

An Introductory Course in Computational Neuroscience—Paul Miller (Notes)

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0.1 xLIF

0.1.1 Modelling the Leaky membrane potential

Nernst Potential

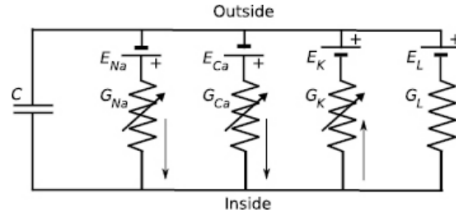
The *Nernst potential* E_A of an ion A of charge z_A with intracellular concentration $[A_{\text{in}}]$ and extracellular concentration $[A_{\text{out}}]$ is given by

$$E_A = \frac{k_B T}{z_A q_e} \ln \left(\frac{[A_{\text{out}}]}{[A_{\text{in}}]} \right)$$

where T is the temperature in Kelvin, k_B the Boltzmann constant ($1.39 \times 10^{-23} JK^{-1}$) (which converts units of temperature to units of thermal energy). q_e is the fundamental electronic charge ($1.60 \times 10^{-19} C$).

Model

Considering this representation of a neuron's membrane:



If all channels with variable conductance are closed, then the current will only flow through the leak channels (subscript L) until the cell membrane is at the leak potential E_L . The current through a channel is given by

$$I_t = G_t(V_m - E_t)$$

Where G_t represents conductance and E_t the nernst potential; t represents the type of channel.

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Equilibrium

When the cell is at equilibrium the different currents balance each other out and sum to zero:

$$I_m = \sum_t I_t = \sum_t G_t(V_m - E_t) = 0$$

In the context of this current model this can be rewritten as

$$G_{Na}(V_m - E_{Na}) + G_{Ca}(V_m - E_{Ca}) + G_K(V_m - E_K) + G_L(V_m - E_L)$$

Solving for V_m we can see that the *resting membrane potential*—where no net current flows, is the weighted average of the individual Nernst potentials:

$$V_m = \frac{G_{Na}E_{Na} + G_{Ca}E_{Ca} + G_KE_K + G_LE_L}{G_{Na} + G_{Ca} + G_K + G_L}$$

The derivation of the resting membrane potential is typically more complicated.

Leaky membrane potential

Here we consider the passive properties of the cell, where the variable conductance of all channels are fixed. With this we treat the circuit as having a single ‘leak’ conductance and potential.

The membrane potential is generated by the charge stored on the membrane; it depends on both the stored charge and the membrane’s capacitance C_m via the equation

$$Q = C_m V_m$$

The current is defined as positive when it flows *out* of the cell; with that we have

$$\frac{dQ}{dt} = -I_m = -G_L(V_m - E_L)$$

Fixing the capacitance we obtain the dynamics of the resting membrane potential as

$$C_m \frac{dV_m}{dt} = G_L(E_L - V_m)$$

This is a linear first order ODE.

0.1.2 Solution for Leaky ODE

We had the dynamics of the resting membrane potential as

$$C_m \frac{dV_m}{dt} = G_L(E_L - V_m)$$

Expressing in standard form and solving by integrating factor:

$$\begin{aligned} \frac{dV_m}{dt} + \frac{G_L}{C_m} V_m &= \frac{G_L}{C_m} E_L \\ V_m &= \frac{1}{\exp\left(\frac{G_L}{C_m} t\right)} \left(\int \exp\left(\frac{G_L}{C_m} t\right) \cdot \frac{G_L}{C_m} E_L dt + c \right) \end{aligned}$$

To simplify we define the *time constant* $\tau_m = C_m/G_L$:

$$\begin{aligned} V_m &= \frac{1}{\exp\left(\frac{t}{\tau_m}\right)} \left(\int \exp\left(\frac{1}{\tau_m} t\right) \cdot \frac{1}{\tau_m} E_L dt + c \right) \\ &= \exp\left(-\frac{t}{\tau_m}\right) \left(\exp\left(\frac{t}{\tau_m}\right) E_L + A \right) \\ &= E_L + \exp\left(-\frac{t}{\tau_m}\right) \cdot A \end{aligned}$$

At initial condition $V_m(0)$:

$$V_m(0) = E_L + A \implies A = V_m(0) - E_L$$

With that we have the solution

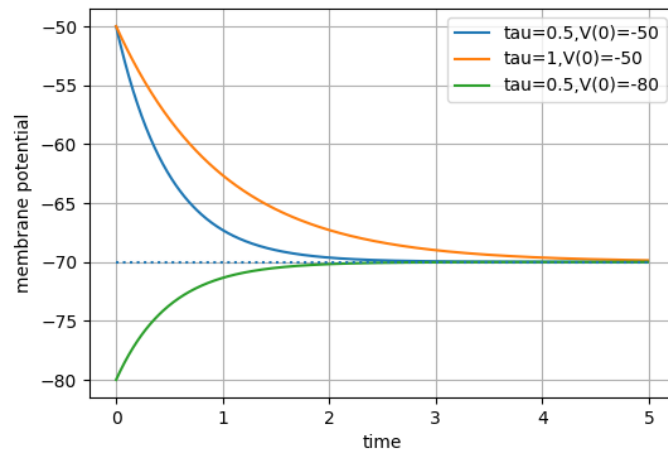
$$V_m = E_L + (V_m(0) - E_L) \exp\left(-\frac{t}{\tau_m}\right)$$

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Illustrated

See that the equation tends to E_L , and that τ_m dictates how fast this decay occurs (thus the name). Illustrated here using code from `leakymembrane.py`:

$$V_m = E_L + (V_m(0) - E_L) \exp\left(-\frac{t}{\tau_m}\right)$$



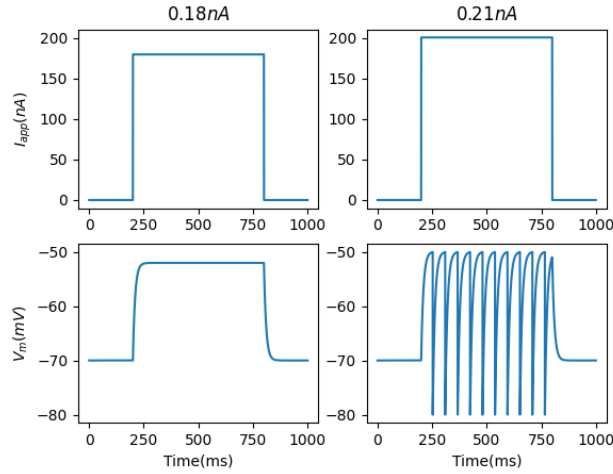
0.1.3 Leaky-Integrate-and-Fire

The LIF introduces a framework on which we can build more realistic models of neurons. It is essentially the initial model for the leaky membrane with an additional term I_{app} modelling an applied current. The spike is modelled by the membrane potential being reset to a low (hyperpolarised) value V_{reset} after the potential reaches some threshold V_{th} (threshold potential):

$$C_m \frac{dV_m}{dt} = G_L(E_L - V_m) + I_{\text{app}}; \text{ if } V_m > V_{th} \text{ then } V_m \mapsto V_{\text{reset}}$$

Notice that an actual ‘spike’ shape hasn’t been modelled, and would have to be put in by hand before V_m is reset; spike times are recorded at the time when the membrane potential crosses the threshold.

The following was simulated using code from `LIFfunc.py`:



Threshold current

See that there is an insufficient level for I_{app} where no ‘firing’ occurs (the membrane potential does not reach threshold and the model does not spike). Setting $dV_m/dt = 0$ allows us to obtain the *steady state* membrane potential as

$$V_m^{ss} = E_L + I_{\text{app}}/G_L$$

If the steady state is below threshold then the model does not fire. By setting $V_m^{ss} = V_{th}$ we can obtain the *threshold current* I_{th} —the minimum applied current required to elicit firing:

$$I_{th} = G_L(V_{th} - E_L)$$

0.1.4 Firing rate-Current (F-I) curve of the LIF model

For the LIF (and fixed applied current), we can come up with a closed-form way of determining the time for the neuron's membrane potential to increase from its reset value to the threshold.

Solution for LIF model with constant I_{app}

Ignoring the hard reset part of the LIF (this doesn't matter since we are trying to find time between a period where the model doesn't reset), we have

$$C_m \frac{dV_m}{dt} = G_L(E_L - V_m) + I_{\text{app}}$$

When the applied current is fixed this equation can be solved:

$$V_m(t) = V_m^{ss} + (V_m(0) - V_m^{ss}) \exp\left(-\frac{t}{\tau}\right)$$

where V_m^{ss} denotes the steady state membrane potential.

0.1.5 Solution for LIF with fixed applied current and no reset

Ignoring the hard reset part of the LIF and fixing I_{app} , we have

$$C_m \frac{dV_m}{dt} = G_L(E_L - V_m) + I_{\text{app}}$$

This can be solved. Where $\tau = C_m/G_L$:

$$\frac{dV_m}{dt} = \frac{G_L}{C_m}(E_L - V_m) + \frac{I_{\text{app}}}{C_m}$$

In standard form:

$$\frac{dV_m}{dt} + \frac{1}{\tau}V_m = \frac{1}{\tau}E_L + \frac{I_{\text{app}}}{C_m}$$

Solving by integrating factor:

$$\begin{aligned} V_m(t) &= \frac{1}{\exp(t/\tau)} \left(\int \exp\left(\frac{t}{\tau}\right) \left(\frac{1}{\tau}E_L + \frac{I_{\text{app}}}{C_m} \right) dt + c \right) \\ &= \exp\left(-\frac{t}{\tau}\right) \left(\tau \exp\left(\frac{t}{\tau}\right) \left(\frac{1}{\tau}E_L + \frac{I_{\text{app}}}{C_m} \right) + A \right) \\ &= \exp\left(-\frac{t}{\tau}\right) \left(\exp\left(\frac{t}{\tau}\right) E_L + \exp\left(\frac{t}{\tau}\right) \tau \frac{I_{\text{app}}}{C_m} + A \right) \\ &= E_L + \frac{C_m}{G_L} \frac{I_{\text{app}}}{C_m} + A \exp\left(-\frac{t}{\tau}\right) \\ &= E_L + \frac{I_{\text{app}}}{G_L} + A \exp\left(-\frac{t}{\tau}\right) \end{aligned}$$

Initial condition:

$$\begin{aligned} V_m(0) &= E_L + \frac{I_{\text{app}}}{G_L} + A \\ A &= V_m(0) - \left(E_L + \frac{I_{\text{app}}}{G_L} \right) \end{aligned}$$

We have the solution

$$V_m(t) = E_L + \frac{I_{\text{app}}}{G_L} + \left(V_m(0) - \left(E_L + \frac{I_{\text{app}}}{G_L} \right) \right) \exp\left(-\frac{t}{\tau}\right)$$

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In relation to steady state potential

We had

$$V_m(t) = E_L + \frac{I_{\text{app}}}{G_L} + \left(V_m(0) - \left(E_L + \frac{I_{\text{app}}}{G_L} \right) \right) \exp \left(-\frac{t}{\tau} \right)$$

Recall that setting $dV_m/dt = 0$ in the LIF equation (again ignoring the hard reset) gave us the steady state membrane potential V_m^{ss} :

$$V_m^{ss} = E_L + I_{\text{app}}/G_L$$

See that this can be slotted into our solution:

$$V_m(t) = V_m^{ss} + (V_m(0) - V_m^{ss}) \exp \left(-\frac{t}{\tau} \right)$$

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