

## Assignment 3

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### Addition of spins $j = \frac{1}{2}$ and $j = 1$

$|j_1, j_2; j, m\rangle$  basis kets are the eigen-kets of  $J^2, J_z, J_1^2, J_2^2$ . The defining equation for the Clebsch-Gordan coefficients is:

$$\begin{aligned} |j_1, j_2; j, m\rangle &= \sum_{m_1 m_2} \langle j_1, j_2; m_1, m_2 | j_1, j_2; j, m \rangle |j_1, j_2; m_1, m_2\rangle \\ &= \sum_{m_1 m_2} C_{m_1 m_2}^{jm} |j_1, j_2; m_1, m_2\rangle \end{aligned}$$

Consider the action of  $J_+$  on the state  $|j_1, j_2; j, m\rangle$ ,

$$\begin{aligned} \sqrt{(j-m)(j+m+1)} |j_1, j_2; j, m+1\rangle &= \sum_{m_1 m_2} \sqrt{(j_1-m_1)(j_1+m_1+1)} C_{m_1 m_2}^{jm} |j_1, j_2; m_1+1, m_2\rangle + \\ &\quad \sum_{m_1 m_2} \sqrt{(j_2-m_2)(j_2+m_2+1)} C_{m_1 m_2}^{jm} |j_1, j_2; m_1, m_2+1\rangle \end{aligned}$$

Using the orthonormality of the basis kets and the definition of the Clebsch-Gordan coefficients,

$$\sqrt{(j-m)(j+m+1)} C_{m'_1 m'_2}^{j m+1} = \sqrt{(j_1-m'_1+1)(j_1+m'_1)} C_{m'_1-1 m'_2}^{j m} + \sqrt{(j_2-m'_2+1)(j_2+m'_2)} C_{m'_1 m'_2-1}^{j m}$$

Similarly,

$$\sqrt{(j+m)(j-m+1)} C_{m'_1 m'_2}^{j m-1} = \sqrt{(j_1+m'_1+1)(j_1-m'_1)} C_{m'_1+1 m'_2}^{j m} + \sqrt{(j_2+m'_2+1)(j_2-m'_2)} C_{m'_1 m'_2+1}^{j m}$$

### Calculation for $j = \frac{3}{2}$

$$m = \pm \frac{3}{2}$$

This state can only be obtained when  $(m_1, m_2) = (\pm 1, \pm \frac{1}{2})$ . Therefore,

$$C_{m_1 m_2}^{\frac{3}{2} \frac{3}{2}} = \delta_{m_1, 1} \delta_{m_2, \frac{1}{2}}$$

$$C_{m_1 m_2}^{\frac{3}{2} - \frac{3}{2}} = \delta_{m_1, -1} \delta_{m_2, -\frac{1}{2}}$$

$$m = \frac{1}{2}$$

Since, the Clebsch-gordan coefficients are known for  $m = \frac{3}{2}$ , using the already obtained relation:

$$\sqrt{\left(\frac{3}{2} + \frac{3}{2}\right)\left(\frac{3}{2} - \frac{3}{2} + 1\right)} C_{m'_1 m'_2}^{\frac{3}{2} \frac{3}{2}-1} = \sqrt{(1+m'_1+1)(1-m'_1)} C_{m'_1+1 m'_2}^{\frac{3}{2} \frac{3}{2}} + \sqrt{\left(\frac{1}{2} + m'_2 + 1\right)\left(\frac{1}{2} - m'_2\right)} C_{m'_1 m'_2+1}^{\frac{3}{2} \frac{3}{2}}$$

Simplifying,

$$\sqrt{3} C_{m'_1 m'_2}^{\frac{3}{2} \frac{1}{2}} = \sqrt{2} \delta_{m'_1, 0} \delta_{m'_2, \frac{1}{2}} + \delta_{m'_1, 1} \delta_{m'_2, -\frac{1}{2}}$$

$$m = -\frac{1}{2}$$

Since, the Clebsch-gordan coefficients are known for  $m = -\frac{3}{2}$ , using the already obtained relation:

$$\sqrt{\left(\frac{3}{2} - -\frac{3}{2}\right)\left(\frac{3}{2} + -\frac{3}{2} + 1\right)} C_{m'_1 m'_2}^{\frac{3}{2}(-\frac{3}{2}+1)} = \sqrt{(1 - m'_1 + 1)(1 + m'_1)} C_{m'_1 - 1 m'_2}^{\frac{3}{2} - \frac{3}{2}} + \sqrt{\left(\frac{1}{2} - m'_2 + 1\right)\left(\frac{1}{2} + m'_2\right)} C_{m'_1 m'_2 - 1}^{\frac{3}{2} - \frac{3}{2}}$$

Simplifying,

$$\sqrt{3} C_{m'_1 m'_2}^{\frac{3}{2} - \frac{1}{2}} = \sqrt{2} \delta_{m'_1, 0} \delta_{m'_2, -\frac{1}{2}} + \delta_{m'_1, -1} \delta_{m'_2, \frac{1}{2}}$$

## Calculation for $j = \frac{1}{2}$

$$m = \frac{1}{2}$$

This state can be constructed by the superposition of  $|1, -\frac{1}{2}\rangle$  and  $|0, \frac{1}{2}\rangle$ . By convention, the Clebsch-Gordan coefficients are taken to be orthonormal, which in turn implies from:

$$|C_{1-\frac{1}{2}}^{\frac{1}{2}\frac{1}{2}}|^2 + |C_{0\frac{1}{2}}^{\frac{1}{2}\frac{1}{2}}|^2 = 1,$$

that

$$C_{1-\frac{1}{2}}^{\frac{1}{2}\frac{1}{2}} = \cos \theta$$

$$C_{0\frac{1}{2}}^{\frac{1}{2}\frac{1}{2}} = \sin \theta$$

Consider the action of  $J_+$  on the state in discussion.

$$\sqrt{\left(\frac{1}{2} - \frac{1}{2}\right)\left(\frac{1}{2} + \frac{1}{2} + 1\right)} C_{m'_1 m'_2}^{\frac{1}{2}\frac{1}{2}+1} = \sqrt{(1 - m'_1 + 1)(1 + m'_1)} C_{m'_1 - 1 m'_2}^{\frac{1}{2}\frac{1}{2}} + \sqrt{\left(\frac{1}{2} - m'_2 + 1\right)\left(\frac{1}{2} + m'_2\right)} C_{m'_1 m'_2 - 1}^{\frac{1}{2}\frac{1}{2}}$$

Simplifying,

$$0 = \sqrt{(1 - m'_1 + 1)(1 + m'_1)} C_{m'_1 - 1 m'_2}^{\frac{1}{2}\frac{1}{2}} + \sqrt{\left(\frac{1}{2} - m'_2 + 1\right)\left(\frac{1}{2} + m'_2\right)} C_{m'_1 m'_2 - 1}^{\frac{1}{2}\frac{1}{2}}$$

Now, let  $m'_1 = 1$  and  $m'_2 = \frac{1}{2}$ . Substituting in the above equation,

$$\begin{aligned} 0 &= \sqrt{(1 - 1 + 1)(1 + 1)} C_{1-1\frac{1}{2}}^{\frac{1}{2}\frac{1}{2}} + \sqrt{\left(\frac{1}{2} - \frac{1}{2} + 1\right)\left(\frac{1}{2} + \frac{1}{2}\right)} C_{1\frac{1}{2}-1}^{\frac{1}{2}\frac{1}{2}} \\ &= \sqrt{2} \sin \theta + \cos \theta \\ \implies \tan \theta &= -\frac{1}{\sqrt{2}} \\ \implies \sin \theta &= \pm \sqrt{\frac{1}{3}} \\ \implies \cos \theta &= \mp \sqrt{\frac{2}{3}} \end{aligned}$$

Without loss of generality, I choose one sign. Therefore,

$$C_{m_1 m_2}^{\frac{1}{2}\frac{1}{2}} = \sqrt{\frac{2}{3}} \delta_{m_1, 1} \delta_{m_2, -\frac{1}{2}} - \sqrt{\frac{1}{3}} \delta_{m_1, 0} \delta_{m_2, \frac{1}{2}}$$

$$m = -\frac{1}{2}$$

Here, although an analogous calculation to the previous sub-section is possible, we choose to use the operation of  $J_-$  on the  $m = \frac{1}{2}$  state, to preserve the sign-convention adopted.

$$\sqrt{(\frac{1}{2} + \frac{1}{2})(\frac{1}{2} - \frac{1}{2} + 1)} C_{m'_1 m'_2}^{\frac{1}{2} \frac{1}{2} - 1} = \sqrt{(1 + m'_1 + 1)(1 - m'_1)} C_{m'_1 + 1 m'_2}^{\frac{1}{2} \frac{1}{2}} + \sqrt{(\frac{1}{2} + m'_2 + 1)(\frac{1}{2} - m'_2)} C_{m'_1 m'_2 + 1}^{\frac{1}{2} \frac{1}{2}}$$

Only non-zero values occur when either  $(m'_1, m'_2) = (-1, \frac{1}{2})$  or  $(m'_1, m'_2) = (0, -\frac{1}{2})$  Therefore,

$$C_{-1 \frac{1}{2}}^{\frac{1}{2} - \frac{1}{2}} = \sqrt{2} C_{0 \frac{1}{2}}^{\frac{1}{2} \frac{1}{2}} + \sqrt{(\frac{1}{2} + \frac{1}{2} + 1)(\frac{1}{2} - \frac{1}{2})} C_{-1 \frac{3}{2}}^{\frac{1}{2} \frac{1}{2}} = -\sqrt{\frac{2}{3}}$$

Similarly,

$$C_{0 - \frac{1}{2}}^{\frac{1}{2} - \frac{1}{2}} = \sqrt{\frac{1}{3}}$$

Hence,

$$C_{m_1 m_2}^{\frac{1}{2} - \frac{1}{2}} = -\sqrt{\frac{2}{3}} \delta_{m_1, -1} \delta_{m_2, \frac{1}{2}} + \sqrt{\frac{1}{3}} \delta_{m_1, 0} \delta_{m_2, -\frac{1}{2}}$$