

Prob. 4.2.

Proof

Suppose $|\tilde{\psi}_i\rangle = \sum_j u_{ij} |\tilde{\varphi}_j\rangle$ for some unitary u_{ij} . Then

$$\sum_i |\tilde{\psi}_i\rangle\langle\tilde{\psi}_i| = \sum_{ijk} u_{ij} u_{ik}^* |\tilde{\varphi}_j\rangle\langle\tilde{\varphi}_k| \quad (2.168)$$

$$= \sum_{jk} \left(\sum_i u_{ki}^\dagger u_{ij} \right) |\tilde{\varphi}_j\rangle\langle\tilde{\varphi}_k| \quad (2.169)$$

$$= \sum_{jk} \delta_{kj} |\tilde{\varphi}_j\rangle\langle\tilde{\varphi}_k| \quad (2.170)$$

$$= \sum_j |\tilde{\varphi}_j\rangle\langle\tilde{\varphi}_j|, \quad (2.171)$$

which shows that $|\tilde{\psi}_i\rangle$ and $|\tilde{\varphi}_j\rangle$ generate the same operator.

Conversely, suppose

$$A = \sum_i |\tilde{\psi}_i\rangle\langle\tilde{\psi}_i| = \sum_j |\tilde{\varphi}_j\rangle\langle\tilde{\varphi}_j|. \quad (2.172)$$

Let $A = \sum_k \lambda_k |k\rangle\langle k|$ be a decomposition for A such that the states $|k\rangle$ are orthonormal, and the λ_k are strictly positive. Our strategy is to relate the states $|\tilde{\psi}_i\rangle$ to the states $|\tilde{k}\rangle \equiv \sqrt{\lambda_k} |k\rangle$, and similarly relate the states $|\tilde{\varphi}_j\rangle$ to the states $|\tilde{k}\rangle$. Combining the two relations will give the result. Let $|\psi\rangle$ be any vector orthonormal to the space spanned by the $|\tilde{k}\rangle$, so $\langle\psi|\tilde{k}\rangle\langle\tilde{k}|\psi\rangle = 0$ for all k , and thus we see that

$$0 = \langle\psi|A|\psi\rangle = \sum_i \langle\psi|\tilde{\psi}_i\rangle\langle\tilde{\psi}_i|\psi\rangle = \sum_i |\langle\psi|\tilde{\psi}_i\rangle|^2. \quad (2.173)$$

Thus $\langle\psi|\tilde{\psi}_i\rangle = 0$ for all i and all $|\psi\rangle$ orthonormal to the space spanned by the $|\tilde{k}\rangle$. It follows that each $|\tilde{\psi}_i\rangle$ can be expressed as a linear combination of the $|\tilde{k}\rangle$, $|\tilde{\psi}_i\rangle = \sum_k c_{ik} |\tilde{k}\rangle$. Since $A = \sum_k |\tilde{k}\rangle\langle\tilde{k}| = \sum_i |\tilde{\psi}_i\rangle\langle\tilde{\psi}_i|$ we see that

$$\sum_k |\tilde{k}\rangle\langle\tilde{k}| = \sum_{kl} \left(\sum_i c_{ik} c_{il}^* \right) |\tilde{k}\rangle\langle\tilde{l}|. \quad (2.174)$$

The operators $|\tilde{k}\rangle\langle\tilde{l}|$ are easily seen to be linearly independent, and thus it must be that

$\sum_i c_{ik} c_{il}^* = \delta_{kl}$. This ensures that we may append extra columns to c to obtain a unitary matrix v such that $|\tilde{\psi}_i\rangle = \sum_k v_{ik} |\tilde{k}\rangle$, where we have appended zero vectors to the list of $|\tilde{k}\rangle$. Similarly, we can find a unitary matrix w such that $|\tilde{\varphi}_j\rangle = \sum_k w_{jk} |\tilde{k}\rangle$. Thus $|\tilde{\psi}_i\rangle = \sum_j u_{ij} |\tilde{\varphi}_j\rangle$, where $u = vw^\dagger$ is unitary. \square

Prob. 4-3.

$$(A) \rho = \frac{3}{4}|0\rangle\langle 0| + \frac{1}{4}|1\rangle\langle 1|$$

$$(B) \rho = \frac{1}{3}\left(\frac{|00\rangle\langle 11|}{\sqrt{2}}\right)\left(\frac{|00\rangle\langle 01|}{\sqrt{2}}\right) + \frac{1}{6}\left(\frac{|00\rangle\langle 11|}{\sqrt{2}}\right)\left(\frac{|00\rangle\langle 11|}{\sqrt{2}}\right) + \frac{1}{2}|0\rangle\langle 0|$$

$$= \frac{3}{4}|00\rangle\langle 00| + \frac{1}{12}|00\rangle\langle 11| + \frac{1}{12}|11\rangle\langle 00| + \frac{1}{4}|11\rangle\langle 11|$$

Prob. 4-4.

\Rightarrow if ρ is pure, $\exists |q_i\rangle$ s.t. $\rho = |q_i\rangle\langle q_i| \Rightarrow \text{tr}(\rho^2) = \text{tr}(|q_i\rangle\langle q_i|q_i\rangle\langle q_i|) = \text{tr}(\rho) = 1$.

\Leftarrow if $\text{tr}(\rho^2) = 1$, $\exists \{|\psi_i\rangle, p_i\}$ s.t. $\rho = \sum_i p_i |\psi_i\rangle\langle \psi_i|$ with $\{|\psi_i\rangle\}$ orthonormal (\because Hermitian)

$$\Rightarrow \rho^2 = \sum_{ij} p_i p_j |\psi_i\rangle\langle \psi_i| |\psi_j\rangle\langle \psi_j| = \sum_i p_i^2 |\psi_i\rangle\langle \psi_i|$$

$$\Rightarrow \text{tr}(\rho^2) = \sum_i p_i^2 \leq \sum_i p_i = 1$$

The equality condition is when $p_i^2 = p_i$ for all i , which means $p_i = 0$ or $1 \Rightarrow$ only allows pure states.

Prob. 4-5.

$$|\Psi^\pm\rangle = \frac{|00\rangle \pm |11\rangle}{\sqrt{2}} \Rightarrow \rho_{\Psi^\pm} = \frac{1}{2}|00\rangle\langle 00| \pm \frac{1}{2}|00\rangle\langle 11| \pm \frac{1}{2}|11\rangle\langle 00| + \frac{1}{2}|11\rangle\langle 11|$$

$$\Rightarrow \text{tr}_B(\rho_{\Psi^\pm}) = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$$

$$|\Psi^\pm\rangle = \frac{|00\rangle \pm |11\rangle}{\sqrt{2}} \Rightarrow \rho_{\Psi^\pm} = \frac{1}{2}|00\rangle\langle 00| \pm \frac{1}{2}|00\rangle\langle 11| \quad \downarrow \text{when } \langle 1| \rho | 1 \rangle$$

$$\Rightarrow \text{tr}_B(\rho_{\Psi^\pm}) = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \quad \leftarrow \text{when } \langle 0| \rho | 0 \rangle$$

Prob. 4-6.

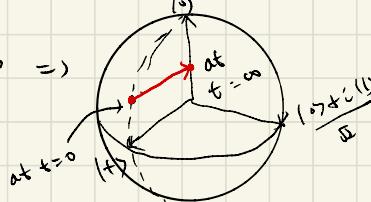
$$(A) \rho = \frac{1}{2}I + (a-\frac{1}{2})\sigma_z + Re(b)\sigma_x + Im(b)\sigma_y$$

$$(B) \Rightarrow \text{Suppose } |\rho|^2 = 1 \Leftrightarrow (a-\frac{1}{2})^2 + (Re(b))^2 + (Im(b))^2 = 1/4 \\ \Rightarrow \text{tr}(\rho^2) = \text{tr}\left(\frac{1}{4}I + \frac{1}{2}(a-\frac{1}{2})\sigma_z + \frac{1}{2}Re(b)\sigma_x + \frac{1}{2}Im(b)\sigma_y + \left[(a-\frac{1}{2})^2 + (Re(b))^2 + (Im(b))^2\right]I\right) \\ = \text{tr}\left(\frac{2}{3}I\right) = 1 \quad (\because \text{tr}(\sigma_i) = 0)$$

$$(\Leftarrow) \text{ Suppose } \text{tr}(\rho) = 1$$

$$\Rightarrow \text{tr}\left[\frac{1}{2}I + (a-\frac{1}{2})\sigma_z + (Re(b))^2 + (Im(b))^2\right]I = \frac{1}{2} + 2\left[(a-\frac{1}{2})^2 + (Re(b))^2 + (Im(b))^2\right] = 1 \Leftrightarrow |\rho|^2 = 1 \quad \square$$

$$(C) \quad n_z = \frac{1}{3}, \quad n_x = \frac{2\sqrt{2}}{3}e^{i\omega t}, \quad n_y = 0 \quad \Rightarrow$$



Prob. 4-9.

by spectral decomposition, $\rho = \sum_j p_j |\psi_j\rangle\langle\psi_j|$ where $\{|\psi_j\rangle\}$ is orthonormal.

$$\Rightarrow \rho \log \rho = \left(\sum_j p_j |\psi_j\rangle\langle\psi_j| \right) \sum_i \log(p_i) |\psi_i\rangle\langle\psi_i| = \sum_i p_i \log(p_i) |\psi_i\rangle\langle\psi_i|$$

$$\Rightarrow \text{tr}(\rho \log \rho) = \text{tr}\left(\sum_i p_i \log(p_i) |\psi_i\rangle\langle\psi_i|\right).$$

$$\Rightarrow \text{Under unitary } U, \quad \tilde{\rho} = \sum_j p_j U |\psi_j\rangle\langle\psi_j| U^\dagger = \sum_j p_j |\tilde{\psi}_j\rangle\langle\tilde{\psi}_j|$$

$$\Rightarrow \text{tr}(\tilde{\rho} \log \tilde{\rho}) = \text{tr}\left(\sum_i p_i \log(p_i) |\tilde{\psi}_i\rangle\langle\tilde{\psi}_i|\right)$$

$$\Rightarrow \text{tr}\left(\sum_i p_i \log(p_i) U |\psi_i\rangle\langle\psi_i| U^\dagger\right) = \text{tr}(U^\dagger U \rho \log \rho) = \underline{\text{tr}(\rho \log \rho)}$$

(check) $at t=0, 0$

$at t=\infty, 1$.

$$\rho(t) = \frac{1+e^{-i\omega t}}{2} |+\rangle\langle+| + \frac{1-e^{-i\omega t}}{2} |- \rangle\langle-|$$

$$\Rightarrow \text{entropy} = -\text{tr}(\rho \log \rho) = -\left[\frac{1+e^{-i\omega t}}{2} \log \frac{1+e^{-i\omega t}}{2} + \frac{1-e^{-i\omega t}}{2} \log \frac{1-e^{-i\omega t}}{2} \right]$$

Prob. 4-8.

$$\begin{aligned}
 & \text{define } |\tilde{\psi}\rangle = U|\psi\rangle \Rightarrow |\psi\rangle = U^+|\tilde{\psi}\rangle \Rightarrow i\hbar \frac{d(U^+|\tilde{\psi}\rangle)}{dt} = H(U^+|\tilde{\psi}\rangle) \\
 & \Rightarrow i\hbar \frac{dU^+}{dt}|\tilde{\psi}\rangle + i\hbar U^+ \frac{d(\tilde{\psi})}{dt} = HU^+|\tilde{\psi}\rangle \\
 & \Rightarrow i\hbar UU^+ \frac{d(\tilde{\psi})}{dt} = i\hbar \frac{d}{dt}|\tilde{\psi}\rangle = UHU^+|\tilde{\psi}\rangle - i\hbar U \frac{dU^+}{dt}|\tilde{\psi}\rangle = \underbrace{(UHU^+ - i\hbar U \frac{dU^+}{dt})}_{= \tilde{H}}|\tilde{\psi}\rangle \\
 & \Rightarrow i\hbar \frac{d(\tilde{\psi})}{dt} = \tilde{H}|\tilde{\psi}\rangle, \quad \tilde{H} = UHU^+ - i\hbar U \frac{dU^+}{dt} \\
 & \quad \& (\tilde{H})^+ = \tilde{H}; \text{ Hamilton.}
 \end{aligned}$$

$$\begin{aligned}
 \left(\frac{d}{dt}(UU^+) \right) &= \left(\frac{dU}{dt}U^+ + U \frac{dU^+}{dt} \right) \\
 \Rightarrow \left(i\hbar U \frac{dU^+}{dt} \right)^+ &= \left[i\hbar \left(-\frac{dU}{dt}U^+ \right) \right]^+ \\
 &= i\hbar U \frac{dU^+}{dt}
 \end{aligned}$$

Prob. 4-9.

$$\begin{aligned}
 (A) \quad G_x \cdot G_y &= iG_y G_x \in P_1, \text{ where } G_{ijk} = \begin{cases} 1 & \rightarrow \text{cyclic} \\ -1 & \rightarrow \text{anti-cyclic} \end{cases} \\
 & \& I \cdot G_z = G_z \in P_1, \quad G_z^2 = I \in P_1 \\
 & \& \pm G_z \in P_1, \quad \pm iG_i \in P_1 \quad (i=0, x, y) \Rightarrow \text{closed under} \\
 & & \text{multiplication.}
 \end{aligned}$$

$$\begin{aligned}
 \text{also, } \quad G_x \cdot G_y &= iG_z \quad G_y \cdot G_x = -iG_z \quad G_z^2 = I \\
 G_y \cdot G_z &= iG_x \quad G_z \cdot G_y = -iG_x \quad (iG_z)^2 = -I \\
 G_z \cdot G_x &= iG_y \quad G_x \cdot G_z = -iG_y \quad (G_x \cdot G_x) = iI \\
 & \quad -I, iI = -iI
 \end{aligned}$$

\Rightarrow All elements in P_1 can be generated by $\langle G_x, G_y, G_z \rangle$

$$\begin{aligned}
 (B) \quad H G_x H^+ &= \frac{1}{\sqrt{2}} \left(\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right) = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = G_x \in P_1 \\
 H G_y H^+ &= \frac{1}{\sqrt{2}} \left(\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right) = \frac{1}{2} \begin{pmatrix} i & -i \\ -i & i \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix} = -G_y \in P_1 \\
 H G_z H^+ &= H \cdot H \cdot G_x \cdot H \cdot H = G_x \in P_1 \\
 S G_x S^+ &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = G_y \in P_1 \\
 S G_y S^+ &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = -G_x \in P_1
 \end{aligned}$$

$$SG_z S^+ = G_z \left(\because [G_z, S] = 0 \right) \in P_1$$

Remaining cases are just phase multiplication

$$(C) CNOT = \begin{pmatrix} 1 & & \\ & 1 & \\ & & 0 & 1 \\ & & & 0 \end{pmatrix} = \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix}$$

$$\Rightarrow CNOT G_x \otimes G_i CNOT = \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix} \begin{pmatrix} 0 & G_i \\ G_i & 0 \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix}$$

$$= \begin{pmatrix} 0 & G_i G_x \\ G_x G_i & 0 \end{pmatrix} = \begin{cases} \begin{pmatrix} 0 & -G_x G_i \\ G_x G_i & 0 \end{pmatrix} = -i G_y \otimes G_x G_i \in P_2 \\ \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} (i = x) = G_x \otimes I \in P_2 \end{cases}$$

$$CNOT G_y \otimes G_i CNOT = \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix} \begin{pmatrix} 0 & -i G_i \\ i G_i & 0 \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix} = \begin{pmatrix} 0 & -i G_i G_x \\ i G_x G_i & 0 \end{pmatrix}$$

$$= \begin{cases} i G_x \otimes G_x G_i \quad (i \neq x) \in P_2 \\ G_y \otimes I \quad (i = x) \end{cases}$$

$$CNOT G_z \otimes G_i CNOT = \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix} \begin{pmatrix} G_i & \\ & -G_i \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & G_x \end{pmatrix} = \begin{pmatrix} G_i & 0 \\ 0 & -G_x G_i G_x \end{pmatrix}$$

$$= \begin{cases} \begin{pmatrix} G_x & \\ & -G_x \end{pmatrix} = G_x \otimes G_x \quad (i = x) \in P_2 \\ \begin{pmatrix} G_i & \\ & G_i \end{pmatrix} = I \otimes G_i \quad (i \neq x) \end{cases}$$

QED