# Assignment 1

Mytraya Gattu, Roshni Singh, Ramakant Pal 180050032, 18B030020, 180260028

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## Problem 1

Given that  $\langle x'|x\rangle = \delta(x'-x)$ , we can write the operator X in the position basis (its own eigenbasis) as:

$$\langle x'|X|x\rangle = x\delta(x'-x) = x'\delta(x'-x)$$

We also take the canonical commutation relation as one of the postulates:

$$[X,P] = i\hbar$$

Now we can write:

$$\langle x'| [X, P] | x \rangle = i\hbar \delta (x' - x)$$

$$\implies \langle x'| XP | x \rangle - \langle x'| PX | x \rangle = i\hbar \delta (x' - x)$$

$$\implies x' \langle x'| P | x \rangle - x \langle x'| P | x \rangle = i\hbar \delta (x' - x)$$

Therefore,

$$\langle x'|P|x\rangle = i\hbar \frac{\delta(x'-x)}{x'-x} = -i\hbar \frac{d}{dx} \left(\delta(x-x')\right)$$

### Problem 2

In the ket representation,

$$\partial_t |\psi(t)\rangle = -\frac{iH(t)}{\hbar} |\psi(t)\rangle$$

The wave function  $\psi(x,t)$  is defined as  $\langle x|\psi(t)\rangle$ . Multiplying the equation by  $\langle x|$  on both sides we get:

$$\partial_t \langle x | \psi(t) \rangle = -i \langle x | \frac{H(t)}{\hbar} | \psi(t) \rangle$$

$$\implies \partial_t \psi(x,t) = -i \langle x | \frac{H(t)}{\hbar} | \psi(t) \rangle$$

We now insert the identity  $\int dx'|x'\rangle\langle x'|=1$ , to get

$$\partial_t \psi(x,t) = -\frac{i}{\hbar} \int dx' \langle x|H(t)|x'\rangle \langle x'|\psi(t)\rangle$$

$$\implies \partial_t \psi(x,t) = -\frac{i}{\hbar} \int dx' \psi(x',t) \langle x|H(t)|x'\rangle$$

Now, we assume that H(t) is given by  $H(t) = \frac{P^2}{2m} + V$ .

$$\langle x|H(t)|x'\rangle = \langle x|\frac{P^2}{2m} + V|x'\rangle = \langle x|\frac{P^2}{2m}|x'\rangle + \langle x|V|x'\rangle$$

We use the position space representation of the momentum and potential operators to get:

$$\langle x|H(t)|x'\rangle = [-\hbar^2 \frac{\partial^2}{\partial x^2} + V(x)]\delta(x - x')$$

Using this, therefore

$$\partial_t \psi(x,t) = -\frac{i}{\hbar} \int dx' \psi(x',t) [-\hbar^2 \frac{\partial^2}{\partial x^2} + V(x)] \delta(x-x')$$

Therefore, we finally get:

$$-\frac{\hbar}{i}\frac{\partial}{\partial t}\psi(x,t) = \left(-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2} + V(x)\right)\psi(x,t)$$

(Schrodinger's Equation)

#### Problem 3

For a free non-relativistic particle,

$$H(t) = \frac{P^2}{2m}$$

Using the operator identity in the momentum space (taking the normalisation of plane wave states to be  $\frac{1}{\sqrt{(2\pi\hbar)^D}}$  in D dimensions), we have

$$\langle x'|e^{-\frac{\imath}{\hbar}Ht}|x\rangle = \int dp\langle x|p\rangle\langle p|e^{-\frac{\imath}{2\hbar m}P^2t}|x'\rangle$$

Simplifying.

$$\langle x'|e^{-\frac{\imath}{\hbar}Ht}|x\rangle = \int dp \frac{1}{2\pi\hbar} e^{\frac{\imath}{\hbar}p(x-x') - \frac{\imath}{2\hbar m}p^2t}$$

(Since, 
$$\langle x|p\rangle = \sqrt{\frac{1}{2\pi\hbar}}e^{\frac{\imath}{\hbar}p\cdot x}$$
)

Resolving the integral, (by completing the square in the exponent):

$$\langle x'|e^{-\frac{\imath}{\hbar}Ht}|x\rangle = \sqrt{\frac{m}{2\imath t\pi\hbar}}e^{\frac{im(x-x')^2}{2\hbar t}}$$

#### Working

The power of e is:

$$-\frac{i}{\hbar} \left[ \frac{p^2 t}{2m} - p \left( x - x' \right) \right]$$

Putting  $\alpha^2 = \frac{t}{2m}$ , we get

$$\frac{-i}{\hbar} \left[ \alpha^2 p^2 - 2(\alpha p) \frac{(x-x')}{2\alpha} + \frac{(x-x')^2}{4\alpha^2} \right] + \frac{i}{\hbar} \left[ \frac{(x-x')^2}{4\alpha^2} \right]$$

which can be written as:

$$-\frac{i}{\hbar} \quad \left(\alpha p - \frac{(x - x')}{2\alpha}\right)^2 + \frac{i}{\hbar} \frac{(x - x')^2}{4\alpha^2}$$

Therefore we can rewrite the integral as:  $\left\langle x' \left| e^{-\frac{i}{\hbar}Ht} \right| x \right\rangle = \int dp \frac{1}{2\pi\hbar} e^{-\frac{i}{\hbar}} \left( \alpha p - \frac{(x-x')}{2\alpha} \right)^2 + \frac{i}{\hbar} \frac{(x-x')^2}{4\alpha^2}$  Therefore, we can write our integral as,

$$\int_{-\infty}^{\infty} dx e^{-ix^2} = \sqrt{-i\pi}$$

Therefore we get,

$$\left\langle x'\left|e^{-\frac{\imath}{\hbar}Ht}\right|x\right\rangle = e^{\frac{i}{\hbar}\frac{(x-x')^2}{4\alpha^2}}\frac{1}{2\pi\hbar}\frac{\sqrt{-\imath\pi}}{\alpha}$$

Putting the value of  $\alpha^2 = \frac{t}{2m}$  , we get the result stated above.