Module 2a: Introduction to Finite Difference Methods: 1-Dimensional FD Expressions

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Announcements

- First homework posted online this week.
- ECHO360 Lecture capture.
- Permission numbers refresh today, come see me after class.
- Project description hopefully posted Thursday evening.

References and Acknowledgements

The following materials were used in the preparation of this lecture:

- Tannehill, Anderson and Pletcher, Computational fluid Mechanics and Heat Transfer.
- 2 16.920 Notes

The author of these slides wishes to thank these sources for making the current lecture.

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How to solve a simple PDE using a computer

• Consider the 1-D Poisson Equation (ODE) between x = 0 and x = 1:

$$\nabla^2 u = \frac{\partial^2 u}{\partial x^2} = f = 1 \tag{1}$$

$$u(x=0) = 0 (2)$$

$$u(x=1) = 0 (3)$$

- For this problem, we are going to assume that u represents the deflection of a string, and f represents some applied transverse force.
- Discuss with your neighbor(s):
 - What is the actual/real solution to this problem?
 - \bigcirc How can you going to represent the solution u, using a computer code?
 - 4 How can you represent the geometry/domain in the computer code?
 - 4 How can you represent the governing ODE in the computer?
 - Ooes the solution at a given location depend on the neighboring solutions?

The Exact Solution

Solution Representation

- We will see in this course that there are two ways to numerically represent a solution:
 - Pointwise
 - ② Functional

Pointwise

Functional

Geometry Representation

- We will initially use point-wise representation of the solution.
 - \bullet Setup the points where the solution is to be determined \to discretization or mesh.

Solution Representation

- The solution representation and solution method can have a direct impact on geometry representation in the computer.
- Even if the solution is needed at one location (max deflection), we usually need to solve the problem where dependency exists:
 - Elliptic
 - Parabolic
 - Hyperbolic

Solution Representation

 This sub-module: Elliptic equations → smooth solutions, infinite domain of influence and dependence.

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• Let's say we wish to approximate a derivative:

$$\frac{du}{dx}$$
 (4)

• How can we approximate this derivative?

- Fundamental definition of the derivative:
 - Take the value for u at two different x— locations, and simply take the
 difference between the u-value and divide by the spatial distance
 (difference in the x-locations.

$$\frac{du}{dx} \simeq \frac{u_{i+1} - u_i}{x_{i+1} - x_i} + error \tag{5}$$

- As the two points get closer together, the error diminishes.
- In the limit as the two points approach each other, we recover the derivative.
- How accurate is this approximation? As $(x_{i+1} x_i) \to 0$?

- The idea is to approximate derivatives using finite not infinitesimal differences in the variables.
- To make this a viable method, we need to:
 - GOAL # 1: Come up with a way to represent a diversity of derivatives, eg:
 - $\frac{\partial^2 u}{\partial x^2}$
 - $\frac{\partial^3 u}{\partial x^3}$
 - GOAL # 2: Quantify and reduce the error of the approximation → better solution.
 - **GOAL** # 3: Develop expressions and solutions to PDEs using these derivatives.

Mathematics: Taylor Series Expansion

 GOAL #1, Method 1: Taylor Series Expansion in positive x-direction:

$$u(x_0 + \Delta x) = u(x_0) + \frac{\partial u}{\partial x}|_0 \Delta x + \frac{\partial^2 u}{\partial x^2}|_0 \frac{(\Delta x)^2}{2!} + \dots + \frac{\partial^n u}{\partial x^n}|_0 \frac{(\Delta x)^n}{n!} + \dots$$

Geometry and Mathematics of the Taylor Series Expansion

Mathematics: Taylor Series Expansion

Taylor Series Expansions in negative x-direction:

$$u(x_0 - \Delta x) = u(x_0) - \frac{\partial u}{\partial x}|_0 \Delta x +$$

$$+ \frac{\partial^2 u}{\partial x^2}|_0 \frac{(\Delta x)^2}{2!} - \dots$$

$$+ (-1)^n \frac{\partial^n u}{\partial x^n}|_0 \frac{(\Delta x)^n}{n!} + \dots$$

Mathematics: Taylor Series Expansion

• If we start with:

$$u(x_0 + \Delta x) = u(x_0) + \frac{\partial u}{\partial x} \|_0 \Delta x + \frac{\partial^2 u}{\partial x^2} \|_0 \frac{(\Delta x)^2}{2!} + \dots + \frac{\partial^n u}{\partial x^n} \|_0 \frac{(\Delta x)^n}{n!} + \dots$$

We can re-arrange the equation so that:

$$\begin{aligned} \frac{\partial u}{\partial x} \|_{0} &= \frac{\left(u(x_{0}) - u(x_{0} + \Delta x)\right)}{\Delta x} + \\ &+ \frac{\frac{\partial^{2} u}{\partial x^{2}} \|_{0} \frac{(\Delta x)^{2}}{2!}}{\Delta x} + \dots \\ &+ \frac{\frac{\partial^{n} u}{\partial x^{n}} \|_{0} \frac{(\Delta x)^{n}}{n!}}{\Delta x} + \dots \end{aligned}$$

This gives us:

$$\frac{\partial u}{\partial x}\|_{0} = \frac{(u(x_{0}) - u(x_{0} + \Delta x))}{\Delta x} + O(\Delta x)$$
 (6)

• Comment: The error decays proportionally to decreases in Δx

• If we start with:

$$u(x_0 - \Delta x) = u(x_0) - \frac{\partial u}{\partial x} \|_0 \Delta x + \frac{\partial^2 u}{\partial x^2} \|_0 \frac{(\Delta x)^2}{2!} - \dots + (-1)^n \frac{\partial^n u}{\partial x^n} \|_0 \frac{(\Delta x)^n}{n!} + \dots$$

• We can re-arrange the equation so that:

$$\begin{aligned} \frac{\partial u}{\partial x} \|_{0} &= \frac{\left(u(x_{0}) - u(x_{0} - \Delta x)\right)}{\Delta x} + \\ &+ \frac{\frac{\partial^{2} u}{\partial x^{2}} \|_{0} \frac{(\Delta x)^{2}}{2!}}{\Delta x} - \dots \\ &+ (-1)^{n} \frac{\frac{\partial^{n} u}{\partial x^{n}} \|_{0} \frac{(\Delta x)^{n}}{n!}}{\Delta x} + \dots \end{aligned}$$

This gives us:

$$\frac{\partial u}{\partial x}\|_0 = \frac{(u(x_0) - u(x_0 - \Delta x))}{\Delta x} + O(\Delta x)$$

ullet Comment: The error again decays proportionally to decreases in Δx

• If we start with:

$$u(x_0 - \Delta x) = u(x_0) - \frac{\partial u}{\partial x} \|_0 \Delta x + \frac{\partial^2 u}{\partial x^2} \|_0 \frac{(\Delta x)^2}{2!} - \dots + (-1)^n \frac{\partial^n u}{\partial x^n} \|_0 \frac{(\Delta x)^n}{n!} + \dots$$

• And subtract the expression:

$$u(x_0 + \Delta x) = u(x_0) + \frac{\partial u}{\partial x} \|_0 \Delta x + \frac{\partial^2 u}{\partial x^2} \|_0 \frac{(\Delta x)^2}{2!} ... + \frac{\partial^n u}{\partial x^n} \|_0 \frac{(\Delta x)^n}{n!} + ...$$

$$\frac{\partial u}{\partial x}\|_{0} = \frac{u(x_{0} + \Delta x) - u(x_{0} - \Delta x)}{2\Delta x} + \frac{1}{2} \frac{\partial^{3} u}{\partial x^{3}} \|_{0} \frac{(\Delta x)^{2}}{3!} + \dots$$

- This is a more accurate approximation to the first derivative
- As we reduce the spacing between points by factor of 2, the error in the derivative changes by a factor of 4

- Let's define a system that uses i (and j) indices to indicate our position on a grid.
- The spatial difference between *i*-points is Δx .

 Let's look at a couple of common first derivative finite difference approximations:

$$\frac{du}{dx} = \frac{u_{i+1} - u_i}{\Delta x} + O(\Delta x) \tag{7}$$

$$\frac{du}{dx} = \frac{u_i - u_{i-1}}{\Delta x} + O(\Delta x) \tag{8}$$

$$\frac{du}{dx} = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O(\Delta x^2)$$
(9)

$$\frac{du}{dx} = \frac{3u_i - 4u_{i-1} + u_{i-2}}{2\Delta x} + O(\Delta x^2)$$
 (10)

$$\frac{du}{dx} = \frac{-3u_i + 4u_{i+1} - u_{i+2}}{2\Delta x} + O(\Delta x^2)$$
 (11)

$$\frac{du}{dx} = \frac{u_{i+1} - u_i}{\Delta x} + O(\Delta x)$$

$$\frac{du}{dx} = \frac{u_i - u_{i-1}}{\Delta x} + O(\Delta x)$$

$$\frac{du}{dx} = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O(\Delta x^2)$$

$$\frac{du}{dx} = \frac{3u_i - 4u_{i-1} + u_{i-2}}{2\Delta x} + O(\Delta x^2)$$

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Thought Experiment

• The deflection of a 1-D string.

Thought Experiment

 Now that we have a governing equation, let's determine how to solve it using Finite Differences:

$$\frac{\partial^2 u}{\partial x^2} = f(x) \tag{12}$$

With boundary conditions on both ends of:

$$u_L = u_R = 0 (13)$$

• How do we solve this?

Thought Experiment

- Step 1: Discretize the domain into n intervals (n + 1 points).
- Step 2: Write finite difference equations for each internal node of the problem
- Step 3: Form a system of linear equations
- Step 4: Enforce the boundary conditions for the end points
- Step 5: Solve the system of equations

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- Method 2: for finding finite difference formula is Lagrange Interpolation. This is a more rigorous approach.
- When deriving new finite difference formula, we want to find an approximation of the form:

$$\frac{d^m u}{dx^m} \simeq \sum_{j=-left}^{right} \delta_j^m u_j \tag{14}$$

ullet We need a way to find the values of the coefficients: δ^m_i

• Example of this notation: The first derivative

• Start by defining a Lagrange Polynomial:

$$L_{j}(x) = \frac{(x - x_{l})...(x - x_{j-1})(x - x_{j+1})...(x - x_{r})}{(x_{j} - x_{l})...(x_{j} - x_{j-1})(x_{j} - x_{j+1})...(x_{j} - x_{r})}$$
(15)

- The values for the above equation are 1 when $x = x_j$ and 0 when $x_i \neq x_i$.
- Here, x_i is a node other than x_j

 We can approximate the solution by adding the appropriate combinations of Lagrange Polynomials together:

$$\hat{u}(x) = \sum_{j=-left}^{right} L_j(x)u_j$$
 (16)

• Rather than the solution \hat{u} , we want to represent the derivatives, $\frac{d^m u}{dx^m}$ up to order m.

$$\frac{d^m u}{dx^m} \simeq \frac{d^m \hat{u}}{dx^m}|_{x=x_0} = \sum_{j=-left}^{right} \frac{d^m L_j}{dx^m}|_{x=x_0} u_j$$
(17)

• What this means, is that the coefficients δ_j^m that we wish to find are simply (by pattern matching):

$$\delta_j^m = \frac{d^m L_j}{dx^m}|_{x=x_0} \tag{18}$$

- Here we will look for the finite difference equations that are based on a 3-point polynomial.
- Start first by expressing the solution as:

$$\hat{u}(x) = \sum_{j=-l}^{r} L_{j}(x)u_{j}$$

$$= \frac{(x-x_{j})(x-x_{j+1})}{(x_{j-1}-x_{j})(x_{j-1}-x_{j+1})}u_{j-1}$$

$$+ \frac{(x-x_{j-1})(x-x_{j+1})}{(x_{j}-x_{j-1})(x_{j}-x_{j+1})}u_{j}$$

$$+ \frac{(x-x_{j-1})(x-x_{j})}{(x_{i+1}-x_{i-1})(x_{i+1}-x_{i})}u_{j+1}$$

- To find the finite difference equation coefficients (the δ_i^m):
 - 1 It helps to first expand the polynomial in each numerator
 - 2 Differentiate the various Lagrangian polynomials with respect to x
 - Oepending upon where the finite difference is centered, insert the appropriate x entry
 - ullet Simplify the expression to get the value for δ_j^m in terms of Δx

• Start with the term δ_{j-1}^m , where the finite difference is centered at x_{j-1} :

$$L_{j-1}(x) = \frac{(x-x_j)(x-x_{j+1})}{(x_{j-1}-x_j)(x_{j-1}-x_{j+1})}$$

$$L_{j-1}(x) = \frac{(x^2-xx_j-xx_{j+1}+x_jx_{j+1})}{(x_{j-1}-x_j)(x_{j-1}-x_{j+1})}$$

$$\frac{dL_{j-1}(x)}{dx} = \frac{(2x-x_j-x_{j+1})}{(x_{j-1}-x_j)(x_{j-1}-x_{j+1})}$$

$$\delta_{j-1}^m(x=x_{j-1}) = \frac{-3\Delta x}{2\Delta x^2} = \frac{-3}{2\Delta x}$$

• Similarly, for Largange polynomials centered at different x points, we can find the following expressions when $x = x_{j-1}$:

$$\delta_{j-1}^m = \frac{-3}{2\Delta x}$$

$$\delta_j^m = \frac{2}{\Delta x} = \frac{4}{2\Delta x}$$

$$\delta_{j+1}^m = \frac{-1}{2\Delta x}$$

 See the example sheet for derivations of the other terms and other finite difference expressions.

• The result is:

$$\frac{du}{dx} = \delta_{j-1}^{m} u_{j} + \delta_{j}^{m} u_{j+1} + \delta_{j+1}^{m} u_{j+2}$$

$$= \frac{-3}{2\Delta x} u_{j} + \frac{2}{\Delta x} u_{j+1} + \frac{-1}{2\Delta x} u_{j+2}$$

$$= \frac{-3u_{j} + 4u_{j+1} - u_{j+2}}{2\Delta x}$$

- We have derived the equation for the one sided first derivative.
- Note: We had 3-points, and as such we represented the solution using a polynomial of order 2.
- The error in the first derivative is therefore second order $(O(\Delta x^2))$.

- Applying the same approach, we can find many other finite difference formulas.
- One formula of particular interest is the second derivative, central difference formula:

$$\frac{d^2u}{dx^2} \simeq \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} + O(h^2)$$