a curve, f(s, t) = 0, and

$$\hat{l}_t = -\frac{f_c + f_s s}{f_t + f_{st} s},\tag{5}$$

enabling the use of variable limits of integration.

Consider integration over the domain,  $R_+$ , where  $\widehat{\psi}_{BCSZ} > 0$ :



$$\int \int_{R_+} M(s,t) = \int_{-1}^1 ds \int_{\hat{I}_t}^1 dt \ M(s,t)$$

## Isn't MATLAB Right?

#### Initial thinking:

- Must respect curved boundary of integration regions
  - Only possible with analytic integration via variable limits of integration.
- MATLAB gives a solution, that solution must be right
  - Boldly [blindly] assumes we live in an analytic, not finite precision world

#### This results in:

- Occasional trouble with non-linear iteration
  - Predominantly with nearly zero solutions
- Maybe blindly trusting MATLAB is bad?

### Quadrature integrate $\psi_{i,M}$

$$\int \int_{R^+} |\mathbf{J}| \, B_i \widehat{\psi}_{BCSZ} \, \, ds dt$$

$$E_{i} = \frac{|\psi_{i,sym} - \psi_{i,num}|}{|\psi_{i,sym}|}$$

$$\widehat{E}_{i} = \frac{|\psi_{i,MAX} - \psi_{i,num}|}{|\psi_{i,MAX}|}$$

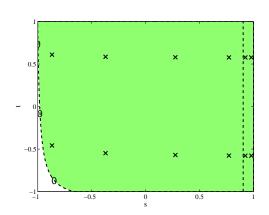


Figure: Quadrature layout with i = 4

### Evidence of Numerical Precision Issues

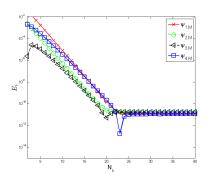


Figure:  $E_i$  for quadrature test.

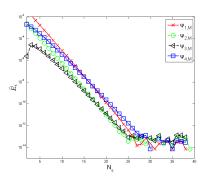


Figure:  $\widehat{E}_i$  for quadrature test.

**Modification**: Adaptive Gauss-Kronrod quadrature to evaluate cell integrals

- Newton iteration with finite difference formed Jacobian
- $\widehat{\psi}_{BCSZ}^{(0)} = \widetilde{\psi}_{UBLD}$
- Damping with restart based on iteration count
- Solve for local unknowns (1 solve per cell, direction, group [if needed])
- Search for scaled  $\widehat{\psi}_{BCSZ}$ . Scale using  $\widehat{\psi}_{BCSZ}^{(0)}$
- $\epsilon = \epsilon_{rel} R^{(0)} + \epsilon_{abs} \| RHS \|_{L_2}$
- $\bullet$   $\epsilon_{rel} = 10^{-10}$ ,  $\epsilon_{abs} = 10^{-12}$
- Usually 7-10 Newton iterations per solve

### Implementation in PDT

- PDT assumes unknowns live at cell vertices, and all methods are linear (not non-linear)
  - But, BCSZ is obviously non-linear
- Work around required
  - Add if statement in sweep. Affects all methods
  - 2 Multiply  $\psi_{i,M}$  of converged  $\psi_{BCSZ}$  by inverse mass matrix
    - ullet Gives a bilinear function with same cell  $\psi_{i,M}$  moments as  $\psi_{BCSZ}$
  - Prepare outflow for downwind cells
    - Calculate linear/nodal values necessary to yield BCSZ edge moments
- Exact BCSZ solution representation cannot be recreated without sweeping again
  - Cell average BCSZ flux retained with work around
  - Spatial moments preserved, does not affect physics coupling

Computational Results

# Methods to Compare

- UBLD: Unlumped Bilinear DFEM- Galerkin DFEM, no explanation necessary
- SCB/FLBLD: Subcell Corner Balance- On rectangles, Equivalent to UBLD with mass matrix lumping, surface matrix lumping, and other manipulations
- BCSZ: Bilinear consistent set-to-zero- Non-linear, Petrov-Galerkin DFEM, satisfies all bilinear spatial moments of the transport equation

Computational Results

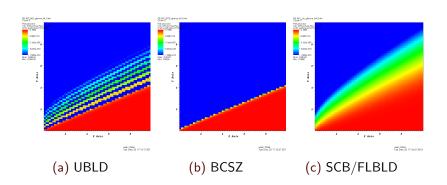
### Test Problem

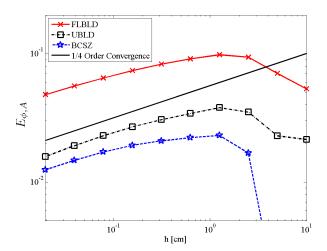
- $10[cm] \times 10[cm]$  square void
- Vacuum BC on left, top, right edges
- Incident flux of 1  $[n/(cm^2 sec ster)]$  in one direction on bottom edge
  - $\mu = 0.868890300722, \ \eta = 0.35002117452$
- L<sup>2</sup> like norm of cell average scalar flux

$$E_{\phi_A} = \sqrt{\sum_{c=1}^{N_{cells}} \Delta x_c \Delta y_c (\widetilde{\phi}_A - \phi_{A,exact})^2},$$

# Orthogonal Mesh

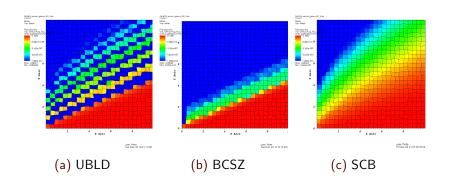
### 625 square mesh cells





### Distorted Mesh

#### $25 \times 25$ cells, with distorted interior vertices



### Conclusions and Future Work

#### Conclusions

Goal

- BCSZ is strictly non-negative
- BCSZ is more accurate than UBLD or SCB for a glancing void
- BCSZ can be applied to non-orthogonal meshes
- BCSZ requires significant local computation, but is computationally feasible

#### Future Work

- Develop a problem large/complex enough that timing can be performed
- Move non-linear iteration out of individual cells to enable preconditioning / DSA
- Consider removing case selection statement in favor of apply GK quad to entire cell

# Acknowledgments

Thanks for your time! Portions of this work were funded by the

Department of Energy CSGF program, administered by the Krell Institute, under grant DE-FG02-97ER25308.

Additional support was provided by the Department of Energy, National Nuclear Security Administration, under Award Number(s) DE-NA0002376.



