

# Magnetic dichroism in darkfield UV photoemission electron microscopy

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Here comes the abstract ...

**Introduction.** Ultrafast spin and magnetization dynamics are exciting and rapidly growing fields in condensed matter physics with promising implications for both future research and device applications. Ultrafast imaging of magnetic domains on the micrometer scale is well established based on all-optical methods, as e.g. Kerr microscopy. On the nanometer scale, however, electron microscopy is the method of choice due to electrons short de Broglie wavelength. For imaging magnetic domains the magnetic circular dichroism (MCD) contrast is used in photoelectron emission microscopy (PEEM). The intensity recorded for a particular domain changes with the helicity of the incident radiation, thereby producing magnetic contrast without the need for an explicit electron spin detection. By tuning the incident X-ray radiation to a magnetic core level absorption edge, substantial and element-specific MCD asymmetries have been reported. With the wide availability of tunable synchrotron radiation, this technique of XMCD-PEEM is well established for magnetic domain imaging on the nanometer scale and for slow dynamics. However, the pulse length of synchrotron radiation of typically 30-50 ps renders XMCD-PEEM unsuitable on ultrafast timescales. Replacing the incident X-ray radiation by ultrashort laser pulses solves this issue straightforwardly and allows for pump-probe experiments on a few femtoseconds timescale. In addition, experiments can be performed in the laboratory with UV laser sources. With UV light, electrons close to the Fermi level are excited to energies slightly above the escape threshold. However, in threshold photoemission the reported MCD contrasts are quite small, especially for in-plane magnetization.

In this letter we introduce the concept of darkfield PEEM for MCD in threshold photoemission, which allows efficient MCD imaging with an order-of-magnitude enhanced MCD contrast for in-plane magnetization. It paves the way for MCD imaging on femtosecond timescales using modern lasers. Darkfield PEEM imaging uses an aperture for photoelectron momentum selection in the back focal plane of the electron imaging column. We apply this approach to magnetic domain imaging of Fe(001). While the general momentum dependence

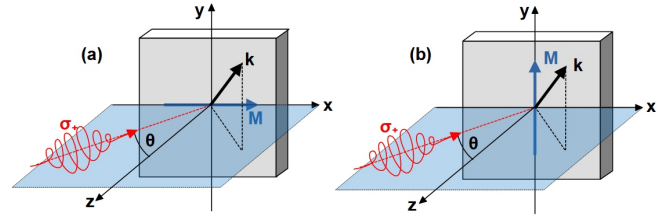


FIG. 1. Symmetry analysis. A circular polarized laser pulse (orange, with helicity  $\sigma_+$  impinges onto a magnetic domain (rectangular solid, with magnetization direction  $\mathbf{M}$ ). The pulse's incidence direction and the surface normal ( $z$ -axis) span the scattering plane (blue;  $xz$ -plane). The off-normal detection of photoelectrons with wavevector  $\mathbf{k}$  results in a chiral setup. If the scattering plane is a mirror plane of the lattice, the photoemission intensities for fixed  $\mathbf{k}$  obey  $I(\sigma_+, +\mathbf{M}) = I(\sigma_-, -\mathbf{M})$  (panel a) or  $I(\sigma_+, +\mathbf{M}) = I(\sigma_-, +\mathbf{M})$  (panel b).

of the MCD contrast will be rationalized by symmetry consideration, the experimental results will be quantitatively compared with state-of-the-art full-relativistic photoemission calculations.

Review experiments

Respective experiments on Ni/Cu(001) with threshold PEEM revealed that the magnetic contrast is very small

Give number

and therefore might not be sufficient for ultrafast imaging

Add reference

**Conceptual basis.** Key to the proposed approach are asymmetries of the photoemission intensities, which follow from symmetry consideration. Since a PEEM detects photoelectrons with (almost) all emission angles, one is concerned with a chiral geometry for photoelectrons with off-normal wavevector  $\mathbf{k}$  (Fig. 1). This chirality results in magnetic dichroism and, hence, in magnetic contrast.

For a conceptual approach let's assume light incidence within one of the high-symmetry [100] direction. The photoemission intensity of electrons detected with wavevector  $\mathbf{k}$  depends then on the helicity,  $\sigma_+$  or  $\sigma_-$ , of

the incident circular polarized laser radiation and on the two orientations  $\pm M$  of the in-plane magnetization in a selected domain, yielding four intensities  $I_{\mathbf{k}}(\sigma_{\pm}, \pm M)$  (shortened  $I_{\pm\pm}$ ). The latter are combined into the total intensity

$$I \equiv I_{++} + I_{+-} + I_{-+} + I_{--}. \quad (1)$$

In order to disentangle the two main contrast mechanisms we define appropriate asymmetries:

$$A_{\text{pol}} \equiv [(I_{++} + I_{+-}) - (I_{-+} + I_{--})] / I, \quad (2a)$$

$$A_{\text{ex}} \equiv [(I_{++} + I_{--}) - (I_{+-} + I_{-+})] / I. \quad (2b)$$

In the polarization asymmetry  $A_{\text{pol}}$  the magnetization's orientation is averaged out; it thus encodes contrast due to the light's helicity (as if the domain were nonmagnetic). Contrast due to the exchange splitting is quantified by the exchange asymmetry  $A_{\text{ex}}$ , in which one averages over the mutual orientations of helicity and magnetization.

*Experimental aspects.* Setup ...

The experimental results as well the conceptual ideas are supported by relativistic photoemission computations, briefly described in the Supplemental Material [1].

*Contrast mechanisms.* In a first step we analyze the two main contrast mechanism: light polarization and exchange splitting.

The polarization asymmetry  $A_{\text{pol}}$ , defined in Eq. (2a) and depicted in Fig. 2, depends on the binding energy of the initial states. Both theoretical (top row) and experimental data (bottom row) show that this contrast mechanism is sizable (with absolute values up to about 40 % in theory and 20 % in experiment) and thus cannot be ignored.

The theoretical patterns (top row in Fig. 2) exhibit two nodal lines ( $k_x = 0$  and  $k_y = 0$ ). Moreover, one finds changes of sign if either  $k_x$  or  $k_y$  is reversed. These features are imposed by the symmetry of the setup. The experimental counterparts (bottom row) display the same features but slightly oblique or off-center, most clearly for binding energies with comparably small asymmetry (cf. 0.15 eV and 0.25 eV). We attribute these deviations to imperfections in experiment, for example a small misalignment of the light incidence with respect to a crystal mirror plane and inhomogeneous magnetization in the surface layers. Moreover, we assume an electron self-energy that is independent of  $k_{\parallel}$  (see [1]). Nevertheless, the agreement of experiment and theory shows that ...

The energy-dependent exchange asymmetry  $A_{\text{ex}}$ , defined in Eq. (2b) and shown in Fig. 3, is smaller than  $A_{\text{pol}}$  (Fig. 2) (with absolute values up to 10 % in theory and 6 % in experiment). Nevertheless it exhibits a clear

odd symmetry upon reversal of  $k_x$ ; it does not change sign upon reversal of  $k_y$ , in contrast to  $A_{\text{pol}}$ . Both features are dictated by the symmetry of the setup (see [1]).

The above findings support that the asymmetries  $A_{\text{pol}}$  and  $A_{\text{ex}}$  are suitable tools to disentangle and to quantify the main contrast mechanisms for domain imaging.

*Domain imaging.* Aperture, selection of wave vectors.

Center position of the aperture shows very weak contrast (Fig. 4).

Sizable contrast in the other off-center positions of the aperture. Allows to image domain structure unequivocally. Size of contrast/asymmetry ...

*Prospects.* The present investigation proves that magnetic domains can be imaged with a laser-equipped photoelectron emission microscope with high contrast using a momentum-selection of the detected photoelectrons. The common belief that a such PEEM provides only weak contrast is thereby refuted.

For a proof of principle we applied the proposed improvement to in-plane magnetized Fe(001) surfaces. However, the approach is general so that it can easily be applied to other ferromagnets, even to those with out-of-plane magnetization. Thus, it well suited for studying magnetic reorientation transitions, as for example observed for Ni/Cu(001). We see its main capabilities in investigations of ultrafast magnetization dynamics using femtosecond laser pulses. As applications one might think of ultrafast motion of domain walls or of large skyrmions.

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[1] See Supplemental Material at [URL will be inserted by publisher] for supporting information and additional results.

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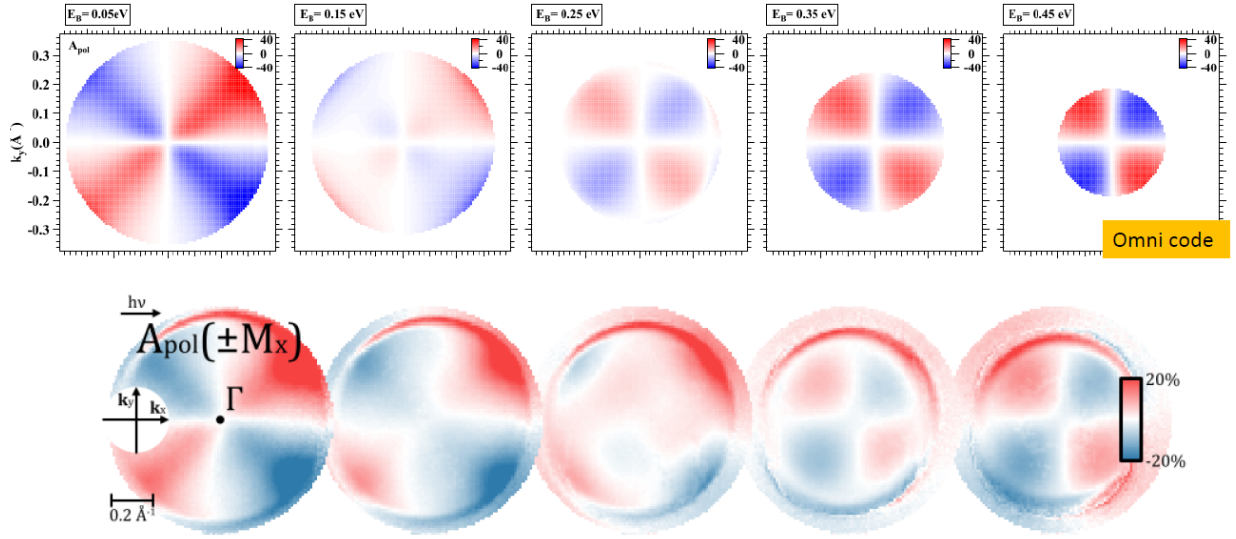


FIG. 2. Polarization asymmetry  $A_{\text{pol}}$  of Fe(001) at selected binding energies versus wave vector  $\mathbf{k}_{\parallel}$  of the photoelectrons. Top row: theoretical results obtained from photoemission calculations. The binding energy is indicated at each panel. The color scale, showing  $A_{\text{pol}}$  as defined in Eq. (2a) in percent, is identical for all panels in this row. Bottom row: respective experimental results. The arrow marked  $h\nu$  indicates the light incidence direction.  $\Gamma$  is the center of the surface Brillouin zone ( $\mathbf{k}_{\parallel} = 0$ ).

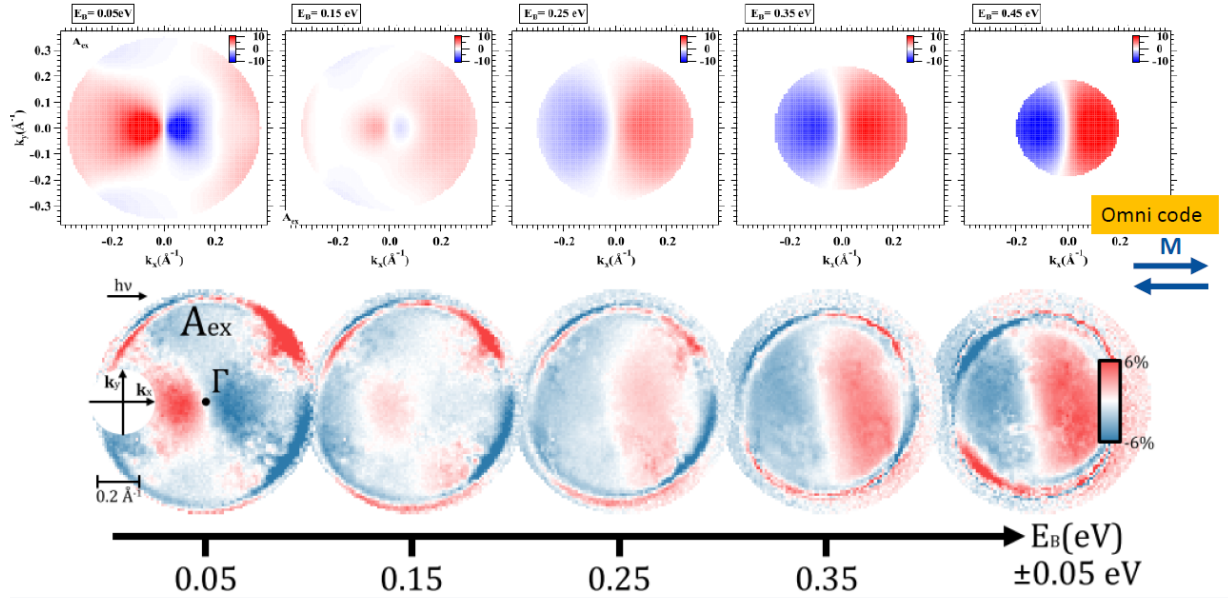


FIG. 3. As Figure 2, but for the exchange asymmetry  $A_{\text{ex}}$  of Fe(001). The orientations of the magnetization are indicated by arrows on the right-hand side.

FIG. 4. Darkfield magnetic-circular dichroism imaging of Fe(100). Left: schematic, nine aperture positions .... Right: domain imaging using the nine aperture positions.