# Math 254A Lecture 28 Notes

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# 1 Variational Principles for the Entropy Rate

### 1.1 Recap

Last time, we showed that

$$s(\mu) := \inf_{W,U \ni \mu} \lim_{B \uparrow \mathbb{Z}^d} \frac{1}{|B|} \log |\{x \in A^B : P_x^W \in U\}|$$
$$= \begin{cases} h(\mu) := \lim_{B} \frac{1}{|B|} H(\mu_B) & \text{if } \mu \in P^T \\ -\infty & \text{otherwise.} \end{cases}$$

Here, we extend h by  $h(\mu) - \infty$  if  $\mu \notin P^T$ . Then  $h: M(A^{\mathbb{Z}^d}) \to [-\infty, \log |A|]$  is concave and upper semicontinuous, and the set  $\{h > -\infty\} = \{h \ge 0\} = P^T$ . The upper bound  $\log |A|$  is achieved when  $\mu = \mathrm{Unif}_A^{\times \mathbb{Z}^d}$ .

Now, we will see two variational principles.

### 1.2 The first variational principle

**Theorem 1.1.** Let  $\psi: A^{\mathbb{Z}^d} \to \mathbb{R}^r$  depend only on coordinates in a finite  $W \subseteq \mathbb{Z}^d$ . For  $x \in \mathbb{R}^r$ , let

$$s(\psi,y) = \inf_{V\ni x} \lim_{B\uparrow\mathbb{Z}^d} \frac{1}{|B|} \log |\{x\in A^B: \frac{1}{|B|} \Psi_B(x)\in V\}|,$$

where the inf is over open, convex neighborhoods of x in  $\mathbb{R}^r$ . Then

$$s(\psi, y) = \sup\{h(\mu) : \mu \in P^T, \langle \psi, \mu \rangle = y\}$$
  
= 
$$\sup\{h(\mu) : \mu \in P, \langle \psi, \mu \rangle = y\}.$$

with the convention that  $\sup \emptyset = -\infty$ .

Proof.

$$\frac{1}{|B|} \Psi_B(x) = \frac{1}{|B|} \sum_{v+W \subseteq B} \psi(T^v x) 
= \frac{1}{|\{v : v + W \subseteq B\}|} \sum_{v+W \subseteq B} \psi(T^v x) + o(|B|) 
= \langle \psi, P_x^W \rangle + o(|B|),$$

This gives  $(\geq)$ : For any  $V \subseteq \mathbb{R}^r$  and finite  $W \subseteq \mathbb{Z}^d$ , we have

$$\frac{1}{|B|}\log|\{x\in A^B: \langle \psi, P^W_x\rangle \in V\}|.$$

The condition  $\langle \psi, P_x^W \rangle \in V$  defines any convex neighborhood of any  $\mu$  such that  $\langle \psi, \mu \rangle = y$ . So taking  $\lim_B$  of the above, we get that it is  $\geq h(\mu)$  for any such  $\mu$ .

Now consider  $(\leq)$ . Let

$$h = \sup\{h(\mu) : \mu \in P, \langle \psi, \mu \rangle = y\}.$$

The set  $\{\mu \in P : \langle \psi, \mu \rangle = y\}$  is compact, so there exists a window W and open convex sets  $U_1, \ldots, U_r$  in  $P(A^W)$  such that  $\{\mu \in P : \langle \psi, \mu \rangle = y\} \subseteq \bigcup_i \{\mu \in P : \mu_W \in U_i\}$  and

$$\frac{1}{|B|}\log|\{x: P_x^W \in U_i\}| \le (h+\varepsilon) + o(1)$$

for all i. Finally, by compactness again, if  $V \subseteq \mathbb{R}^r$  is a small enough neighborhood of y, then

$$\bigcup_{i} \{ \mu \in P : \mu_W \in U_i \} \supseteq \{ \mu : \langle \psi, \mu \rangle \in V \}$$

So

$$\frac{1}{|B|}\log|\{x:\langle\psi,P_x^W\rangle\in V\}| \le \max_i \frac{1}{|B|}\log|\{x:P_x^W\in U_i\}| + \frac{\log s}{|B|} \le h + \varepsilon$$

as  $B \uparrow \mathbb{Z}^d$ . Since  $\varepsilon > 0$  is arbitrary, we get  $s(\psi, y) = h$ , as desired.

Corollary 1.1. For any convex, open  $V \subseteq \mathbb{R}^r$ ,

$$\begin{split} s(\psi, v) &= \lim_{B} \frac{1}{|B|} \log |\{x : \langle \psi, P_x^W \rangle \in V\}| \\ &= \sup_{y \in V} s(\psi, y) \\ &= \sup\{h(\mu) : \mu \in P^T, \langle \psi, \mu \rangle \in V\} \\ &= \sup\{h(\mu) : \mu \in P, \langle \psi, \mu \rangle \in V\}. \end{split}$$

From this, we can return to interactions giving the total potential energy  $\varphi = (\varphi_F)_F$ , assumed (for simplicity) to be a finite range interaction. Look at

$$|\{x \in A^B : \frac{1}{|B|}\Phi(x) \in I\},$$

where I is a small open interval, and  $\Phi_B(x) = \sum_{F \subseteq B} \varphi_F(x_F) = \sum_{F'} |B| \langle \varphi_{F'}, P_x^W \rangle + o(|B|)$ . Here, W is a big enough window to see all nonzero translates, and F' runs over one copy of each finite set  $\subseteq W$  up to translation. So this set is

$$\left| \left\{ x \in A^B : \sum_{F'} \langle \varphi_{F'}, P_x^W \rangle \in I \right\} \right|.$$

 $\sum_{F'} \langle \varphi_{F'}, P_x^W \rangle \in I$  is an open, convex condition in  $\mathbb{R}^r$ , so

$$\frac{1}{|B|} \log \left| \left\{ x \in A^B : \sum_{F'} \langle \varphi_{F'}, P_x^W \rangle \in I \right\} \right| \xrightarrow{B \uparrow \mathbb{Z}^d} \sup \left\{ h(\mu) : \mu \in P^T \sum_{F'} \langle \varphi_{F'}, \mu \rangle \in I \approx y \right\}.$$

We can use this result to predict the most likely values of any other observable if there is a unique measure  $\mu$  that maximizes  $h(\mu)$  subject to the constraint  $\sum_{F'} \langle \varphi_{F'}, \mu \rangle = y$ .

**Corollary 1.2.** If there is a unique measure  $\mu$  that maximizes  $h(\mu)$  subject to the constraint  $\sum_{F'} \langle \varphi_{F'}, \mu \rangle = y$ .

**Remark 1.1.** There always exists a  $\mu$  achieving the supremum if the set  $\{\mu : \sum_{F'} \langle \varphi_{F'}, \mu \rangle = y\} \neq \emptyset$ . by upper semicontinuity of h on the above weak\* compact set.

So the key question is when we get uniqueness of that maximizer. We will discuss this next time.

### 1.3 A variational principle for the Fenchel-Legendre transform of h

To understand the second variational principle, we need to extend the first version from  $\varphi: A^{\mathbb{Z}^d} \to \mathbb{R}^r$  to any  $\psi \in C(A^{\mathbb{Z}^d})$ . To apply  $\psi$  "inside a box," given  $x \in A^B$ , let  $\widehat{x}$  be any element of  $A^{\mathbb{Z}^d}$  such that  $\widehat{x}_B = x$ . Given B and  $\psi \in C(A^{\mathbb{Z}^d})$ , let

$$s_B \psi(x) = \sum_{v \in B} \psi(T^v \widehat{x}).$$

**Lemma 1.1.** If  $\hat{x}$ ,  $\check{x}$  are two choices of extension, then

$$\left| \sum_{v \in B} \psi(T^v \widehat{x}) - \sum_{v \in B} \psi(T^v \check{x}) \right| = o(|B|).$$

Now a fiddly extension of the first variational principle gives

$$\frac{1}{|B|} \log |\{x \in A^B : \frac{1}{|B|} s_B \psi(x) \in V\}| = \sup \{h(\mu) : \mu \in P^T, \langle \psi, \mu \rangle \in V\}.$$

This version is good because we can now handle the whole Banach space  $C(A^{\mathbb{Z}^d})$ , which is the dual of  $M(A^{\mathbb{Z}^d})$ , equipped with the weak\* topology. This leads to a description of the Fenchel-Lengendre transform of h:

**Theorem 1.2** (2nd variational principle). On  $C(A^{\mathbb{Z}^d})$ ,

$$h^*(f) := \sup\{h(\mu) - \langle f, \mu \rangle : \mu \in M(A^{\mathbb{Z}^d})\}$$
$$= \lim_{B \uparrow \mathbb{Z}^d} \frac{1}{|B|} \log \sum_{x \in A^B} e^{-s_B f(x)}.$$

The  $e^{-s_B f(x)}$  are the Gibbs weights that define the canonical distribution on  $A^B$ . In ergodic theory and much of mathematical physics, this limit is called the *pressure* of f (denoted p(f)). Caution: this is not always the physical pressure of the system.