

Homework 8

Philip Westphal

UW NetID: philw16@uw.edu

Problem 1

Solve the initial-boundary value problem

$$\begin{aligned} u_{tt} &= 9u_{xx} + 2\delta(x) \text{ for all } t > 0 \text{ and } -\infty < x < \infty \\ u(0, x) &= e^{-x^2} \text{ and } u_t(0, x) = \frac{1}{1+x^2} \text{ for all } -\infty < x < \infty \end{aligned}$$

Plot your solution at several representative times.

Solution.

Considering the following general problem

$$\begin{aligned} u_{tt} &= c^2 u_{xx} + h(x) \text{ for all } t > 0 \text{ and } -\infty < x < \infty \\ u(0, x) &= f(x) \text{ and } u_t(0, x) = g(x) \text{ for all } -\infty < x < \infty \end{aligned}$$

we will solve three cases, one with each of $h(x)$, $f(x)$, and $g(x)$ non-zero while the other functions are left as 0. We can then superimpose the solutions together to get the full solution. First, we consider the case where $f(x) = \exp(-x^2)$ and $g(x) = h(x) = 0$. From derivations in class, we know the general solution for a problem of this type is given by

$$u(t, x) = \frac{1}{2}[f(x+ct) + f(x-ct)]$$

Thus, for our problem we have

$$u(t, x) = \frac{1}{2}[\exp(-(x+3t)^2) + \exp(-(x-3t)^2)]$$

Next, we consider the case where $g(x) = \frac{1}{1+x^2}$ and $f(x) = h(x) = 0$. From Homework 5, we know the general solution to this problem can be found as

$$u(t, x) = \frac{1}{2c} \int_{x-ct}^{x+ct} g(y) dy$$

For our problem, we then have

$$\begin{aligned}
u(t, x) &= \frac{1}{6} \int_{x-3t}^{x+3t} \frac{1}{1+y^2} dy \\
&= \frac{1}{6} \arctan(y) \Big|_{x-3t}^{x+3t} \\
u(t, x) &= \frac{1}{6} [\arctan(x+3t) - \arctan(x-3t)]
\end{aligned}$$

Finally, we consider the case where $h(x) = 2\delta(x)$ and $f(x) = g(x) = 0$. From class, the general solution to this problem will be of the form

$$u(t, x) = \int_0^\infty \int_{-\infty}^\infty G(t, x; \tau, y) h(\tau, y) dy d\tau$$

We observe that as we have a delta function for $h(x)$, our integral can be simplified before considering $G(t, x; \tau, y)$.

$$\begin{aligned}
u(t, x) &= 2 \int_0^\infty \int_{-\infty}^\infty G(t, x; \tau, y) \delta(y) dy d\tau \\
&= 2 \int_0^\infty G(t, x; \tau, 0) d\tau \\
&= 2 \int_0^t G(t, x; \tau, 0) d\tau \leftarrow \text{Noting integral will be 0 for } \tau \leq t
\end{aligned}$$

From class, we know $G(t, x; \tau, 0)$ can be given by

$$\begin{aligned}
G(t, x; \tau, 0) &= \begin{cases} \frac{1}{2c} & x > -c(t-\tau) \text{ and } x < c(t-\tau) \text{ and } \tau < t \\ 0 & \text{otherwise} \end{cases} \\
&= \begin{cases} \frac{1}{2c} & \tau < t + x/c \text{ and } \tau < t - x/c \text{ and } \tau < t \\ 0 & \text{otherwise} \end{cases}
\end{aligned}$$

For our problem, this will be equivalent to

$$G(t, x; \tau, 0) = \begin{cases} \frac{1}{6} & \tau < t + x/3 \text{ and } \tau < t - x/3 \text{ and } \tau < t \\ 0 & \text{otherwise} \end{cases}$$

We can consider this problem in two parts to find where G is non-zero. First, if $x > 0$, we note that $\tau < t + \frac{x}{3}$ will always be satisfied as the right hand side of the inequality is growing. Thus, we need to consider the more restrictive inequality $\tau < t - \frac{x}{3}$. We find the boundary where $\tau = 0$:

$$0 = t - \frac{x}{3} \rightarrow x = 3t$$

Thus, if $x \geq 3t$ then $G = 0$. For $x > 0$ & $x < 3t$, we have:

$$u(t, x) = 2 \int_0^{t-\frac{x}{3}} \frac{1}{6} d\tau = \frac{t}{3} - \frac{x}{9}$$

We use a similar logic if $x < 0$, however in this case the more restrictive inequality will be $\tau < t + \frac{x}{3}$.

We again find the boundary where $\tau = 0$

$$0 = t + \frac{x}{3} \rightarrow x = -3t$$

Therefore if $x \leq -3t$ than $G = 0$. For $x < 0$ & $x > -3t$, we have:

$$u(t, x) = 2 \int_0^{t+\frac{x}{3}} \frac{1}{6} d\tau = \frac{t}{3} + \frac{x}{9}$$

If $x = 0$, we have

$$u(t, x) = 2 \int_0^t \frac{1}{6} d\tau = \frac{t}{3}$$

Combining all our solutions, we have the following full solution:

$$u(t, x) = \frac{1}{2} [\exp(-(x+3t)^2) + \exp(-(x-3t)^2)] + \frac{1}{6} [\arctan(x+3t) - \arctan(x-3t)] + \begin{cases} \frac{t}{3} - \frac{x}{9} & 0 \leq x < 3t \\ \frac{t}{3} + \frac{x}{9} & -3t < x \leq 0 \end{cases}$$

Figure Some plots of $u(t, x)$ are shown in figure 1

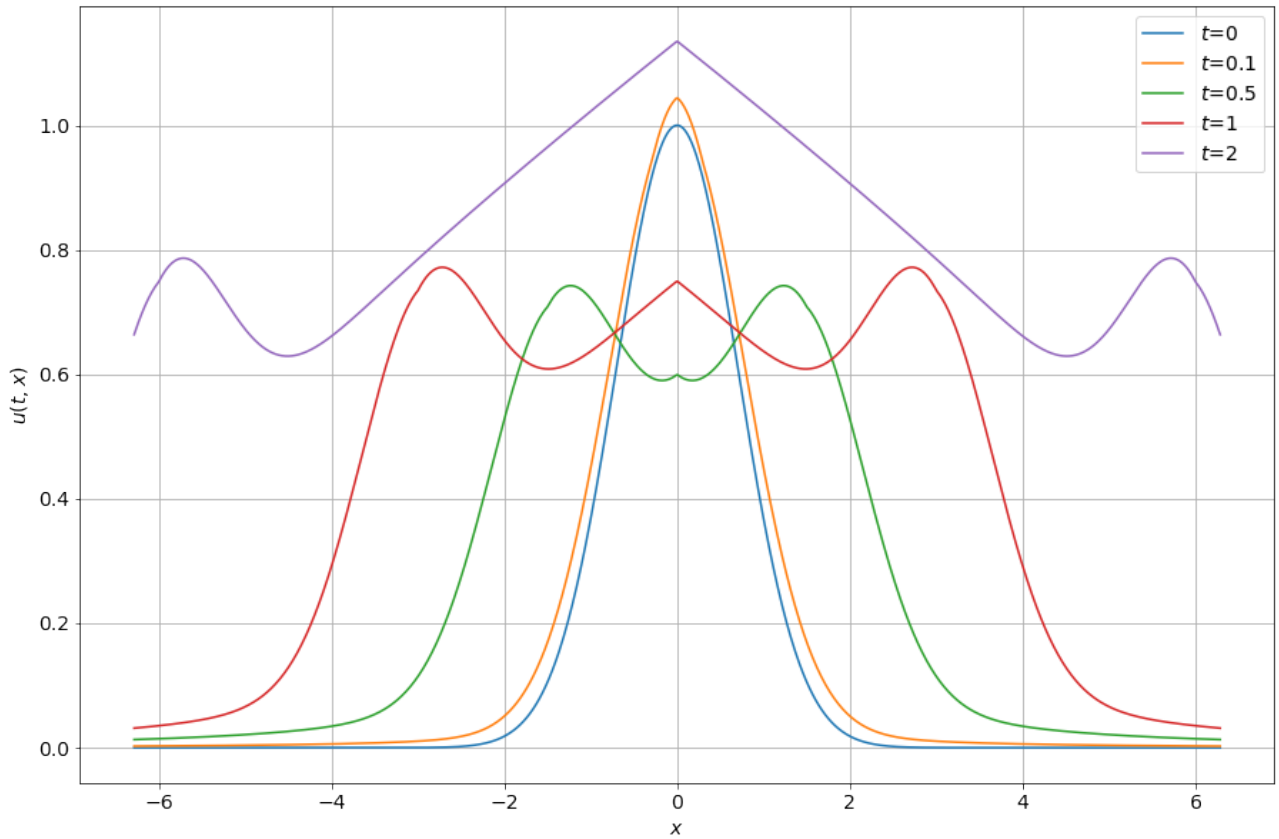


Figure 1: Solution $u(t, x)$ at various t for $x \in [-2\pi, 2\pi]$

Problem 2

Define the operator $L : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ by

$$L[\mathbf{x}] = \begin{pmatrix} 1 & 2 \\ -1 & 3 \end{pmatrix} \mathbf{x}$$

Find the adjoint L^* for each of the following inner products. (In each case, the same inner product is used on both the domain and range.)

Part 1

The standard Euclidean inner product.

Solution.

We can write the inner product of two vectors in the following form, where A is a matrix

$$\langle x, y \rangle = x^T A y$$

For the standard Euclidean inner product, A is equal to the identity matrix. We know an adjoint operator can be found using the following definition

$$\langle L[u], v \rangle = \langle u, L^*[v] \rangle$$

As the transform given in the problem is a matrix multiplication, we can define the transformation as Lx , where L is the matrix

$$L = \begin{bmatrix} 1 & 2 \\ -1 & 3 \end{bmatrix}$$

The corresponding adjoint operator L^* will also be a matrix multiplication. Substituting our inner product definition from above, we can rewrite the adjoint equality in the following form.

$$\begin{aligned} \langle L[u], v \rangle &= \langle u, L^*[v] \rangle \\ (Lu)^T Av &= u^T AL^*v \\ u^T L^T Av &= u^T AL^*v \\ 0 &= u^T L^T Av - u^T AL^*v \\ 0 &= u^T (L^T A - AL^*)v \end{aligned}$$

In order for this to always hold, we require

$$L^T A - AL^* = 0$$

Therefore, we know that

$$L^* = A^{-1} L^T A$$

For the Euclidean inner product, we know A will be the identity matrix. Thus

$$L^* = L^T = \begin{bmatrix} 1 & -1 \\ 2 & 3 \end{bmatrix}$$

Part 2

$$\langle \mathbf{x}, \mathbf{y} \rangle = 2x_1y_1 + 3x_2y_2$$

Solution.

We observe that this inner product can be written as

$$\langle x, y \rangle = x^T \begin{bmatrix} 2 & 0 \\ 0 & 3 \end{bmatrix} y$$

Thus for this inner product, our A matrix is given by

$$A = \begin{bmatrix} 2 & 0 \\ 0 & 3 \end{bmatrix}$$

We can find A^{-1} as follows

$$A^{-1} = \frac{1}{6-0} \begin{bmatrix} 3 & 0 \\ 0 & 2 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{3} \end{bmatrix}$$

And $L^T A$ as

$$L^T A = \begin{bmatrix} 1 & -1 \\ 2 & 3 \end{bmatrix} \begin{bmatrix} 2 & 0 \\ 0 & 3 \end{bmatrix} = \begin{bmatrix} 2 & -3 \\ 4 & 9 \end{bmatrix}$$

Therefore

$$L^* = \begin{bmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{3} \end{bmatrix} \begin{bmatrix} 2 & -3 \\ 4 & 9 \end{bmatrix} = \begin{bmatrix} 1 & -\frac{3}{2} \\ \frac{4}{3} & 3 \end{bmatrix}$$

Part 3

$$\langle \mathbf{x}, \mathbf{y} \rangle = \mathbf{x}^T \begin{pmatrix} 2 & -1 \\ -1 & 4 \end{pmatrix} \mathbf{y}$$

Solution.

In this problem, we are given the matrix A in the inner product definition. We can find A^{-1} as

$$A^{-1} = \frac{1}{8-1} \begin{bmatrix} 4 & 1 \\ 1 & 2 \end{bmatrix} = \begin{bmatrix} \frac{4}{7} & \frac{1}{7} \\ \frac{1}{7} & \frac{2}{7} \end{bmatrix}$$

And $L^T A$ as

$$L^T A = \begin{bmatrix} 1 & -1 \\ 2 & 3 \end{bmatrix} \begin{bmatrix} 2 & -1 \\ -1 & 4 \end{bmatrix} = \begin{bmatrix} 3 & -5 \\ 1 & 10 \end{bmatrix}$$

Therefore

$$L^* = \begin{bmatrix} \frac{4}{7} & \frac{1}{7} \\ \frac{1}{7} & \frac{2}{7} \end{bmatrix} \begin{bmatrix} 3 & -5 \\ 1 & 10 \end{bmatrix} = \begin{bmatrix} \frac{13}{7} & -\frac{10}{7} \\ \frac{5}{7} & \frac{15}{7} \end{bmatrix}$$

Problem 3

Define the operator $L : U \rightarrow U$, where

$$U = \{C^1([0, 1]) \mid u(0) = u(1) = 0\}$$

and $L[u] = u_x$. Use the weighted inner product

$$\langle u_1, u_2 \rangle = \int_0^1 u_1(x) u_2(x) e^x dx$$

on both the domain and range.

Part 1

Find the adjoint operator L^* .

Solution.

To find the adjoint operator, we utilize the inner product equality and convert it to an integral form

$$\begin{aligned} \langle L[u], v \rangle &= \langle u, L^*[v] \rangle \\ \int_0^1 \frac{du}{dx} v(x) \exp(x) dx &= \int_0^1 u(x) L^*[v] \exp(x) dx \end{aligned}$$

Our goal is to get the left hand side of the equation into a form that matches the right hand side. We can evaluate the left hand side using integration by parts. Let

$$u = v(x) \exp(x) \quad du = \frac{d}{dx}(v(x) \exp(x)) \quad dv = \frac{du}{\exp(x)} \quad v = u(x)$$

$$\begin{aligned} LHS &= u(x)v(x) \exp(x) \Big|_0^1 - \int_0^1 u(x) \frac{d}{dx}(v(x) \exp(x)) dx \\ &= [u(1)v(1)\cancel{e} - \cancel{u(0)v(0)}] + \int_0^1 u(x) \left(-\frac{1}{\exp(x)}\right) \frac{d}{dx}(v(x) \exp(x)) \exp(x) dx \\ &= \int_0^1 u(x) \left[-\frac{1}{\exp(x)} \frac{d}{dx}(v(x) \exp(x))\right] \exp(x) dx \end{aligned}$$

We see this is now in a pattern that matches the right hand side of our inner product definition. Thus

$$L^*[v] = -\frac{1}{\exp(x)} \frac{d}{dx}(v(x) \exp(x)) = -\frac{1}{\exp(x)} \left(v(x) \exp(x) + \exp(x) \frac{dv}{dx} \right) = v(x) - \frac{dv}{dx}$$

Therefore, the operator L^* can be defined as

$$\boxed{L^* = -1 - \partial_x}$$

Part 2

Find the operator $S = L^* \circ L$.

Solution.

We can define $S[u]$ as

$$\begin{aligned} S[u] &= L^*[L[u]] \\ &= (-1 - \partial_x)(\partial_x(u)) \\ &= (-1 - \partial_x)(u_x) \\ &= -u_x - u_{xx} \end{aligned}$$

Therefore

$$\boxed{S = -\partial_x - \partial_{xx}}$$

Part 3

Solve the boundary value problem $S[u] = 2e^x$.

Solution.

Substituting for $S[u]$ from Part 2, we have the following relation

$$\begin{aligned} -u_x - u_{xx} &= 2\exp(x) \\ u_{xx} + u_x &= -2\exp(x) \end{aligned}$$

This is a second order ODE that we can solve with the the principle of superposition. First, we find the solution to the homogeneous problem using the characteristic equation

$$\begin{aligned} u_{xx} + u_x &= 0 \\ r^2 + r &= 0 \\ r(r + 1) &= 0 \rightarrow r = 0, -1 \\ \therefore u_{gen} &= c_1 + c_2 \exp(-x) \end{aligned}$$

Next, we look for a particular solution using the method of undetermined coefficients. We assume our particular solution is of the form

$$u_p(x) = A \exp(x)$$

Substituting into our ODE, we have

$$A \exp(x) + A \exp(x) = -2 \exp(x) \rightarrow \therefore A = -1$$

By the principle of superposition, our full solution is thus

$$u(x) = c_1 + c_2 \exp(-x) - \exp(x)$$

We can utilize the given boundary conditions in the problem to solve for the coefficients in the problem

$$u(0) = 0 = c_1 + c_2 - 1 \rightarrow c_1 = 1 - c_2$$

$$u(1) = 0 = c_1 + \frac{c_2}{e} - e$$

$$0 = 1 - c_2 + \frac{c_2}{e} - e$$

$$e - 1 = c_2\left(\frac{1}{e} - 1\right)$$

$$\therefore c_2 = -e, \quad c_1 = 1 + e$$

Our final solution can then be given as

$$u(x) = 1 + e - \exp(1 - x) - \exp(x)$$

Problem 4

Consider the operator $L : U \rightarrow U$ where $L[u] = -u_{xx}$ and

$$U = \{C^2([0, 1]) \mid u(0) = 0 \text{ and } u'(1) + \beta u(1) = 0\}$$

for some real constant β . Use the standard L^2 inner product for both the domain and range.

Is L self-adjoint? Is L positive definite? For positive-definiteness, you should find that your answer depends on β . It is enough to find a range of β 's where L is positive definite - you do not need to prove that L is not positive definite for other values of β . Show your work.

Solution.

To determine if L is self adjoint, we use the following equality

$$\langle L[u], v \rangle = \langle u, L^*[v] \rangle$$

We express this in integral form as

$$\int_0^1 -u''(x)v(x) \, dx = \int_0^1 u(x)L^*[v(x)] \, dx$$

We aim to make the left hand side of the equation be in the same form of the right hand side. We can integrate by parts with the following substitutions

$$u = v(x) \quad du = v'(x) \quad dv = -u''(x) \quad v = -u'(x)$$

$$\begin{aligned} LHS &= -v(x)u'(x)\Big|_0^1 + \int_0^1 u'(x)v'(x) \, dx \\ &= [-v(1)u'(1) + \cancel{v(0)u'(0)}] + \int_0^1 u'(x)v'(x) \, dx \\ &= -v(1)u'(1) + \int_0^1 u'(x)v'(x) \, dx \end{aligned}$$

We can integrate by parts again with the following substitution

$$u = v'(x) \quad du = v''(x) \quad dv = u'(x) \quad v = u(x)$$

$$\begin{aligned}
LHS &= -v(1)u'(1) + u(x)v'(x)|_0^1 - \int_0^1 u(x)v''(x) dx \\
&= -v(1)u'(1) + [u(1)v'(1) + \cancel{u(0)v'(0)}] + \int_0^1 u(x)[-v''(x)] dx \\
&= u(1)v'(1) - u'(1)v(1) + \int_0^1 u(x)[-v''(x)] dx \\
&= u(1)v'(1) + \beta u(1)v(1) + \int_0^1 u(x)[-v''(x)] dx \leftarrow \text{Apply BCs} \\
&= u(1)[\cancel{v'(1)} + \beta v(1)] + \int_0^1 u(x)[-v''(x)] dx \\
&= \int_0^1 u(x)[-v''(x)] dx
\end{aligned}$$

We see this matches the form of the right hand side of our initial inner product equation. Therefore, we know that

$$L^*[v(x)] = -\partial_{xx}[v(x)]$$

This is the same as the the given operator L , indicating that L is self adjoint. To find where L is positive definite, we need to indentify the conditions for which the following relation holds

$$\langle u, L[u] \rangle > 0$$

Using the integral definition of the inner product, we have

$$\int_0^1 -u(x)u''(x) dx > 0$$

We use the following substitution to integrate by parts

$$u = u(x) \quad du = u'(x) \quad dv = -u''(x) \quad v = -u'(x)$$

$$\begin{aligned}
&-u(x)u'(x)|_0^1 + \int_0^1 [u'(x)]^2 dx > 0 \\
&[-u(1)u'(1) + \cancel{u(0)u'(0)}] + \int_0^1 [u'(x)]^2 dx > 0 \\
&-u(1)u'(1) + \int_0^1 [u'(x)]^2 dx > 0
\end{aligned}$$

Using our initial conditions, we make the substitution $u'(1) = -\beta u(1)$

$$\begin{aligned}
&-u(1)u'(1) + \int_0^1 [u'(x)]^2 dx > 0 \\
&\beta u(1)^2 + \int_0^1 [u'(x)]^2 dx > 0 \\
&\beta u(1)^2 > -\int_0^1 [u'(x)]^2 dx \\
&\beta > \frac{-\int_0^1 [u'(x)]^2 dx}{u(1)^2}
\end{aligned}$$

Thus, we have that L is positive definite for

$$\beta > \frac{-\int_0^1 [u'(x)]^2 dx}{u(1)^2}$$

Problem 5

Let U be the space of continuously differentiable, complex-valued, 2π -periodic functions $u : \mathbb{R} \rightarrow \mathbb{C}$. Define the linear operator $L : U \rightarrow U$ such that

$$L[u] = i \frac{du}{dx}$$

Use the inner product

$$\langle u_1, u_2 \rangle = \int_{-\pi}^{\pi} u_1(x) \overline{u_2(x)} dx$$

for both the domain and range. Show that L is self-adjoint.

Solution.

With the given definition $L[u]$, the resulting operator L can be defined as

$$L = i\partial_x$$

A function $f : \mathbb{R} \rightarrow \mathbb{C}$ can be written in the following form of real and complex components

$$f(x) = a(x) + ib(x)$$

Thus, it follows that

$$f'(x) = a'(x) + ib'(x)$$

The complex conjugate of $f'(x)$ is given as

$$\overline{f'(x)} = a'(x) - ib'(x)$$

We also can calculate the derivative of the conjugate of $f(x)$

$$\frac{d}{dx}(\overline{f(x)}) = \frac{d}{dx}(a(x) - ib(x)) = a'(x) - ib'(x)$$

Thus, we see that

$$\overline{f'(x)} = \frac{d}{dx}(\overline{f(x)})$$

To show the operator L is self adjoint, we use the equality relation of the inner product definition

$$\begin{aligned} \langle L[u], v \rangle &= \langle u, L^*[v] \rangle \\ \int_{-\pi}^{\pi} i \frac{du}{dx} \overline{v(x)} dx &= \int_{-\pi}^{\pi} u(x) \overline{L^*[v(x)]} dx \end{aligned}$$

We attempt to match the left hand side of the equation to the right hand side. We can use integration by parts with the following substitutions on the left hand side

$$u = \overline{v(x)} \quad du = \overline{\frac{d}{dx}(v(x))} \quad dv = \frac{du}{dx} \quad v = u(x)$$

$$\begin{aligned}
LHS &= i \left[u(x) \overline{v(x)} \Big|_{-\pi}^{\pi} - \int_{-\pi}^{\pi} u(x) \overline{\frac{d}{dx}(v(x))} dx \right] \\
&= i \left[(u(\pi) \overline{v(\pi)} - u(-\pi) \overline{v(-\pi)}) - \int_{-\pi}^{\pi} u(x) \overline{\frac{d}{dx}(v(x))} dx \right]
\end{aligned}$$

As from the problem definitions the functions in U are 2π periodic, we know that $u(\pi) = u(-\pi)$. By the same logic, $v(\pi) = v(-\pi)$, and if these values are equal than it follows that $\overline{v(\pi)} = \overline{v(-\pi)}$. Therefore, we know that

$$u(\pi) \overline{v(\pi)} - u(-\pi) \overline{v(-\pi)} = 0$$

Thus

$$LHS = \int_{-\pi}^{\pi} u(x) \left[-i \frac{d}{dx}(v(x)) \right] dx$$

We know by complex conjugate rules that $\overline{z\overline{w}} = \overline{z} w$, and as a result the following relationship holds true

$$\overline{i \frac{d}{dx}(v(x))} = \overline{i} \overline{\frac{d}{dx}(v(x))} = -i \frac{d}{dx}(v(x))$$

We now have a pattern that matches the right hand side of our initial integral relation

$$\int_{-\pi}^{\pi} u(x) i \frac{d}{dx}(v(x)) dx = \int_{-\pi}^{\pi} u(x) \overline{L^*[v(x)]} dx$$

From inspection, we have that

$$L^*[v] = i \frac{dv}{dx} \rightarrow L^* = i \partial_x$$

As $L^* = L$, we have proved that the operator is self adjoint.