Simulation and High-Performance Computing Part 7: Iterative Methods

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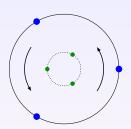
October 1st, 2020

Example: Lagrange points

Setting: We have bodies with masses m_1, \ldots, m_n in points $x_1, \ldots, x_n \in \mathbb{R}^2$ in a coordinate system rotating at an angular velocity of α .

Lagrange points: There are points $z \in \mathbb{R}^2$ where the gravitational forces and the centrifugal force cancel each other out.

Stable Lagrange points tend to accumulate trojan asteroids of interest to astronomers.

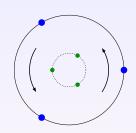


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Equilibrium of forces described by

$$\alpha^2 z + \gamma \sum_{i=1}^n m_i \frac{x_i - z}{\|x_i - z\|^3} = 0.$$

Nonlinear equations like this are generally solved by numerical algorithms.

Iterations

Nonlinear systems can always be written in the form

$$f(x) = 0$$

with a function $f: \mathbb{R}^n \to \mathbb{R}^n$.

Problem: Frequently impossible to compute exact solution.

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Problem: Frequently impossible to compute exact solution.

Idea: Starting with an initial guess $x_0 \in \mathbb{R}^n$, improve it by applying an iteration function $\Phi \colon \mathbb{R}^n \to \mathbb{R}^n$.

$$x_{m+1} := \Phi(x_m).$$

This procedure may never yield the exact solution, but it can provide us with arbitrarily accurate approximations.

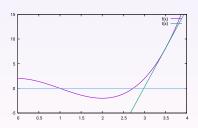
Newton iteration

Idea: Approximate f by its tangent

$$f(x) \approx f(x_m) + f'(x_m)(x - x_m),$$

look for the tangent's zero.

$$0 = f(x_m) + f'(x_m)(x_{m+1} - x_m)$$
 \iff $x_{m+1} = x_m - \frac{f(x_m)}{f'(x_m)}$



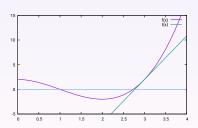
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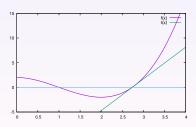
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Taylor expansion around x_m yields $\eta \in \mathbb{R}$ with

$$0 = f(x) = f(x_m) + f'(x_m)(x - x_m) + \frac{f''(\eta)}{2}(x - x_m)^2,$$

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$$0 = \frac{f(x_m)}{f'(x_m)} + x - x_m + \frac{f''(\eta)}{2f'(x_m)}(x - x_m)^2,$$

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Convergence: We assume $\frac{|f''(y)|}{2|f'(z)|} \le c$ and obtain $|x_{m+1} - x| \le c |x_m - x|^2$.

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Example: If
$$c|x_0 - x| \leq \frac{1}{2}$$
, we find

$$c |x_1 - x| \le c^2 |x_0 - x|^2 \le \frac{1}{4},$$
 $c |x_2 - x| \le c^2 |x_1 - x|^2 \le \frac{1}{16},$ $c |x_3 - x| \le c^2 |x_2 - x|^2 \le \frac{1}{256},$

Taylor expansion around x_m yields $\eta \in \mathbb{R}$ with

$$0 = f(x) = f(x_m) + f'(x_m)(x - x_m) + \frac{f''(\eta)}{2}(x - x_m)^2,$$

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$$x_{m+1} - x = \frac{f''(\eta)}{2f'(x_m)}(x - x_m)^2.$$

Convergence: We assume $\frac{|f''(y)|}{2|f'(z)|} \le c$ and obtain $|x_{m+1} - x| \le c |x_m - x|^2$.

Example: If $c|x_0 - x| \leq \frac{1}{2}$, we find

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 $c |x_2 - x| \le c^2 |x_1 - x|^2 \le \frac{1}{16},$ $c |x_3 - x| \le c^2 |x_2 - x|^2 \le \frac{1}{256},$ $c |x_m - x| \le 2^{-2^m}.$

Experiment: Square root

Approach: Square root of $a \in \mathbb{R}_{>0}$ is a zero of $f(x) = x^2 - a$.

$$x_{m+1} = x_m - \frac{f(x_m)}{f'(x_m)} = x_m - \frac{x_m^2 - a}{2x_m} = \frac{1}{2} \left(x_m + \frac{a}{x_m} \right).$$

m	$x_0 = 2$	$x_0 = 10$	$x_0 = -5$
1	8.58_{-2}	3.69+0	1.29+0
2	2.45_{-3}	1.33_{+0}	3.06_{-1}
3	2.12_{-6}	3.23_{-1}	2.72_{-2}
4	1.59_{-12}	3.00_{-2}	2.57_{-4}
5		3.12_{-4}	2.34_{-8}
6		3.44_{-8}	2.22_{-16}
7		2.22_{-16}	

Observation: Convergence quadratic once accurate enough.

Experiment: Reciprocal square root

Approach: Reciprocal square root of $a \in \mathbb{R}_{>0}$ is a zero of $f(x) = \frac{1}{x^2} - a$.

$$x_{m+1} = x_m - \frac{f(x_m)}{f'(x_m)} = x_m - \frac{x_m^{-2} - a}{-2 x_m^{-3}} = \frac{x_m}{2} (3 - a x_m^2).$$

m	$x_0 = 1$	$x_0 = 1.5$	$x_0 = 2$
-111		-	<i>∧</i> 0 − ∠
1	2.07_{-1}	4.18_{-1}	4.29_{+0}
2	8.21_{-2}	4.43_{-1}	1.18_{+2}
3	1.37_{-2}	3.30_{-1}	1.62_{+6}
4	3.98_{-4}	1.95_{-1}	4.27_{+18}
5	3.36_7	7.32_{-2}	7.77_{+55}
6	2.40_{-13}	1.10_{-2}	4.70_{+167}
7	1.11_{-16}	2.55_{-4}	∞
8	0	1.38_{-7}	

Observation: Convergence quadratic once accurate enough, no convergence for unsuitable initial guesses x_0 .

Multidimensional Newton

Idea: Replace the tangent by tangential hyperplanes.

$$0 = f(x) \approx f(x_m) + Df(x_m)(x - x_m)$$

with the Jacobian matrix

$$Df(z) = \begin{pmatrix} \frac{\partial f_1}{\partial z_1}(z) & \cdots & \frac{\partial f_1}{\partial z_n}(z) \\ \vdots & \ddots & \vdots \\ \frac{\partial f_n}{\partial z_1}(z) & \cdots & \frac{\partial f_n}{\partial z_n}(z) \end{pmatrix}.$$

Result: Solving the approximated problem yields

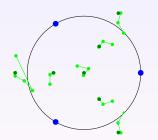
$$x_{m+1} = x_m - Df(x_m)^{-1}f(x_m).$$

We have to solve a linear system of equations in each step.

Experiment: Lagrange points

Goal: Find solutions of

$$\alpha^2 x + \gamma \sum_{i=1}^n \frac{y_i - x}{\|y_i - x\|^3} = 0.$$



Observation: Different starting points lead to different results.

Example: Resonance frequencies

Goal: Find non-trivial solution of the wave equation

$$\frac{\partial^2 u}{\partial t^2}(t,x) = c\Delta u(t,x).$$

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Approach: Separate variables $u(t,x) = \cos(\omega t) e(x)$.

$$\omega^2 e(x) = -c\Delta e(x).$$

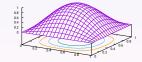
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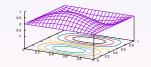
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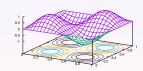
$$\frac{\partial^2 u}{\partial t^2}(t,x) = c\Delta u(t,x).$$

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$$\omega^2 e(x) = -c\Delta e(x).$$







Eigenvalue problems

Task: Given a matrix $A \in \mathbb{R}^{n \times n}$, find $\lambda \in \mathbb{R}$ and $e \in \mathbb{R}^n$ with

$$Ae = \lambda e$$
,

$$e \neq 0$$

 λ is called an eigenvalue, e an eigenvector.

Variations:

- Find the smallest or largest eigenvalue and a matching eigenvector.
- Find an entire basis of eigenvectors.
- Find a basis for an invariant subspace.

Power iteration

Approach: If a vector $x \in \mathbb{R}^n$ can be written as

$$x = e_1 + e_2 + \ldots + e_k$$

with eigenvectors e_1, \ldots, e_k for eigenvalues $\lambda_1, \ldots, \lambda_k$, we have

$$A^m x = A^m e_1 + A^m e_2 + \dots + A^m e_k$$

= $\lambda_1^m e_1 + \lambda_2^m e_2 + \dots + \lambda_k^m e_k$.

Idea: If $|\lambda_1| > |\lambda_2|, \ldots, |\lambda_k|$, the first component in the sum will dominate if m is large enough, i.e., we will have an approximation of e_1 .

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Practice: To avoid an overflow, compute the normalized sequence

$$\hat{x}_{m+1} := Ax_m, \qquad x_{m+1} := \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$$

Approach: If e_1, \ldots, e_k is an orthogonal basis of eigenvectors with eigenvalues $\lambda_1, \ldots, \lambda_k$ satisfying $|\lambda_1| \geq |\lambda_2| \geq \ldots \geq |\lambda_k|$, we can bound the angle between x_m and e_1 by

$$\tan^2 \angle (x_m, e) = \frac{\|\lambda_2^m e_2 + \ldots + \lambda_k^m e_k\|^2}{\|\lambda_1^m e_1\|^2}$$

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$$= \frac{|\lambda_{2}|^{2m} \|e_{2}\|^{2} + \ldots + |\lambda_{k}|^{2m} \|e_{k}\|^{2}}{|\lambda_{1}|^{2m} \|e_{1}\|^{2}}$$

$$\leq \frac{|\lambda_{2}|^{2m} \|e_{2}\|^{2} + \ldots + |\lambda_{2}|^{2m} \|e_{k}\|^{2}}{|\lambda_{1}|^{2m} \|e_{1}\|^{2}}$$

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$$\begin{split} \tan^2 \angle (x_m, e) &= \frac{\|\lambda_2^m e_2 + \ldots + \lambda_k^m e_k\|^2}{\|\lambda_1^m e_1\|^2} \\ &= \frac{|\lambda_2|^{2m} \|e_2\|^2 + \ldots + |\lambda_k|^{2m} \|e_k\|^2}{|\lambda_1|^{2m} \|e_1\|^2} \\ &\leq \frac{|\lambda_2|^{2m} \|e_2\|^2 + \ldots + |\lambda_2|^{2m} \|e_k\|^2}{|\lambda_1|^{2m} \|e_1\|^2} \\ &= \frac{|\lambda_2|^{2m}}{|\lambda_1|^{2m}} \frac{\|e_2\|^2 + \ldots + \|e_k\|^2}{\|e_1\|^2} \end{split}$$

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$$\begin{split} \tan^2 \angle(x_m,e) &= \frac{\|\lambda_2^m e_2 + \ldots + \lambda_k^m e_k\|^2}{\|\lambda_1^m e_1\|^2} \\ &= \frac{|\lambda_2|^{2m} \|e_2\|^2 + \ldots + |\lambda_k|^{2m} \|e_k\|^2}{|\lambda_1|^{2m} \|e_1\|^2} \\ &\leq \frac{|\lambda_2|^{2m} \|e_2\|^2 + \ldots + |\lambda_2|^{2m} \|e_k\|^2}{|\lambda_1|^{2m} \|e_1\|^2} \\ &= \frac{|\lambda_2|^{2m}}{|\lambda_1|^{2m}} \frac{\|e_2\|^2 + \ldots + \|e_k\|^2}{\|e_1\|^2} = \left(\frac{|\lambda_2|}{|\lambda_1|}\right)^{2m} \tan^2 \angle(x_0,e). \end{split}$$

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Result: Convergence rate $\frac{|\lambda_2|}{|\lambda_1|}$ for the tangent.

Inverse iteration

Task: Frequently we need not the largest, but the smallest eigenvalue, e.g., when computing the fundamental frequency of a system.

Idea: Work with A^{-1} instead of A.

$$A^{-m}x = A^{-m}e_1 + A^{-m}e_2 + \dots + A^{-m}e_k$$

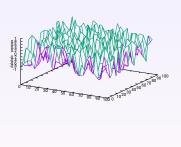
= $\frac{1}{\lambda_1^m}e_1 + \frac{1}{\lambda_2^m}e_2 + \dots + \frac{1}{\lambda_k^m}e_k$.

If $|\lambda_1| < |\lambda_2|, \dots, |\lambda_k|$, the first component in the sum will dominate if m is large enough.

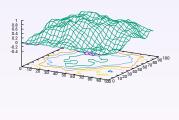
Practice: We have to solve a linear system in each step.

$$A\hat{x}_{m+1} = x_m,$$
 $x_{m+1} := \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$

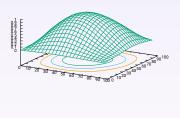
m	$\sin \alpha_m$	$\tan \alpha_{\it m}$	ratio
0	1.00_{+0}	6.97 ₊₁	



m	$\sin lpha_{\it m}$	$ an lpha_{\it m}$	ratio
0	1.00_{+0}	6.97_{+1}	
1	3.87_{-1}	4.19_{-1}	166.3

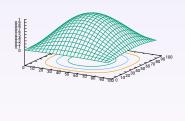


m	$\sin \alpha_m$	$ an lpha_{\it m}$	ratio
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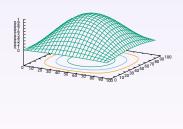
Goal: Find the smallest eigenvalue of the wave equation.

m	$\sin lpha_{\it m}$	$ an lpha_{\it m}$	ratio
0	1.00_{+0}	6.97_{+1}	
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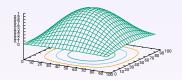
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1	3.87_{-1}	4.19_{-1}	166.3
2	1.20_{-1}	1.20_{-1}	3.5
3	4.46_{-2}	4.46_{-2}	2.7
4	1.73_{-2}	1.73_{-2}	2.6



Experiment: Inverse iteration

Goal: Find the smallest eigenvalue of the wave equation.

m	$\sin lpha_{\it m}$	$\tan lpha_{\it m}$	ratio
0	1.00_{+0}	6.97_{+1}	
1	3.87_{-1}	4.19_{-1}	166.3
2	1.20_{-1}	1.20_{-1}	3.5
3	4.46_{-2}	4.46_{-2}	2.7
4	1.73_{-2}	1.73_{-2}	2.6
5	6.86_{-3}	6.86_{-3}	2.5
6	2.73_{-3}	2.73_{-3}	2.5
7	1.09_{-3}	1.09_{-3}	2.5
8	4.36_{-4}	4.36_4	2.5



Inverse iteration with shift

Task: Sometimes we need the eigenvalue closest to a number μ , e.g., to find resonances close to the frequency of an incident wave.

Idea: Work with $(A - \mu I)^{-1}$ instead of A.

$$(A - \mu I)^{-m} x = (A - \mu I)^{-m} e_1 + (A - \mu I)^{-m} e_2 + \ldots + (A - \mu I)^{-m} e_k$$

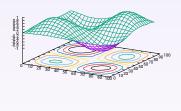
= $\frac{1}{(\lambda_1 - \mu)^m} e_1 + \frac{1}{(\lambda_2 - \mu)^m} e_2 + \ldots + \frac{1}{(\lambda_k - \mu)^m} e_k$.

If $|\lambda_1 - \mu| < |\lambda_2 - \mu|, \dots, |\lambda_k - \mu|$, the first component in the sum will dominate if m is large enough.

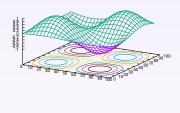
Practice: We have to solve a linear system in each step.

$$(A - \mu I)\hat{x}_{m+1} = x_m,$$
 $x_{m+1} := \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$

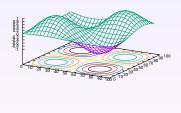
m	$\sin \alpha_m$	$\tan \alpha_{\it m}$	ratio
0	1.00_{+0}	8.30_{+1}	
1	1.73_{-1}	1.76_{-1}	472.6



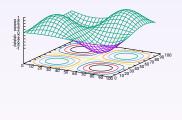
m	$\sin \alpha_m$	$\tan \alpha_{\it m}$	ratio
0	1.00_{+0}	8.30_{+1}	
1	1.73_{-1}	1.76_{-1}	472.6
2	2.08_{-2}	2.09_{-2}	8.4



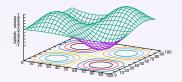
m	$\sin lpha_{\it m}$	$ an lpha_{\it m}$	ratio
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2	2.08_{-2}	2.09_{-2}	8.4
3	3.04_{-3}	3.04_{-3}	6.9



$\sin lpha_{\it m}$	$\tan \alpha_{\it m}$	ratio
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1.73_{-1}	1.76_{-1}	472.6
2.08_{-2}	2.09_{-2}	8.4
3.04_{-3}	3.04_3	6.9
4.64_{-4}	4.64_{-4}	6.6
	$ \begin{array}{r} 1.00_{+0} \\ 1.73_{-1} \\ 2.08_{-2} \\ 3.04_{-3} \end{array} $	1.00 ₊₀ 8.30 ₊₁ 1.73 ₋₁ 1.76 ₋₁ 2.08 ₋₂ 2.09 ₋₂ 3.04 ₋₃ 3.04 ₋₃



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2	2.08_{-2}	2.09_{-2}	8.4
3	3.04_{-3}	3.04_3	6.9
4	4.64_{-4}	4.64_{-4}	6.6
5	7.15_{-5}	7.15_{-5}	6.5
6	1.11_{-5}	1.11_{-5}	6.5
7	1.72_{-5}	1.72_{-6}	6.4
8	2.67_{-7}	2.67_7	6.4



Rayleigh quotient

Task: Assuming that we have an approximation x of an eigenvector e, how do we find an approximation of the corresponding eigenvalue λ ?

Idea: Use the inner product. For the exact eigenvector e, we have

$$Ae = \lambda e, \qquad \langle e, Ae
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Rayleigh quotient: For the approximated eigenvector, we find

$$\frac{\langle x, Ax \rangle}{\langle x, x \rangle} \approx \lambda.$$

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Stopping criterion: The Rayleigh quotient can be used to estimate the accuracy of an approximation by checking

$$\|Ax - \tilde{\lambda}x\| \leq \epsilon \, \|x\| \qquad \text{ with } \qquad \tilde{\lambda} := \frac{\langle x, Ax \rangle}{\langle x, x \rangle}.$$

Rayleigh iteration

Task: Find good shift parameters μ for the inverse iteration.

Idea: Use the Rayleigh quotient, since it provides approximate eigenvalues.

$$\mu_m := \frac{\langle x_m, Ax_m \rangle}{\langle x_m, x_m \rangle}, \qquad (A - \mu_m I) \hat{x}_{m+1} = x_m, \qquad x_{m+1} := \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$$

Rayleigh iteration

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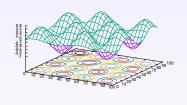
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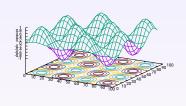
Consequences:

- In each step, a different matrix $A \mu_m I$ has to be used.
- If the initial vector x_0 is close enough to the desired eigenvector, we obtain quadratic or even cubic convergence.
- If the initial vector x_0 is not close enough to the desired eigenvector, we may observe convergence to another eigenvector or even no convergence at all.

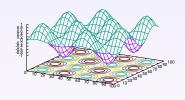
μ	$\sin lpha_{\it m}$	$ an lpha_{\it m}$	ratio
3.10_{+2}	1.00_{+0}	5.43 ₊₁	
3.20_{+2}	2.03_{-1}	2.07_{-1}	261.9
	,	3.10_{+2} 1.00_{+0}	3.10_{+2} 1.00_{+0} 5.43_{+1}



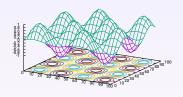
m	μ	$\sin lpha_{m}$	$ an lpha_{\it m}$	ratio
0	3.10_{+2}	1.00_{+0}	5.43 ₊₁	
1	3.20_{+2}	2.03_{-1}	2.07_{-1}	261.9
2	3.15_{+2}	4.10_{-2}	4.10_{-2}	5.1



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1	3.20_{+2}	2.03_{-1}	2.07_{-1}	261.9
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3	3.15_{+2}	3.85 ₋₅	3.84 ₋₅	1067.6



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2	3.15_{+2}	4.10_{-2}	4.10_{-2}	5.1
3	3.15_{+2}	3.85 ₋₅	3.84 ₋₅	1067.6
4	3.15_{+2}	3.06_{-13}	3.06_{-13}	1.25_{+8}



Summary

Newton iteration can be used to solve nonlinear equations f(x) = 0.

$$x_{m+1} = x_m - \frac{f(x_m)}{f'(x_m)},$$
 $x_{m+1} = x_m - Df(x_m)^{-1}f(x_m).$

If x_m is sufficiently close to the solution, we have quadratic convergence.

Power iteration can be used to solve eigenvalue problems $Ae = \lambda e$.

$$\hat{x}_{m+1} = Ax_m,$$
 $x_{m+1} = \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$

Inverse iteration can be used to choose specific eigenvalues.

$$(A - \mu I)\hat{x}_{m+1} = x_m,$$
 $x_{m+1} = \frac{\hat{x}_{m+1}}{\|\hat{x}_{m+1}\|}.$

Rayleigh quotient provides an approximation of an eigenvalue.