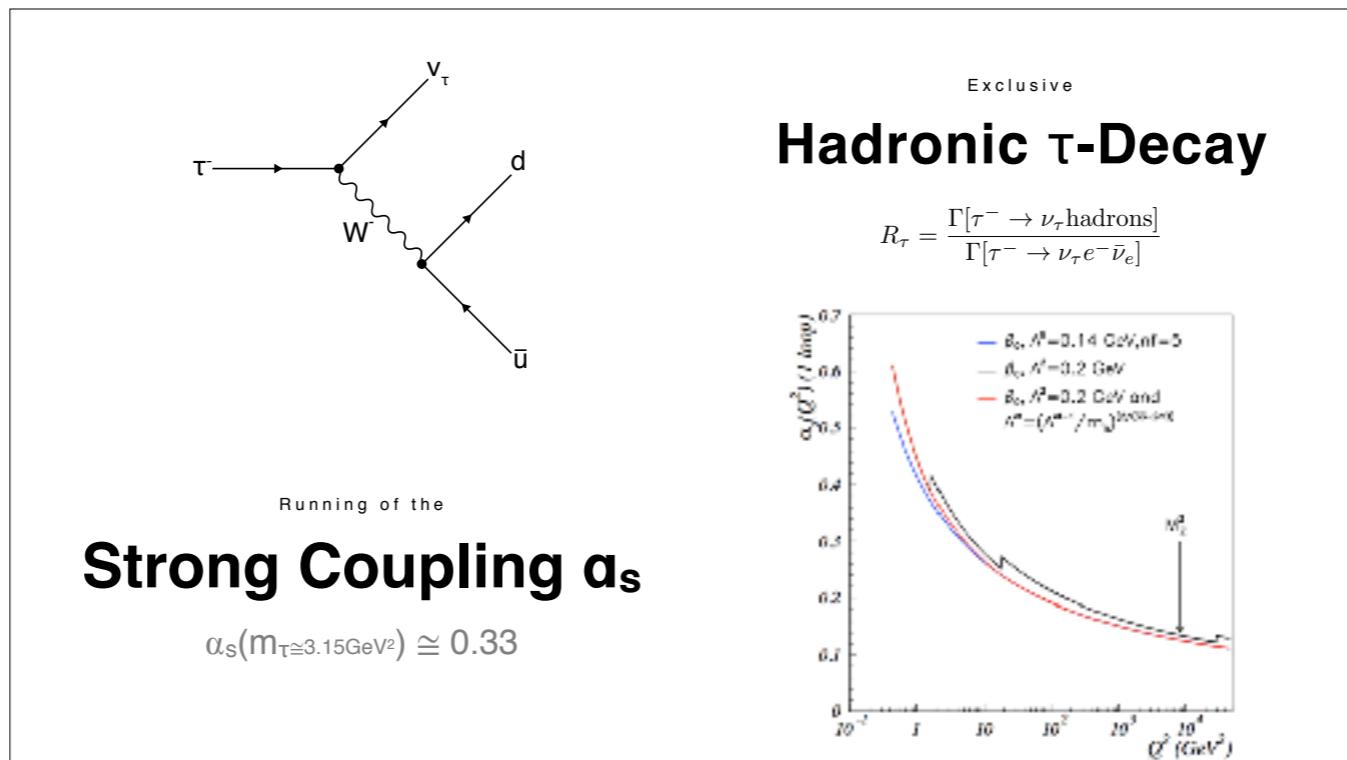


ALEPH

Determination of the QCD Coupling from ALEPH τ Decay Data

Dirk Hornung

τ Decay



- Hadronic decay width of τ most precise determinations of α_s
- R_τ = inclusive hadronic decay ratio can be calculated within QCD & sensitive α_s
- τ decays into a ν_τ & W boson, decaying into quark anti-quark (lepton-neutrino) pairs.
- Running α_s . $\alpha_s(m_\tau = 3.15 \text{ GeV}^2) = 0.33 \Rightarrow R_\tau$ sensitive to α_s
- $\alpha_s(m_\tau)$ rel. err: 4%, but $\alpha_s(M_Z)$ rel. err. : 1%

Contents

- Theoretical Framework**
- ALEPH Data**
- Fitting Methodology**
- Duality Violation**
- Determination of α_s**
- Program**
- Summary**

Theory

Framework

$$\text{Inclusive Ratio} \quad R_\tau = 12\pi S_{EW} \int_0^{m_\tau} \frac{ds}{m_\tau^2} \left(1 - \frac{s}{m_\tau^2}\right) \left[\left(1 + 2\frac{s}{m_\tau^2}\right) \text{Im } \Pi^{(1)}(s) + \text{Im } \Pi^{(0)}(s) \right]$$

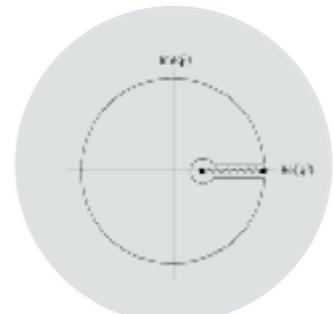
$$\Pi^{(J)}(s) \equiv |V_{uq}|^2 \left(\Pi_{ud,V}^{(J)} + \Pi_{ud,A}^{(J)}(s) \right)$$

Two-Point Correlation Function

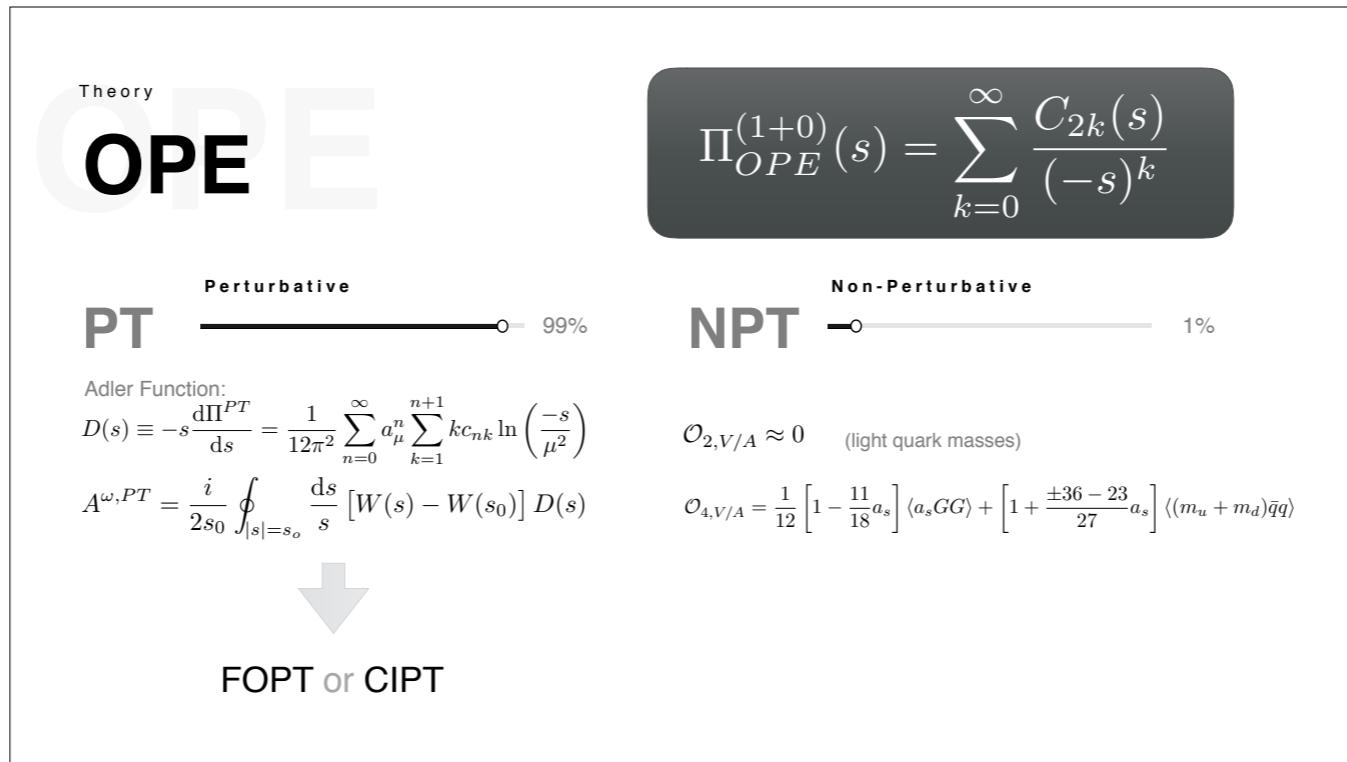
$$\begin{aligned} \Pi_{\mu\nu}(q) &= i \int d^4x e^{iqx} \langle 0 | T \left\{ \mathcal{J}_{ij}^\mu(x) \mathcal{J}_{ij}^{\nu\dagger}(0) \right\} \rangle \\ &= (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi_{ij,\mathcal{J}}^{(1)}(q^2) + q^\mu q^\nu \Pi_{ij,\mathcal{J}}^{(0)}(q^2) \end{aligned}$$

$$(i, j = u, d; \mathcal{J} = V, A) \quad V_{ij}^\mu = \bar{q}_j \gamma^\mu q_i \quad A_{ij}^\mu = \bar{q}_j \gamma^\mu \gamma_5 q_i$$

$$\int_{s_{th}}^{s_0} \frac{ds}{s_0} \omega(s) \text{Im } \Pi_{V/A}(s) = \frac{i}{2} \oint_{|s|=s_0} \frac{ds}{s_0} \omega(s) \Pi_{V/A}(s)$$



- Main ingredient: R_τ (depends on the two-point function).
- Π split up in transversal (spin 0)/ longitudinal (spin 1) & V/A part.
- Experimentally: measure spectral-function ρ
- Theoretically: employ OPE (not valid on positive real axis, axis of exp. data)
- Use of Cauchy's theorem, connect exp. with th.
- $\omega(s)$ is an analytic function, called the weight function



- Use OPE for scalar Π
- The OPE consist PT ($k=0$) & NPT part ($k \geq 1$)
- NPT-corrections contribute to only 1% (kinematic weight)
- PT \rightarrow Adler function (a_s , c_{nk} , $\log(s)$). $D(s)$ “independent” of μ , but higher order \Rightarrow th. err.
- Partial integration, include $D(s)$ \Rightarrow different weights
- NPT terms contain “condensates”.
- Omit $D=2$ term (proportional to the m_q^2)
- $D=4$ mainly Gluon Condensate
- FOPT vs CIPT
- Two operations in momentum: integration & sum, which first? sum \Rightarrow FOPT, int. \Rightarrow CIPT
- FOPT: bigger $R_T \Rightarrow$ smaller a_s

Theory

OPE D=6

Wilson coefficients:

$$\begin{aligned} C_6^{V-A}(Q^2) \langle O_6 \rangle &= 4\pi^2 a_s \left\{ \left[2 + \left(\frac{25}{6} - L \right) a_s \right] \langle Q_-^o \rangle - \left(\frac{11}{18} - \frac{2}{3}L \right) a_s \langle Q_-^s \rangle \right\} \\ C_6^{V+A}(Q^2) \langle O_6 \rangle &= -4\pi^2 a_s \left\{ \left[2 + \left(\frac{155}{24} - \frac{7}{2}L \right) a_s \right] \langle Q_+^o \rangle + \left(\frac{11}{18} - \frac{2}{3}L \right) a_s \langle Q_+^s \rangle + \right. \\ &\quad \left. \left[\frac{4}{9} + \left(\frac{37}{36} - \frac{95}{162}L \right) a_s \right] \langle Q_3 \rangle + \left(\frac{35}{108} - \frac{5}{18}L \right) a_s \langle Q_4 \rangle + \right. \\ a_s &\equiv \frac{\alpha_s}{\pi} \qquad \qquad \qquad \left(\frac{14}{81} - \frac{4}{27}L \right) a_s \langle Q_6 \rangle - \left(\frac{2}{81} + \frac{4}{27}L \right) a_s \langle Q_7 \rangle \\ L &\equiv \log \frac{Q^2}{\mu^2} \end{aligned}$$

L.E. Adam, K.G. Chetyrkin, Phys. Lett. B329, 129 (1994)

- Master (end 2015)
- Anomalous dimension of four-quark operators for D=6 term to 1st order
- Contributions only from 4-quark condensates (3-gluon condensate at 2nd order)
- Contribution to the $\Pi_{\mu\nu}$ computed at next-to-leading order (Adam, Chetyrkin, 1994)
- Results for V-A and V+A (V-A the penguin diagrams cancelled, $n_f=3$ & $n_c=3$)

Theory

OPE D=6

Basis

Singlet:

$$Q_{\pm}^S \equiv Q_V^S \pm Q_A^S$$

Octet:

$$Q_{\pm}^O \equiv Q_V^O \pm Q_A^O$$

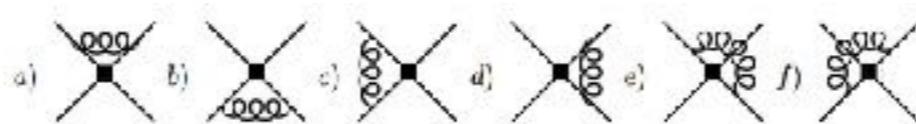
$$\begin{aligned}
Q_V^O &= (\bar{u}\gamma_\mu t^a d\bar{d}\gamma^\mu t^a u), \quad Q_A^O = (\bar{u}\gamma_\mu \gamma_5 t^a d\bar{d}\gamma^\mu \gamma_5 t^a u), \\
Q_V^S &= (\bar{u}\gamma_\mu d\bar{d}\gamma^\mu u), \quad Q_A^S = (\bar{u}\gamma_\mu \gamma_5 d\bar{d}\gamma^\mu \gamma_5 u), \\
Q_3 &\equiv (\bar{u}\gamma_\mu t^a u + \bar{d}\gamma_\mu t^a d) \sum_{q=u,d,s} (\bar{q}\gamma^\mu t^a q), \\
Q_4 &\equiv (\bar{u}\gamma_\mu \gamma_5 t^a u + \bar{d}\gamma_\mu \gamma_5 t^a d) \sum_{q=u,d,s} (\bar{q}\gamma^\mu \gamma_5 t^a q), \\
Q_5 &\equiv (\bar{u}\gamma_\mu u + \bar{d}\gamma_\mu d) \sum_{q=u,d,s} (\bar{q}\gamma^\mu q), \\
Q_6 &\equiv (\bar{u}\gamma_\mu \gamma_5 u + \bar{d}\gamma_\mu \gamma_5 d) \sum_{q=u,d,s} (\bar{q}\gamma^\mu \gamma_5 q), \\
Q_7 &\equiv \sum_{q=u,d,s} (\bar{q}\gamma_\mu t^a q) \sum_{q'=u,d,s} (\bar{q}'\gamma^\mu t^a q'), \\
Q_8 &\equiv \sum_{q=u,d,s} (\bar{q}\gamma_\mu \gamma_5 t^a q) \sum_{q'=u,d,s} (\bar{q}'\gamma^\mu \gamma_5 t^a q'), \\
Q_9 &\equiv \sum_{q=u,d,s} (\bar{q}\gamma_\mu q) \sum_{q'=u,d,s} (\bar{q}'\gamma^\mu q'), \\
Q_{10} &\equiv \sum_{q=u,d,s} (\bar{q}\gamma_\mu \gamma_5 q) \sum_{q'=u,d,s} (\bar{q}'\gamma^\mu \gamma_5 q').
\end{aligned}$$

- The appearing four-quark operators are a subset which belong to the complete basis required for the one-loop renormalisation.
- $Q_{V,A}$: current-current operators
- $Q_3 - Q_{10}$ penguin operators

Theory

OPE D=6

Current-current diagrams



Penguin diagrams:



- Lots of penguin diagrams

Theory

OPE D=6

Flavor non-diagonal

$$\gamma_{Q+}^{(1)} = \begin{pmatrix} -\frac{3}{N_c} & \frac{3C_F}{2N_c} & -\frac{1}{3N_c} & 0 & 0 \\ 3 & 0 & \frac{3}{2} & 0 & 0 \\ 0 & 0 & \frac{N_f}{4} - \frac{3N_c}{4} - \frac{1}{3N_c} & -\frac{3N_c}{4} - \frac{3}{N_c} & 0 \\ \frac{3}{2} + \frac{3}{2N_c} & -\frac{3C_F}{2N_c} & \frac{3N_c}{4} + \frac{3}{2} - \frac{11}{6N_c} & -\frac{3N_c}{4} + \frac{3}{2} + \frac{3}{2N_c} & 0 \\ 0 & 0 & \frac{11}{3} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\hat{\gamma}_{Q-}^{(1)} = \begin{pmatrix} -\frac{3N_c}{2} + \frac{3}{N_c} & -\frac{3C_F}{2N_c} \\ -3 & 0 \end{pmatrix}$$

Anomalous dimension matrix

$$\hat{\gamma}(a_s) = a_s \hat{\gamma}^{(1)} + a_s^2 \hat{\gamma}^{(2)} + \dots$$

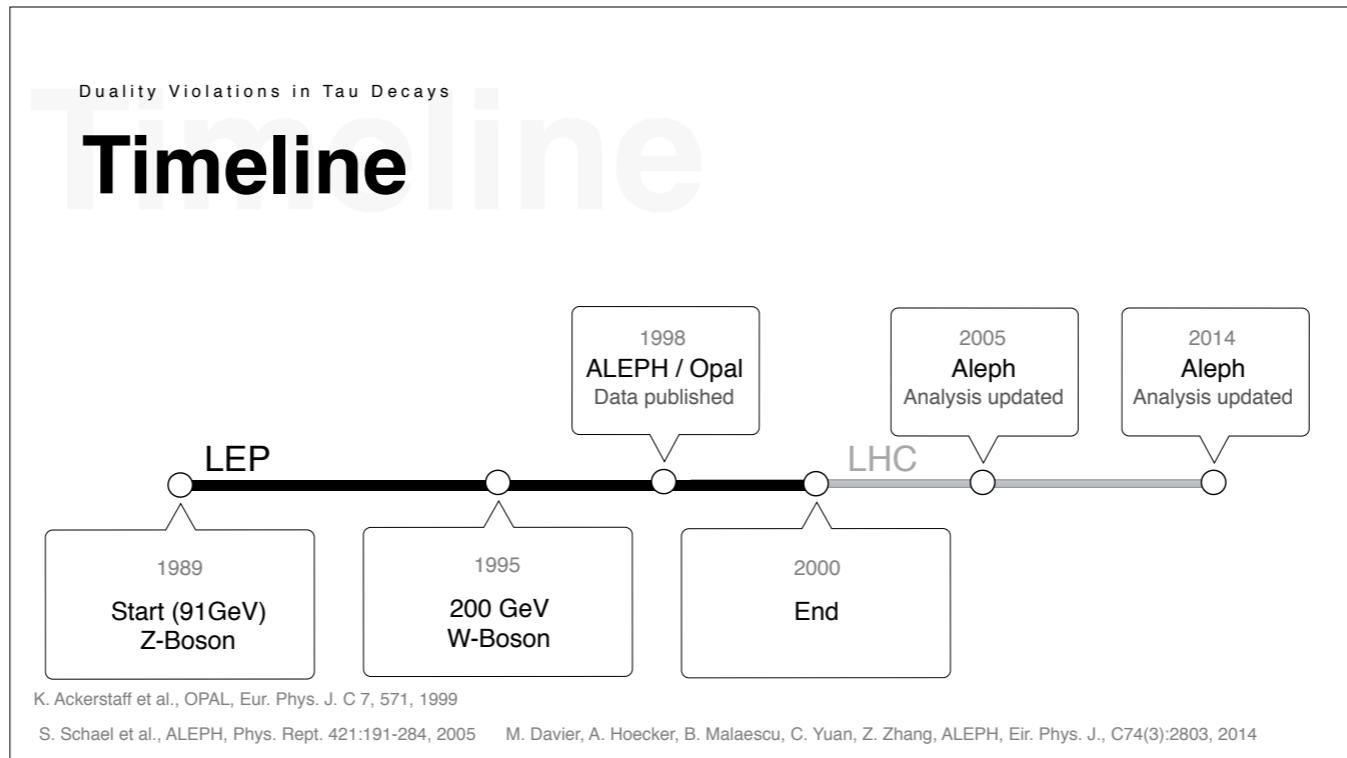
V/A currents:

$$j_\mu^V(x) = (\bar{u}\gamma_\mu d)(x)$$

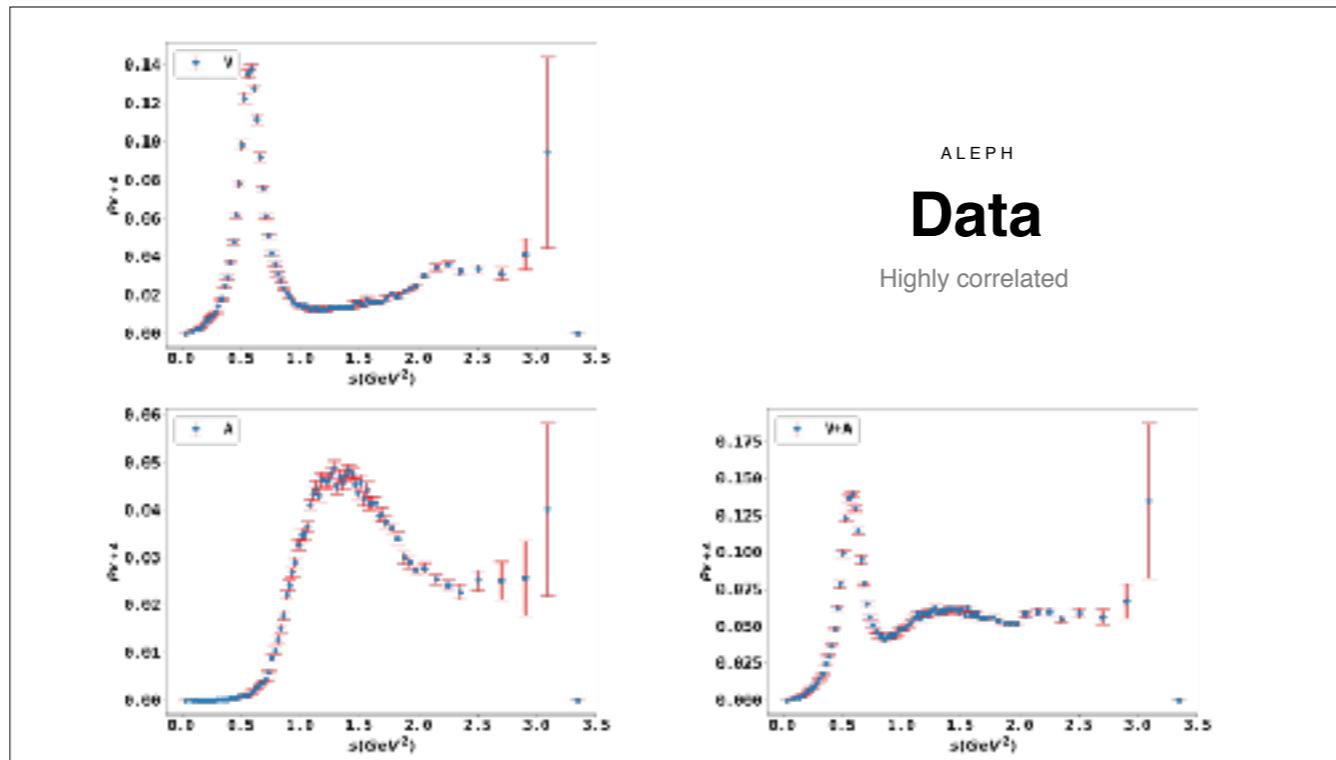
$$j_\mu^A(x) = (\bar{u}\gamma_\mu\gamma_5 d)(x)$$

Diogo Boito, Dirk Hornung, Matthias Jamin, JHEP 1512 (2015)

- The anomalous dimension matrix defines running of the Operators
- D=6 contribution at different s (used as const. in our fits)
- Flavor non-diagonal V/A correlators
- Q- no penguin diagrams (2x2)
- Same for Flavor diagonal
- used γ to analyse the renormalon structure of dimension 6



- LEP (Large Electron-Positron Collider, same tunnel as LHC)
- Detectors ALEPH & OPAL
- Analysis update (new Parameters, unfolding)
- Opal data have bigger errors



- Updated Aleph data (2014).
- The data channels: V, A & V+A.
- Contains correlation matrix, but highly correlated, needed for χ^2
- V: ρ resonance visible & large errors last bins.
- A: a_1 resonance visible & large errors last 4 bins.
- V+A: quick convergence, after ρ resonance, agreement with naive Parton model.
- V+A: D6 cancels (V,A have oposite signs)

Data

ALEPH

$$\rho(s) \equiv \frac{1}{\pi} \operatorname{Im} \Pi(s)$$

Spectral Function

Normalized invariant mass-squared distribution

$$\left(\frac{1}{N_{V/A}} \right) \left(\frac{dN_{V/A}}{ds} \right)$$

$$\operatorname{Im} \Pi_{\bar{u}d,V}^{(1)}(s) = \frac{1}{2\pi} v_1(s)$$

$$\operatorname{Im} \Pi_{\bar{u}d,A}^{(1)}(s) = \frac{1}{2\pi} a_1(s)$$

$$\operatorname{Im} \Pi_{\bar{u}d,A}^{(0)}(s) = \frac{1}{2\pi} a_0(s)$$

$$v_1(s) \equiv \frac{m_\tau^2}{6|V_{ud}|^2 S_{EW}} \frac{B(\tau^- \rightarrow V^- \nu_\tau)}{B(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)} \frac{dN_V}{N_V ds} \left[\left(1 - \frac{s}{m_\tau^2}\right)^2 \left(1 + \frac{2s}{m_\tau^2}\right) \right]^{-1}$$

$$a_1(s) \equiv \frac{m_\tau^2}{6|V_{ud}|^2 S_{EW}} \frac{B(\tau^- \rightarrow A^- \nu_\tau)}{B(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)} \frac{dN_A}{N_A ds} \left[\left(1 - \frac{s}{m_\tau^2}\right)^2 \left(1 + \frac{2s}{m_\tau^2}\right) \right]^{-1}$$

$$a_0(s) \equiv \frac{m_\tau^2}{6|V_{ud}|^2 S_{EW}} \frac{B(\tau^- \rightarrow \pi^- \nu_\tau)}{B(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)} \frac{dN_A}{N_A ds} \left(1 - \frac{s}{m_\tau^2}\right)^2$$

The ALEPH Collaboration, Phys. Rept., 421:191-284, 2005

- Exp. measured Spectral function
- Normalized invariant mass-squared distribution divided by kinematic factor
- No $J = 0$ contribution (vector current conserved)
- Only contribution to a_0 pion-pole (why do we add this manually?)

Methodology

Fitting

Chi squared

$$\chi^2(\alpha) = (I_i^{exp} - I_i^{th}(\alpha))C_{ij}^{-1}(I_j^{exp} - I_j^{th}(\alpha))$$

$$I_{i=kl}^{exp}(s_k, \omega_l) = \int_{s_{th}}^{s_k} \frac{ds}{s_k} \omega_l(s) \operatorname{Im} \Pi_{V/A}(s)$$

$$I_{i=kl}^{th}(s_k, \omega_l) = \frac{i}{2s_k} \oint_{|s|=s_k} \frac{ds}{s} [W_l(s) - W_l(s_k)] D(s)$$

Parameters

$$a_s \quad \langle a_s GG \rangle \quad \mathcal{O}_{6,V+A} \quad \mathcal{O}_{8,V+A}$$

Different s_0

# (k, l)	2 Moments	
1 (1,1)	s_1	w_1
2 (2,1)	s_2	w_1

Different weights

# (k, l)	3 Moments	
1 (1,1)	s_1	w_1
2 (1,2)	s_1	w_2
2 (1,3)	s_1	w_3

E.g.

Extract max. 2 parameters

max. 3 parameters

- χ^2 : difference exp. th, averaged with correlation matrix inverse
- Integral moments
- Fitted parameters: 4, 12 with DV
- Different combinations => different results
- Different s_0
- Different weights (high correlations!)
- Different s_0 & weights

Fitting

Weights

$$\frac{1}{2\pi i s_0} \oint_{|s|=s_0} ds \left(\frac{s}{s_0}\right)^n \frac{C_{2k}}{(-s)^k} = (-1)^{n+1} \frac{C_{2(n+1)}}{s_0^{n+1}} \delta_{k,n+1}$$

$$\omega_{kl}(s) = \left(1 - \frac{s}{m_\tau^2}\right)^{2+k} \left(\frac{s}{m_\tau^2}\right)^l \left(1 + \frac{2s}{m_\tau^2}\right)$$

$(k, l) = (0, 0), (1, 0), (1, 1), (1, 2), (1, 3)$ Double Pinched

n-th degreee monomial in the weight selects the $D = 2(n+1)$

$$\frac{s}{s_0} + \left(\frac{s}{s_0}\right)^2 + \left(\frac{s}{s_0}\right)^3 + \dots$$

\downarrow \downarrow \downarrow

$\mathcal{O}_{4,V/A}$ $\mathcal{O}_{6,V/A}$ $\mathcal{O}_{8,V/A}$

$$A_{00,V/A}^{ALEPH} = A_{00,V/A}^{ALEPH}(a_s, \mathcal{O}_{6,V/A})$$

$$A_{10,V/A}^{ALEPH} = A_{10,V/A}^{ALEPH}(a_s, \langle a_s GG \rangle, \mathcal{O}_{6,V/A}, \mathcal{O}_{10,V/A})$$

$$A_{11,V/A}^{ALEPH} = A_{11,V/A}^{ALEPH}(a_s, \langle a_s GG \rangle, \mathcal{O}_{6,V/A}, \mathcal{O}_{10,V/A} \mathcal{O}_{12,V/A})$$

$$A_{12,V/A}^{ALEPH} = A_{12,V/A}^{ALEPH}(a_s, \mathcal{O}_{6,V/A}, \mathcal{O}_{10,V/A}, \mathcal{O}_{12,V/A}, \mathcal{O}_{14,V/A})$$

$$A_{13,V/A}^{ALEPH} = A_{13,V/A}^{ALEPH}(a_s, \mathcal{O}_{10,V/A}, \mathcal{O}_{12,V/A}, \mathcal{O}_{14,V/A}, \mathcal{O}_{16,V/A})$$

- weights used for “activation” of condensates
- Standard definition w_{kl}
- Pinched weight: $1-s/m_\tau$ reduce DV (the closer to mtau the less? Get DV stronger with energy?)
- Some possible weights and their parameters

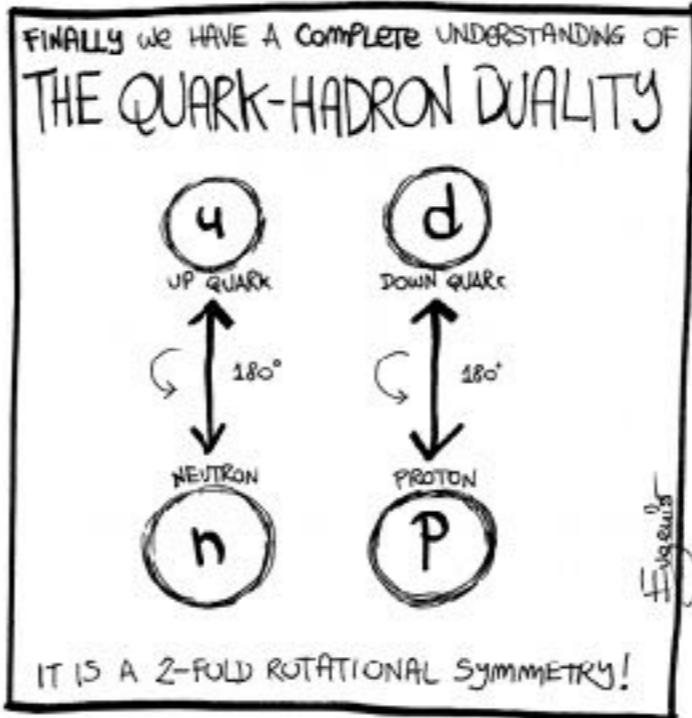
Violations

Duality

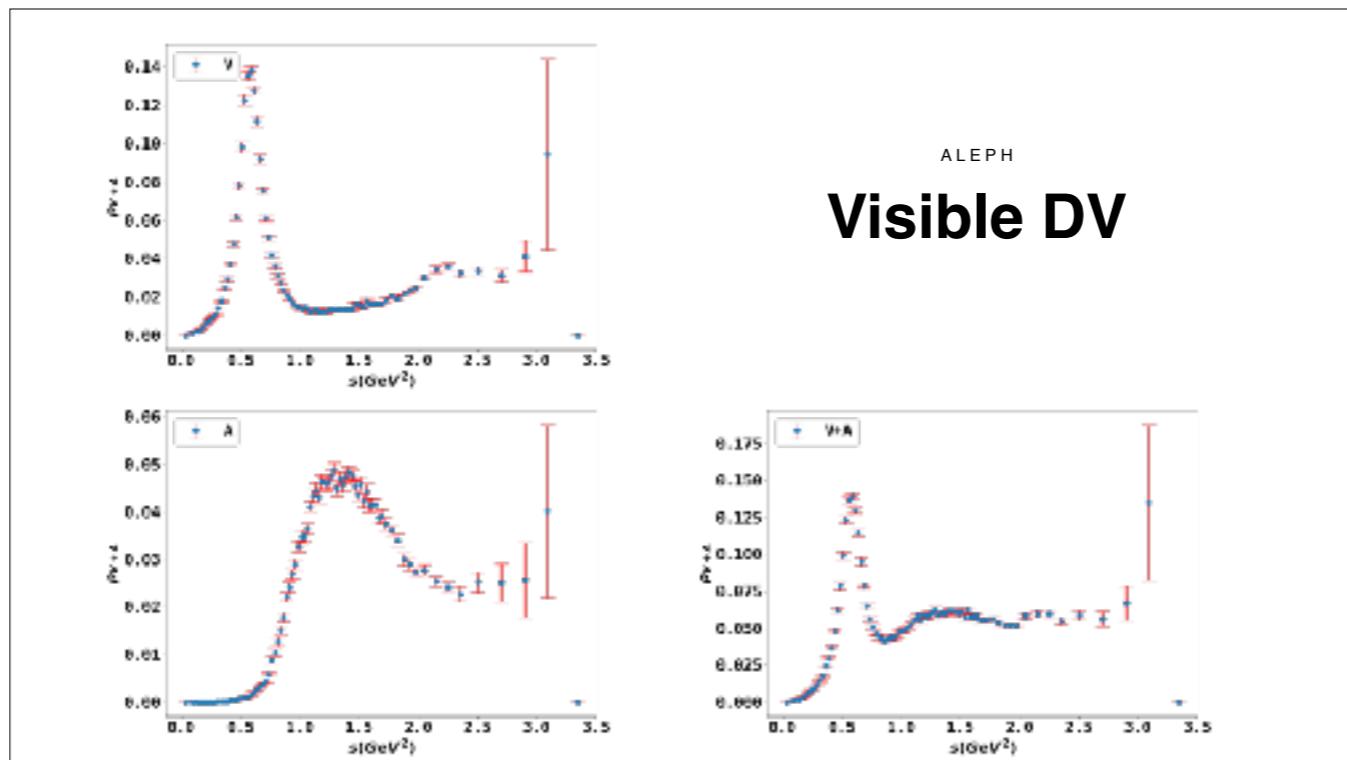
Experiment
Hadrons Theory
Quark/ Gluons
 $\rho_{\text{exp}} = \rho_{\text{th}}$

- What energy is high enough for the quark-hadron duality to set in?

M. Shiftman, hep-th/0009131, 2000



- measure hadrons but calculate with quarks/ gluons?
- Are they really the same?
- Adhoc: $\rho_{\text{exp}} = \rho_{\text{th}}$
- n, p as Hadrons, flip quarks and they are the same



- V,A: DV clearly visible!
- PT cannot reproduce resonances
- Exponential suppressed
- Sinusoidal
- V+A: less DV

Duality

Model

Model:

8 Parameters

$$\Delta\rho_{V/A}^{DV}(s) = e^{-\delta_{V/A} + \gamma_{V/A}s} \sin(\alpha_{V/A} + \beta_{V/A}s)$$

$$\Delta A_{V/A}^{\omega, DV}(s_0) \equiv \frac{i}{2} \oint_{|s_0|=s_0} \frac{ds}{s_0} \omega(s) \left\{ \Pi_{V/A}(s) - \Pi_{V/A}^{OPE}(s) \right\} = -\pi \int_{s_0}^{\infty} \frac{ds}{s_0} \omega(s) \Delta\rho_{V/A}^{DV}(s)$$

O. Catà, Maarten Golterman, Santiago Peris, arxiv:0812.2285v2 (2008)

- V: 4 params + A: 4 params $A = 8$ params
- Add DV to theory side (numerically not easy)

Timeline

- D. Boito
- M. Goltermann
- S. Peris
- K Maltman



- A. Pich
- A. Rodríguez-Sánchez



Philosophy

$\alpha_s = 0.322$ (CIPT), $\alpha_s = 0.307$ (FOPT)

Boito, Goltermann, Peris, Maltman

Pro

- Theoretically well motivated
- Systematic errors have been underestimated
- Assuming Duality a priori leads to worse model

$\alpha_s = 0.348$ (CIPT), $\alpha_s = 0.324$ (FOPT)

Pich, Rodríguez-Sánchez

Contra

- Cannot be derived from first principles
- Far too many parameters (8)
- Poor statistical quality (low p-value, large uncertainties)
- Sufficiently suppressed (pinched weights)

D. Boito, O. Catà, M. Golterman, M. Jamin, K. Maltman, J. Osborne, S. Peris, Phys. Rev., D84:113006, 2011

Architecture

Program

Important Frameworks:

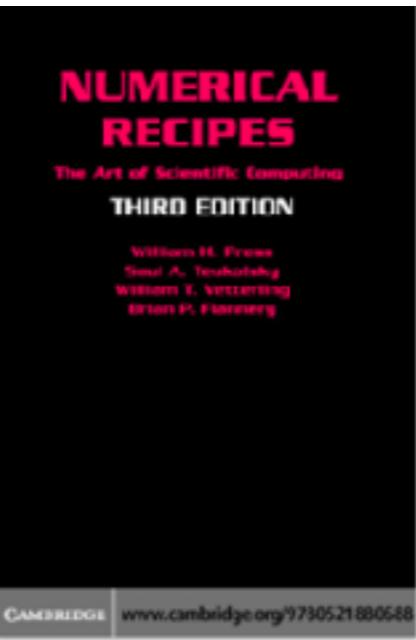
- GSL
- BOOST
- ROOT

Numerics:

- Integration: GSL - QAG adaptive integration
- Inverse Matrix: BOOST - UBLAS
- Non-linear Equation Solver: GSL - Multidimensional Root-Finding
- Minimization: ROOT - MINUIT2

Infos:

- Editor: Emacs
- Parallelization: use multiple cores
- GitHub

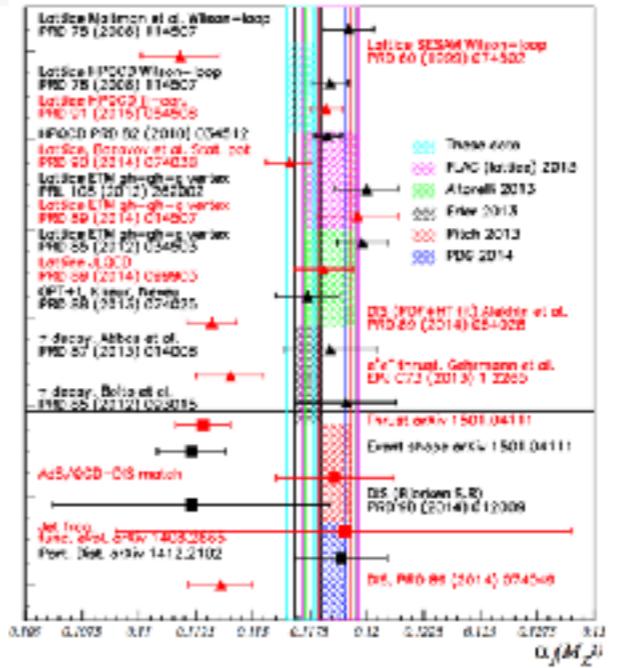


Determination of α_s

Channel	$\alpha_s(m_\tau^2)$	$\langle a_s GG \rangle$
V+A (FOPT)	0.319	-3
V+A (CIPT)	0.339	-16

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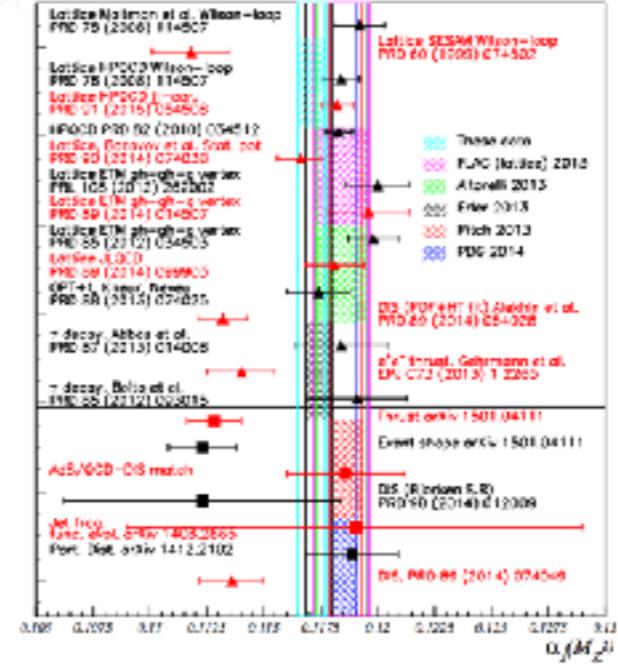


Summary

Channel	$\alpha_s(m_\tau^2)$	$\langle a_s GG \rangle$
V+A (FOPT)	0.319	-3
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V+A (FOPT)	0.319	-3
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- try to fit without DV