

# Nuclear Reactor Theory Project #1

## Group #3

Lee, Seungsup

Miller, Dory

Payant, Andrew

Powers-Luhn, Justin

Zhang, Fan

Material	$\Sigma_{tr}(\text{cm}^{-1})$	$\Sigma_a(\text{cm}^{-1})$	$\nu\Sigma_f(\text{cm}^{-1})$	Relative Absorption
H	$1.79 \times 10^{-2}$	$8.08 \times 10^{-3}$	0	0.053
O	$7.16 \times 10^{-3}$	$4.90 \times 10^{-6}$	0	0
Zr	$2.91 \times 10^{-3}$	$7.01 \times 10^{-4}$	0	0.005
Fe	$9.46 \times 10^{-4}$	$3.99 \times 10^{-3}$	0	0.026
$^{235}\text{U}$	$3.08 \times 10^{-4}$	$9.24 \times 10^{-2}$	0.145	0.602
$^{238}\text{U}$	$6.95 \times 10^{-3}$	$1.39 \times 10^{-2}$	$1.20 \times 10^{-2}$	0.091
$^{10}\text{B}$	$8.77 \times 10^{-6}$	$3.41 \times 10^{-2}$	0	0.223
	$3.62 \times 10^{-2}$	0.1532	0.1570	1.000

TABLE I  
MACROSCOPIC CROSS SECTIONS

## Abstract

THIS IS THE ABSTRACT

## I. INTRODUCTION & BACKGROUND

Proving the capabilities and safety of a reactor design requires effective modeling of the neutron flux in the core (expressed in equation 1). For real cores, however, this is impossible, and must be first simplified, then discretized to provide the solution for a representative mesh.

For this project we have analyzed a simplified, monoenergetic, non-multiplying medium in one dimension. The flux originates from a single source at  $x = 0$  with a strength of  $S = 1 \times 10^8 \text{ s}^{-1}$ . These assumptions simplify the transport equation to that presented in equation 2.

In the following sections, we will first describe the terms in equation 2, then provide both an analytical and a discrete solution. We will also provide an analysis of the accuracy of the analysis as a function of the number of nodes. Finally, we will analyze the solution for different coordinate systems to equation 2.

$$\frac{\partial n}{\partial t} + v\hat{\Omega} \cdot \nabla n + v\Sigma_t n(\mathbf{r}, E', \hat{\Omega}, t) = \int_{4\pi} d\hat{\Omega}' \int_0^\infty dE' v' \Sigma_s(E' \rightarrow E, \hat{\Omega}' \rightarrow \hat{\Omega}) n(\mathbf{r}, E', \hat{\Omega}', t) + s(\mathbf{r}, E, \hat{\Omega}, t) \quad (1)$$

$$-D_m \frac{d^2 \phi}{dx^2} + \Sigma_a^m \phi = \begin{cases} S & (x = 0) \\ 0 & (x > 0) \end{cases} \quad (2)$$

## II. METHODOLOGY

Equation 2 is a simplified description of neutron diffusion through a finite medium, similar to a point source travelling through a shielding material to a detector. The flux, therefore, depends on the transport cross section. This is accounted for in the term  $D_m$ , which is related to the transport coefficient by  $D_m = 3\Sigma_{tr}^{-1}$ . Values for  $\Sigma_{tr}$  for typical reactor materials are found in table I.

### A. Analytic Solution

In the slab, equation 2 is equal to 0,  $-D_m \frac{\partial^2 \phi}{\partial x^2} + \Sigma_a^m \phi = 0$ . In order to better group constants, specify a diffusion length,  $L = \sqrt{D_m/\Sigma_a}$ . We can then solve for  $\phi(x)$ :

$$\begin{aligned} \frac{\partial^2 \phi}{\partial x^2} - \frac{\phi}{L} &= 0 \\ \phi(x) &= Ae^{-x/L} + Ce^{x/L} \end{aligned} \quad (3)$$

First use the boundary condition  $\phi(w) = 0$  to solve for  $C$

$$\begin{aligned} 0 &= Ae^{-w/L} + Ce^{w/L} \\ C &= -Ae^{2w/L} \end{aligned}$$

Next, use the fact that  $-D_m \phi'(0) = J(0) = S/2$  to solve for  $A$

$$\begin{aligned} J(0) &= \frac{S}{2} = -\frac{A}{L} (1 + e^{-2w/L}) \\ A &= \frac{SL}{2D_m} (1 + e^{-2w/L})^{-1} \end{aligned} \quad (4)$$

$$\begin{aligned} -D_m \frac{\phi_1 - \phi_0}{\Delta x} + D_m \left. \frac{d\phi}{dx} \right|_0 + \int_0^{\Delta x/2} \Sigma_a \phi dx &= 0 \\ -D_m \frac{\phi_1 - \phi_0}{\Delta x} + D_m \frac{S}{2} + \int_0^{\Delta x/2} \Sigma_a \phi dx &= 0 \\ -D_m \frac{\phi_1 - \phi_0}{\Delta x} + D_m \frac{S}{2} + \Sigma_a \phi_0 \int_0^{\Delta x/2} dx &= 0 \end{aligned}$$

### Flux vs. Distance for $n$ Nodes

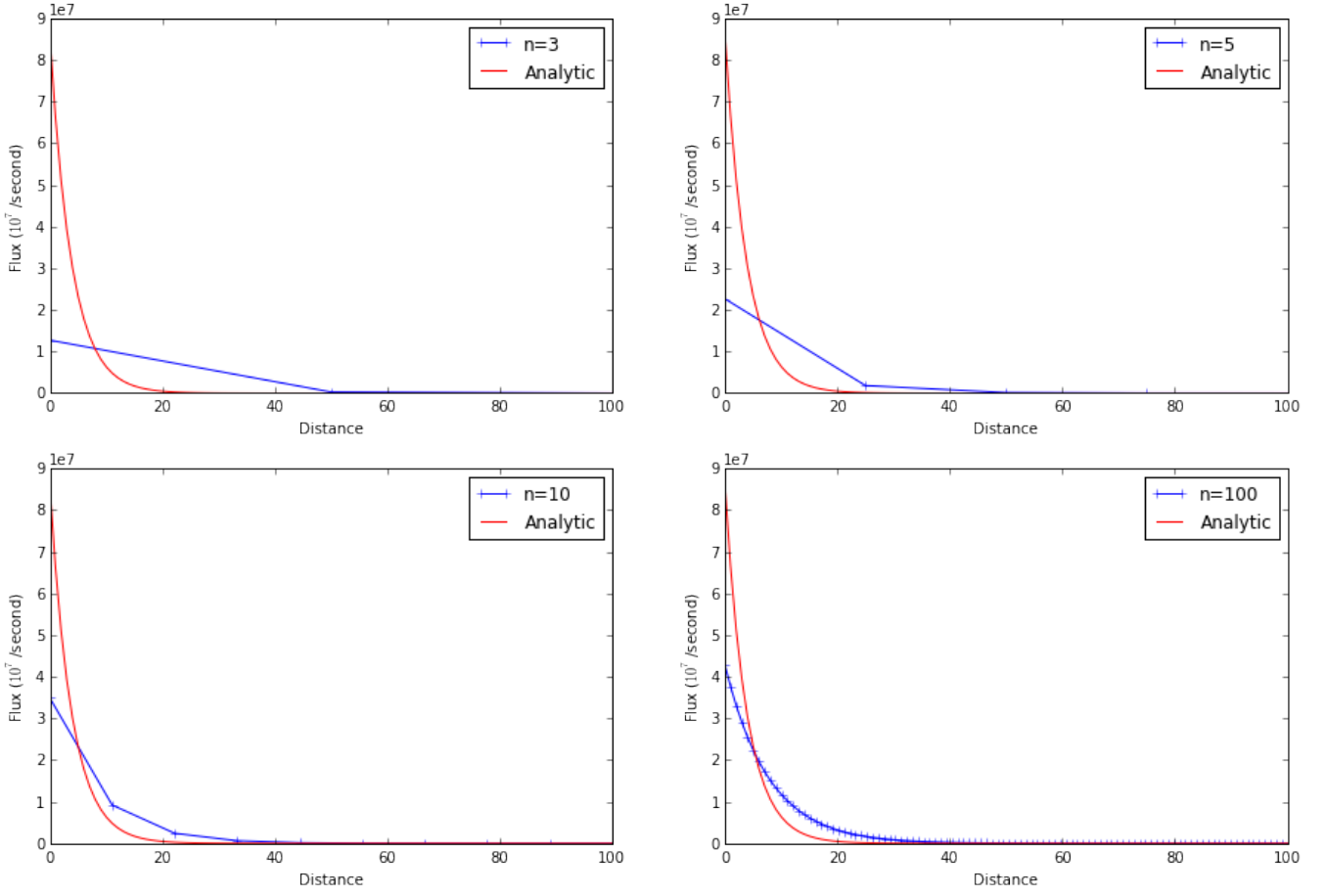


Fig. 1. Comparison of Numerical and Analytical Solutions

Divide by  $\Delta x$  on both sides:

$$\begin{aligned} -D_m \frac{\phi_1 - \phi_0}{\Delta x^2} - \frac{S}{2\Delta x} + \frac{1}{2}\Sigma_a \phi_0 &= 0 \\ \frac{-D_m}{\Delta x^2} \phi_1 + \left( \frac{D}{\Delta x^2} + \frac{1}{2}\Sigma_a \right) \phi_0 &= 0 \end{aligned}$$

This gives a final matrix  $\mathbf{A}$  (for  $n = 5$  nodes):

$$\begin{bmatrix} \frac{D_m}{\Delta x^2} + \frac{1}{2}\Sigma_a & \frac{-D_m}{\Delta x^2} & 0 & 0 \\ \frac{-D_m}{\Delta x^2} & \frac{2D_m}{\Delta x^2} + \Sigma_a & \frac{-D_m}{\Delta x^2} & 0 \\ 0 & \frac{-D_m}{\Delta x^2} & \frac{2D_m}{\Delta x^2} + \Sigma_a & \frac{-D_m}{\Delta x^2} \\ 0 & 0 & \frac{-D_m}{\Delta x^2} & \frac{2D_m}{\Delta x^2} + \Sigma_a \end{bmatrix}$$

### III. RESULTS

A Fortran program was developed to implement the discrete solution to the transport equation from section II-B. The source code is included as an attachment to this report. The numerical solution produced values that converged with the analytical solution as the number of nodes increased, showing clearly exponential behavior as  $n \rightarrow \infty$ . This behavior can be seen in figure 1.

### IV. CONCLUSIONS

The conclusions go here