## Chapter 1

## τ decays into hadrons

$$R_{\tau} = \frac{\Gamma(\tau \to \nu_{\tau} + Hadrons)}{\Gamma(\tau \to \nu_{\tau} e^{+} e^{-})} \tag{1.1}$$

The theoretical expression of the hadronic  $\tau$ -decay ratio was first derived by [**Tsai1971**] (using current algebra, a more recent derivation making use of the \*optical theorem\* can be taken from [**Schwab2002**]):

$$R_{\tau} = 12\pi \int_{0}^{m_{\tau}} = \frac{ds}{m_{\tau}^{2}} \left( 1 - \frac{s}{m_{\tau}^{2}} \right) \left[ \left( 1 + 2\frac{s}{m_{\tau}^{2}} \right) \operatorname{Im} \Pi^{(T)}(s) + \operatorname{Im} \Pi^{(L)} \right]. \quad (1.2)$$

 $R_{\tau}$  introduces a problematic integral over the real axis of  $\Pi(s)$  from 0 up to  $m_{\tau}$  for two reasons:

- The *perturbative Quantum Chromodynamcs* (**pQCD**) and the OPE breaks down for low energies (over which we have to integrate).
- The positive euclidean axis of  $\Pi(s)$  has a discontinuity cut and can theoretically not be evaluated.

Cauchy's Theorem

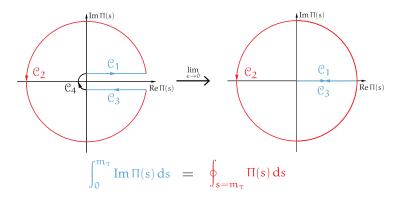
$$\int_{\mathcal{C}} f(z) dz = 0 \tag{1.3}$$

$$\begin{split} \oint_{s=m_{\tau}} \Pi(s) &= \int_{0}^{m_{\tau}} \Pi(s+i\varepsilon) + \int_{\mathcal{C}_{2}} \Pi(s) \, ds + \int_{m_{\tau}}^{0} \Pi(s-i\varepsilon) \, ds + \int_{\mathcal{C}_{4}} \Pi(s) \, ds \\ &= \int_{0}^{m_{\tau}} \Pi(s+i\varepsilon) - \Pi(s-i\varepsilon) \, ds + \int_{\mathcal{C}_{2}} \Pi(s) \, ds + \int_{\mathcal{C}_{4}} \Pi(s) \, ds \\ &= \int_{0}^{m_{\tau}} \Pi(s+i\varepsilon) - \overline{\Pi(s+i\varepsilon)} + \int_{\mathcal{C}_{2}} \Pi(s) \, ds + \int_{\mathcal{C}_{4}} \Pi(s) \, ds \end{split} \tag{1.4}$$

$$\lim_{\epsilon \to 0} 2i \int_{0}^{m_{\tau}} \operatorname{Im} \Pi(s) \, ds + \oint_{s=m_{\tau}} \Pi(s) \, ds$$

where we made use of  $\Pi(z) = \overline{\Pi(\overline{z})}$ , because  $\Pi(s)$  is analytic and  $\Pi(z) - \overline{\Pi(z)} = 2i \operatorname{Im} \Pi(z)$ 

$$\int_0^{m_\tau} \Pi(s) \, \mathrm{d}s = \frac{\mathrm{i}}{2} \oint_{s=m_\tau} \Pi(s) \, \mathrm{d}s \tag{1.5}$$



$$R_{\tau} = 6\pi i \oint_{s=m_{\tau}} \frac{ds}{m_{\tau}^2} \left( 1 - \frac{s}{m_{\tau}^2} \right) \left[ \left( 1 + 2\frac{s}{m_{\tau}^2} \right) \Pi^{(T)}(s) + \Pi^{(L)} \right]$$
 (1.6)

$$\begin{split} R_{\tau} &= 6\pi i \oint_{|s| = m_{\tau}^2} \frac{ds}{m_{\tau}^2} \left(1 - \frac{s}{m_{\tau}^2}\right)^2 \left[ \left(1 + 2\frac{s}{m_{\tau}^2}\right) \Pi^{(L+T)}(s) - \left(\frac{2s}{m_{\tau}^2}\right) \Pi^{(L)}(s) \right] \\ D^{(L+T)}(s) &\equiv -s \frac{d}{ds} \Pi^{(L+T)}(s), \qquad D^{(L)}(s) \equiv \frac{s}{m_{\tau}^2} \frac{d}{ds} (s \Pi^{(L)}(s)) \end{aligned} \tag{1.8}$$

Integration by parts

$$\int_{a}^{b} u(x)V(x) dx = \left[ U(x)V(x) \right]_{a}^{b} - \int_{a}^{b} U(x)v(x) dx$$
 (1.9)

$$\begin{split} R_{\tau}^{(1)} &= \frac{6\pi i}{m_{\tau}^{2}} \oint_{|s|=m_{\tau}^{2}} \underbrace{\left(1 - \frac{s}{m_{\tau}^{2}}\right)^{2} \left(1 + 2\frac{s}{m_{\tau}^{2}}\right)}_{=u(x)} \underline{\Pi^{(L+T)}(s)} \\ &= \frac{6\pi i}{m_{\tau}^{2}} \left\{ \left[ -\frac{m_{\tau}^{2}}{2} \left(1 - \frac{s}{m_{\tau}^{2}}\right)^{3} \left(1 + \frac{s}{m_{\tau}^{2}}\right) \underline{\Pi^{(L+T)}(s)} \right]_{|s|=m_{\tau}^{2}} \\ &+ \oint_{|s|=m_{\tau}^{2}} \underbrace{-\frac{m_{\tau}^{2}}{2} \left(1 - \frac{s}{m_{\tau}^{2}}\right)^{3} \left(1 + \frac{s}{m_{\tau}^{2}}\right) \underbrace{\frac{d}{ds}}_{=v(x)} \underline{\Pi^{(L+T)}(s)} \right\}}_{=u(x)} \\ &= -3\pi i \oint_{|s|=m_{\tau}^{2}s} \frac{ds}{s} \left(1 - \frac{s}{m_{\tau}^{2}}\right)^{3} \left(1 + \frac{s}{m_{\tau}^{2}}\right) \underbrace{\frac{d}{ds}}_{=v(x)} \underline{\Pi^{(L+T)}(s)} \\ &= -3\pi i \oint_{|s|=m_{\tau}^{2}s} \underbrace{\frac{ds}{s} \left(1 - \frac{s}{m_{\tau}^{2}}\right)^{3} \left(1 + \frac{s}{m_{\tau}^{2}}\right) \underbrace{\frac{d}{ds}}_{=v(x)} \underline{\Pi^{(L+T)}(s)}}_{=v(x)} \end{split}$$

where we fixed the integration constant to  $C=-\frac{m_{\tau}^2}{2}$  in the second line and left the antiderivatives contained in the squared brackets untouched. Parametrizing the expression in the squared brackets

$$\left[ -\frac{m_{\tau}^2}{2} \left( 1 - e^{-i\phi} \right)^3 \left( 1 + e^{-i\phi} \right) \Pi^{(L+T)}(m_{\tau}^2 e^{-i\phi}) \right]_0^{2\pi} = 0 \tag{1.11}$$

where  $s\to m_\tau^2 e^{-\mathfrak{i}\,\varphi}$  and  $(1-e^{-\mathfrak{i}\cdot 0})=(1-e^{-\mathfrak{i}\cdot 2\pi})=0.$ 

$$\begin{split} R_{\tau}^{(2)} &= \oint_{|s| = m_{\tau}^2} ds \left( 1 - \frac{s}{m_{\tau}^2} \right)^2 \left( -\frac{2s}{m_{\tau}^2} \right) \Pi^{(L)}(s) \\ &= -4\pi i \oint \frac{ds}{s} \left( 1 - \frac{s}{m_{\tau}^2} \right)^3 D^{(L)}(s) \end{split}$$
(1.12)

$$R_{\tau} = -\pi i \oint_{|s|=m_{\tau}^2} \frac{d}{s} \left( 1 - \frac{s}{m_{\tau}^2} \right)^3 \left[ 3 \left( 1 + \frac{s}{m_{\tau}^2} D^{(L+T)}(s) + 4D^{(L)}(s) \right) \right]$$
 (1.13)

$$R_{\tau} = -\pi i \oint_{|s|=m_{\tau}^2} \frac{d}{x} (1-x)^3 \left[ 3(1+x) D^{(L+T)}(s) + 4 D^{(L)}(s) \right), \tag{1.14}$$

where  $x = s/m_{\tau}^2$ .