
Quantum Metrology with Photoelectrons Vol. 3 *Analysis methodologies*

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CONTENTS

I	Frontmatter	5
1	Overview	7
1.1	General overview	7
1.2	Provisional contents	7
II	Theory & software	9
2	Introduction	11
2.1	Topical introduction: from quantum metrology to a generalised bootstrapping protocol	11
2.2	Context & aims for Vol. 3	14
3	Quantum metrology software platform/ecosystem overview	17
3.1	Analysis components	17
3.2	Additional tools	19
3.3	Python ecosystem (backends, libraries and packages)	19
3.4	Docker deployments	20
3.5	General discussion	20
4	Theory	21
4.1	Observables: photoelectron flux in the LF and MF	21
4.2	Photoionization dynamics	27
4.3	Tensor formulation of photoionization	29
4.4	Density matrix representation	50
4.5	Information content	56
5	Numerical implementation	59
III	Backmatter	61
6	Bibliography	63
7	Glossary	65
IV	Test pages	67
8	Build versions and config tests	69
8.1	Versions	69
8.2	Docker build env	70

8.3	Book versions	70
8.4	Github pkg versions	70
8.5	Full conda env	71
Bibliography		81
Index		85

Quantum Metrology with Photoelectrons Volume 3: *Analysis methodologies*, an open source executable book. This repository contains the source documents (mainly Jupyter Notebooks in Python) and notes for the book, as of Jan 2022 writing is in progress, and the [current HTML build can be found online](#). The book is due to be finished in 2023, and will be published by IOP Press - see below for more details.

Series abstract

Photoionization is an interferometric process, in which multiple paths can contribute to the final continuum photoelectron wavefunction. At the simplest level, interferences between different final angular momentum states are manifest in the energy and angle resolved photoelectron spectra: metrology schemes making use of these interferograms are thus phase-sensitive, and provide a powerful route to detailed understanding of photoionization. In these cases, the continuum wavefunction (and underlying scattering dynamics) can be characterised. At a more complex level, such measurements can also provide a powerful probe for other processes of interest, leading to a more general class of quantum metrology built on phase-sensitive photoelectron imaging. Since the turn of the century, the increasing availability of photoelectron imaging experiments, along with the increasing sophistication of experimental techniques, and the availability of computational resources for analysis and numerics, has allowed for significant developments in such photoelectron metrology.

About the books

- Volume I covers the core physics of photoionization, including a range of computational examples. The material is presented as both reference and tutorial, and should appeal to readers of all levels. ISBN 978-1-6817-4684-5, <http://iopscience.iop.org/book/978-1-6817-4684-5> (IOP Press, 2018)
- Volume II explores applications, and the development of quantum metrology schemes based on photoelectron measurements. The material is more technical, and will appeal more to the specialist reader. ISBN 978-1-6817-4688-3, <http://iopscience.iop.org/book/978-1-6817-4688-3> (IOP Press, 2018)

Additional online resources for Vols. I & II can be found on [OSF](#) and [Github](#).

- Volume III in the series will continue this exploration, with a focus on numerical analysis techniques, forging a closer link between experimental and theoretical results, and making the methodologies discussed directly accessible via new software. The book is due for publication by IOP due in 2023; this volume is also open-source, with a live HTML version at <https://phockett.github.io/Quantum-Metrology-with-Photoelectrons-Vol3/> and source available at <https://github.com/phockett/Quantum-Metrology-with-Photoelectrons-Vol3>.

For some additional details and motivations (including topical video), see [the ePSdata project](#).

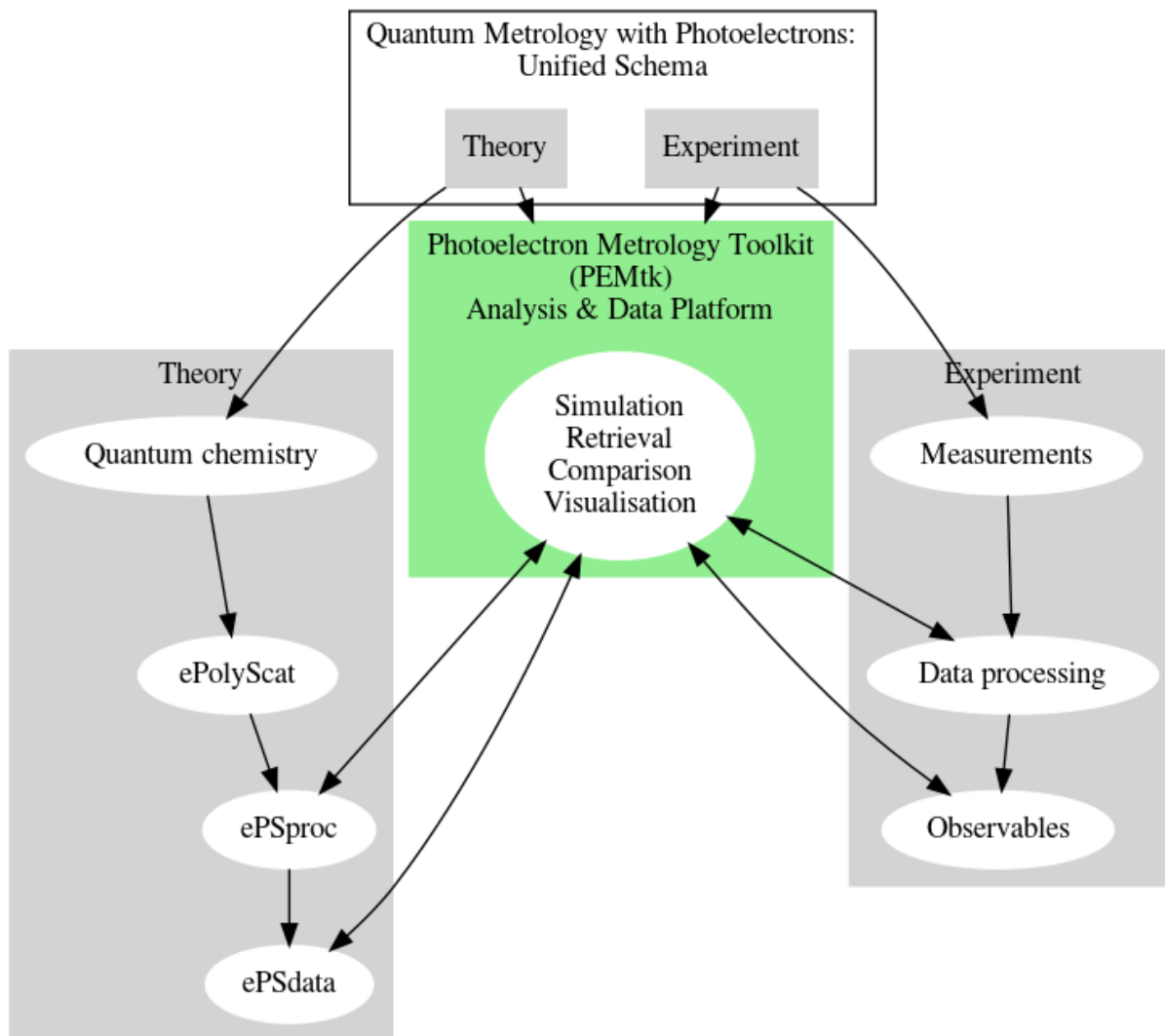
Technical details

This repository contains:

- `doc-source`: the source documents (mainly Jupyter Notebooks in Python)
- `notes`: additional notes for the book,
- the `gh-pages` branch contains the current HTML build, also available at <https://phockett.github.io/Quantum-Metrology-with-Photoelectrons-Vol3/>

The project has been setup to use the [Jupyter Book](#) build-chain (which uses Sphinx on the back-end) to generate HTML and Latex outputs for publication from source Jupyter notebooks & markdown files.

The work *within* the book will make use of the [Photoelectron Metrology Toolkit](#) platform for working with experimental & theoretical data.



Running code examples

Each Jupyter notebook (*.ipynb) can be treated as a stand-alone computational document. These can be run/used/modified independently with an appropriately setup python environment (details to follow).

Building the book

The full book can also be built from source:

1. Clone this repository
2. Run `pip install -r requirements.txt` (it is recommended you do this within a virtual environment)

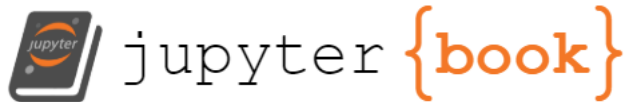
3. (Optional) Edit the books source files located in the `doc-source/` directory
4. Run `jupyter-book clean doc-source/` to remove any existing builds
5. For an HTML build:
 - Run `jupyter-book build doc-source/`
 - A fully-rendered HTML version of the book will be built in `doc-source/_build/html/`.
6. For a LaTeX & PDF build:
 - Run `jupyter-book build doc-source/ --builder pdflatex`
 - A fully-rendered HTML version of the book will be built in `doc-source/_build/latex/`.

See <https://jupyterbook.org/basics/building/index.html> for more information.

Credits

This project is created using the open source [Jupyter Book project](#) and the [executablebooks/cookiecutter-jupyter-book template](#).

To add: build env & main software packages (see automation for this...)



Part I

Frontmatter

OVERVIEW

1.1 General overview

Vol. 3. will focus on analysis techniques for quantum metrology with photoelectrons, including:

- Interpreting experimental data.
- Extraction/reconstruction/determination of quantum mechanical properties (matrix elements, wavefunctions, density matrices) from experimental data.
- Comparison of experimental and theoretical data.
- New analysis methodologies & techniques.
- Introduction to newly-developed software platform (see below).

1.2 Provisional contents

1.2.1 Part 1: theory & software

General review & update of the topic, including recent theory developments.

1. Introduction
 - a. Topic overview.
 - b. Context of vol. 3 (following vols. 1 & 2).
 - c. Aims: Vol. 3 in the series will continue the exploration of quantum metrology with photoelectrons, with a focus on numerical analysis techniques, forging a closer link between experimental and theoretical results, and making the methodologies discussed directly accessible via a new software platform/ecosystem.
2. Quantum metrology software platform/ecosystem overview
 - a. Introduction to python packages for simulation, data analysis, and open-data.
 - b. Photoelectron metrology toolkit (PEMtk) package/platform for experimental data processing & analysis. (See pemtk.readthedocs.io.)
 - c. ePSproc package for theory & simulation. (See epsproc.readthedocs.io.)
 - d. ePSdata platform for data/results library (see [ePSdata motivations](#)).
3. General method development: geometric tensor treatment of photoionization, fitting & matrix-inversion techniques
 - a. Theory development overview - tensor methods (e.g. [ePSproc tensor methods](#))

b. Direct molecular frame reconstruction via matrix-inversion methods (see Gregory, Margaret, Paul Hockett, Albert Stolow, and Varun Makhija. “Towards Molecular Frame Photoelectron Angular Distributions in Polyatomic Molecules from Lab Frame Coherent Rotational Wavepacket Evolution.” *Journal of Physics B: Atomic, Molecular and Optical Physics* 54, no. 14 (July 2021): 145601. DOI: 10.1088/1361-6455/ac135f.)

4. Numerical implementation & analysis platform tools

- a. Tensor methods implementation in ePSproc/PEMtk.
- b. Information content analysis (inc. basis-set exploration, e.g. *PEMtk fitting demo*), see also vol. 2, sect. 12.1.
- c. Density matrix analysis. (e.g. *ePSproc density matrix method dev notes*)
- d. Generalised bootstrapping implementation in PEMtk (see vol. 2, sects. 11.3 & 12.3)

1.2.2 Part 2: numerical examples

Open-source worked examples using the new software platform.

1. Quantum metrology example: generalised bootstrapping for a homonuclear diatomic scattering system (N₂)*
 - a. Experimental data overview & simulation.
 - b. Matrix element extraction (bootstrap protocol, see vol. 2, sects. 11.3 & 12.3) & statistical analysis.
 - c. Direct molecular frame reconstruction via matrix-inversion methods.
 - d. Comparison of methods.
 - e. Information content/quantum information analysis. (See vol. 2, sect. 12.1.)
2. Quantum metrology example: generalised bootstrapping for a heteronuclear scattering system (CO)*
 - a. Experimental data overview & simulation.
 - b. Matrix element extraction (bootstrap protocol, see vol. 2, sects. 11.3 & 12.3) & statistical analysis.
 - c. Direct molecular frame reconstruction via matrix-inversion methods.
 - d. Comparison of methods.
 - e. Information content/quantum information analysis. (See vol. 2, sect. 12.1.)
3. Quantum metrology example: generalised bootstrapping and matrix-inversion methods for a complex/general asymmetric top scattering system (C₂H₄ (ethylene))*
 - a. Experimental data overview & simulation.
 - b. Matrix element extraction (bootstrap protocol, see vol. 2, sects. 11.3 & 12.3) & statistical analysis.
 - c. Direct molecular frame reconstruction via matrix-inversion methods.
 - d. Comparison of methods.
 - e. Information content/quantum information analysis.
4. Future directions & outlook
5. Summary & conclusions

* Exact choice of “simple” and “complex” systems may change, but should include a homonuclear diatomic and/or heteronuclear diatomic, and symmetric and asymmetric top polyatomic systems. May also include an atomic example.

Part II

Theory & software

INTRODUCTION

The overall aim of *Quantum Metrology with Photoelectrons Vol. 3* is to expand, explore, and illustrate, new computational developments in quantum metrology with photoelectrons: specifically, the application of new python-based tools to tackle problems in matrix element retrieval. The book itself is written as a set of Jupyter Notebooks, hence all the material herein is available directly to readers, and can be run locally to further explore the topic interactively, and provide a foundation which can be adapted to apply the methodology to new problems.

Whilst this volume aims to provide a self-contained text, and computational examples which may be used without extensive background knowledge, a brief introduction to the core physics and some recent extensions is also presented herein (Sect. XX BELOW). The unfamiliar reader is referred to *Quantum Metrology Vol. 1* [1] for a more detailed introduction to the physics, and as a more general gateway to the literature. Following the topical introduction, the remainder of Part I introduces the main computational and software tools ([Chapter 3](#)), recent theory developments ([Chapter 4](#)), and concludes with a general overview for approaching matrix element retrieval numerically.

Part II details the application of these tools to a few specific cases, starting with a (relatively) simple homonuclear diatomic example, then escalating in complexity to a the most general polyatomic asymmetric top case.

2.1 Topical introduction: from quantum metrology to a generalised bootstrapping protocol

There are two core topics at the heart of this work, specifically photoelectron spectroscopy (and associated experimental, theoretical and analysis methodologies) and quantum metrology in general. To briefly (re)introduce these topics, and contextually frame the work discussed herein, some brief comments from *Quantum Metrology Vol. 1* [1] are reproduced below; the reader is referred to the introductory chapters of *Quantum Metrology Vol. 1* [1] for a lengthier treatment, and an introductory video to *Phase-sensitive Photoelectron Metrology* can be found online.

2.1.1 Quantum metrology with photoelectrons

To set the general context, consider quantum metrology in general...

Quantum metrology can be loosely defined as any class of experiment which provides detailed information on quantum mechanical properties (phases, coherences, entanglement etc) of a system. To stay with the spirit of modern metrology, this definition can further be refined to measurements which provide high-resolution quantum information; a clear contemporary example is therefore experimental methodologies which provide full quantum state reconstruction (e.g. quantum tomography), and/or make use of quantum mechanical properties as a tool for measurement (e.g. atom interferometry). Traditional high-resolution spectroscopies may also fit within this definition in some cases, although in the majority of cases high-resolution spectroscopic measurements provide transition line-strengths and energies, but lack sufficient information for a full determination or reconstruction of the underlying quantum state.

[...]

...at what point does a measurement of a quantum mechanical system become quantum metrology? A pragmatic view on this is that the complete quantum state of the system must be capable of *unique definition from the experimental measurement(s)*. This is pragmatic in the sense that it leaves the door open for both *inferred* and *direct* reconstruction techniques. In the former case, the experimental data informs the theory and analysis, but is not directly ‘analysed’ or ‘inverted’ to provide or reconstruct the full quantum information; in the latter case one obtains the desired quantum mechanical information from the measurement in a more ‘direct’ fashion (which may, admittedly, still remain as a rather convoluted process, depending on the level of theoretical input required). Traditional spectroscopies again provide a touchstone here - high-resolution spectroscopy measurements can be compared with models or *ab initio* computations to provide quantum mechanical details of a system, but typically do not directly provide this information from a set of measurements alone. In this sense they fit a pragmatic definition of quantum metrology, but not a more specific definition of quantum metrology as a (somewhat) direct empirical technique.

[...]

In summary, while quantum metrology can come in many flavours, at heart it might be considered as any set of measurements (and associated analysis methodologies) which provide detailed (quantitative) quantum mechanical information on a given system of interest - ideally with little or no restriction on the complexity of the system - and it is discussed in this spirit herein.

---Quantum Metrology Vol. 1 [1], Chpt. 1

And, for the specific case of photoionization...

... both *ab initio* methods and *high-dimensionality measurements* (combined with detailed *analysis methodologies*) can nonetheless provide detailed information on the photoionization dynamics. Although the simple analogy with Young’s double-slit [i.e. basic two-path interferometry] fails, the resulting photoelectron flux, measured spatially, remains, in essence, a self-referencing angular interferogram of the continuum wavefunction. In a more abstract sense, the basic interferometer paradigm can be extended to the general ‘photoionization interferometer’, one just has to keep in mind that there are now potentially many, many channels. In the most basic sense, the energy and angle resolved interferograms - the photoelectron flux as a function of energy and angle $I(E, \theta, \phi)$ - which may be measured, are nothing more than an interferometric measurement sensitive to the relative phases of the different angular momentum components.

[...]

In the photoionization community, the angular interferograms (which will usually be considered at a single energy E) are photoelectron angular distributions (*PADs*), and have long been used as a means to learn about the process of photoionization. In this context, *PADs* measured for a range of experimental parameters can provide a dataset with sufficient information content to determine the magnitudes and phases of the photoelectron wavefunction, hence the photoionization dynamics may be reconstructed from the measurements in favourable cases. This class of measurement is traditionally termed a *complete photoionization experiment*, although the exact nature of the completeness may vary. The phase-sensitivity of photoelectron interferograms have also been used in complementary fashions in other contexts, including as a means to probe the phase-shift induced by a specific prepared pathway, and control in multipath schemes, and in many other regimes.

[...]

... the combination of a phase-sensitive quantum mechanical observable - photoelectron interferograms - with modern experimental and computational techniques provides the tools required for a full quantum metrology based on this class of measurement. Following the above discussion and definitions, a full metrology technique is one which allows both the *intrinsic* and *extrinsic/dynamic* quantum mechanical properties of the system under study to be obtained/reconstructed from a measurement, or set of measurements. In the simplest case, one might seek to understand just the intrinsic photoionization dynamics of a scattering system (e.g. the magnitudes and phases of the various pathways [...]), while in more complex cases the intrinsic properties are part of a probe process for additional properties or dynamics of the system [...]. In all cases, the key is measurement (and possibly control) with a high information-content technique, and a

detailed understanding of the processes involved.

---*Quantum Metrology* Vol. 1 [1], Chpt. 1

2.1.2 Generalised geometric metrology protocols

In order to develop a quantitative form of photoelectron spectroscopy, hence analyse photoelectron interferograms in the context of quantum metrology in general, a number of techniques have previously been investigated (see *Quantum Metrology* Vols. 1 & 2 [1, 2]). In general, any applicable technique involves the manipulation or control of parameters which affect the observables in analytically-defined (or otherwise well-characterised) ways; measurements over a set of suitable experimental or control parameters then provide the high information-content dataset required for a full characterisation of the system at hand. Typically, “geometric” (angular-momentum) properties of the system provide a suitable set of control parameters, and a number of experimental methodologies with different flavours of these parameters have been demonstrated (see *Quantum Metrology* Vol. 2 [2]). The main, outstanding, issue with previous techniques was the system-specific nature of many of the applications: ideally, one would like to make use of a generalised protocol, which is independent of the particulars of the system under study, hence does not require, for example, specific spectroscopic properties to be known and/or be experimentally accessible.

The main aim of the work in the current volume is the further development, deployment and demonstrations of, such a scheme.

The focus is on one specific *high information-content* technique: the *generalised bootstrapping* protocol, which makes use of experiments using rotational wavepackets as a (geometric) control dimension, and time-resolved photoelectron measurements as a high-dimensionality, phase-sensitive observable; the combination of these measurements with a quantitative analysis methodology provides a (relatively) general route to a full quantum metrology with photoelectrons (a.k.a. complete photoionization experiments). A brief introduction to the technique is given below, with theoretical and numerical techniques and demonstrations forming the remainder of this book; interested readers can find a longer topical introduction in *Quantum Metrology* Vol. 2 [2] (in particular Chpt. 11), and see also Ref. [3] for an experimental demonstration, and Ref. [4] for a recent review in the context of molecular frame reconstruction.

As defined in *Quantum Metrology* Vol. 2 [2]:

For the analysis of the data [time-resolved photoelectron images from a rotationally-excited system], a ‘bootstrapping’ fitting approach was developed. This methodology [...is illustrated in figure 8.3, and...] is comprised of two stages (potentially split into multiple steps) which allow for the separation of the two sets of unknowns (rotational and ionization dynamics), and provides a way to gradually bootstrap to the complete *MF* results via stages of analysis of increasing complexity. The nature of the fitting at each stage also provides a flexible methodology which can be used to carefully sample the solution hyperspace in order to ensure unique results, and fit with variable information content (experimental measurements) based on computational time and desired precision, based on a similar Monte-Carlo sampling manner to the methodologies already discussed [...]. In all cases, the underlying physics provides stringent limits on the form of the fitting functions, hence the fitting procedure at each stage is expected to be somewhat reliable by construction. Further analysis of the results, including comparison with experimental parameters, additional data not used in the analysis, and *ab initio* calculations all provide additional means of cross-checking and verifying the extracted physical parameters.

In terms of information content, the bootstrapping procedure gradually increases both the experimental information content - the number of geometric configurations of the photoionization interferometer - and the level of physical information included (hence fitted/extracted) in the analysis. In the first step, *ADMs* [i.e. molecular alignment properties] are determined without the need for accurate treatment of the ionization probe [7]; in the second step this information is used as part of the calculation to determine the ionization dynamics. In the sub-steps to determine the ionization dynamics, the experimental information content included in the analysis is gradually increased: the initial coarse steps in this procedure provide a base-line high information content, without the necessity for many temporal points, via the selection of highly distinct molecular axis distributions, while latter sub-steps allow for fine-tuning of the data by gradually coupling additional time-steps into the analysis.

The protocol as presented relies on certain steps to be experimentally realisable, and theoretically calculable:

1. Molecular alignment. Experimentally, this can be induced in any system with a strong (typically $> 10^{12} \sim \text{Wcm}^{-2}$), short (few hundred femtosecond timescale or shorter) infra-red laser pulse, which (impulsively) creates a rotational wavepacket in the system. The exact nature of the wavepacket is laser pulse(s) and system dependent, but the technique is general.
2. Time-resolved photoelectron measurements. Experimentally, this requires - at minimum - a pump-probe type configuration, with the alignment pulse as the pump, and a time-delayed ionization pulse. This is a typical experimental configuration in many ultrafast laser labs, with pulses typically in the atto- or femto-second regime. Measurements may be made by any angle-resolved technique; photoelectron imaging is currently the most accessible and widespread method.
3. Data analysis. This provides the bridge from high information-content measurements to a full quantum metrology (system characterisation). For the generalised bootstrapping approach this requires: a. In order to characterise the rotational wavepacket created, alignment calculations of the system must be possible - such computations are increasingly tractable, if not already (somewhat) routine for a number of groups. These calculations are required in order to determine the rotational wavepacket quantitatively, and in order to determine the corresponding *ADMs* from/for the experiment. b. To characterise the *intrinsic* photoionization dynamics, a set of appropriate geometric basis functions must be computed, and combined with a sufficiently large dataset to enable extraction of the photoionization matrix elements via a fitting procedure. c. (Optional) In cases with *extrinsic* dynamics, these may further be analysed once the *intrinsic* dynamics have been characterised (or as part of that characterisation); this may, however, remain qualitative or semi-quantitative, depending on the system dynamics and complexity.
4. (Optional) *Ab initio* computations may also be performed to compare with any or all of the previous steps; comparison with step 3 is particularly powerful, since one can compare fundamental quantum mechanical properties, as opposed to comparison of measured and simulated observables which may be integrated over many degrees of freedom of the system.

As detailed in the following section (Sect. 2.2), the main aims herein are the development of the methodology and toolkit to address the data analysis requirement (step 3), and to test this methodology for a range of example cases.

2.2 Context & aims for Vol. 3

2.2.1 Scientific aims

The work in the current volume primarily addresses recent developments towards a generalised bootstrapping protocol (i.e. the analysis of the data time-resolved photoelectron images from a rotationally-excited system), as previously outlined in *Quantum Metrology* Vol. 2 [2] Sect. 12.3; in particular the new *Photoelectron Metrology Toolkit* [4] has been built with the aim of making the protocol easy to use and apply to any given problem (as distinct from a bespoke/per-experiment analysis methodology and/or non-open-source codebase).

Part I herein includes a full precis of the new codebase, along with the theory and numerics implemented towards this end; Part II provides multiple demonstrations of the new codebase, including the use of the toolkit to investigate more complex systems beyond the simple homonuclear diatomic case demonstrated to date.

Although the analysis herein focuses on the rotational wavepacket case, the techniques and codebase developed are equally applicable to *any methodology or protocol making use of geometric properties as a variable*, and are built with all such problems in mind - although minor modifications or extensions may be required for specific cases. Examples include other cases discussed in *Quantum Metrology* Vols. 1 & 2 [1, 2], e.g. use of shaped laser pulses, use of narrow-band, state-selected rotational excitation; in all cases the fitting/retrieval of matrix elements is carried out in the same manner, and the only changes required to the methodology are the choice of control variable and the corresponding input experimental or theoretical parameters - this is discussed further in Chpt. XX.

2.2.2 Technical context and notes

As noted previously, Vol. 3 is somewhat distinct from the previous volumes in the series; although involving computational elements, *Quantum Metrology* Vols. 1 & 2 [1, 2] % Vols. 1 & 2 [1, 2] are more traditional publications. The material presented in this volume aims to continue the exploration of quantum metrology with photoelectrons, with a focus on numerical analysis techniques, forging a closer link between experimental and theoretical results, and making the methodologies discussed directly accessible via a new software platform/ecosystem, *Photoelectron Metrology Toolkit* [4]. %, introduced in more detail in [Chapter 3](#). In order to fulfill this aim, Vol. 3 is a computational/computable document, with code directly available to readers to facilitate code transparency and reuse. This can be broken down as follows:

1. The book itself is written as a set of Jupyter Notebooks.¹
 - These are .ipynb files, usually running a python kernel, each of which is designed such that it can be modified and used independently.
 - The full book is compiled from these sections using the Jupyter Book project platform,² which includes build tools and specifications for the specific flavour of Markdown (MyST) used for the written text.
 - The book source code is available via a Github repository, *Quantum Metrology Vol. 3 (Github repo)*, which includes all the notebooks (in the `doc-source` directory), as well as installation and build notes for building the book itself.
 - An HTML version is also available at *Quantum Metrology Vol. 3 (HTML version)*, which includes interactive figures.
2. The code examples *within* the book make use the new *Photoelectron Metrology Toolkit* [4].
 - In order to run code examples, a specific python environment (with various additional python packages) is required.
 - A full introduction to the relevant software tool-chain, including installation instructions for the codes used *within* the book, can be found in [Chapter 3: Quantum metrology software platform/ecosystem overview](#).
 - For a quick and easy installation, including all requirements, a Docker build of the platform can also be used, see [Sect. 3.4: Docker deployments](#). % [Open Photoionization Docker Stacks](#) [5]).
 - Once configured, any code examples from the book can be executed locally by the user/reader, and modified as desired.
3. The book can be regarded as, essentially, a manual and introduction to the *Photoelectron Metrology Toolkit* [4], as well as a foundation for those wishing to use (and potentially extend) the platform.
 - Part I covers all required background material, including details of the theory and numerical methods implemented.
 - Part II contains various examples of usage for a range of problems, and possible extensions.
 - Since no specific knowledge of the underlying physics should be required to use the software tools, they will hopefully also provide a suitable platform for new researchers wishing to learn about photoionization in general.
 - It is, of course, also hoped that established researchers in the field will find the tools useful, and readily adaptable, to related problems of interest.

Finally, it is of note that whilst readers unfamiliar with the Jupyter and Python ecosystem may find that there is somewhat of a barrier to entry for making use of the platform, it is one that may be worth surmounting given the ubiquity of these tools, and general usefulness in modern scientific/data-science workflows; readers already making use of these tools in their work should have no difficulty, and the platform adheres to standard practice wherever possible.

¹ For more information on the [Jupyter Project and ecosystem](#), see [jupyter.org](#) and Refs. [6, 7, 8].)

² For more information see [jupyterbook.org](#) and Refs. [9, 10].

2.2.3 Open science, open source software and reproducibility

QUANTUM METROLOGY SOFTWARE PLATFORM/ECOSYSTEM OVERVIEW

STUB

In recent years, a unified Python codebase/ecosystem/platform has been in development to tackle various aspects of photoionization problems, including *ab initio* computations and experimental data handling, and (generalised) matrix element retrieval methods. The eponymous *Quantum Metrology with Photoelectrons* platform is introduced here, and is used for the analysis herein. The main aim of the platform is to provide a unifying data platform, and analysis routines, for photoelectron metrology, including new methods and tools, as well as a unifying bridge between these and existing tools. [Fig. 3.1](#) provides a general overview of some of the main tools and tasks/layers.

As of late 2022, the new parts of the platform - primarily the [Photoelectron Metrology Toolkit](#) [4] library - implement general data handling (although not a full experimental analysis toolchain), matrix element handling and retrieval, which will be the main topic of this volume. In the future, it is hoped that the platform will be extended to other theoretical and experimental methods, including full experimental data handling.

3.1 Analysis components

The two main components of the platform for analysis tasks, as used herein, are:

- The [Photoelectron Metrology Toolkit](#) [4] (PEMtk) codebase aims to provide various general data handling routines for photoionization problems. At the time of writing, simulation of observables and fitting routines are implemented, along with some basic utility functions. Much of this is detailed herein, and more technical details and ongoing documentation can be found in the [PEMtk documentation](#) [11].
- The [ePSproc](#) codebase [12, 13, 14] aims to provide methods for post-processing with *ab initio* radial dipole matrix elements from [ePolyScat](#) (ePS) [15, 16, 17, 18], or equivalent matrix elements from other sources (dedicated support for R-matrix results from the [RMT suite](#) [19, 20] is in development). The core functionality includes the computation of AF and MF observables. Manual computation without known matrix elements is also possible, e.g. for investigating limiting cases, or data analysis and fitting - hence these routines also provide the backend functionality for PEMtk fitting routines. Again more technical details can be found in the [ePSproc documentation](#) [14].

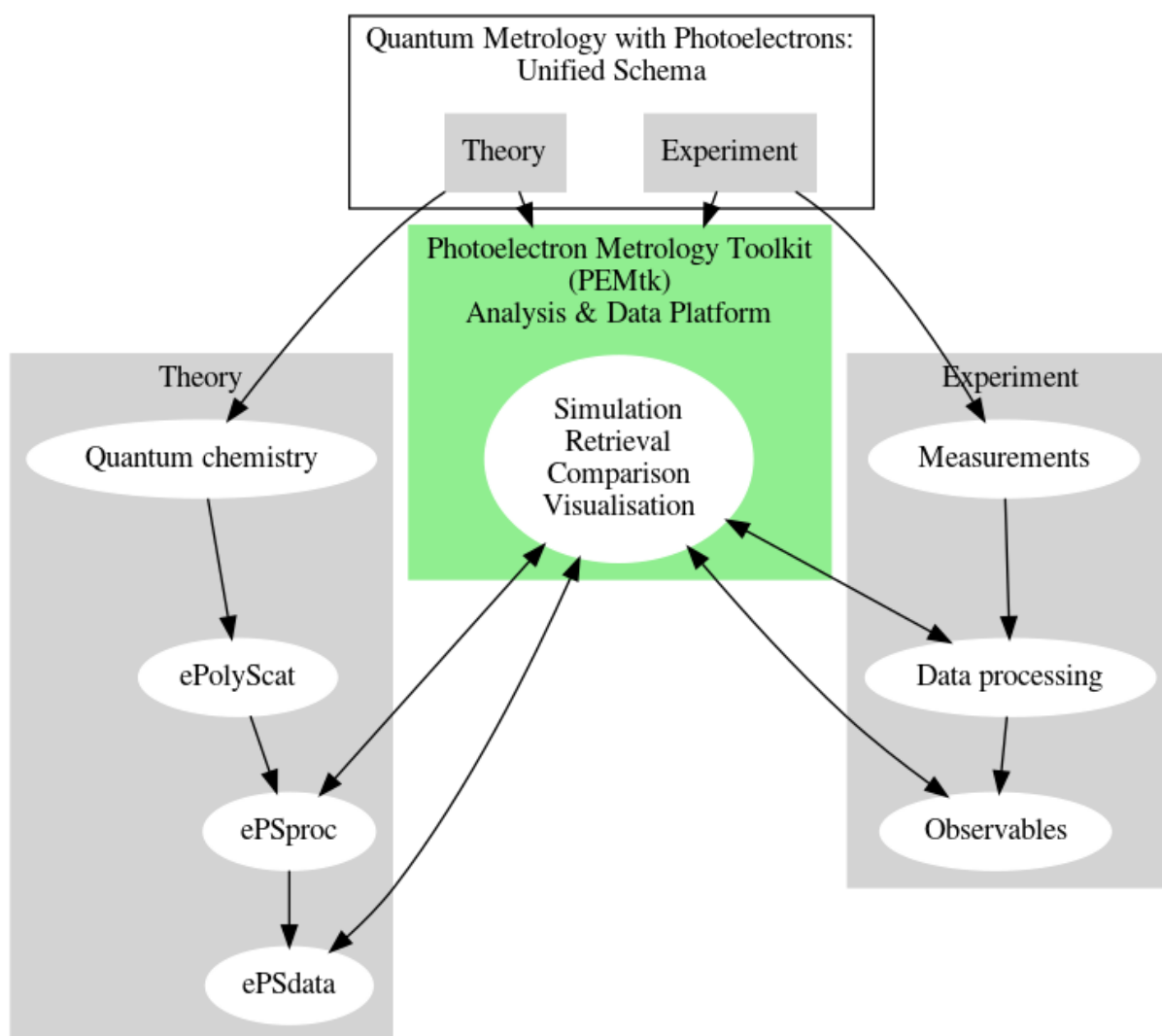


Fig. 3.1: Quantum metrology with photoelectrons ecosystem overview.

3.2 Additional tools

Other tools listed in Fig. 3.1 include:

- Quantum chemistry layer. The starting point for *ab initio* computations. Many tools are available, but for the examples herein, all computations made use of Gamess (“The General Atomic and Molecular Electronic Structure System”) [21, 22] for electronic structure computations, and inputs to ePolyScat.
 - For a python-based approach, various packages are available, e.g. PySCF, PyQuante, Psi can be used for electronic structure calculation, although note that some ePSproc [13] routines currently require Gamess files (specifically for visualisation of orbitals).
 - A range of other python tools are available, including cclib for file handling and conversion, Chemlab for molecule wavefunction visualisations, see further notes below.
- ePolyScat (ePS) [15, 16, 17, 18] is an open-source tool for numerical computation of electron-molecule scattering & photoionization by Lucchese & coworkers. All matrix elements used herein were obtained via ePS calculations. For more details see ePolyScat website and manual [18] and Refs. [15, 16, 17].
- ePSdata [23] is an open-data/open-science collection of ePS + ePSproc results.
 - ePSdata collects ePS datasets, post-processed via ePSproc (Python) in Jupyter notebooks, for a full open-data/open-science transparent pipeline.
 - Source notebooks are available on the ePSdata [23] Github project repository, and notebooks + datasets via ePSdata Zenodo [24]. Each notebook + dataset is given a Zenodo DOI for full traceability, and notebooks are versioned on Github.
 - Note: ePSdata may also be linked or mirrored on the existing ePolyScat Collected Results OSF project, but will effectively supercede those pages.
 - All results are released under Creative Commons Attribution-NonCommercial-ShareAlike 4.0 (CC BY-NC-SA 4.0) license, and are part of an ongoing Open Science initiative.

3.3 Python ecosystem (backends, libraries and packages)

The core analysis tools, which constitute the Photoelectron Metrology Toolkit [4] platform, are themselves built with the aid of a range of open-source python packages/libraries which handle various backend functionality. Notably, they make use of the following key packages:

- General functionality makes use of the usual Scientific Python stack, in particular:
 - Numpy for general numerical methods and data types.
 - pandas for statistical methods, and various tabulation and sorting tasks.
 - Scipy for some special functions and computational routines, particularly spherical harmonics and fitting routines (see below).
- General tensor handling and manipulation makes use of the Xarray library [25, 26].
- Angular momentum functions (Wigner D and 3js) are currently implemented directly, or via the Spherical Functions library [27], and have been tested for consistency with the definitions in Zare (for details see the ePSproc docs [14]). The Spherical Functions library also uses numpy_quaternion which implements a quaternion datatype in Numpy.
- Spherical harmonics are defined with the usual physics conventions: orthonormalised, and including the Condon-Shortley phase. Numerically they are implemented directly or via SciPy’s sph_harm function (see the SciPy docs

for details [28]. Further manipulation and conversion between different normalisations can be readily implemented with the SHtools library [29, 30].

- Non-linear optimization (fitting) is handled via the `lmfit` library, which implements and/or wraps a range of non-linear fitting routines in Python [31, 32]; for the Levenberg-Marquardt least-squares minimization method used herein this wraps `Scipy's least_squares` functionality, which therefore constituted the core minimization routine [28] for the demonstration cases.
- Symmetry functionality, specifically computing symmetrized harmonics $X_{hl}^{\Gamma\mu*}(\theta, \phi)$ (see (4.1)), makes use of `libmsym` [33, 34] (symmetry coefficients) and `SHtools` [29, 30] (general spherical harmonic handling and conversion).
- Some specialist (optional) tools also make use of additional libraries, although these are not required for basic use; in particular:
 - For 3D orbital visualizations with `ePSproc` [13]: `pyvista` for 3D plotting (which itself is built on VTK), `cclib` for electronic structure file handling and conversion, and methods based on `Chemlab` for molecule wavefunction (orbital) computation from electronic structure files are all used on the backend.
 - For general plotting a range of tools are used, or can be used, including `Matplotlib` (basic plotting, including `Xarray` plotters), `Holoviews` (data handling and interactive plotting, wraps various backends), `Bokeh` (implemented via `Holoviews`), `Plotly` (mainly used for spherical polar plotting), and `Seaborn` (for statistical and some specialist plots).
 - `Numba` is used for numerical acceleration in some routines, although remains mainly experimental in `eP-Sproc` at the time of writing (an exception to this is the Spherical Functions library, which does make full use of `Numba` acceleration).

Further comments, including conventions and numerical examples, can be found in Chpt. XX.

3.4 Docker deployments

A Docker-based distribution of various codes for tackling photoionization problems is also available from the `Open Photoionization Docker Stacks` [5] project, which aims to make a range of these tools more accessible to interested researchers, and fully cross-platform/portable. The project currently includes Docker builds for `ePS`, `ePSproc` and `PEMtk`.

3.5 General discussion

Note that, at the time of writing, rotational wavepacket simulation is not yet implemented in the `PEMtk` suite, and these must be obtained via other codes. An initial build of the `limapack` suite for rotational wavepacket simulations is currently part of the `Open Photoionization Docker Stacks` [5], but has yet to be tested.

- *Observables: photoelectron flux in the LF and MF*
- *Photoionization dynamics*
- *Tensor formulation of photoionization*
- *Density matrix representation*
- *Information content*

4.1 Observables: photoelectron flux in the LF and MF

The observables of interest - the photoelectron flux as a function of energy, ejection angle, and time - can be written quite generally as expansions in radial and angular basis functions. Various types and definitions are given in this section, including worked numerical examples.

4.1.1 Spherical harmonics

The photoelectron flux as a function of energy, ejection angle, and time, can be written generally as an expansion in spherical harmonics:

$$\bar{I}(\epsilon, t, \theta, \phi) = \sum_{L=0}^{2n} \sum_{M=-L}^L \bar{\beta}_{L,M}(\epsilon, t) Y_{L,M}(\theta, \phi) \quad (4.1)$$

Here the flux in the laboratory frame (*LF*) or aligned frame (*AF*) is denoted $\bar{I}(\epsilon, t, \theta, \phi)$, with the bar signifying ensemble averaging, and the molecular frame flux by $I(\epsilon, t, \theta, \phi)$. Similarly, the expansion parameters $\bar{\beta}_{L,M}(\epsilon, t)$ include a bar for the LF/AF case. These observables are generally termed photoelectron angular distributions (*PADs*), often with a prefix denoting the reference frame, e.g. LFPADs, MFPADs, and the associated expansion parameters $\bar{\beta}_{L,M}(\epsilon, t)$ are generically termed *anisotropy parameters*. The polar coordinate system (θ, ϕ) is referenced to an experimentally-defined axis in the *LF/AF* case (usually defined by the laser polarization), and the molecular symmetry axis in the *MF*. Some arbitrary examples are given in Fig. 4.1, which illustrates both a range of distributions of increasing complexity, and some basic code to set $\beta_{L,M}$ parameters and visualise them; the values used as tabulated in Fig. 4.2.

```
# Plot some distributions from specified BLMs

# Set specific LM coeffs by list with setBLMs, items are [l,m,value]
from epsproc.sphCalc import setBLMs

# BLM = setBLMs([[0,0,1],[1,1,1],[2,2,1]])
# BLM = setBLMs([[0,0,1,1,1],[1,1,1,0.5,0.2],[2,2,1,1,0.2]]) # Note different index
```

(continues on next page)

(continued from previous page)

```

BLM = setBLMs([[0,0,1,1,1,1],[1,1,0,0.5,0.8,1],[2,0,1,0.5,0,0],
              [4,2,0,0,0,0.5],[4,-2,0,0,0,0.5]])

# Set the backend to 'pl' for an interactive surface plot with Plotly
# NOTE PL FIG RETURN BROKEN FOR THIS CASE (ePSproc v1.3.1), so run sphSumPlotX too.
dataPlot, figObj = ep.sphFromBLMPlot(BLM, facetDim='t', plotFlag = False, backend = _
    ↪plotBackend);
figObj = ep.sphSumPlotX(dataPlot,facetDim='t', plotFlag = False, backend = _
    ↪plotBackend);

# And GLUE for display later with caption
# from myst_nb import glue
# glue("padExamplePlot", figObj[0], display=False);
# Glue with Plotly wrapper.
# gluePlotly("padExamplePlot", figObj[0]) # Working in Render test notebook, but _
    ↪not here? Issue with subplots?

# Test in separate cell...
gluePlotly("padExamplePlot", figObj[0]) # Working in Render test notebook, but not _
    ↪here? Issue with subplots?

```

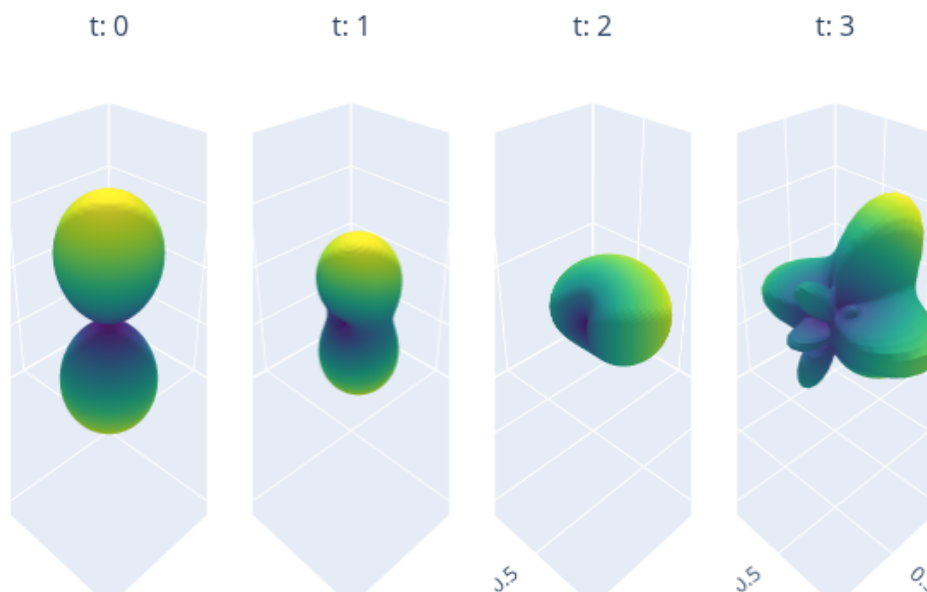


Fig. 4.1: Examples of angular distributions (expansions in spherical harmonics $Y_{L,M}$), for a range of cases. Note that up-down asymmetry is associated with odd- l contributions, and breaking of cylindrical symmetry with $m \neq 0$ terms.

In general, the spherical harmonic rank and order (L, M) of Eq. (4.1) are constrained by experimental factors in the *LF* or *AF*, and n is effectively limited by the molecular alignment (which is correlated with the photon-order for gas

	t	0	1	2	3
l	m				
0	0	1.0	1.0	1.0	1.0
1	0	0.0	0.5	0.8	1.0
2	0	1.0	0.5	0.0	0.0
4	-2	0.0	0.0	0.0	0.5
	2	0.0	0.0	0.0	0.5

Fig. 4.2: Values used for the plots in Fig. 4.1.

phase experiments, or conservation of angular momentum in the *LF* more generally [35]), but in the *MF* is defined by the maximum continuum angular momentum $n = l_{max}$ imparted by the scattering event [36] (note lower-case l here refers specifically to the continuum photoelectron wavefunction, see Eq. (4.7)).

For basic cases these limits may be low: for instance, a simple 1-photon photoionization event ($n = 1$) from an isotropic ensemble (zero net ensemble angular momentum) defines $L_{max} = 2$; for cylindrically symmetric cases (i.e. $D_{\infty h}$ symmetry) $M = 0$ only. For *MF* cases, $l_{max} = 4$ is often given as a reasonable rule-of-thumb for the continuum - hence $L_{max} = 8$ - although in practice higher- l may be populated. Some realistic example cases are discussed later (**PART II**), see also ref. [1] for more discussion and complex examples.

In general, these observables may also be dependent on various other parameters; in Eq. (4.1) two such parameters, (ϵ, t) , are included, as the usual variables of interest. Usually ϵ denotes the photoelectron energy, and t is used in the case of time-dependent (usually pump-probe) measurements. As discussed below (Sect. 4.2), the origin of such dependencies may be complicated but, in general, the associated photoionization matrix elements are energy-dependent, and time-dependence may also appear for a number of intrinsic or extrinsic (experimental) reasons, e.g. electronic or nuclear dynamics, rotational (alignment) dynamics, electric field dynamics etc. In many cases only one particular aspect may be of interest, so t can be used as a generic label to index changes as per Fig. 4.1.

4.1.2 Symmetrized harmonics

Symmetrized (or generalised) harmonics, which essentially provide correctly symmetrized expansions of spherical harmonics (Y_{LM}) functions for a given irreducible representation, Γ , of the molecular point-group can be defined by linear combinations of spherical harmonics (Refs. [37, 38, 39] as below):

$$X_{hl}^{\Gamma\mu*}(\theta, \phi) = \sum_{\lambda} b_{hl\lambda}^{\Gamma\mu} Y_{l,\lambda}(\theta, \phi) \quad (4.1)$$

where:

- Γ is an irreducible representation;
- (l, λ) define the usual spherical harmonic indices (rank, order), but note the use of (l, λ) by convention, since these harmonics are usually referenced to the *MF*;
- $b_{hl\lambda}^{\Gamma\mu}$ are symmetrization coefficients;
- index μ allows for indexing of degenerate components (note here the unfortunate convention that the label μ is also used for photon projection terms in general, as per Sect. 4.3.2 - in ambiguous cases the symmetrization term will instead be labelled as μ_X , although in many cases may actually be redundant and safely dropped from the symmetrization coefficients);
- h indexes cases where multiple components are required with all other quantum numbers identical.

Analogously to Eq. (4.1), a general expansion of an observable in the symmetrized harmonic basis set can then be defined as:

$$\bar{I}^{\Gamma}(\epsilon, t, \theta, \phi) = \sum_{\Gamma \mu h l} \bar{\beta}_{hl}^{\Gamma \mu}(\epsilon, t) X_{hl}^{\Gamma \mu*}(\theta, \phi) \quad (4.2)$$

Alternatively, by substitution into Eq. (4.1), and assigning $l = L$ and $\lambda = M$, a general symmetrized expansion may also be defined as:

$$\bar{I}(\epsilon, t, \theta, \phi) = \sum_{\Gamma \mu h} \sum_{L=0}^{2n} \sum_{M=-L}^L b_{hLM}^{\Gamma \mu} \bar{\beta}_{L,M}^{\Gamma}(\epsilon, t) Y_{L,M}(\theta, \phi) \quad (4.3)$$

However, in many cases the symmetrization coefficients are subsumed into the $\beta_{L,M}^{\Gamma}$ terms (or underlying matrix elements); in this case a simplified symmetrized expansion can be defined as:

$$\bar{I}^{\Gamma}(\epsilon, t, \theta, \phi) = \sum_{L=0}^{2n} \sum_{M=-L}^L \bar{\beta}_{L,M}^{\Gamma}(\epsilon, t) Y_{L,M}(\theta, \phi) \quad (4.3)$$

Where the expansion is defined for a given symmetry and irreducible representation with the shorthand Γ ; in many systems a single label may be sufficient here, since allowed (L, M) terms will be defined uniquely by irreducible representation, although multiple quantum numbers may be required for unique definition in the most general cases as per Eq. (4.1) (e.g. for cases with degenerate components). Further details and usage in relation to channel functions are also discussed in Sect. 4.3 (see, in particular, Eq. (4.12) for a similar general case), and in relation to fitting for specific cases in **PART II**.

The exact form of these coefficients will depend on the point-group of the system, see, e.g. Refs. [39, 40]. Numerical routines for the generation of symmetrized harmonics are implemented in [Photoelectron Metrology Toolkit](#) [4]: point-groups, character table generation and symmetrization (computing $b_{hLM}^{\Gamma \mu}$ parameters) is handled by [libmsym](#) [33, 34]; additional handling also makes use of [pySHtools](#) [29, 30].

A brief numerical example is given below, for Td symmetry ($l_{max} = 6$), and more details can be found in the [PEMtk documentation](#) [11]. In this case, full tabulations of the parameters lists all $b_{hLM}^{\Gamma \mu}$ for each irreducible representation, and the corresponding PADs are illustrated in Fig. 4.4.

Note: Full tabulations of the parameters available in HTML or notebook formats only.

```
# Import class
from pemtk.sym.symHarm import symHarm

# Compute hamronics for Td, lmax=4
sym = 'Td'
lmax=6

symObj = symHarm(sym, lmax)

# Character tables can be displayed - this will render directly in a notebook.
symObj.printCharacterTable()

# Glue items for later
glue("symHarmPG", sym, display=False)
glue("symHarmLmax", lmax, display=False)
glue("charTab", symObj.printCharacterTable(returnPD=True), display=False) # As above, ↵
↵but with PD object return and glue.
```

Character	dim	E	C2 ¹	S4 ¹	σ _d	C3 ¹
A1	1	1.0	1.0	1.0	1.0	1.0
A2	1	1.0	1.0	-1.0	-1.0	1.0
E	2	2.0	2.0	0.0	0.0	-1.0
T1	3	3.0	-1.0	1.0	-1.0	0.0
T2	3	3.0	-1.0	-1.0	1.0	0.0

Fig. 4.3: Example character table for Td symmetry generated with the Photoelectron Metrology Toolkit [4] wrapper for libmsym [33, 34].

```
# The full set of expansion parameters can be tabulated

# pd.set_option('display.max_rows', 100)

symObj.displayXlm() # Display values (note this defaults to REAL harmonics)
# symObj.displayXlm(YlmType='comp') # Display values for COMPLEX harmonic expansion.

# To plot using ePSproc/PEMtk class, these values can be converted to ePSproc BLM_
# data type...

# Run conversion - the default is to set the coeffs to the 'BLM' data type
symObj.toePSproc()

# Set to new key in data class
data.data['symHarm'] = {}

for dataType in ['BLM']: #['matE', 'BLM']:
    data.data['symHarm'][dataType] = symObj.coeffs[dataType]['b (comp)'] # Select_
    # expansion in complex harmonics
    data.data['symHarm'][dataType].attrs = symObj.coeffs[dataType].attrs

# Plot full harmonics expansions, plots by symmetry
# Note 'squeeze=True' to force drop of singleton dims may be required.
# data.padPlot(keys='symHarm', dataType='BLM', facetDims = ['Cont'], squeeze = True,
# backend=plotBackend)

data.padPlot(keys='symHarm', dataType='BLM', facetDims = ['Cont'], squeeze = True,
backend=plotBackend, plotFlag=False, returnFlag=True) # Working
figObj = data.data['symHarm']['plots']['BLM']['polar'][0]

# And GLUE for display later with caption
# from myst_nb import glue
# glue("padExamplePlot2", figObj, display=False);
gluePlotly("symHarmPADs", figObj)
```

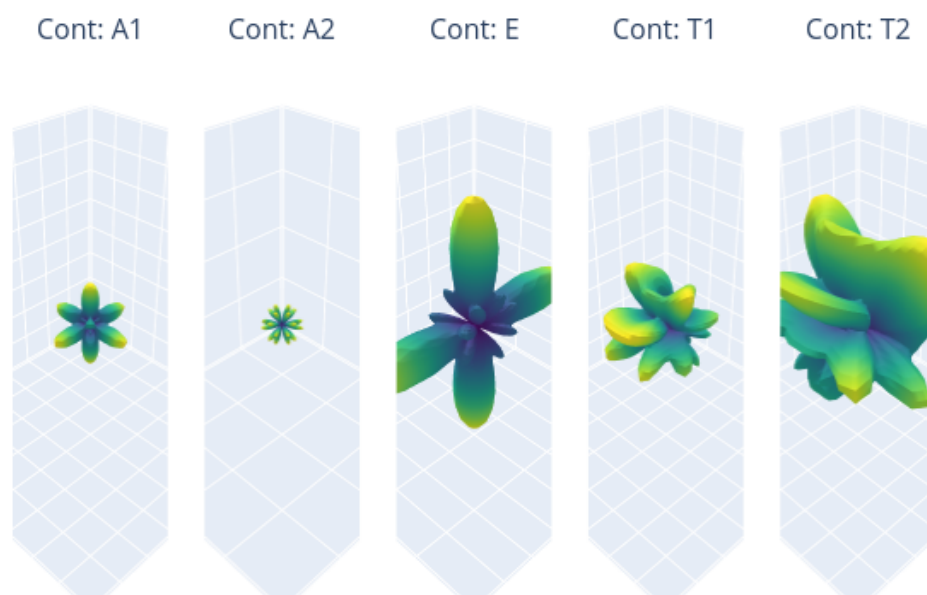


Fig. 4.4: Examples of angular distributions from expansions in symmetrized harmonics $X_{hl}^{\Gamma\mu*}(\theta, \phi)$, for all irreducible representations in Td symmetry ($l_{max}=6$). (Note A_2 only has components for $l \geq 6$.)

4.1.3 Real harmonics

4.1.4 Legendre polynomials

4.2 Photoionization dynamics

The core physics of photoionization has been covered extensively in the literature, and only a very brief overview is provided here with sufficient detail to introduce the metrology/reconstruction/retrieval problem; the reader is referred to Vol. 1 [1] (and refs. therein) for further details and general discussion.

Photoionization can be described by the coupling of an initial state of the system to a particular final state (photoion(s) plus free photoelectron(s)), coupled by an electric field/photon. Very generically, this can be written as a matrix element $\langle \Psi_i | \hat{\Gamma}(\mathbf{E}) | \Psi_f \rangle$, where $\hat{\Gamma}(\mathbf{E})$ defines the light-matter coupling operator (depending on the electric field \mathbf{E}), and Ψ_i, Ψ_f the total wavefunctions of the initial and final states respectively.

There are many flavours of this fundamental light-matter interaction, depending on system and coupling. For metrology, the focus is currently on the simplest case of single-photon absorption, in the weak field (or perturbative), dipolar regime, resulting in a single photoelectron. (For more discussion of various approximations in photoionization, see Refs. [41, 42].) In this case the core physics is well defined, and tractable (albeit non-trivial), via the separation of matrix elements into radial (energy) and angular-momentum (geometric) terms pertaining to couplings between various elements of the problem; the retrieval of such matrix elements is a well-defined problem, making use of analytic terms in combination with fitting methodologies as explored herein. Again, more extensive background and discussion can be found in *Quantum Metrology* Vol. 1 [1], and references therein. % [TODO: Add some more refs here?]

The basic case also provides a strong foundation for extension into more complex light-matter interactions, in particular cases with shaped laser-fields (i.e. a time-dependent coupling $\hat{\Gamma}(\mathbf{E}, \mathbf{t})$) and multi-photon processes (which require multiple matrix elements, and/or different approximations). Note, however, that non-perturbative (strong field) light-matter interactions are, typically, not amenable to description in a separable picture in this manner. In such cases the laser field, molecular and continuum properties are strongly coupled, and are typically treated numerically in a fully time-dependent manner (although some separation of terms may work in some cases, depending on the system and interaction(s) at hand).

Underlying the photoelectron observables is the photoelectron continuum state $|\mathbf{k}\rangle$, prepared via photoionization. The photoelectron momentum vector is denoted generally by $\mathbf{k} = k\hat{\mathbf{k}}$, in the molecular frame (*MF*). The ionization matrix elements associated with this transition provide the set of quantum amplitudes completely defining the final continuum scattering state,

$$|\Psi_f\rangle = \sum \int |\Psi_+; \mathbf{k}\rangle \langle \Psi_+; \mathbf{k} | \Psi_f \rangle d\mathbf{k}, \quad (4.3)$$

where the sum is over states of the molecular ion $|\Psi_+\rangle$. The number of ionic states accessed depends on the nature of the ionizing pulse and interaction. For the dipolar case,

$$\hat{\Gamma}(\mathbf{E}) = \hat{\mu} \cdot \mathbf{E} \quad (4.4)$$

Hence,

$$\langle \Psi_+; \mathbf{k} | \Psi_f \rangle = \langle \Psi_+; \mathbf{k} | \hat{\mu} \cdot \mathbf{E} | \Psi_i \rangle \quad (4.5)$$

Where the notation implies a perturbative photoionization event from an initial state i to a particular ion plus electron state following absorption of a photon $h\nu$, $|\Psi_i\rangle + h\nu \rightarrow |\Psi_+; \mathbf{k}\rangle$, and $\hat{\mu} \cdot \mathbf{E}$ is the usual dipole interaction term [43], which includes a sum over all electrons s defined in position space as \mathbf{r}_s :

$$\hat{\mu} = -e \sum_s \mathbf{r}_s \quad (4.6)$$

The position space photoelectron wavefunction is typically expressed in the “partial wave” basis, expanded as (asymptotic) continuum eigenstates of orbital angular momentum, with angular momentum components (l, m) (note lower case notation

for the partial wave components, distinct from upper-case for the similar terms (L, M) in the observables),

$$\Psi_{\mathbf{k}}(r) \equiv \langle r | \mathbf{k} \rangle = \sum_{lm} Y_{lm}(\hat{\mathbf{k}}) \psi_{lm}(r, k) \quad (4.7)$$

where r are MF electronic coordinates and $Y_{lm}(\hat{\mathbf{k}})$ are the spherical harmonics.

Similarly, the ionization dipole matrix elements can be separated generally into radial (energy-dependent or ‘dynamical’ terms) and geometric (angular momentum) parts (this separation is essentially the Wigner-Eckart Theorem, see Ref. [44] for general discussion), and written generally as (using notation similar to [45]):

$$\langle \Psi_+; \mathbf{k} | \hat{\mu} \cdot \mathbf{E} | \Psi_i \rangle = \sum_{lm} \gamma_{l,m} \mathbf{r}_{k,l,m} \quad (4.8)$$

Provided that the geometric part of the matrix elements $\gamma_{l,m}$ - which includes the geometric rotations into the LF arising from the dot product in Eq. (4.8) and other angular-momentum coupling terms - are known, knowledge of the so-called radial (or reduced) dipole matrix elements, at a given k thus equates to a full description of the system dynamics (and, hence, the observables).

For the simplest treatment, the radial matrix element can be approximated as a 1-electron integral involving the initial electronic state (orbital), and final continuum photoelectron wavefunction:

$$\mathbf{r}_{k,l,m} = \int \psi_{lm}(r, k) r \Psi_i(r) dr \quad (4.9)$$

As noted above, the geometric terms $\gamma_{l,m}$ are analytical functions which can be computed for a given case - minimally requiring knowledge of the molecular symmetry and polarization geometry, although other factors may also play a role (see Sect. 4.3.2 for details).

The photoelectron angular distribution (PAD) at a given (ϵ, t) can then be determined by the squared projection of $|\Psi_f\rangle$ onto a specific state $|\Psi_+; \mathbf{k}\rangle$; very generally this can be written in terms of the energy and angle-resolved observable, which arises as the coherent square:

$$I(\epsilon, \theta, \phi) = \langle \Psi_f | \Psi_f \rangle \quad (4.10)$$

Expansion in terms of the components of the matrix elements as detailed above then yields a separation into radial and angular components (see *Quantum Metrology* Vol. 1 [1], Sect. 2.1 for a full derivation), which can be written (at a single energy) as (following Eq. 2.45 of *Quantum Metrology* Vol. 1 [1]):

$$I(\theta, \phi; k) = \sum_{l'l'} \sum_{\lambda\lambda'} \sum_{mm'} \gamma_{\alpha\alpha_+ l\lambda m l' \lambda' m'} r_{kl\lambda} r_{kl'\lambda'} e^{i(\eta_{l\lambda}(k) - \eta_{l'\lambda'}(k))} Y_{lm}(\hat{k}) Y_{l'm'}^*(\hat{k}) \quad (4.11)$$

In this form α denotes all other quantum numbers required to define the initial state, and α_+ the final state of the molecular ion. The radial matrix elements $r_{kl\lambda}$, denote an integral over the radial part of the wavefunctions, in this case labelled by the MF quantum numbers, and the associated scattering phase is given by $\eta_{l\lambda}(k)$ (i.e. the matrix elements are written in magnitude-phase form, rather than complex form). The γ terms denotes a general set of geometric parameters arising from the coherent square. A tensor form is also given herein, see Sect. 4.3.2, including a full breakdown of these terms and numerical implementation. Comparison with Eq. (4.1) then indicates that the amplitudes in Eq. (4.8) also determine the observable anisotropy parameters $\beta_{L,M}(\epsilon, t)$ (Eqn. (4.1)), which basically collect all the terms in Eq. (4.11) and the product over spherical harmonics, into a result set of (L, M) . (Note that the photoelectron energy ϵ and momentum k are used somewhat interchangeably herein, with the former usually preferred in reference to observables.)

Note, also, that in the treatment above there is no time-dependence incorporated in the notation; however, a time-dependent treatment readily follows, and may be incorporated either as explicit time-dependent modulations in the expansion of the wavefunctions for a given case, or implicitly in the radial matrix elements. Examples of the former include, e.g. a rotational or vibrational wavepacket, or a time-dependent laser field. The rotational wavepacket case is discussed herein (see Sect. 4.3.2). The radial matrix elements are a sensitive function of molecular geometry and electronic configuration in general, hence may be considered to be responsive to molecular dynamics, although they are formally time-independent

in a Born-Oppenheimer basis - for further general discussion and examples see Ref. [46] and *Quantum Metrology* Vol. 1 [1]; discussions of more complex cases with electronic and nuclear dynamics can be found in Refs. [42, 47, 48, 49].

Typically, for reconstruction experiments, a given measurement will be selected to simplify this as much as possible by, e.g., populating only a single ionic state (or states for which the corresponding observables are experimentally energetically-resolvable), and with a bandwidth $d\mathbf{k}$ which is small enough such that the matrix elements can be assumed constant over the observation window. Importantly, the angle-resolved observables are sensitive to the magnitudes and (relative) phases of these matrix elements - as emphasised in the magnitude-phase form of Eq. (4.11) - and can be considered as angular interferograms.

4.3 Tensor formulation of photoionization

A number of authors have treated MFPADs and related problems [REFS]; herein, a geometric tensor based formalism is developed, which is close in spirit to the treatments given by Underwood and co-workers [49, 50, 51], but further separates various sets of physical parameters into dedicated tensors; this allows for a unified theoretical and numerical treatment, where the latter computes properties as tensor variables which can be further manipulated and investigated. Furthermore, the tensors can readily be converted to a density matrix representation [44, 52], which is more natural for some quantities, and also emphasizes the link to quantum state tomography and other quantum information techniques. Much of the theoretical background, as well as application to aspects of the current problem, can be found in the textbooks of Blum [52] and Zare [44].

Within this treatment, the observables can be defined in a series of simplified forms, emphasizing the quantities of interest for a given problem. Some details are defined in the following subsections,

4.3.1 Channel functions

A simple form of the equations (cf. the general form of Eq. (4.11), see also *Quantum Metrology* Vol. 2 [2] Chpt. 12), amenable to fitting and numerical implementation, is to write the observables in terms of “channel functions”, which define the ionization continuum for a given case and set of parameters u (e.g. defined for the MF, or defined for a specific experimental configuration),

$$\beta_{L,M}^u = \sum_{\zeta, \zeta'} \gamma_{L,M}^{u, \zeta \zeta'} \mathbb{I}^{\zeta \zeta'} \quad (4.12)$$

Where ζ, ζ' collect all the required quantum numbers, and define all (coherent) pairs of components. The term $\mathbb{I}^{\zeta \zeta'}$ denotes the coherent square of the ionization matrix elements:

$$\mathbb{I}^{\zeta, \zeta} = I^\zeta(\epsilon) I^{\zeta'}(\epsilon) \quad (4.13)$$

This is effectively a convolution equation (cf. refs. [50, 53]) with channel functions $\gamma_{L,M}^{u, \zeta \zeta'}$, for a given “experiment” u , summed over all terms ζ, ζ' . Aside from the change in notation (which is here chosen to match the formalism of Refs. [15, 16, 17]),

these matrix elements are essentially identical to the simplified (radial) forms $\mathbf{r}_{k,l,m}$ defined in Eqn. (4.8), in the case where $\zeta = k, l, m$. Note, also, that the matrix elements used herein are usually assumed to be symmetrized (unless explicitly stated), i.e. expanded in spherical harmonics per Eq. (4.1) with any additional terms $b_{hl\lambda}^{\Gamma\mu}$ incorporated into the value of the matrix elements.

These complex matrix elements can also be equivalently defined in a magnitude, phase form:

$$I^\zeta(\epsilon) \equiv \mathbf{r}_\zeta \equiv r_\zeta e^{i\phi_\zeta} \quad (4.14)$$

This tensorial form is numerically implemented in the **ePSproc** [13] codebase, and is in contradistinction to standard numerical routines in which the requisite terms are usually computed from vectorial and/or nested summations, which

can be somewhat opaque to detailed interpretation, and typically implement the full computation of the observables in one monolithic computational routine. The [Photoelectron Metrology Toolkit](#) [4] codebase implements matrix element retrieval based on the tensor formalism, with pre-computation of all the geometric tensor components (channel functions) prior to a fitting protocol for matrix element analysis, essentially a fit to Eqn. (4.12), with terms $I^\zeta(\epsilon)$ as the unknowns (in magnitude, phase form per (4.14)). The main computational cost of a tensor-based approach is that more RAM is required to store the full set of tensor variables; however, the method is computationally efficient since it is inherently parallel (as compared to a traditional, serial loop-based solution), hence may lead to significantly faster evaluation of observables. Furthermore, the method allows for the computational routines to match the formalism quite closely, and investigation of the properties of the channel functions for a given problem in general terms, as well as for specific experimental cases.

TODO: numerical examples here or below. TODO: benchmarks, or link to them (see test fitting notebooks...?).

4.3.2 Full tensor expansion

In more detail, the channel functions $\mathcal{T}_{L,M}^{u,\zeta\zeta'}$ can be given as a set of tensors, defining each aspect of the problem. The following equations illustrate this for the *MF* and *LF/AF* cases, fully expanding the general form of Eq. (4.12) in terms of the relevant tensors. Further details and numerical examples are given in the following sub-sections.

For the MF:

$$\begin{aligned} \beta_{L,-M}^{\mu_i,\mu_f}(\epsilon) &= (-1)^M \sum_{P,R',R} [P]^{\frac{1}{2}} E_{P-R}(\hat{e}; \mu_0) \\ &\times \sum_{l,m,\mu} \sum_{l',m',\mu'} (-1)^{(\mu'-\mu_0)} \Lambda_{R',R}(R_{\hat{n}}; \mu, P, R, R') B_{L,-M}(l, l', m, m') \\ &\times I_{l,m,\mu}^{p_i\mu_i,p_f\mu_f}(\epsilon) I_{l',m',\mu'}^{p_i\mu_i,p_f\mu_f^*}(\epsilon) \end{aligned} \quad (4.15)$$

And the LF/AF as:

$$\begin{aligned} \bar{\beta}_{L,-M}^{\mu_i,\mu_f}(E, t) &= (-1)^M \sum_{P,R',R} [P]^{\frac{1}{2}} E_{P-R}(\hat{e}; \mu_0) \\ &\times \sum_{l,m,\mu} \sum_{l',m',\mu'} (-1)^{(\mu'-\mu_0)} \bar{\Lambda}_{R'}(\mu, P, R') B_{L,S-R'}(l, l', m, m') \\ &\times I_{l,m,\mu}^{p_i\mu_i,p_f\mu_f}(\epsilon) I_{l',m',\mu'}^{p_i\mu_i,p_f\mu_f^*}(\epsilon) \sum_{K,Q,S} \Delta_{L,M}(K, Q, S) A_{Q,S}^K(t) \end{aligned} \quad (4.16)$$

In both cases a set of geometric tensor terms are required,

these terms provide details of:

- $E_{P-R}(\hat{e}; \mu_0)$: polarization geometry & coupling with the electric field.
- $B_{L,M}(l, l', m, m')$: geometric coupling of the partial waves into the $\beta_{L,M}$ terms (spherical tensors). Note for the *AF* case the terms may be reindexed by $M = S - R'$, which allows for the projection dependence on the *ADMs* (see below).
- $\Lambda_{R',R}(R_{\hat{n}}; \mu, P, R, R')$, $\bar{\Lambda}_{R'}(\mu, P, R')$: frame couplings and rotations (note slightly different terms for *MF* and *AF*).
- $\Delta_{L,M}(K, Q, S)$: alignment frame coupling (*AF* only).
- $A_{Q,S}^K(t)$: ensemble alignment described as a set of *axis distribution moments* (*ADMs*, *AF* only).
- Square-brackets indicate degeneracy terms, e.g. $[P]^{\frac{1}{2}} = (2P + 1)^{\frac{1}{2}}$.

And $I_{l,m,\mu}^{p_i\mu_i,p_f\mu_f}(\epsilon)$ are the (radial) dipole ionization matrix elements, as a function of energy ϵ . These matrix elements are essentially identical to the simplified forms $r_{k,l,m}$ defined in Eqn. (4.8), except with additional indices to label symmetry and polarization components defined by a set of partial-waves $\{l, m\}$, for polarization component μ (denoting the photon angular momentum components) and channels (symmetries) labelled by initial and final state indexes $(p_i\mu_i, p_f\mu_f)$.

The notation here follows that used by [ePolyScat \(ePS\)](#) [15, 16, 17, 18], and these matrix elements again represent the quantities to be obtained numerically from data analysis, or from an [ePolyScat \(or similar\)](#) calculation.

Following the tensor components detailed above, the full form of the channel functions of Eq. (4.12) for the *AF* and *MF* can be written as:

$$\begin{aligned} \Upsilon_{L,M}^{u,\zeta\zeta'} = & (-1)^M [P]^{\frac{1}{2}} E_{P-R}(\hat{e}; \mu_0) (-1)^{(\mu' - \mu_0)} \Lambda_{R',R}(R_{\hat{n}}; \mu, P, R, R') \\ & \times B_{L,-M}(l, l', m, m') \end{aligned} \quad (4.17)$$

$$\begin{aligned} \bar{\Upsilon}_{L,M}^{u,\zeta\zeta'} = & (-1)^M [P]^{\frac{1}{2}} E_{P-R}(\hat{e}; \mu_0) (-1)^{(\mu' - \mu_0)} \bar{\Lambda}_{R'}(\mu, P, R') \\ & \times B_{L,S-R'}(l, l', m, m') \Delta_{L,M}(K, Q, S) A_{Q,S}^K(t) \end{aligned} \quad (4.18)$$

Note that, in this case as given, time-dependence arises purely from the $A_{Q,S}^K(t)$ terms in the *AF* case, and the electric field term currently describes only the photon angular momentum coupling, although can in principle also describe time-dependent/shaped fields. Similarly, a time-dependent initial state (e.g. a vibrational wavepacket) could also describe a time-dependent *MF* case.

It should be emphasized, however, that the underlying physical quantities are essentially identical in all the theoretical approaches, with a set of coupled angular-momenta defining the geometrical part of the photoionization problem, despite these differences in the details of the theory and notation.

The various tensors defined above are implemented as functions in the [ePSproc codebase](#) [12, 13, 14], and further wrapped for fitting cases in the [Photoelectron Metrology Toolkit](#) [4]. In the remainder of this section, numerical examples using these codes are illustrated and explored. Full computational details can be found in the [ePSproc documentation](#) [14], including [extended discussion of each tensor](#) and complete function references in the [geomCalc submodule documentation](#).

4.3.3 Numerical aside: symmetry-defined channel functions

In the following sub-sections, each component is defined in detail, including numerical examples. For illustration purposes, the numerical example uses a minimal set of assumptions, and is defined initially purely by symmetry, although further terms may be required for some of the geometric terms and are discussed where required.

For this example, the D_{2h} point group is used, representing a fairly general case of a planar asymmetric top system, e.g. ethylene (C_2H_4). Note that, in this case, the symmetrization coefficients ($b_{hl\lambda}^{\Gamma\mu}$, see (4.1)) have the property that $\mu = 0$ only, and the h index is redundant, since it maps uniquely to l - see [Fig. 4.5](#) - so these indexes can be dropped. Note, also, the unfortunate convention that the label μ is used for multiple indexes; to avoid ambiguity this term is remapped to μ_X in the numerics below. However, in this case, since μ can be dropped from the symmetrization coefficients, there is actually no ambiguity in later usage.

Note: Full tabulations of the parameters available in HTML or notebook formats only.

```
# Setup symmetry-defined matrix elements using PEMtk

# Import class
from pemtk.sym.symHarm import symHarm

# Compute hamronics for Td, lmax=4
sym = 'D2h'
lmax=4
```

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```

lmaxPlot = 2 # Set lmaxPlot for subselection on plots later.

# Glue items for later
glue("symHarmPGmatE", sym, display=False)
glue("symHarmLmaxmatE", lmax, display=False)
glue("symHarmBasislmaxPlot", lmaxPlot, display=False)

# TODO: consider different labelling here, can set at init e.g. dims = ['C', 'h', 'muX', 'l', 'm'] - 25/11/22 code currently fails for mu mapping, remap below instead
symObj = symHarm(sym, lmax)
# symObj = symHarm(sym, lmax, dims = ['Cont', 'h', 'muX', 'l', 'm'])

# To plot using ePSproc/PEMtk class, these values can be converted to ePSproc BLM_
# data type...

# Run conversion - the default is to set the coeffs to the 'BLM' data type
dimMap = {'C': 'Cont', 'mu': 'muX'}
symObj.toePSproc(dimMap=dimMap)

# Run conversion with a different dimMap & dataType
dataType = 'matE'
# symObj.toePSproc(dimMap = {'C': 'Cont', 'h': 'it', 'mu': 'muX'}, dataType=dataType)
symObj.toePSproc(dimMap = dimMap, dataType=dataType)
# symObj.toePSproc(dimMap = {'C': 'Cont', 'h': 'it'}, dataType=dataType) # Drop mu >_
# muX mapping for now
# symObj.coeffs[dataType]

# Example using data class (setup in init script)
data = pemtkFit()

# Set to new key in data class
dataKey = sym
data.data[dataKey] = {}

for dataType in ['matE', 'BLM']:
    data.data[dataKey][dataType] = symObj.coeffs[dataType]['b (comp)'].sum(['h', 'muX', 'l']) # Select expansion in complex harmonics, and sum redundant dims
    data.data[dataKey][dataType].attrs = symObj.coeffs[dataType].attrs

```

```

# Display results (real harmonics)
symObj.displayXlm(setCols='h') #, dropLevels='mu')

```

```

# Glue version for JupyterBook output
glue("D2hXlm", symObj.displayXlm(setCols='h', returnPD=True), display=False) # As_
# above, but with PD object return and glue.

```

```

# ep.matEleSelector(data.data[dataKey]['matE'], thres = None, inds = {}, dims = 'Eke',
# sq = True, drop = True)

```

```

# Compute basis functions for given matrix elements

# Set data
data.subKey = dataKey

```

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Character (Γ)	PREFIX (μ)	1	h m	b					
				0	1	2	3	4	5
A1g	0	0	0	1.0					
		2	0		1.0				
			2			1.0			
		4	0				1.0		
			2					1.0	
A1u	0		4						1.0
		3	-2	1.0					
		2	-2	1.0					
		4	-4		1.0				
			-2			1.0			
B1u	0	1	0	1.0					
		3	0		1.0				
			2			1.0			
		2	1	1.0					
		4	1		1.0				
B2g	0		3				1.0		
		1	-1	1.0					
		3	-3		1.0				
			-1			1.0			
		2	-1	1.0					
B3g	0	4	-3		1.0				
			-1				1.0		
		1	1	1.0					
		3	1		1.0				
			3			1.0			

Fig. 4.5: Symmetrized harmonics coefficients ($b_{hl\lambda}^{\Gamma\mu}$, see (4.1)) for D2h symmetry ($l_{max}=4$) generated with the [Photoelectron Metrology Toolkit](#) [4] wrapper for [libmsym](#) [33, 34]. Note that, in this case, the coefficients have the property that $\mu = 0$ only, and the h index is redundant (maps uniquely to l).

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```
# Using PEMtk - this only returns the product basis set as used for fitting
BetaNormX, basisProduct = data.afblmMatEfit(selDims={}, sqThres=False)

# Using ePSproc directly - this includes full basis return if specified
BetaNormX2, basisFull = ep.geomFunc.afblmXprod(data.data[data.subKey]['matE'],
basisReturn = 'Full', selDims={}, sqThres=False) #, BLMRenorm = BLMRenorm,
**kwargs)

# The basis dictionary contains various numerical parameters, these are investigated
below.
# See also the ePSproc docs at https://epsproc.readthedocs.io/en/latest/methods/
geometric_method_dev_260220_090420_tidy.html
print(f"Product basis elements: {basisProduct.keys()}")
print(f"Full basis elements: {basisFull.keys()}")

# Use full basis for following sections
basis = basisFull
```

```
Product basis elements: dict_keys(['BLMtableResort', 'polProd', 'phaseConvention',
BLMRenorm'])
Full basis elements: dict_keys(['QNs', 'EPRX', 'lambdaTerm', 'BLMtable',
BLMtableResort', 'AFTERm', 'AKQS', 'polProd', 'phaseConvention', 'BLMRenorm'])
```

4.3.4 Matrix element geometric coupling term $B_{L,M}$

The coupling of the partial wave pairs, $|l, m\rangle$ and $|l', m'\rangle$, into the observable set of $\{L, M\}$ is defined by a tensor contraction with two $3j$ terms.

$$B_{L,M} = (-1)^m \left(\frac{(2l+1)(2l'+1)(2L+1)}{4\pi} \right)^{1/2} \begin{pmatrix} l & l' & L \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} l & l' & L \\ -m & m' & M \end{pmatrix} \quad (4.19)$$

Note that this term is equivalent, effectively, to a triple integral over spherical harmonics (e.g. Eq. 3.119 in Zare [44]):

$$\int_0^{2\pi} \int_0^\pi Y_{J_3 M_3}(\theta, \phi) Y_{J_2 M_2}(\theta, \phi) Y_{J_1 M_1}(\theta, \phi) \sin \theta d\theta d\phi = \left(\frac{(2J_1+1)(2J_2+1)(2J_3+1)}{4\pi} \right)^{1/2} \\ \times \begin{pmatrix} J_1 & J_2 & J_3 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} J_1 & J_2 & J_3 \\ M_1 & M_2 & M_3 \end{pmatrix}$$

And a similar term appears in the contraction over a pair of harmonics into a resultant harmonic (e.g. Eqs. C.21, C.22 in Blum [52]) - this is how the term arises in the derivation of the observables.

$$Y_{J_1 M_1}(\theta, \phi) Y_{J_2 M_2}(\theta, \phi) = \sum_{J_3 M_3} \left(\frac{(2J_1+1)(2J_2+1)(2J_3+1)}{4\pi} \right)^{1/2} \\ \times \begin{pmatrix} J_1 & J_2 & J_3 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} J_1 & J_2 & J_3 \\ M_1 & M_2 & M_3 \end{pmatrix} Y_{J_3 M_3}^*(\theta, \phi)$$

Note also some definitions use conjugate spherical harmonics, which can be converted as, e.g., Eq. C.21 in Blum [52]:

$$\beta_{L,M}^{\mu_i, \mu_f} Y_{LM}^*(\theta_k, \phi_k) = \beta_{L,-M}^{\mu_i, \mu_f} (-1)^M Y_{L,-M}(\theta_k, \phi_k) \quad (4.20)$$

In the current **Photoelectron Metrology Toolkit** [4] codebase, the relevant basis item can be inspected as below, in order to illustrate the sensitivity of different (L, M) terms to the matrix element products. Note for the **AF** case the terms

may be reindexed by $M = S - R'$ - this allows for all MF projections to contribute, rather than just a single specified polarization geometry. However, in many typical cases, this term is nonetheless restricted to only $M = 0$ components overall by other geometric factors (see below).

The code cells below illustrate this for the current example case, and Fig. 4.6 offers a general summary. In general, this is a convenient way to visualize the selection rules into the observable: for instance, only terms $l = l'$ and $m = -m'$ contribute to the overall photoionization cross-section term ($L = 0, M = 0$), and the maximum observable $L_{max} = 2l_{max}$. However, since these terms are fairly simply followed algebraically in this case, via the rules inherent in the $3j$ product (Eq. (4.19)), this is not particularly insightful (although useful pedagogically). These visualizations will become more useful when dealing with real sets of matrix elements, and specific polarization geometries, which will further modulate or restrict the $B_{L,M}$ terms.

Numerically, various standard functions may be used to quickly gain deeper insight, for example min/max, averages etc. Such considerations may provide a quick sanity-check for a given case, and may prove useful when planning experiments to investigate particular aspects or channels of a given system. Other properties of the basis functions may also be interrogated numerically; for instance, correlation maps provide an alternative way to check which terms are strongly correlated or coupled, or will dominate a given aspect of the observable.

```
# Tabulate basis (use ep.multiDimXrToPD, or other? Should have wrappers...?)

basisKey = 'BLMtableResort' # Key for BLM basis set

# Reformat basis for display (optional)
stackDims = {'LM':['L','M']}
basisPlot = basis[basisKey].rename({'S-Rp':'M'}).stack(stackDims)

# Convert to Pandas
pd, _ = ep.multiDimXrToPD(basisPlot, colDims=stackDims)

# Summarise properties and tabulate via Pandas Describe
pd.describe().T
```

```
*** Plot BLM terms for basis set - basic
basisKey = 'BLMtableResort'

# Basic plot
ep.lmPlot(basisPlot, xDim=stackDims); # Basic plot with all terms
```

```
*** Plot BLM terms for basis set - Plot with some additional figure formatting,
↳options

cmap=None # cmap = None for default. 'vlag' good?
# cmap = 'vlag'

labelRound = 1
catLegend=False
titleString=''
titleDetails=False
labelCols = [1,1]

# With global formatting args
# ep.lmPlot(basis[basisKey].where((basis[basisKey].l<=lmaxPlot) & (basis[basisKey].lp
↳<=lmaxPlot)).rename({'S-Rp':'M'}).stack({'LM':['L','M']}), xDim={'LM':['L','M']}},
#          cmap=cmap, labelRound = labelRound, catLegend=catLegend,
↳titleString=titleString, titleDetails=titleDetails, labelCols = labelCols);
```

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```

# With fig return - NEEDS WORK, displays but doesn't return fig object?
# UPDATE 28/11/22: still displays, but now have glue() working correctly.
# daPlot, daPlotpd, legendList, gFig = ep.lmPlot(basis[basisKey].
  ↳ where((basis[basisKey].l<=lmaxPlot) & (basis[basisKey].lp<=lmaxPlot)).rename({'S-Rp
  ↳ ':'M'}).stack({'LM':['L','M']}), xDim={'LM':['L','M']},
#
  pType = 'r', cmap=cmap, labelRound = labelRound, catLegend=catLegend,
  ↳ titleString=titleString, titleDetails=titleDetails, labelCols = labelCols);

# 30/11/22 simplified version
daPlot, daPlotpd, legendList, gFig = ep.lmPlot(basisPlot.where((basisPlot.l
  ↳ <=lmaxPlot) & (basisPlot.lp<=lmaxPlot)),
  xDim=stackDims, pType = 'r', cmap=cmap,
  ↳ labelRound = labelRound,
  catLegend=catLegend,
  ↳ titleString=titleString, titleDetails=titleDetails,
  labelCols = labelCols);

# For glue
glue("lmPlot_BLM_basis_D2h", gFig.fig, display=False)

```

4.3.5 Electric field geometric coupling term $E_{P,R}(\hat{e}; \mu_0)$

The coupling of two 1-photon terms (which arises in the square of the ionization matrix element as per Eq. (4.11)) can be written as a tensor contraction:

$$E_{PR}(\hat{e}) = [e \otimes e^*]_R^P = [P]^{\frac{1}{2}} \sum_p (-1)^R \begin{pmatrix} 1 & 1 & P \\ p & R-p & -R \end{pmatrix} e_p e_{R-p}^* \quad (4.21)$$

Where:

- e_p and e_{R-p} define the field strengths for the polarizations p and $R-p$, which are coupled into the spherical tensor E_{PR} ;
- square-brackets indicate degeneracy terms, e.g. $[P]^{\frac{1}{2}} = (2P+1)^{\frac{1}{2}}$;
- the symbol μ_0 is conventionally used to denote the *LF* field projection definition, given in full as $(1, \mu_0)$ for the 1-photon case, **in general for the LF/AF case $\mu_0 = p$, while for the MF case all projection terms are allowed, and are usually labelled by μ or q . TO FIX**

(To derive this result, one can start from, e.g., Eq. 5.40 in Zare

$$[A^{(1)} \otimes B^{(1)}]_q^k = \sum_m \langle 1m, 1q-m | kq \rangle A(1, m) B(1, q-m) \quad (4.22)$$

Convert to $3j$ form:

$$[A^{(1)} \otimes B^{(1)}]_q^k = \sum_m (-1)^q [k]^{1/2} \begin{pmatrix} 1 & 1 & k \\ m & q-m & -q \end{pmatrix} A(1, m) B(1, q-m) \quad (4.23)$$

And substitute in appropriate terms.)

As before, we can visualise these values...

```

# For illustration, recompute EPR term for default case.
EPRX = ep.geomCalc.EPR(form = 'xarray')

```

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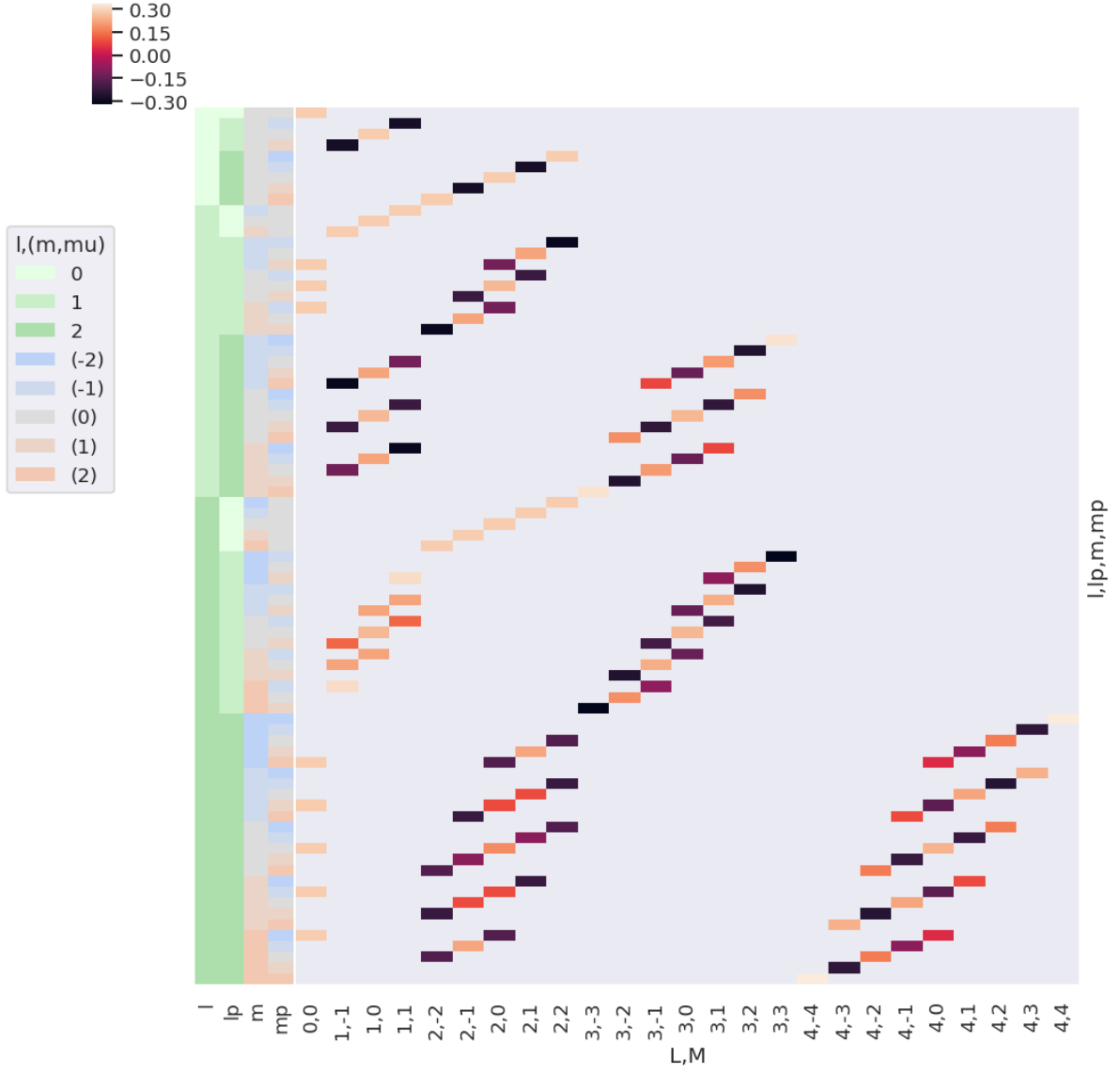


Fig. 4.6: Example $B_{L,M}$ basis functions for D_{2h} symmetry. Note figure is truncated to $l_{max} = l'_{max} = 2$ for clarity.

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```

# Set parameters to restack the Xarray into (L,M) pairs
# plotDimsRed = ['l', 'p', 'lp', 'R-p']
plotDimsRed = ['p', 'R-p']
xDim = {'PR':['P','R']}

# Plot with ep.lmPlot(), real values
daPlot, daPlotpd, legendList, gFig = ep.lmPlot(EPRX, plotDims=plotDimsRed, xDim=xDim,
pType = 'r')
# Version summed over l,m
# daPlot, daPlotpd, legendList, gFig = ep.lmPlot(EPRX.unstack().sum(['l', 'lp', 'R-p']),
xDim=xDim, pType = 'r')

# For glue
glue("lmPlot_EPR_basis", gFig.fig, display=False)

```

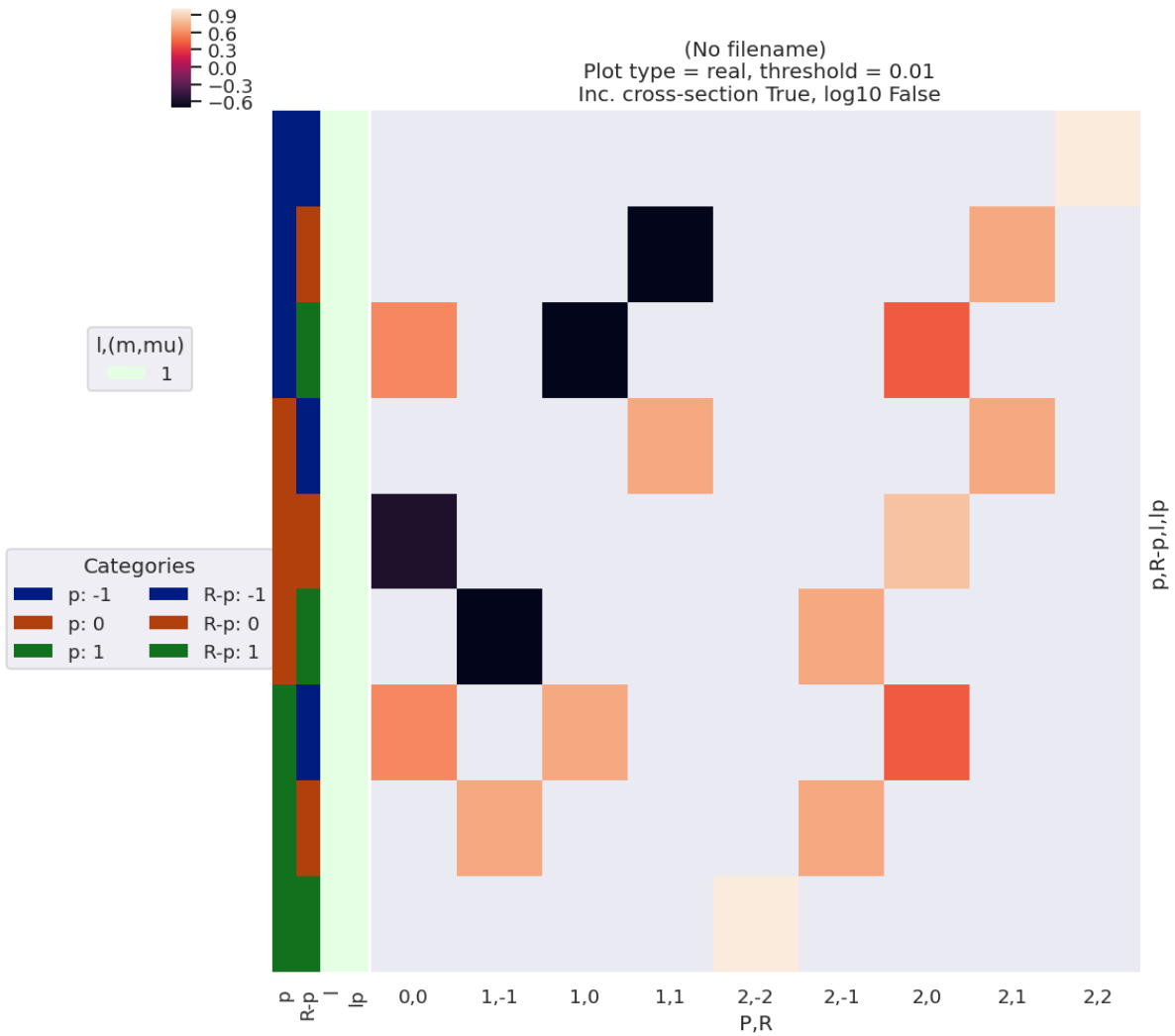


Fig. 4.7: Example $E_{P,R}$ basis functions. Note that for linearly polarised light $p = R - p = 0$ only, hence only the terms $E_{0,0}$ and $E_{2,0}$ are non-zero in this case. For non-linearly polarised cases many other terms are allowed.

4.3.6 Molecular frame projection term Λ

For the molecular frame case, the coupling between the *LF* and *MF* can be defined by a projection term, $\Lambda_{R',R}(R_{\hat{n}})$:

$$\Lambda_{R',R}(R_{\hat{n}}) = (-1)^{(R')} \begin{pmatrix} 1 & 1 & P \\ \mu & -\mu' & R' \end{pmatrix} D_{-R',-R}^P(R_{\hat{n}}) \quad (4.24)$$

This is similar to the E_{PR} term, and essentially rotates it into the *MF*, defining the projections of the polarization vector (photon angular momentum) μ into the *MF* for a given molecular orientation (frame rotation) defined by $R_{\hat{n}}$.

For the *LF/AF* case, the same term appears but in a simplified form:

$$\bar{\Lambda}_{R'} = (-1)^{(R')} \begin{pmatrix} 1 & 1 & P \\ \mu & -\mu' & R' \end{pmatrix} \equiv \Lambda_{R',R'}(R_{\hat{n}} = 0) \quad (4.25)$$

This form pertains since - in the *LF/AF* case - there is no specific frame transformation defined (i.e. there is no single molecular orientation defined in relation to the light polarization, rather a distribution as defined by the *ADMs*), but the total angular momentum coupling of the photon terms is still required in the equations.

Numerically, the function is calculated for a specified set of orientations, which default to the standard set of (x, y, z) MF polarization cases. For the *LF/AF* case, this term is still used, but restricted to $R_{\hat{n}} = (0, 0, 0) = z$, i.e. no frame rotation relative to the *LF* E_{PR} definition.

TODO: frame rotation illustration as well as term plots? Cf. QM1?

4.3.7 Alignment tensor $\Delta_{L,M}(K, Q, S)$

Finally, for the *LF/AF* case, the alignment tensor couples the molecular axis ensemble (defined as a set of *ADMs*) and the photoionization multipole terms into the final observable.

$$\Delta_{L,M}(K, Q, S) = (2K+1)^{1/2} (-1)^{K+Q} \begin{pmatrix} P & K & L \\ R & -Q & -M \end{pmatrix} \begin{pmatrix} P & K & L \\ R' & -S & S-R' \end{pmatrix} \quad (4.26)$$

In the full equations for the observable, this term appears in a summation with the *ADMs*, as:

$$\tilde{\Delta}_{L,M}(t) = \sum_{K,Q,S} \Delta_{L,M}(K, Q, S) A_{Q,S}^K(t) \quad (4.27)$$

This summed alignment term can be considered, essentially, as a (coherent) geometric averaging of the *MF* observable weighted by the axis distribution in the *AF* (for more on the axis averaging as a convolution, see Refs. [1, 51]); equivalently, the averaging can be considered as a purely angular-momentum coupling effect, which accounts for all contributing moments of the various aspects of the system, and defines the allowed projections onto the final observables in the *LF*.

Mappings of these terms are investigated numerically below, for some exemplar cases.

Basic cases

Fig. 4.8 illustrates the alignment tensor $\Delta_{L,M}(K, Q, S)$ for some basic cases, and values are also tabulated in Fig. 4.9. Note that for illustration purposes the term is subselected with $K = 0, Q = 0, S = 0$ and $R' = 0; R \neq 0$ terms are included to illustrate the elliptically-polarized case, which can give rise to non-zero M terms.

For the simplest case of an unaligned ensemble, this term is restricted to $K = Q = S = 0$, i.e. $\Delta_{L,M}(0, 0, 0)$; for single-photon ionization with linearly-polarized light ($p = 0$, hence $P = 0, 2$ and $R = R' = 0$), this has non-zero values for $L = 0, 2$ and $M = 0$ only. Typically, this simplest case is synonymous with standard *LF* results, and maintains cylindrical and up-down symmetry in the observable.

For circularly polarized light ($p = \pm 1$, hence $P = 0, 1, 2$ and $R = R' = 0$), odd- L is allowed, signifying up/down symmetry breaking in the observable (where up/down pertains to the propagation direction of the light, conventionally

the z -axis). For elliptically polarized light, mixing of terms with different p allows for non-zero R terms, hence non-zero M is allowed, signifying breaking of cylindrical symmetry in the observable.

```

**** Set range of ADMs for test as time-dependent values - single call
# AKQS = ep.setADMs (ADMs = [[0,0,0,1,1,1,1],
#                             [2,0,0,0,0.5,1,1],
#                             [4,0,0,0,0,0.3]]) #, t=[0,1,2]) # Nested list or np.
# array OK.
# AKQS = ep.setADMs (ADMs = np.array([[0,0,0,1,1],[2,0,0,0,0.5]]), t=[0,1])

**** Alternative form with ADM inds separate
# ADMinds = np.array([[0,0,0],[2,0,0],[4,0,0]])
# AKQS = ep.setADMs (ADMs = np.stack((np.ones(4),np.linspace(0,1,4),np.linspace(0,0.2,
# 4))),
#                    KQSLabels=ADMinds)

# For computation set ADMs and EPR term first.
# Note that delta term is independent of the absolute values of the ADMs(t), but does
# use this term to define limits on some quantum numbers.

**** Neater version
# Set ADMs for increasing alignment...
tPoints = 10
inputADMs = [[0,0,0, *np.ones(tPoints)], # * np.sqrt(4*np.pi)], # Optional
# multiplier for normalisation
#           [2,0,0, *np.linspace(0,1,tPoints)],
#           [4,0,0, *np.linspace(0,0.5,tPoints)],
#           [6,0,0, *np.linspace(0,0.3,tPoints)],
#           [8,0,0, *np.linspace(0,0.2,tPoints)]]

AKQS = ep.setADMs (ADMs = inputADMs) # TODO WRAP TO CLASS (IF NOT ALREADY!)

**** Plot ADMs
# daPlot, daPlotpd, legendList, gFig = ep.lmPlot(AKQS, xDim = 't', pType = 'r',
# squeeze = False, thres=None, cmap='vlag') # Note squeeze = False required for 1D
# case (should add this to code!)
# daPlotpd

# Example alignment term, assuming all other terms allow/unity valued.

# Use default EPR term - note this computes for all pol states, p=[-1,0,1]
EPR = ep.geomCalc.EPR(form='xarray')

# Compute alignment terms
Afterm, DeltaTerm = ep.geomCalc.deltaLMKQS (EPR, AKQS)

**** Plot Delta term with subselections
# xDim = {'KQ':['K','Q']}
xDim = {'LM':['L','M']}
daPlot, daPlotpd, legendList, gFig = ep.lmPlot(DeltaTerm.sel(K=0,Q=0,S=0,Rp=0).sel({
# 'S-Rp':0}), xDim = xDim, pType = 'r', squeeze = False, thres=None) #, cmap='vlag')
# # , fillna=True) # Note squeeze = False required for 1D case (should add this to
# code!)
# daPlotpd.fillna('')

# Glue versions for JupyterBook output

```

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```

glue("deltaTerm000-lmPlot", gFig.fig, display=False)
glue("deltaTerm000-tab", daPlotpd.fillna(''), display=False) # As above, but with PD
↪object return and glue.

```

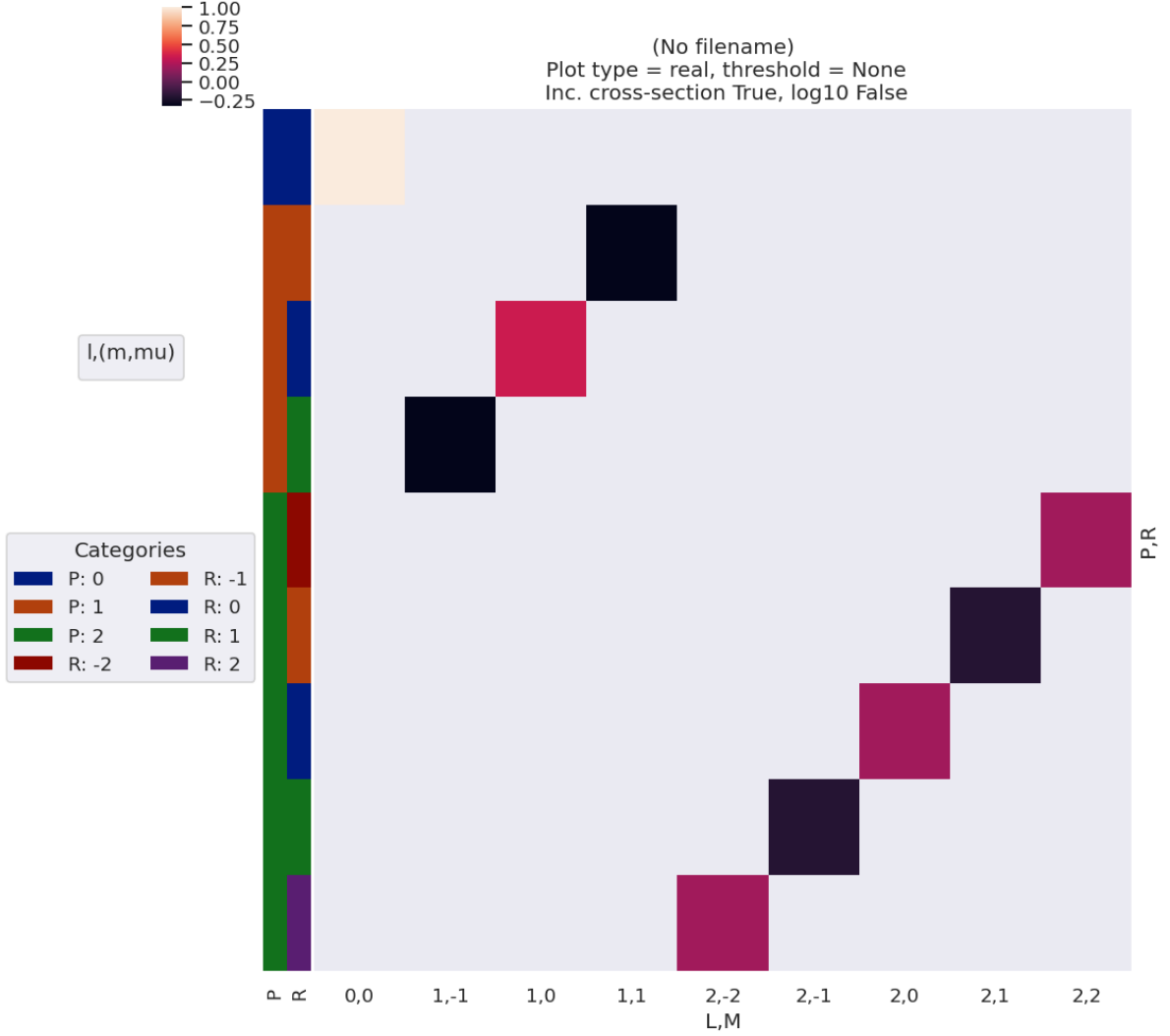


Fig. 4.8: Example $\Delta_{L,M}(0,0,0)$ basis functions (see also Fig. 4.9). For illustration purposes, the plot only shows terms for $R' = 0$. See main text for discussion.

```

*** Plot ADMs
daPlot, daPlotpd, legendList, ADMFig = ep.lmPlot(AKQS, xDim = 't', pType = 'r',
↪squeeze = False, cmap='vlag') # Note squeeze = False required for 1D case (should
↪add this to code!)
# daPlotpd

*** Plot subsection
daPlot, daPlotpd, legendList, AFFig = ep.lmPlot(Afterm.sel(R=0).sel(Rp=0).sel({'S-Rp
↪':0})),
                                xDim = 't', pType = 'r', squeeze =
↪False, cmap='vlag') # Note squeeze = False required for 1D case (should add this
↪to code!)

```

(continues on next page)

	L	0	1		2					
	M	0	-1	0	1	-2	-1	0	1	2
P	R									
0	0	1.0								
1	-1				-0.333					
	0			0.333						
	1		-0.333							
2	-2								0.2	
	-1							-0.2		
	0						0.2			
	1					-0.2				
	2				0.2					

Fig. 4.9: Example $\Delta_{L,M}(0,0,0)$ basis functions (see also Fig. 4.8). For illustration purposes, the table only shows terms for $R' = 0$. See main text for discussion.

(continued from previous page)

```
# Glue versions for JupyterBook output
glue("ADMs-linearRamp-lmPlot", ADMFig.fig, display=False)
glue("AFterm-linearRamp-lmPlot", AFFig.fig, display=False)
```

For cases with aligned molecular ensembles, additional terms can similarly appear depending on the alignment as well as the properties of the ionizing radiation. Again, the types of terms follow some typical patterns dependent on the symmetry of the ensemble, as well as the order of the terms allowed. For instance, $L_{max} = P_{max} + K_{max} = 2 + K_{max}$, and K_{max} represents the overall degree of alignment of the ensemble; hence an aligned ensemble may be signified by higher-order terms in the observable (if allowed by other terms in the overall expansion) or, equivalently, aligning an ensemble prior to ionization can be used as a way to control which terms contribute to the alignment tensor.

Since this is a coherent averaging, additional interferences can also appear in the AF - or be restricted in the AF - depending on these geometric parameters and the contributing matrix elements. Additionally, any effects modulating these terms, for instance a time-dependent alignment (rotational wavepacket), vibronic dynamics (vibrational and/or electronic wavepacket), time-dependent laser field (control field) may be anticipated to lead to both changes in these terms and, potentially, interesting effects in the observable. Such effects have been discussed in more detail in *Quantum Metrology* Vol. 2 [2], and in the current case the focus is purely on rotational wavepackets.

Fig. 4.11 shows $\tilde{\Delta}_{L,M}(t)$ for various choices of alignment (as per the *ADMs* shown in Fig. 4.10), and illustrates some of the general features discussed. Note, for example:

- L_{max} varies with alignment; in the demonstration case $K_{max} = 8$ at later times, resulting in $L_{max} = 10$, whilst at $t = 0$ $K_{max} = 0$, thus restricting terms to $L_{max} = 2$.
- Odd- L values are correlated with $P = 1$ terms.
- Only $M = 0$ terms are allowed in this case ($Q = S = 0$).

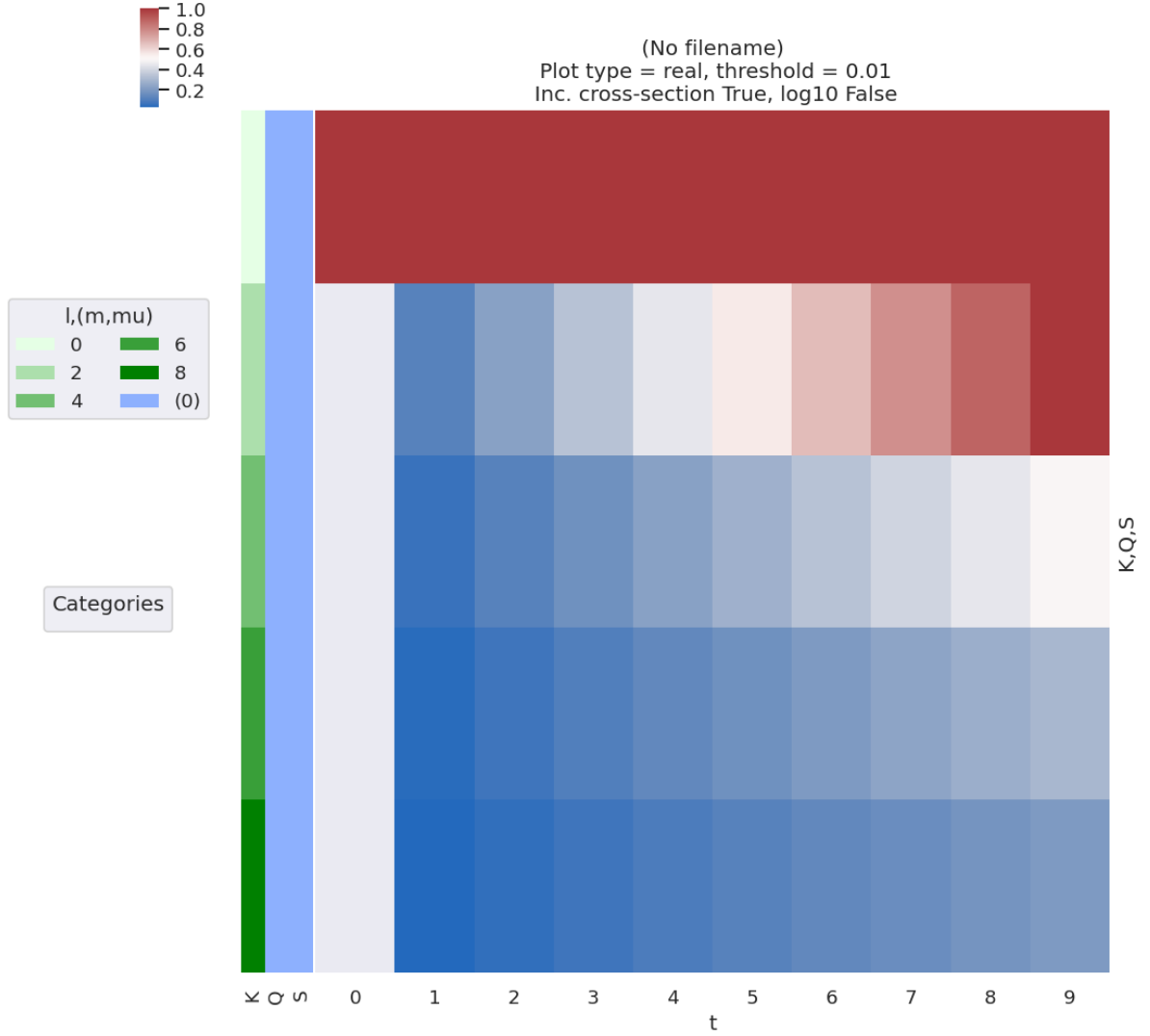


Fig. 4.10: Example ADMs used for AF basis function example (see Fig. 4.11). These ADMs essentially show an increasing degree of alignment with the t parameter, with high-order terms increasing at later t .

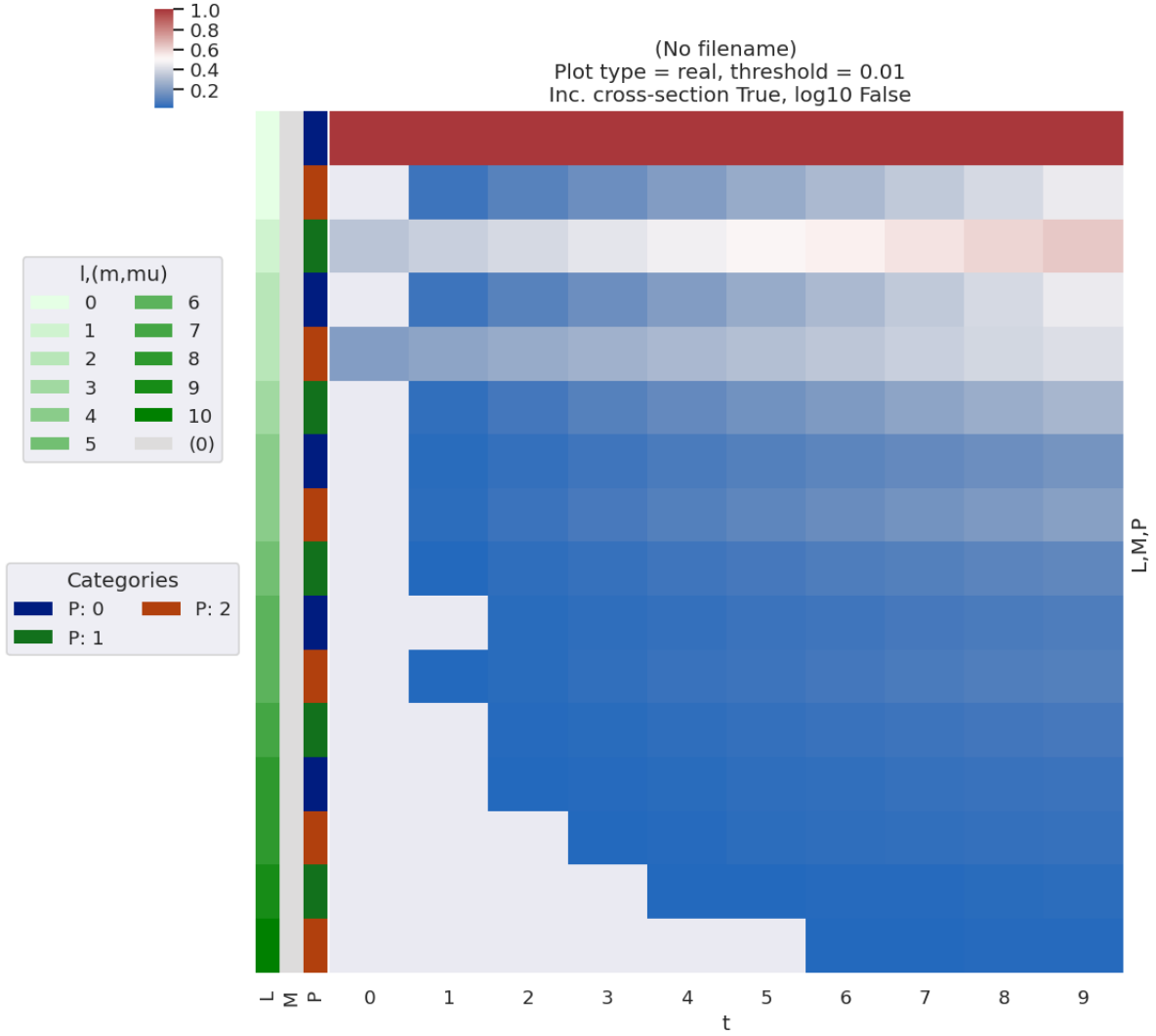


Fig. 4.11: Example of $\tilde{\Delta}_{L,M}(t)$ basis values for various choices of alignment (as per Fig. 4.10). The ADMs essentially show an increasing degree of alignment with the t parameter, with high-order terms increasing at later t , and this is reflected in the $\tilde{\Delta}_{L,M}(t)$ terms with higher-order L appearing at later t .

3D alignments and symmetry breaking

As discussed above, for the case where $Q \neq 0$ and/or $S \neq 0$ additional symmetry breaking can occur. It is simple to examine these effects numerically via changing the trial *ADMs* used to determine $\tilde{\Delta}_{L,M}(t)$ (Eq. (4.27)).

```

**** Neater version
# Set ADMs for increasing alignment...
tPoints = 10
inputADMs3D = [[0,0,0, *np.ones(tPoints)],      # * np.sqrt(4*np.pi)], # Optional
               ↪ multiplier for normalisation
               [2,0,0, *np.linspace(0,1,tPoints)],
               [2,0,2, *np.linspace(0,0.5,tPoints)],
               [2,2,0, *np.linspace(0,0.5,tPoints)],
               [2,2,2, *np.linspace(0,0.8,tPoints)]]

AKQS = ep.setADMs(ADMs = inputADMs3D)  # TODO WRAP TO CLASS (IF NOT ALREADY!)

# Compute alignment terms
AFTERM, DeltaTerm = ep.geomCalc.deltaLMKQS(EPR, AKQS)

**** Plot all
# daPlot, daPlotpd, legendList, gFig = ep.lmPlot(AFTERM,      #.sel(R=0).sel(Rp=0),
#                                                 xDim = 't', pType = 'r', squeeze =
↪ False, cmap='vlag') # Note squeeze = False required for 1D case (should add this
↪ to code!)

**** Plot subsection, L<=2
# .sel(R=0).sel(Rp=0) gives M=0 terms only for R=Rp=0
# .sel(R=0) gives M=0 terms only
# .sel(Rp=0) gives M!=0 terms, S-Rp=0 terms only

daPlot, daPlotpd, legendList, gFig = ep.lmPlot(AFTERM.where(AFTERM.L<=2).sel(Rp=0),
↪ # .sel(R=0).sel(Rp=0), #M=0 for R=Rp=0
                                                xDim = 't', pType = 'r', squeeze =
↪ False, cmap='vlag') # Note squeeze = False required for 1D case (should add this
↪ to code!)

# Glue versions for JupyterBook output
glue("ADMs-3DlinearRamp-lmPlot", ADMFig.fig, display=False)

```

For illustration purposes, Fig. 4.12 shows a subsection of the $\tilde{\Delta}_{L,M}(t)$ basis values, indicating some of the key features in the full 3D case, subselected for $L \leq 2$ and $R' = 0$ terms only.

Note, in particular, the presence of $M \neq 0$ terms in general, and a complicated dependence of the allowed terms on the alignment, which may increase, decrease, or even change sign. % TODO: give more explicit examples here, may want to also reduce and consolidate illustrations further? As previously, these behaviours are generally useful for understanding specific cases or planning experiments for specific systems; this is explored further in Part II which focuses on the results for particular molecules (hence symmetries and sets of matrix elements).

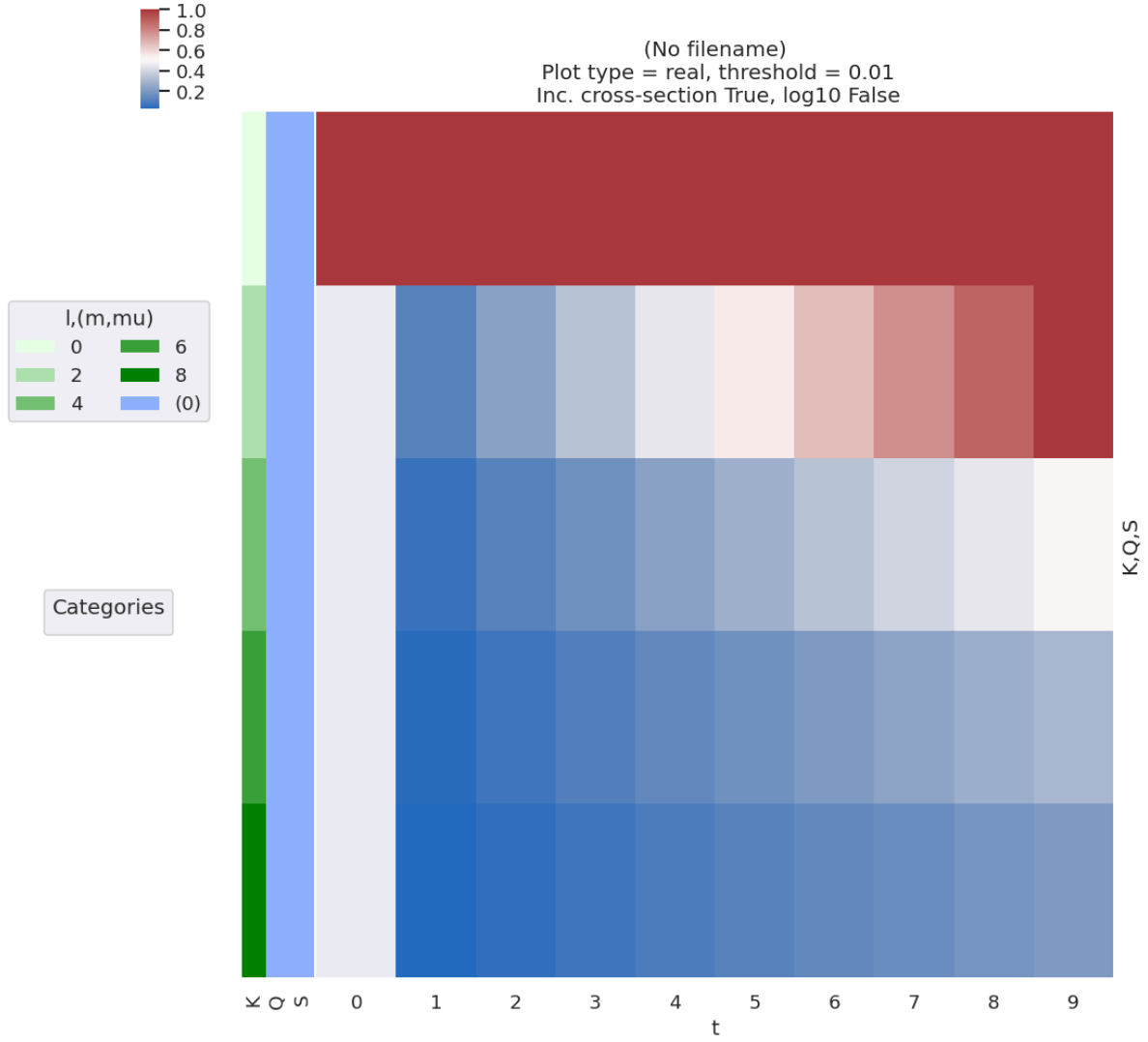


Fig. 4.12: Example $\tilde{\Delta}_{L,M}(t)$ basis values for various choices of “3D” alignment, i.e. including some $K \neq 0$ and $S \neq 0$ terms. Note, in particular, the presence of $M \neq 0$ terms in general, and a complicated dependence of the allowed terms on the alignment (*ADMs*), which may increase, decrease, or even change sign for a given L, M .

4.3.8 Tensor product terms

Following the above, further resultant terms can also be examined, up to and including the full channel functions $\Upsilon_{L,M}^{u,\zeta\zeta'}$ (see (4.12)) for a given case. Numerically these are all implemented in the main `ePSproc` codebase [12, 13, 14], and can be returned by these functions for inspection - the full basis set already defined includes some of these product. Custom tensor product terms are also readily computed with the codebase, with tensor multiplications handled natively by the Xarray data objects.

Polarisation & ADM product term: the main product basis returned, `polProd`, contains the tensor product $\Lambda_R \otimes E_{PR}(\hat{e}) \otimes \Delta_{L,M}(K, Q, S) \otimes A_{Q,S}^K(t)$, expanded over all quantum numbers (see [full definition here](#)). This term, therefore, incorporates all of the dependence (or response) of the AF- β_{LM} s on the polarisation state, and the axis distribution. Example results, making use of the linear-ramp *ADM*s of Sect. 4.3.7 are illustrated in Fig. 4.13.

The full channel (response) functions $\Upsilon_{L,M}^{u,\zeta\zeta'}$ as defined in (4.17) and (4.18) can be determined by the product of this term with the $B_{L,M}$ tensor. This is essentially the complete geometric basis set, % (give or take a phase term or two), hence equivalent to the AF- β_{LM} if the ionization matrix elements were set to unity. This illustrates not only the coupling of the geometric terms into the observable L, M , but also how the partial wave $|l, m\rangle$ terms map to the observables. Example results, making use of the linear-ramp *ADM*s of Sect. 4.3.7 are illustrated in Fig. 4.14.

```
# Set data - set example ADMs to data structure & subset for calculation
data.setADMs(ADMs = inputADMs)
data.setSubset(dataKey = 'ADM', dataType = 'ADM')

# Using PEMtk - this only returns the product basis set as used for fitting
BetaNormX, basisProduct = data.afblmMatEfit(selDims={}, sqThres=False)
```

```
Subselected from dataset 'ADM', dataType 'ADM': 50 from 50 points (100.00%)
```

```
basisKey = 'polProd' # Key for BLM basis set

# Reformat basis for display (optional)
# stackDims = {'LM':['L', 'M']}
# basisPlot = basis[basisKey].rename({'S-Rp':'M'}).stack(stackDims)

# Convert to Pandas
# pd, _ = ep.multiDimXrToPD(basisPlot, colDims=stackDims)

# Subselect on pol state
daPlot, daPlotpd, legendList, gFig = ep.lmPlot(basisProduct[basisKey].sel(Labels='A'),
    xDim='t', cmap = cmap, mDimLabel='mu'); # , cmap='vlag'); # Subselect on pol state
# ep.lmPlot(basisProduct[basisKey], xDim='t', cmap = cmap); # , cmap='vlag'); #
    Subselect on pol state

# Glue versions for JupyterBook output
glue("polProd-linearRamp-lmPlot", gFig.fig, display=False)
```

```
# Channel functions
daPlot, daPlotpd, legendList, gFig = ep.lmPlot((basisProduct['BLMtableResort'] *
    basisProduct['polProd']).sel(Labels='A').sel({'S-Rp':0}).sel(L=2), xDim='t',
    cmap=cmap, mDimLabel='m'); # Basic case
# ep.lmPlot((basis['BLMtableResort'] * basis['polProd']), xDim='t', cmap='vlag');

# Glue versions for JupyterBook output
glue("channelFunc-linearRamp-lmPlot", gFig.fig, display=False)
```

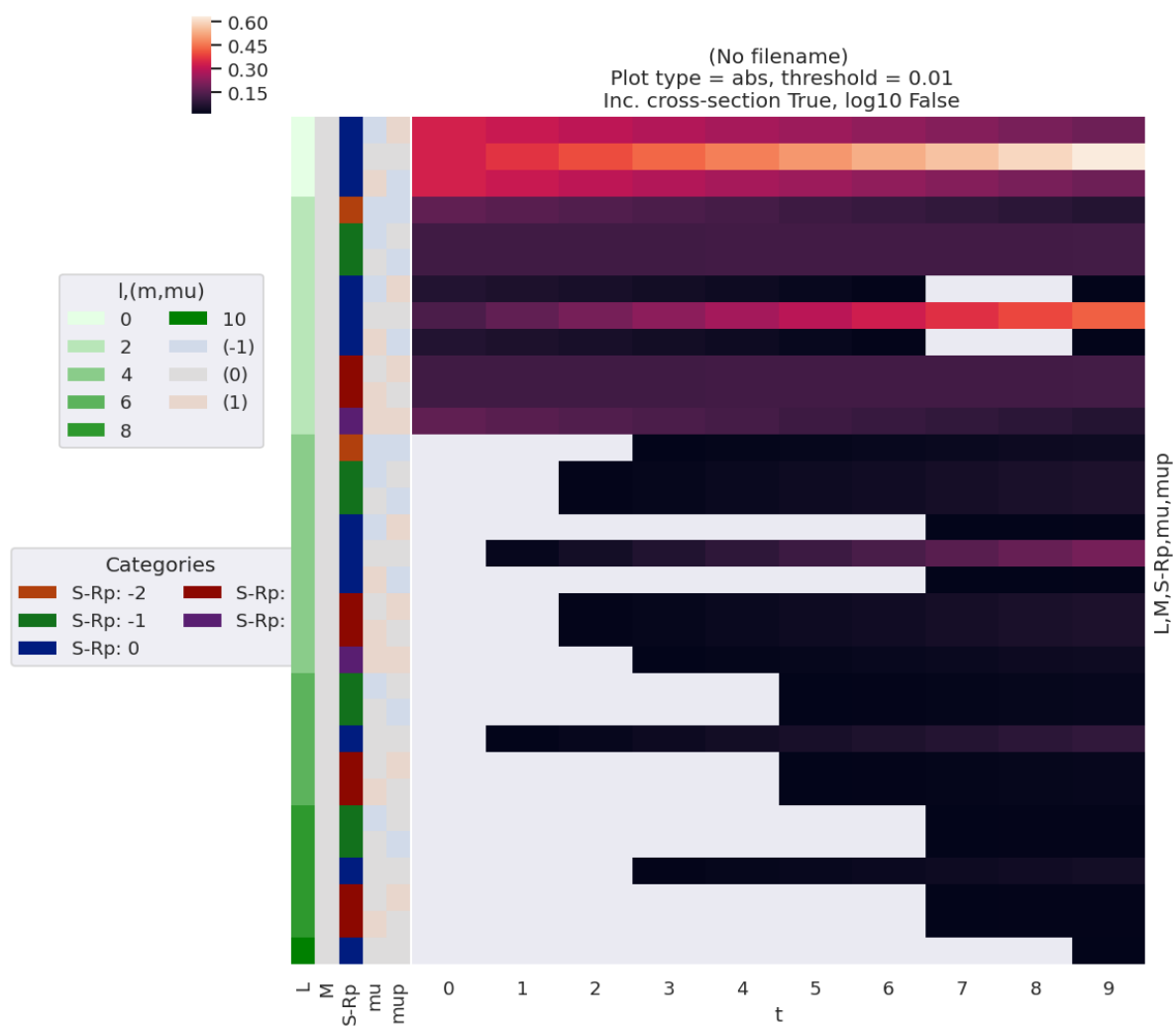


Fig. 4.13: Example product basis function for the polarisation and ADM terms.

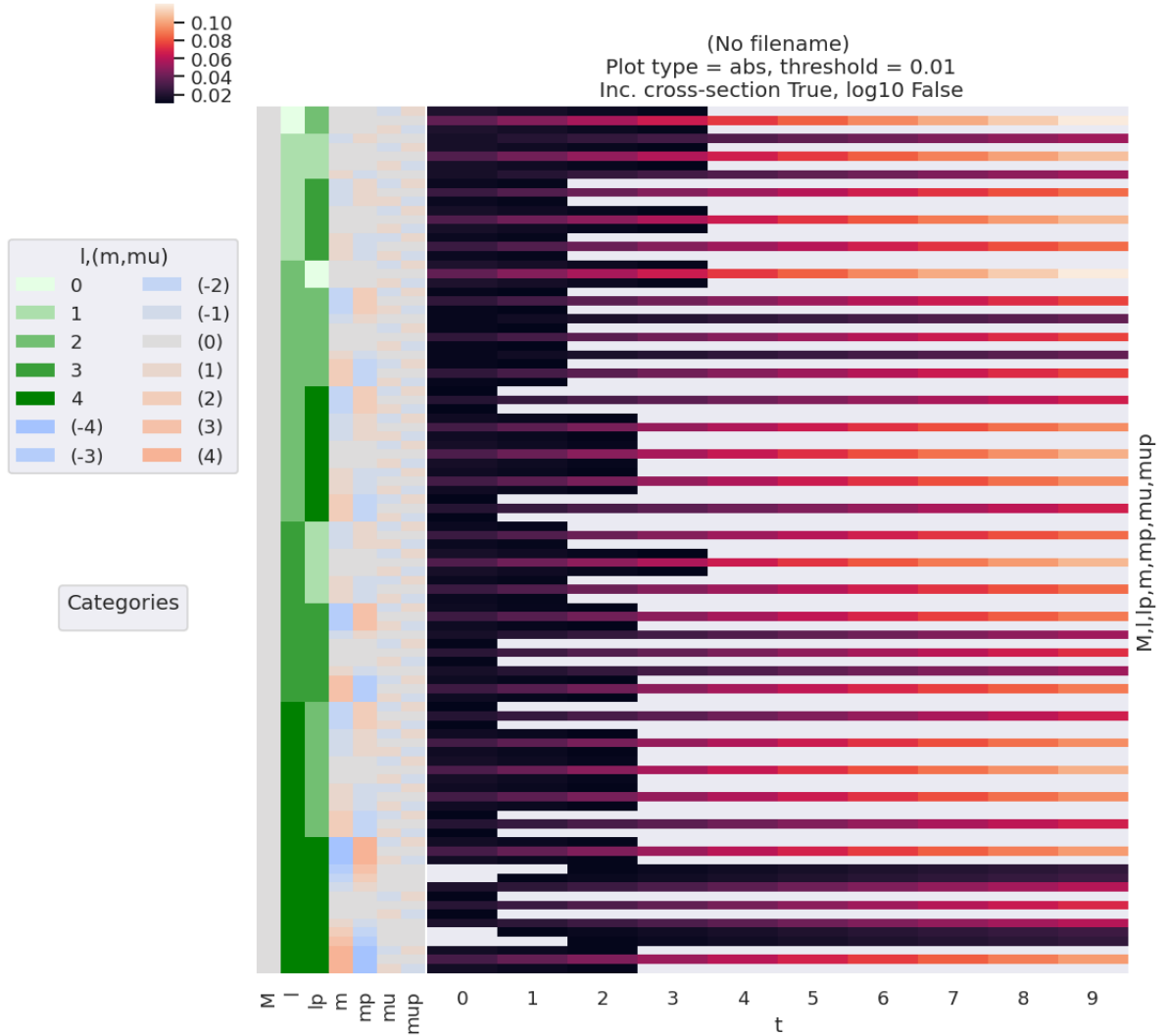


Fig. 4.14: Example of $\tilde{Y}_{L,M}^{u,\zeta\zeta'}$ basis values for various choices of alignment (as per Fig. 4.10), shown for $L = 2$ only. The basis essentially shows the observable terms if the ionization matrix elements are neglected, hence the sensitivity of the configuration to each pair of partial wave terms. Note, in general, that the sensitivity to any given pair $\langle l', m' | l, m \rangle$, increases with alignment (hence with t in this example) for the linear polarisation case ($\mu = \mu' = 0$), but typically decreases with alignment for cross-polarised terms.

4.4 Density matrix representation

The density operator associated with the continuum state in Eq. (4.3) is easily written as $\hat{\rho} = |\Psi_c\rangle\langle\Psi_c| \equiv |\mathbf{k}\rangle\langle\mathbf{k}|$.¹ The full final continuum state as a density matrix in the $\zeta\zeta'$ representation (with the observable dimensions L, M explicitly included in the density matrix), which will also be dependent on the choice of channel functions (u), can then be given as

$$\rho_{L,M}^{u,\zeta\zeta'} = \mathcal{T}_{L,M}^{u,\zeta\zeta'} \mathbb{I}^{\zeta,\zeta'} \quad (4.28)$$

Here the density matrix can be interpreted as the final, LF/AF or MF density matrix (depending on the channel functions used), incorporating both the intrinsic and extrinsic effects (i.e. all channel couplings and radial matrix elements for the given measurement), with dimensions dependent on the unique sets of quantum numbers required - in the simplest case, this will just be a set of partial waves $\zeta = (l, m)$.

In the channel function basis, this leads to a (radial or reduced) density matrix given by the radial matrix elements:

$$\rho^{\zeta\zeta'} = \mathbb{I}^{\zeta,\zeta'} \quad (4.29)$$

This form encodes purely intrinsic (molecular scattering) photoionization dynamics (thus characterises the scattering event), whilst the full form $\rho_{L,M}^{u,\zeta\zeta'}$ of Eq. (4.28) includes any additional effects incorporated via the channel functions. For reconstruction problems, it is usually the reduced form of Eq. (4.29) that is of interest, since the remainder of the problem is already described analytically by the channel functions $\mathcal{T}_{L,M}^{u,\zeta\zeta'}$. In other words, the retrieval of the radial matrix elements $\mathbb{I}^{\zeta,\zeta'}$ and the radial density matrix $\rho^{\zeta\zeta'}$ are equivalent, and both can be viewed as completely describing the photoionization dynamics.

The L, M notation for the full density matrix $\rho_{L,M}^{u,\zeta\zeta'}$ (Eq. (4.28)) indicates here that these dimensions should not be summed over, hence the tensor coupling into the $\beta_{L,M}^u$ parameters can also be written directly in terms of the density matrix:

$$\beta_{L,M}^u = \sum_{\zeta,\zeta'} \rho_{L,M}^{u,\zeta\zeta'} \quad (4.30)$$

In fact, this form arises naturally since the $\beta_{L,M}^u$ terms are the state multipoles (geometric tensors) defining the system, which can be thought of as a coupled basis equivalent of the density matrix representations (see, e.g., Ref. [52], Chpt. 4.).

In a more traditional notation (following Eq. (4.3), see also Ref. [54]), the density operator can be expressed as:

$$\rho(t) = \sum_{LM} \sum_{KQS} A_{QS}^K(t) \sum_{\zeta\zeta'} \mathcal{T}_{L,M}^{u,\zeta\zeta'} |\zeta, \Psi_+\rangle\langle\zeta, \Psi_+| \mu_q \rho_i \mu_{q'}^* |\zeta', \Psi_+\rangle\langle\zeta', \Psi_+| \quad (4.31)$$

with $\rho_i = |\Psi_i\rangle\langle\Psi_i|$. This is, effectively, equivalent to an expansion in the various tensor operators defined above, in a standard state-vector notation.

The main benefit of a density matrix representation in the current work is as a rapid way to visualize the phase relations between the photoionization matrix elements (the off-diagonal density matrix elements), and the ability to quickly check the overall pattern of the elements, hence confirm that no phase-relations are missing and orthogonality relations are fulfilled - some examples are given below. Since the method for computing the density matrices is also numerically equivalent to a tensor outer-product, density matrices and visualizations can also be rapidly composed for other properties of interest, e.g. the various channel functions defined herein, providing another complementary methodology and tool for investigation. (Further examples can be found in the [ePSproc documentation](#) [14], as well as in the literature, see, e.g., Ref. [52] for general discussion, Ref. [45] for application in pump-probe schemes.)

Furthermore, as noted above, the density matrix elements provide a complete description of the photoionization event, and hence make clear the equivalence of the “complete” photoionization experiments (and associated continuum reconstruction methods) discussed herein, with general quantum tomography schemes [55]. The density matrix can also be

¹ For general discussion of density matrix techniques and applications in AMO physics, see Blum’s textbook *Density Matrix Theory and Applications* [52], which is referred to extensively herein.

used as the starting point for further analysis based on standard density matrix techniques - this is discussed, for instance, in Ref. [52], and can also be viewed as a bridge between traditional methods in spectroscopy and AMO physics, and more recent concepts in the quantum information sciences (see, e.g., Refs. [56, 57] for recent discussions in this context).

4.4.1 Setup

This follows the setup in Sect. 4.3 *Tensor formulation of photoionization*, using a symmetry-based set of basis functions for demonstration purposes.

```
# Setup symmetry-defined matrix elements using PEMtk

# Import class
from pemtk.sym.symHarm import symHarm

# Compute hamronics for Td, lmax=4
sym = 'D2h'
lmax=4

lmaxPlot = 2 # Set lmaxPlot for subselection on plots later.

# Glue items for later
# glue("symHarmPGmatE", sym, display=False)
# glue("symHarmLmaxmatE", lmax, display=False)
# glue("symHarmBasislmaxPlot", lmaxPlot, display=False)

# TODO: consider different labelling here, can set at init e.g. dims = ['C', 'h', 'muX', 'l', 'm'] - 25/11/22 code currently fails for mu mapping, remap below instead
symObj = symHarm(sym, lmax)
# symObj = symHarm(sym, lmax, dims = ['Cont', 'h', 'muX', 'l', 'm'])

# To plot using ePSproc/PEMtk class, these values can be converted to ePSproc BLM_
# data type...

# Run conversion - the default is to set the coeffs to the 'BLM' data type
dimMap = {'C':'Cont', 'mu':'muX'}
symObj.toePSproc(dimMap=dimMap)

# Run conversion with a different dimMap & dataType
dataType = 'matE'
# symObj.toePSproc(dimMap = {'C':'Cont', 'h':'it', 'mu':'muX'}, dataType=dataType)
symObj.toePSproc(dimMap = dimMap, dataType=dataType)
# symObj.toePSproc(dimMap = {'C':'Cont', 'h':'it'}, dataType=dataType) # Drop mu >_
# muX mapping for now
# symObj.coeffs[dataType]

# Example using data class (setup in init script)
data = pemtkFit()

# Set to new key in data class
dataKey = sym
data.data[dataKey] = {}

for dataType in ['matE', 'BLM']:
    data.data[dataKey][dataType] = symObj.coeffs[dataType]['b (comp)'].sum(['h', 'muX', 'l']) # Select expansion in complex harmonics, and sum redundant dims
    data.data[dataKey][dataType].attrs = symObj.coeffs[dataType].attrs
```

```

# Compute basis functions for given matrix elements

# Set data
data.subKey = dataKey

# Using PEMtk - this only returns the product basis set as used for fitting
BetaNormX, basisProduct = data.afblmMatEfit(selDims={}, sqThres=False)

# Using ePSproc directly - this includes full basis return if specified
BetaNormX2, basisFull = ep.geomFunc.afblmXprod(data.data[data.subKey]['matE'],
↪ basisReturn = 'Full', selDims={}, sqThres=False) #, BLMRenorm = BLMRenorm,
↪ **kwargs)

# The basis dictionary contains various numerical parameters, these are investigated
↪ below.
# See also the ePSproc docs at https://epsproc.readthedocs.io/en/latest/methods/
↪ geometric_method_dev_260220_090420_tidy.html
print(f"Product basis elements: {basisProduct.keys()}")
print(f"Full basis elements: {basisFull.keys()}")

# Use full basis for following sections
basis = basisFull

```

```

Product basis elements: dict_keys(['BLMtableResort', 'polProd', 'phaseConvention',
↪ 'BLMRenorm'])
Full basis elements: dict_keys(['QNs', 'EPRX', 'lambdaTerm', 'BLMtable',
↪ 'BLMtableResort', 'AFterm', 'AKQS', 'polProd', 'phaseConvention', 'BLMRenorm'])

```

Compute density matrix and plot.

```

# DEMO CODE FROM http://jake/jupyter/user/paul/doc/tree/code-share/stimpy-docker-
↪ local/MFPADs_recon_manuscript_dev_April_2022/MFrecon_manuscript_fig_generation_
↪ 170422-Stimpy_MAIN-oldPkgs.ipynb
# SEE ALSO DOCS, https://epsproc.readthedocs.io/en/dev/methods/density_mat_notes_demo_
↪ 300821.html#Density-Matrices

# Import routines
from epsproc.calc import density

# Compose density matrix

# Set dimensions/state vector/representation
# These must be in original data, but will be restacked as necessary to define the
↪ effective basis space.
denDims = 'LM' #, 'mu']
selDims = None #{'Type': 'L'}
pTypes=['r', 'i']
thres = 1e-4 # 0.2 # Threshold out 1>3 terms if using full 'orb5' set.
normME = False
normDen = 'max'

# Calculate - Ref case
# matE = data.data['subset']['matE']
# Set data from master class
# k = 'orb5' # N2 orb5 (SG) dataset
# k = 'subset'

```

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```

k = sym
matE = data.data[k]['matE']
if normME:
    matE = matE/matE.max()

daOut, *_ = density.densityCalc(matE, denDims = denDims, selDims = selDims, thres = _
    ↪thres) # OK

if normDen=='max':
    daOut = daOut/daOut.max()
elif normDen=='trace':
    daOut = daOut/(daOut.sum('Sym').pipe(np.trace)**2) # Need sym sum here to get 2D
    ↪trace

# daPlot = density.matPlot(daOut.sum('Sym'))
daPlot = density.matPlot(daOut.sum('Sym'), pTypes=pTypes)

# # Retrieved
# matE = data.data['agg']['matE']['compC']
# if normME:
#     matE = matE/matE.max()

# daOut2, *_ = density.densityCalc(matE, denDims = denDims, selDims = selDims, thres = _
    ↪thres) # OK

# if normDen=='max':
#     daOut2 = daOut2/daOut2.max()
# elif normDen=='trace':
#     daOut2 = daOut2/(daOut2.sum('Sym').pipe(np.trace)**2)

# daPlot2 = density.matPlot(daOut2.sum('Sym'), pTypes=pTypes) #.sel(Eke=slice(0.5,1.
    ↪5,1)))

# # Compute difference
# daDiff = daOut.sum('Sym') - daOut2.sum('Sym')
# daDiff.name = 'Difference'
# daPlotDiff = density.matPlot(daDiff, pTypes=pTypes)

# #***** Plot
# daLayout = (daPlot.layout('pType') + daPlot2.opts(show_title=False).layout('pType').
    ↪opts(show_title=False) + daPlotDiff.opts(show_title=False).layout('pType')).cols(1)
    ↪ # No cols? AH - set to 1 works.
# # daLayout.opts(width=300, height=300) # Doesn't work?
# daLayout.opts(hvPlotters.opts.HeatMap(width=300, frame_width=300, aspect='square',
    ↪tools=['hover'], colorbar=True, cmap='coolwarm')) # .opts(show_title=False) # .
    ↪opts(title="Custom Title") #OK

```

Set plot kdims to ['LM', 'LM_p']; pass kdims = [dim1,dim2] for more control.

```

# Test with glue

glue("glueTest", daPlot, display=False)
glue("realOnly", daPlot.select(pType='Real'), display=False)

```

```
WARNING:param.HeatMapPlot03456: Due to internal constraints, when aspect and width/
↳height is set, the bokeh backend uses those values as frame_width/frame_height.
↳instead. This ensures the aspect is respected, but means that the plot might be
↳slightly larger than anticipated. Set the frame_width/frame_height explicitly to
↳suppress this warning.
WARNING:param.HeatMapPlot03530: Due to internal constraints, when aspect and width/
↳height is set, the bokeh backend uses those values as frame_width/frame_height.
↳instead. This ensures the aspect is respected, but means that the plot might be
↳slightly larger than anticipated. Set the frame_width/frame_height explicitly to
↳suppress this warning.
```

```
:HoloMap      [pType]
  :HeatMap    [LM, LM_p]    (b (comp))
```

Fig. 4.15: Raw output to glue.

```
:HeatMap      [LM, LM_p]    (b (comp))
```

Fig. 4.16: Real only output to glue.

```
# Test with glueHV()

glueHV("glueTestGHV", daPlot)
glueHV("realOnlyGHV", daPlot.select(pType='Real'))
```

```
WARNING:param.HeatMapPlot03568: Due to internal constraints, when aspect and width/
↳height is set, the bokeh backend uses those values as frame_width/frame_height.
↳instead. This ensures the aspect is respected, but means that the plot might be
↳slightly larger than anticipated. Set the frame_width/frame_height explicitly to
↳suppress this warning.
```

```
-----
ModuleNotFoundError                                Traceback (most recent call last)
/opt/conda/lib/python3.9/site-packages/bokeh/util/dependencies.py in import_
↳required(mod_name, error_msg)
    79     try:
--> 80         return import_module(mod_name)
    81     except ImportError as e:

/opt/conda/lib/python3.9/importlib/__init__.py in import_module(name, package)
    126         level += 1
--> 127     return _bootstrap._gcd_import(name[level:], package, level)
    128

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _gcd_import(name, package,
↳level)

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _find_and_load(name, import_)

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _find_and_load_unlocked(name,
↳import_)
```

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```

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _call_with_frames_removed(f,
↳ *args, **kwargs)

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _gcd_import(name, package,
↳ level)

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _find_and_load(name, import_)

/opt/conda/lib/python3.9/importlib/_bootstrap.py in _find_and_load_unlocked(name,
↳ import_)

ModuleNotFoundError: No module named 'selenium'

The above exception was the direct cause of the following exception:

RuntimeError                                Traceback (most recent call last)
/tmp/ipykernel_9213/890998523.py in <module>
      1 # Test with glueHV()
      2
----> 3 glueHV("glueTestGHV", daPlot)
      4 glueHV("realOnlyGHV", daPlot.select(pType='Real'))

~/jake-home/buildTmp/_latest_build/pdf/doc-source/scripts/setup_notebook.py in
↳ glueHV(name, fig, **kwargs)
    231     imgFile = f'{name}.{imgFormat}'
    232     imgFile = os.path.join(imgPath, imgFile)
--> 233     hv.save(fig, imgFile, fmt=imgFormat)
    234     # fig.write_image(imgFile, format=imgFormat) # See https://plotly.
↳ com/python/static-image-export/
    235

/opt/conda/lib/python3.9/site-packages/holoviews/util/__init__.py in save(obj,
↳ filename, fmt, backend, resources, toolbar, title, **kwargs)
    786     if formats[-1] in supported:
    787         filename = '.'.join(formats[:-1])
--> 788     return renderer_obj.save(obj, filename, fmt=fmt, resources=resources,
    789                             title=title)
    790

/opt/conda/lib/python3.9/site-packages/holoviews/plotting/renderer.py in save(self_
↳ or_cls, obj, basename, fmt, key, info, options, resources, title, **kwargs)
    593         return
    594
--> 595     rendered = self_or_cls(plot, fmt)
    596     if rendered is None: return
    597     (data, info) = rendered

/opt/conda/lib/python3.9/site-packages/holoviews/plotting/renderer.py in __call__
↳ (self, obj, fmt, **kwargs)
    197     return self.static_html(plot), info
    198     else:
--> 199     data = self._figure_data(plot, fmt, **kwargs)
    200     data = self._apply_post_render_hooks(data, obj, fmt)
    201     return data, info

```

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```

/opt/conda/lib/python3.9/site-packages/holoviews/plotting/bokeh/renderer.py in _
↳figure_data(self, plot, fmt, doc, as_script, **kwargs)
    125         elif fmt == 'png':
    126             from bokeh.io.export import get_screenshot_as_png
--> 127             img = get_screenshot_as_png(plot.state, driver=state.webdriver)
    128             imgByteArr = BytesIO()
    129             img.save(imgByteArr, format='PNG')

/opt/conda/lib/python3.9/site-packages/bokeh/io/export.py in get_screenshot_as_
↳png(obj, driver, timeout, resources, width, height)
    227
    228     '''
--> 229     from .webdriver import webdriver_control
    230
    231     with _tmp_html() as tmp:

/opt/conda/lib/python3.9/site-packages/bokeh/io/webdriver.py in <module>
    22
    23 from ..util.dependencies import import_required # isort:skip
--> 24 import_required("selenium.webdriver",
    25                 "To use bokeh.io image export functions you need selenium "
    26                 "('conda install selenium' or 'pip install selenium')")

/opt/conda/lib/python3.9/site-packages/bokeh/util/dependencies.py in import_
↳required(mod_name, error_msg)
    80         return import_module(mod_name)
    81     except ImportError as e:
--> 82         raise RuntimeError(error_msg) from e
    83
    84 #-----

↳---

RuntimeError: To use bokeh.io image export functions you need selenium ('conda_
↳install selenium' or 'pip install selenium')

```

Fig. 4.17: Raw output to glueHV.

Fig. 4.18: Real only output to glueHV.

4.5 Information content

As discussed in *Quantum Metrology* Vol. 2 [2], the information content of a single observable might be regarded as simply the number of contributing $\beta_{L,M}$ parameters. In set notation:

$$M = n\{\beta_{L,M}\} \quad (4.32)$$

where M is the information content of the measurement, defined as $n\{\dots\}$ the cardinality (number of elements) of the set of contributing parameters. A set of measurements, made for some experimental variable u , will then have a total information content:

$$M_u = \sum_u n\{\beta_{L,M}^u\}$$

In the case where a single measurement contains multiple $\beta_{L,M}$, e.g. as a function of energy ϵ or time t , the information content will naturally be larger:

$$\begin{aligned} M_u &= \sum_{u,k,t} n\{\beta_{L,M}^u(\epsilon, t)\} \\ &= M_u \times M \end{aligned}$$

where the second line pertains if each measurement has the same native information content, independent of u . It may be that the variable k is continuous (e.g. photoelectron energy), but in practice it will usually be discretized in some fashion by the measurement.

In terms of purely experimental methodologies, a larger M_u clearly defines a richer experimental measurement which explores more of the total measurement space spanned by the full set of $\{\beta_{L,M}^u(k, t)\}$. However, in this basic definition a larger M_u does not necessarily indicate a higher information content for quantum retrieval applications. The reason for this is simply down to the complexity of the problem (cf. Eq. (4.12)), in which many couplings define the sensitivity of the observable to the underlying system properties of interest. In this sense, more measurements, and larger M , may only add redundancy, rather than new information.

A more complete accounting of information content would, therefore, also include the channel couplings, i.e. sensitivity/dependence of the observable to a given system property, in some manner. For the case of a time-dependent measurement, arising from a rotational wavepacket, this can be written as:

$$M_u = n\{\mathcal{Y}_{L,M}^u(\epsilon, t)\}$$

In this case, each (ϵ, t) is treated as an independent measurement with unique information content, although there may be redundancy as a function of t depending on the nature of the rotational wavepacket and channel functions.

(Note this is in distinction to previously demonstrated cases where the time-dependence was created from a shaped laser-field, and was integrated over in the measurements, which provided a coherently-multiplexed case, see refs. [58, 59, 60] for details.)

NUMERICAL IMPLEMENTATION

Further details of the numerical implementations - base packages, conventions and refs - to go here or in appendix?

Part III

Backmatter

BIBLIOGRAPHY

GLOSSARY

MF Molecular frame (MF) - coordinate system referenced to the molecule, usually with the z-axis corresponding to the highest symmetry axis.

LF Laboratory or lab frame (LF) - coordinate system referenced to the laboratory frame, usually with the z-axis corresponding to the laser field polarization. For circularly or elliptically polarized light the propagation direction is conventionally used for the z-axis. In some cases a different z-axis may be chosen, e.g. as defined by a detector.

AF Aligned frame (AF) - coordinate system referenced to molecular alignment axis or axes. For 1D alignment, the z-axis usually corresponds to the alignment field polarization, and hence may be identical to the standard **LF** definition. For high degrees of (3D) alignment the AF may approach the **MF** in the ideal case, although will usually be limited by the symmetry of the system.

PADs Photoelectron angular distributions (PADs), often with a prefix denoting the reference frame, e.g. LFPADs, MFPADs (sometimes also hyphenated, e.g. LF-PADs). Usage is often synonymous with the associated *anisotropy parameters* (or “betas”).

anisotropy parameters Expansion parameters $\beta_{L,M}$ for an expansion in spherical harmonics (or similar basis sets of angular momentum functions in polar coordinates), e.g. Eq. (4.1). Often referred to simply as “beta parameters”, and may be dependent on various properties, e.g. $\beta_{L,M}(\epsilon, t \dots)$. Herein upper-case L, M usually refer to observables or the general case, whilst lower-case (l, m) usually refer specifically to the photoelectron wavefunction partial waves, and (l, λ) usually denote these terms referenced specifically to the molecular frame.

ADMs Expansion parameters $A_{Q,S}^K(t)$ for describing a molecular ensemble alignment described as a set of axis distribution moments, usually expanded as Wigner rotation matrix element, spherical harmonics or Legendre polynomial functions.

Axis distribution moments See *ADMs*.

Part IV

Test pages

BUILD VERSIONS AND CONFIG TESTS

8.1 Versions

```
import scooby
scooby.Report(additional=['pemtk', 'epsproc', 'xarray', 'pandas', 'scipy', 'matplotlib',
↳ 'jupyterlab', 'plotly', 'holoviews'])
```

```
* sparse not found, sparse matrix forms not available.
* natsort not found, some sorting functions not available.
* Setting plotter defaults with epsproc.basicPlotters.setPlotters(). Run directly_
↳ to modify, or change options in local env.
* Set Holoviews with bokeh.
* pyevtk not found, VTK export not available.
```

```
OMP: Info #271: omp_set_nested routine deprecated, please use omp_set_max_active_
↳ levels instead.
```

```
-----
Date: Tue Feb 07 12:32:51 2023 EST
```

```
      OS : Linux
      CPU(s) : 64
      Machine : x86_64
      Architecture : 64bit
      RAM : 62.8 GiB
      Environment : Jupyter
      File system : btrfs
```

```
Python 3.9.7 | packaged by conda-forge | (default, Sep 29 2021, 19:20:46)
[GCC 9.4.0]
```

```
      pemtk : 0.0.1
      epsproc : 1.3.2-dev
      xarray : 2022.3.0
      pandas : 1.3.4
      scipy : 1.7.1
      matplotlib : 3.4.3
      jupyterlab : 3.2.1
      plotly : 5.11.0
      holoviews : 1.15.2
      numpy : 1.20.3
```

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```
IPython : 7.28.0
scooby  : 0.7.0
-----
```

```
!jupyter-book --version
```

```
Jupyter Book      : 0.13.1
External ToC      : 0.2.4
MyST-Parser       : 0.15.2
MyST-NB           : 0.13.2
Sphinx Book Theme : 0.3.3
Jupyter-Cache     : 0.4.3
NbClient          : 0.5.4
```

8.2 Docker build env

To do

8.3 Book versions

```
QMpath = '/home/jovyan/QM3'
!git -C {QMpath} branch
!git -C {QMpath} log --format="%H" -n 1
```

```
gh-pages
* main
77f5cf1734bd9e53f3d804539ea151435ede8d9d
```

```
# Check current remote commits
!git ls-remote --heads https://github.com/phockett/Quantum-Metrology-with-
Photoelectrons-Vol3
```

d4b9479728c9ecab2c6a332864bb105e555f8e7f	refs/heads/gh-pages
77f5cf1734bd9e53f3d804539ea151435ede8d9d	refs/heads/main

8.4 Github pkg versions

Note - can't get versions for local pip installs from repo(?).

```
from pathlib import Path
import epsproc as ep
ep.__file__
```

```
'/home/jovyan/.local/lib/python3.9/site-packages/epsproc/__init__.py'
```

```
# Check current Git commit for local ePSproc version - NOTE THIS ONLY WORKS FOR_
↳INSTALLED FROM GIT CLONES
# from pathlib import Path
# import epsproc as ep
!git -C {Path(ep.__file__).parent} branch
!git -C {Path(ep.__file__).parent} log --format="%H" -n 1
```

```
fatal: not a git repository (or any of the parent directories): .git
fatal: not a git repository (or any of the parent directories): .git
```

```
# Check current remote commits
!git ls-remote --heads https://github.com/phockett/ePSproc
```

897d73392a7b32ffba4ca6b6b4755c61e7c1c8d7	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/certifi-2022.12.7	
92c661789a7d2927f2b53d7266f57de70b3834fa	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/mistune-2.0.3	
fe1e9540c7b91fe571f60562acd31d8e489d491e	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/nbconvert-6.5.1	
70b80a1e3a54de91c2bfe3b6be82d611fcfd5f43	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/pillow-9.3.0	
92fc79b09aafedadcb645f88bb7ed771c96d5b52	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/setuptools-65.5.1	
fa33ed8d63a5c4a4043cc4c261059cc09e4c2bf7	refs/heads/dependabot/pip/notes/
↳envs/envs-versioned/wheel-0.38.1	
41cdfe43750e08c510f98b05e024a9c62da42771	refs/heads/dependabot/pip/
↳setuptools-65.5.1	
dec4e5902adc562ef2bc9baeae19da1781b7e98b	refs/heads/dev
1c0b8fd409648f07c85f4f20628b5ea7627e0c4e	refs/heads/master
69cd89ce5bc0ad6d465a4bd8df6fba15d3fd1aee	refs/heads/numba-tests
ea30878c842f09d525fbf39fa269fa2302a13b57	refs/heads/revert-9-master
baf0be0c962e8ab3c3df57c8f70f0e939f99cbd7	refs/heads/testDev

```
# Check current remote commits
!git ls-remote --heads https://github.com/phockett/PEMtk
```

7b3bcb92ce5262cfc46562366c2ce024dd4509d5	refs/heads/master
3f4686dffdbb310f15692f978ba36d6a3d15e8d3	refs/heads/mfFittingDev

8.5 Full conda env

```
!conda list
```

```
# packages in environment at /opt/conda:
#
# Name                               Version                               Build Channel
```

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_libgcc_mutex	0.1	conda_forge	conda-forge
_openmp_mutex	4.5	1_llvm	conda-forge
alabaster	0.7.12	pypi_0	pypi
alembic	1.7.4	pyhd8ed1ab_0	conda-forge
alsa-lib	1.2.3.2	h166bdaf_0	conda-forge
altair	4.1.0	py_1	conda-forge
ansi2html	1.8.0	py39hf3d152e_1	conda-forge
anyio	3.3.4	py39hf3d152e_0	conda-forge
appdirs	1.4.4	pyh9f0ad1d_0	conda-forge
argon2-cffi	21.1.0	py39h3811e60_0	conda-forge
asteval	0.9.27	pyhd8ed1ab_0	conda-forge
astropy	5.0.4	py39hd257fcd_0	conda-forge
async_generator	1.10	py_0	conda-forge
atk-1.0	2.36.0	h3371d22_4	conda-forge
attrs	21.2.0	pyhd8ed1ab_0	conda-forge
babel	2.9.1	pyh44b312d_0	conda-forge
backcall	0.2.0	pyh9f0ad1d_0	conda-forge
backports	1.0	py_2	conda-forge
backports.functools_lru_cache	1.6.4	pyhd8ed1ab_0	conda-forge
beautifulsoup4	4.10.0	pyha770c72_0	conda-forge
blas	2.112	openblas	conda-forge
blas-devel	3.9.0	12_linux64_openblas	conda-forge
bleach	4.1.0	pyhd8ed1ab_0	conda-forge
blinker	1.4	py_1	conda-forge
blosc	1.21.0	h9c3ff4c_0	conda-forge
bokeh	2.4.1	py39hf3d152e_1	conda-forge
boost-cpp	1.74.0	h312852a_4	conda-forge
bottleneck	1.3.2	py39hce5d2b2_4	conda-forge
brtoli	1.0.9	h7f98852_5	conda-forge
brtoli-bin	1.0.9	h7f98852_5	conda-forge
brtoli-python	1.0.9	py39h5a03fae_7	conda-forge
brtolipty	0.7.0	py39h3811e60_1001	conda-forge
brunsl	0.1	h9c3ff4c_0	conda-forge
bzip2	1.0.8	h7f98852_4	conda-forge
c-ares	1.18.1	h7f98852_0	conda-forge
c-blosc2	2.0.4	h5f21a17_1	conda-forge
ca-certificates	2022.9.24	ha878542_0	conda-forge
cached-property	1.5.2	hd8ed1ab_1	conda-forge
cached_property	1.5.2	pyha770c72_1	conda-forge
cairo	1.16.0	h6cf1ce9_1008	conda-forge
cartopy	0.20.1	py39he7aa91e_2	conda-forge
certifi	2022.9.24	pyhd8ed1ab_0	conda-forge
certipy	0.1.3	py_0	conda-forge
cffi	1.14.6	py39h4bc2ebd_1	conda-forge
cfitsio	4.0.0	h9a35b8e_0	conda-forge
cftime	1.6.0	py39hd257fcd_1	conda-forge
chardet	4.0.0	py39hf3d152e_1	conda-forge
charls	2.2.0	h9c3ff4c_0	conda-forge
charset-normalizer	2.0.0	pyhd8ed1ab_0	conda-forge
chrpath	0.16	h7f98852_1002	conda-forge
click	8.0.3	py39hf3d152e_0	conda-forge
cloudpickle	2.0.0	pyhd8ed1ab_0	conda-forge
colorama	0.4.4	pyh9f0ad1d_0	conda-forge
colorcet	3.0.1	pyhd8ed1ab_0	conda-forge
conda	4.10.3	py39hf3d152e_2	conda-forge
conda-package-handling	1.7.3	py39h3811e60_0	conda-forge

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configurable-http-proxy	4.5.0	node15_he6ea98c_0	conda-forge
cryptography	35.0.0	py39h95dcef6_1	conda-forge
curl	7.79.1	h2574ce0_1	conda-forge
cycler	0.10.0	py_2	conda-forge
cython	0.29.24	py39he80948d_0	conda-forge
cytoolz	0.11.0	py39h3811e60_3	conda-forge
dash	2.7.0	pyhd8ed1ab_0	conda-forge
dask	2021.10.0	pyhd8ed1ab_0	conda-forge
dask-core	2021.10.0	pyhd8ed1ab_0	conda-forge
dbus	1.13.6	h5008d03_3	conda-forge
dcw-gmt	2.1.1	ha770c72_0	conda-forge
debugpy	1.4.1	py39he80948d_0	conda-forge
decorator	5.1.0	pyhd8ed1ab_0	conda-forge
defusedxml	0.7.1	pyhd8ed1ab_0	conda-forge
dill	0.3.4	pyhd8ed1ab_0	conda-forge
distributed	2021.10.0	py39hf3d152e_0	conda-forge
docutils	0.17.1	pypi_0	pypi
ducc0	0.23.0	py39h5a03fae_0	conda-forge
entrypoints	0.3	py39hde42818_1002	conda-forge
epsproc	1.3.2.dev0	pypi_0	pypi
expat	2.4.8	h27087fc_0	conda-forge
ffmpeg	4.4.0	h6987444_4	conda-forge
fftw	3.3.10	nompi_h77c792f_102	conda-forge
flask	2.1.3	pyhd8ed1ab_0	conda-forge
flask-compress	1.13	pyhd8ed1ab_0	conda-forge
font-ttf-dejavu-sans-mono	2.37	hab24e00_0	conda-forge
font-ttf-inconsolata	3.000	h77eed37_0	conda-forge
font-ttf-source-code-pro	2.038	h77eed37_0	conda-forge
font-ttf-ubuntu	0.83	hab24e00_0	conda-forge
fontconfig	2.14.0	h8e229c2_0	conda-forge
fonts-conda-ecosystem	1	0	conda-forge
fonts-conda-forge	1	0	conda-forge
freetype	2.10.4	h0708190_1	conda-forge
freexl	1.0.6	h7f98852_0	conda-forge
fribidi	1.0.10	h36c2ea0_0	conda-forge
fsspec	2021.10.1	pyhd8ed1ab_0	conda-forge
future	0.18.2	pyhd8ed1ab_6	conda-forge
gdal	3.3.3	py39h0494519_2	conda-forge
gdk-pixbuf	2.42.6	h04a7f16_0	conda-forge
geos	3.10.0	h9c3ff4c_0	conda-forge
geotiff	1.7.0	hcfb7246_3	conda-forge
gettext	0.19.8.1	h73d1719_1008	conda-forge
ghostscript	9.54.0	h27087fc_2	conda-forge
ghp-import	2.1.0	pypi_0	pypi
giflib	5.2.1	h36c2ea0_2	conda-forge
gitdb	4.0.9	pypi_0	pypi
gitpython	3.1.29	pypi_0	pypi
glib	2.70.2	h780b84a_4	conda-forge
glib-tools	2.70.2	h780b84a_4	conda-forge
gmp	6.2.1	h58526e2_0	conda-forge
gmpy2	2.1.0b5	py39h78fa15d_0	conda-forge
gmt	6.2.0	h7e416ca_3	conda-forge
gnuplot	5.4.3	hedd7fda_0	conda-forge
gnutls	3.6.13	h85f3911_1	conda-forge
graphicsmagick	1.3.37	hdc87540_0	conda-forge
graphite2	1.3.13	h58526e2_1001	conda-forge

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greenlet	1.1.2	py39he80948d_0	conda-forge
gshhg-gmt	2.3.7	ha770c72_1003	conda-forge
gst-plugins-base	1.18.5	hf529b03_3	conda-forge
gstreamer	1.18.5	h9f60fe5_3	conda-forge
gtk2	2.24.33	h539f30e_1	conda-forge
h5py	3.4.0	nompi_py39h7e08c79_101	conda-forge
harfbuzz	3.1.1	h83ec7ef_0	conda-forge
hdf4	4.2.15	h10796ff_3	conda-forge
hdf5	1.12.1	nompi_h2750804_101	conda-forge
heapdict	1.0.1	py_0	conda-forge
holoviews	1.15.2	pyhd8ed1ab_0	conda-forge
hvplot	0.8.1	py_0	pyviz
icu	68.2	h9c3ff4c_0	conda-forge
idna	3.1	pyhd3deb0d_0	conda-forge
imagecodecs	2021.8.26	py39h571908b_2	conda-forge
imageio	2.9.0	py_0	conda-forge
imagesize	1.4.1	pypi_0	pypi
importlib-metadata	4.8.1	py39hf3d152e_0	conda-forge
importlib_metadata	4.8.1	hd8ed1ab_1	conda-forge
importlib_resources	5.3.0	pyhd8ed1ab_0	conda-forge
ipykernel	6.4.2	py39hef51801_0	conda-forge
ipympl	0.8.2	pyhd8ed1ab_0	conda-forge
ipython	7.28.0	py39hef51801_0	conda-forge
ipython_genutils	0.2.0	py_1	conda-forge
ipywidgets	7.6.5	pyhd8ed1ab_0	conda-forge
itsdangerous	2.1.2	pyhd8ed1ab_0	conda-forge
jbig	2.1	h7f98852_2003	conda-forge
jedi	0.18.0	py39hf3d152e_2	conda-forge
jinja2	3.0.2	pyhd8ed1ab_0	conda-forge
joblib	1.1.0	pyhd8ed1ab_0	conda-forge
jpeg	9d	h36c2ea0_0	conda-forge
json-c	0.15	h98cffda_0	conda-forge
json5	0.9.5	pyh9f0ad1d_0	conda-forge
jsonschema	3.2.0	pypi_0	pypi
jupyter-book	0.13.1	pypi_0	pypi
jupyter-cache	0.4.3	pypi_0	pypi
jupyter-dash	0.4.2	pyhd8ed1ab_1	conda-forge
jupyter-server-mathjax	0.2.6	pypi_0	pypi
jupyter-sphinx	0.3.2	pypi_0	pypi
jupyter_client	7.0.6	pyhd8ed1ab_0	conda-forge
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jupyter_telemetry	0.1.0	pyhd8ed1ab_1	conda-forge
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jupyterhub-base	1.4.2	py39hf3d152e_0	conda-forge
jupyterlab	3.2.1	pyhd8ed1ab_0	conda-forge
jupyterlab-spellchecker	0.7.2	pypi_0	pypi
jupyterlab_pygments	0.1.2	pyh9f0ad1d_0	conda-forge
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jupyterlab_widgets	1.0.2	pyhd8ed1ab_0	conda-forge
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kealib	1.4.14	h87e4c3c_3	conda-forge
kiwisolver	1.3.2	py39h1a9c180_0	conda-forge
krb5	1.19.2	hcc1bbae_2	conda-forge

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lame	3.100	h7f98852_1001	conda-forge
latexcodec	2.0.1	pypi_0	pypi
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lerc	3.0	h9c3ff4c_0	conda-forge
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libarchive	3.5.2	hccf745f_1	conda-forge
libblas	3.9.0	12_linux64_openblas	conda-forge
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libcurl	7.79.1	h2574ce0_1	conda-forge
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libdeflate	1.8	h7f98852_0	conda-forge
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libnetcdf	4.8.1	nompi_hb3fd0d9_101	conda-forge
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linkify-it-py	1.0.3	pypi_0	pypi
llvm-openmp	12.0.1	h4bd325d_1	conda-forge
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loket	0.2.0	py_2	conda-forge
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lzo	2.10	h516909a_1000	conda-forge
mako	1.1.5	pyhd8ed1ab_0	conda-forge
mamba	0.17.0	py39h951de11_0	conda-forge
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markdown-it-py	1.1.0	pypi_0	pypi
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mathjax	2.7.7	ha770c72_3	conda-forge
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ncurses	6.2	h58526e2_4	conda-forge
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param	1.12.2	pyh6c4a22f_0	conda-forge
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partd	1.2.0	pyhd8ed1ab_0	conda-forge
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poppler	21.09.0	ha39eefc_3	conda-forge
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postgresql	13.5	h2510834_1	conda-forge
proj	8.1.1	h277dcde_2	conda-forge
prometheus_client	0.11.0	pyhd8ed1ab_0	conda-forge
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protobuf	3.18.1	py39he80948d_0	conda-forge
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pybtex-docutils	1.0.2	pypi_0	pypi
pycosat	0.6.3	py39h3811e60_1006	conda-forge
pycparser	2.20	pyh9f0ad1d_2	conda-forge
pyct	0.4.6	py_0	conda-forge
pyct-core	0.4.6	py_0	conda-forge
pycurl	7.44.1	py39h72e3413_0	conda-forge
pydata-sphinx-theme	0.8.1	pypi_0	pypi
pyerfa	2.0.0.1	py39hd257fcd_2	conda-forge
pygments	2.10.0	pyhd8ed1ab_0	conda-forge
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pyrsistent	0.17.3	py39h3811e60_2	conda-forge
pyshp	2.3.1	pyhd8ed1ab_0	conda-forge
pyshtools	4.10	py39hd777142_0	conda-forge
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pytables	3.6.1	py39h2669a42_4	conda-forge
python	3.9.7	hb7a2778_3_cpython	conda-forge
python-dateutil	2.8.2	pyhd8ed1ab_0	conda-forge
python-json-logger	2.0.1	pyh9f0ad1d_0	conda-forge
python-kaleido	0.2.1	pyhd8ed1ab_0	conda-forge
python_abi	3.9	2_cp39	conda-forge
pytz	2021.3	pyhd8ed1ab_0	conda-forge
pyviz_comms	2.2.1	pyhd8ed1ab_1	conda-forge

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pyyaml	6.0	py39h3811e60_0	conda-forge
pyzmq	22.3.0	py39h37b5a0c_0	conda-forge
qt	5.12.9	hda022c4_4	conda-forge
quaternion	2022.4.2	py39hd257fcd_0	conda-forge
readline	8.1	h46c0cb4_0	conda-forge
reproc	14.2.3	h7f98852_0	conda-forge
reproc-cpp	14.2.3	h9c3ff4c_0	conda-forge
requests	2.26.0	pyhd8ed1ab_0	conda-forge
requests-unixsocket	0.2.0	py_0	conda-forge
retrying	1.3.3	py_2	conda-forge
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ruamel.yaml.clib	0.2.2	py39h3811e60_2	conda-forge
ruamel.yaml	0.15.80	py39h3811e60_1004	conda-forge
scikit-image	0.18.3	py39hde0f152_0	conda-forge
scikit-learn	1.0.1	py39h7c5d8c9_1	conda-forge
scipy	1.7.1	py39hee8e79c_0	conda-forge
scooby	0.7.0	pyhd8ed1ab_1	conda-forge
seaborn	0.9.0	py_2	conda-forge
seaborn-base	0.11.2	pyhd8ed1ab_0	conda-forge
send2trash	1.8.0	pyhd8ed1ab_0	conda-forge
setuptools	58.2.0	py39hf3d152e_0	conda-forge
setuptools-scm	6.4.2	pyhd8ed1ab_0	conda-forge
shapely	1.8.0	py39hc7dd4e9_2	conda-forge
six	1.16.0	pyh6c4a22f_0	conda-forge
smmap	5.0.0	pypi_0	pypi
snappy	1.1.8	he1b5a44_3	conda-forge
sniffio	1.2.0	py39hf3d152e_1	conda-forge
snowballstemmer	2.2.0	pypi_0	pypi
sortedcontainers	2.4.0	pyhd8ed1ab_0	conda-forge
soupsieve	2.0.1	py_1	conda-forge
spherical_functions	2022.4.1	pyhd8ed1ab_1	conda-forge
sphinx	4.5.0	pypi_0	pypi
sphinx-book-theme	0.3.3	pypi_0	pypi
sphinx-comments	0.0.3	pypi_0	pypi
sphinx-copybutton	0.5.0	pypi_0	pypi
sphinx-design	0.1.0	pypi_0	pypi
sphinx-external-toc	0.2.4	pypi_0	pypi
sphinx-jupyterbook-latex	0.4.7	pypi_0	pypi
sphinx-multitoc-numbering	0.1.3	pypi_0	pypi
sphinx-thebe	0.1.2	pypi_0	pypi
sphinx-togglebutton	0.3.2	pypi_0	pypi
sphinxcontrib-applehelp	1.0.2	pypi_0	pypi
sphinxcontrib-bibtex	2.5.0	pypi_0	pypi
sphinxcontrib-devhelp	1.0.2	pypi_0	pypi
sphinxcontrib-htmlhelp	2.0.0	pypi_0	pypi
sphinxcontrib-jsmath	1.0.1	pypi_0	pypi
sphinxcontrib-qthelp	1.0.3	pypi_0	pypi
sphinxcontrib-serializinghtml	1.1.5	pypi_0	pypi
spinsfast	2022.4.1	py39hc8e5db4_2	conda-forge
sqlalchemy	1.4.26	py39h3811e60_0	conda-forge
sqlite	3.36.0	h9cd32fc_2	conda-forge
statsmodels	0.13.0	py39hce5d2b2_0	conda-forge
sympy	1.9	py39hf3d152e_0	conda-forge
tblib	1.7.0	pyhd8ed1ab_0	conda-forge
tenacity	8.1.0	pyhd8ed1ab_0	conda-forge

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terminado	0.12.1	py39hf3d152e_0	conda-forge
testpath	0.5.0	pyhd8ed1ab_0	conda-forge
threadpoolctl	3.0.0	pyh8a188c0_0	conda-forge
tifffile	2021.10.12	pyhd8ed1ab_0	conda-forge
tiledb	2.3.4	he87e0bf_0	conda-forge
tk	8.6.11	h27826a3_1	conda-forge
toml	0.10.2	pyhd8ed1ab_0	conda-forge
tomli	2.0.1	pyhd8ed1ab_0	conda-forge
toolz	0.11.1	py_0	conda-forge
tornado	6.1	py39h3811e60_1	conda-forge
tqdm	4.62.3	pyhd8ed1ab_0	conda-forge
traitlets	5.1.1	pyhd8ed1ab_0	conda-forge
typing_extensions	3.10.0.2	pyha770c72_0	conda-forge
tzcode	2022a	h166bdaf_0	conda-forge
tzdata	2021e	he74cb21_0	conda-forge
uc-micro-py	1.0.1	pypi_0	pypi
uncertainties	3.1.7	pyhd8ed1ab_0	conda-forge
urllib3	1.26.7	pyhd8ed1ab_0	conda-forge
wcwidth	0.2.5	pyh9f0ad1d_2	conda-forge
webencodings	0.5.1	py_1	conda-forge
websocket-client	0.57.0	py39hf3d152e_4	conda-forge
werkzeug	2.1.2	pyhd8ed1ab_1	conda-forge
wget	3.2	pypi_0	pypi
wheel	0.37.0	pyhd8ed1ab_1	conda-forge
widgetsnbextension	3.5.1	py39hf3d152e_4	conda-forge
x264	1!161.3030	h7f98852_1	conda-forge
x265	3.5	h924138e_3	conda-forge
xarray	2022.3.0	pyhd8ed1ab_0	conda-forge
xerces-c	3.2.3	h9d8b166_3	conda-forge
xlrd	2.0.1	pyhd8ed1ab_3	conda-forge
xorg-kbproto	1.0.7	h7f98852_1002	conda-forge
xorg-libice	1.0.10	h7f98852_0	conda-forge
xorg-libsm	1.2.3	hd9c2040_1000	conda-forge
xorg-libx11	1.7.2	h7f98852_0	conda-forge
xorg-libxau	1.0.9	h7f98852_0	conda-forge
xorg-libxdmcp	1.1.3	h7f98852_0	conda-forge
xorg-libxext	1.3.4	h7f98852_1	conda-forge
xorg-libxrender	0.9.10	h7f98852_1003	conda-forge
xorg-libxt	1.2.1	h7f98852_2	conda-forge
xorg-renderproto	0.11.1	h7f98852_1002	conda-forge
xorg-xextproto	7.3.0	h7f98852_1002	conda-forge
xorg-xproto	7.0.31	h7f98852_1007	conda-forge
xz	5.2.5	h516909a_1	conda-forge
yaml	0.2.5	h516909a_0	conda-forge
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zipp	3.6.0	pyhd8ed1ab_0	conda-forge
zlib	1.2.11	h36c2ea0_1013	conda-forge
zstd	1.5.0	ha95c52a_0	conda-forge

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INDEX

A

ADMs, [65](#)

AF, [65](#)

anisotropy paramters, [65](#)

Axis distribution moments, [65](#)

L

LF, [65](#)

M

MF, [65](#)

P

PADs, [65](#)