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Improving genome assemblies using multi-platform sequence data

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Keywords: computational geometry, graph theory, Hamilton cycles

1 Fixed-Period Problems: The Sublinear Case

With this chapter, the preliminaries are over, and we begin the search for periodic solutions to Hamiltonian systems. All this will be done in the convex case; that is, we shall study the boundary-value problem

$$\dot{x} = JH'(t, x)$$
$$x(0) = x(T)$$

with $H(t,\cdot)$ a convex function of x, going to $+\infty$ when $||x|| \to \infty$.

1.1 Autonomous Systems

In this section, we will consider the case when the Hamiltonian H(x) is autonomous. For the sake of simplicity, we shall also assume that it is C^1 .

We shall first consider the question of nontriviality, within the general framework of (A_∞,B_∞) -subquadratic Hamiltonians. In the second subsection, we shall look into the special case when H is $(0,b_\infty)$ -subquadratic, and we shall try to derive additional information.

The General Case: Nontriviality. We assume that H is (A_{∞}, B_{∞}) -subquadratic at infinity, for some constant symmetric matrices A_{∞} and B_{∞} , with $B_{\infty} - A_{\infty}$ positive definite. Set:

$$\gamma := \text{smallest eigenvalue of } B_{\infty} - A_{\infty}$$
(1)

$$\lambda := \text{largest negative eigenvalue of } J \frac{d}{dt} + A_{\infty} .$$
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Theorem 1 tells us that if $\lambda + \gamma < 0$, the boundary-value problem:

$$\dot{x} = JH'(x)
x(0) = x(T)$$
(3)

has at least one solution \overline{x} , which is found by minimizing the dual action functional:

$$\psi(u) = \int_{0}^{T} \left[\frac{1}{2} \left(\Lambda_{o}^{-1} u, u \right) + N^{*}(-u) \right] dt \tag{4}$$

on the range of Λ , which is a subspace $R(\Lambda)_L^2$ with finite codimension. Here

$$N(x) := H(x) - \frac{1}{2} \left(A_{\infty} x, x \right) \tag{5}$$

is a convex function, and

$$N(x) \le \frac{1}{2} \left(\left(B_{\infty} - A_{\infty} \right) x, x \right) + c \quad \forall x . \tag{6}$$

Proposition 1. Assume H'(0) = 0 and H(0) = 0. Set:

$$\delta := \liminf_{x \to 0} 2N(x) \|x\|^{-2} . \tag{7}$$

If $\gamma < -\lambda < \delta$, the solution \overline{u} is non-zero:

$$\overline{x}(t) \neq 0 \quad \forall t \ .$$
 (8)

Proof. Condition (7) means that, for every $\delta' > \delta$, there is some $\varepsilon > 0$ such that

$$||x|| \le \varepsilon \Rightarrow N(x) \le \frac{\delta'}{2} ||x||^2$$
 (9)

It is an exercise in convex analysis, into which we shall not go, to show that this implies that there is an $\eta>0$ such that

$$f \|x\| \le \eta \Rightarrow N^*(y) \le \frac{1}{2\delta'} \|y\|^2$$
 (10)

Since u_1 is a smooth function, we will have $||hu_1||_{\infty} \leq \eta$ for h small enough, and inequality (10) will hold, yielding thereby:

$$\psi(hu_1) \le \frac{h^2}{2} \frac{1}{\lambda} \|u_1\|_2^2 + \frac{h^2}{2} \frac{1}{\delta'} \|u_1\|^2 . \tag{11}$$

Fig. 1. This is the caption of the figure displaying a white eagle and a white horse on a snow field

If we choose δ' close enough to δ , the quantity $\left(\frac{1}{\lambda} + \frac{1}{\delta'}\right)$ will be negative, and we end up with

$$\psi(hu_1) < 0 \quad \text{for } h \neq 0 \text{ small}. \tag{12}$$

On the other hand, we check directly that $\psi(0) = 0$. This shows that 0 cannot be a minimizer of ψ , not even a local one. So $\overline{u} \neq 0$ and $\overline{u} \neq \Lambda_o^{-1}(0) = 0$.

Corollary 1. Assume H is C^2 and (a_{∞}, b_{∞}) -subquadratic at infinity. Let ξ_1, \dots, ξ_N be the equilibria, that is, the solutions of $H'(\xi) = 0$. Denote by ω_k the smallest eigenvalue of $H''(\xi_k)$, and set:

$$\omega := \operatorname{Min} \left\{ \omega_1, \dots, \omega_k \right\} . \tag{13}$$

If:

$$\frac{T}{2\pi}b_{\infty} < -E\left[-\frac{T}{2\pi}a_{\infty}\right] < \frac{T}{2\pi}\omega\tag{14}$$

then minimization of ψ yields a non-constant T-periodic solution \overline{x} .

We recall once more that by the integer part $E[\alpha]$ of $\alpha \in \mathbb{R}$, we mean the $a \in \mathbb{Z}$ such that $a < \alpha \le a+1$. For instance, if we take $a_{\infty} = 0$, Corollary 2 tells us that \overline{x} exists and is non-constant provided that:

$$\frac{T}{2\pi}b_{\infty} < 1 < \frac{T}{2\pi} \tag{15}$$

or

$$T \in \left(\frac{2\pi}{\omega}, \frac{2\pi}{b_{\infty}}\right) . \tag{16}$$

Proof. The spectrum of Λ is $\frac{2\pi}{T}Z\!\!\!/ = a_{\infty}$. The largest negative eigenvalue λ is given by $\frac{2\pi}{T}k_o + a_{\infty}$, where

$$\frac{2\pi}{T}k_o + a_\infty < 0 \le \frac{2\pi}{T}(k_o + 1) + a_\infty.$$
 (17)

Hence:

$$k_o = E\left[-\frac{T}{2\pi}a_{\infty}\right] . {18}$$

The condition $\gamma < -\lambda < \delta$ now becomes:

$$b_{\infty} - a_{\infty} < -\frac{2\pi}{T} k_o - a_{\infty} < \omega - a_{\infty} \tag{19}$$

which is precisely condition (14).

Lemma 1. Assume that H is C^2 on $\mathbb{R}^{2n} \setminus \{0\}$ and that H''(x) is non-degenerate for any $x \neq 0$. Then any local minimizer \widetilde{x} of ψ has minimal period T.

Proof. We know that \widetilde{x} , or $\widetilde{x} + \xi$ for some constant $\xi \in \mathbb{R}^{2n}$, is a T-periodic solution of the Hamiltonian system:

$$\dot{x} = JH'(x) \,. \tag{20}$$

There is no loss of generality in taking $\xi=0$. So $\psi(x)\geq \psi(\widetilde{x})$ for all \widetilde{x} in some neighbourhood of x in $W^{1,2}\left(\mathbb{R}/T\mathbb{Z};\mathbb{R}^{2n}\right)$.

But this index is precisely the index $i_T(\widetilde{x})$ of the T-periodic solution \widetilde{x} over the interval (0,T), as defined in Sect. 2.6. So

$$i_T(\widetilde{x}) = 0. (21)$$

Now if \tilde{x} has a lower period, T/k say, we would have, by Corollary 31:

$$i_T(\widetilde{x}) = i_{kT/k}(\widetilde{x}) \ge ki_{T/k}(\widetilde{x}) + k - 1 \ge k - 1 \ge 1.$$
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This would contradict (21), and thus cannot happen.

Notes and Comments. The results in this section are a refined version of [1]; the minimality result of Proposition 14 was the first of its kind.

To understand the nontriviality conditions, such as the one in formula (16), one may think of a one-parameter family $x_T, T \in (2\pi\omega^{-1}, 2\pi b_\infty^{-1})$ of periodic solutions, $x_T(0) = x_T(T)$, with x_T going away to infinity when $T \to 2\pi\omega^{-1}$, which is the period of the linearized system at 0.

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Theorem 1 (Ghoussoub-Preiss). Assume H(t,x) is $(0,\varepsilon)$ -subquadratic at infinity for all $\varepsilon>0$, and T-periodic in t

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 is convex $\forall t$ (23)

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$$H(t,x) \ge n(\|x\|)$$
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$$\forall \varepsilon > 0 \;, \quad \exists c \; : \; H(t,x) \le \frac{\varepsilon}{2} \|x\|^2 + c \;.$$
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Assume also that H is C^2 , and H''(t,x) is positive definite everywhere. Then there is a sequence x_k , $k \in \mathbb{N}$, of kT-periodic solutions of the system

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such that, for every $k \in \mathbb{N}$, there is some $p_o \in \mathbb{N}$ with:

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where the Hamiltonian H is $(0, b_{\infty})$ -subquadratic, and the forcing term is a distribution on the circle:

$$f = \frac{d}{dt}F + f_o \quad \text{with } F \in L^2\left(\mathbb{R}/T\mathbb{Z}; \mathbb{R}^{2n}\right) ,$$
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where $f_o := T^{-1} \int_0^T f(t) dt$. For instance,

$$f(t) = \sum_{k \in \mathbb{N}} \delta_k \xi , \qquad (31)$$

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Hamiltonian Mechanics2

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