

# Searches for invisibly decaying Higgs bosons with the CMS detector

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# Abstract

## Declaration

This dissertation is the result of my own work, except where explicit reference is made to the work of others, and has not been submitted for another qualification to this or any other university. This dissertation does not exceed the word limit for the respective Degree Committee.

Patrick Dunne

## Acknowledgements

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# Chapter 1

## Introduction and theory

### 1.1 The standard model of particle physics

The standard model of particle physics (SM) is a quantum field theory describing the interaction of the particles currently thought to be fundamental with the strong, weak and electromagnetic forces. Its predictions which come solely from specifying the particles in the theory, the symmetries the theory respects, how they are broken and 18 free parameters have been tested in many different experiments in some cases up to one part in a trillion [1]. The Higgs boson was the final SM particle to be discovered, and in this section I will describe the SM focussing on the Higgs mechanism, before outlining the motivations behind and some candidates for physics beyond the SM (BSM).

#### 1.1.1 Symmetries and particles

The Lagrangian of the SM is globally invariant under Poincaré group transformations and locally gauge invariant under the group  $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$ . For any theory Poincaré invariance guarantees that the theory respects the observed energy and momentum conservation in the universe, whilst local symmetries give interactions between the particles in the theory mediated by vector gauge bosons.

The observed particles in the universe fall into two types, fermions and bosons. Fermions in the SM are spin half representations of the above symmetry groups. The vector bosons in the standard model emerge as a result of the gauge invariance described above



## 1.2 Dark matter

## 1.3 Some extensions of the standard model incorporating dark matter

# Chapter 2

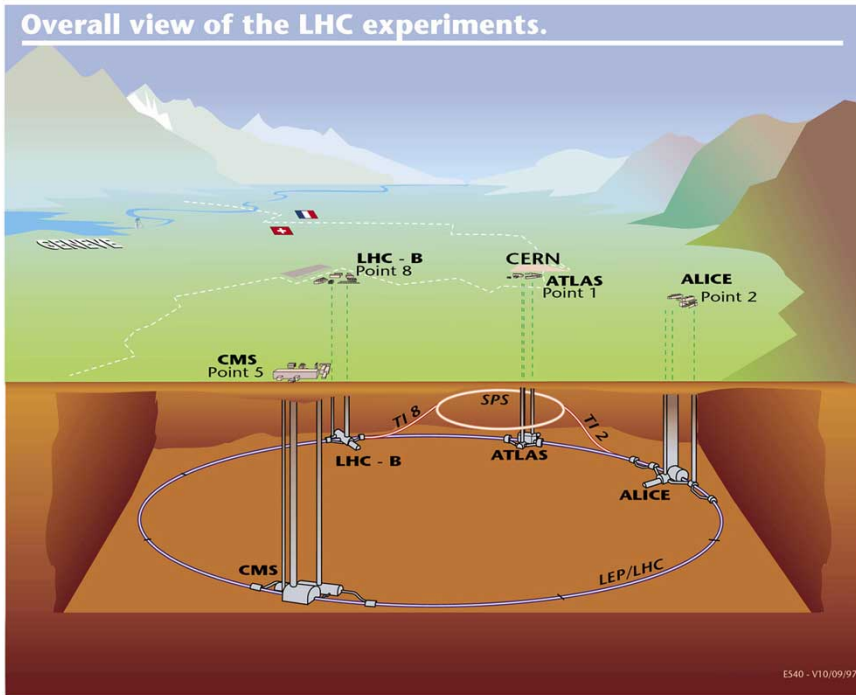
## The LHC and the CMS experiment

The purpose of this chapter is to introduce the CMS experiment and the LHC [2]. Without both of these apparatus the analyses performed for this thesis would, of course, not have been possible. In Section 2.1 an overview of the LHC and the chain of accelerators which feed into it will be given. This will then be followed in section Section ?? by a description of the CMS experiment focussing on the aspects most relevant to the search for invisibly decaying Higgs bosons.

### 2.1 The LHC

The LHC is situated 100m underground in a tunnel formerly built for the LEP accelerator [3] at CERN near Geneva, Switzerland. It is a 27km storage ring which accelerates both protons and heavy ions and collides them at the highest centre of mass energies of any collider built to date. The work contained in this thesis uses data from proton-proton collisions. These protons are obtained by taking hydrogen gas and stripping its atoms of their electrons with an electric field. The first accelerator in the LHC accelerator sequence, Linac 2, then accelerates the protons to 50 MeV. The protons are then accelerated to 1.4 GeV by the next accelerator, the Proton Synchrotron Booster (PSB), which is followed by the Proton Synchrotron (PS) where they reach 25 GeV. The beam energy is then increased to 450 GeV in the Super Proton Synchrotron (SPS). The protons are then injected into the LHC where, at time of writing, the maximum energy the beams have been accelerated to is 6.5 TeV, close to the design maximum of 7 TeV.

When fully filled the LHC contains two counter-rotating beams which are formed of up to 2808 bunches spaced either 25 or 50 ns apart and each containing  $\mathcal{O}(10^{11})$  protons. The



**Figure 2.1:** The layout of the LHC accelerator chain, showing the position of the four main detectors.

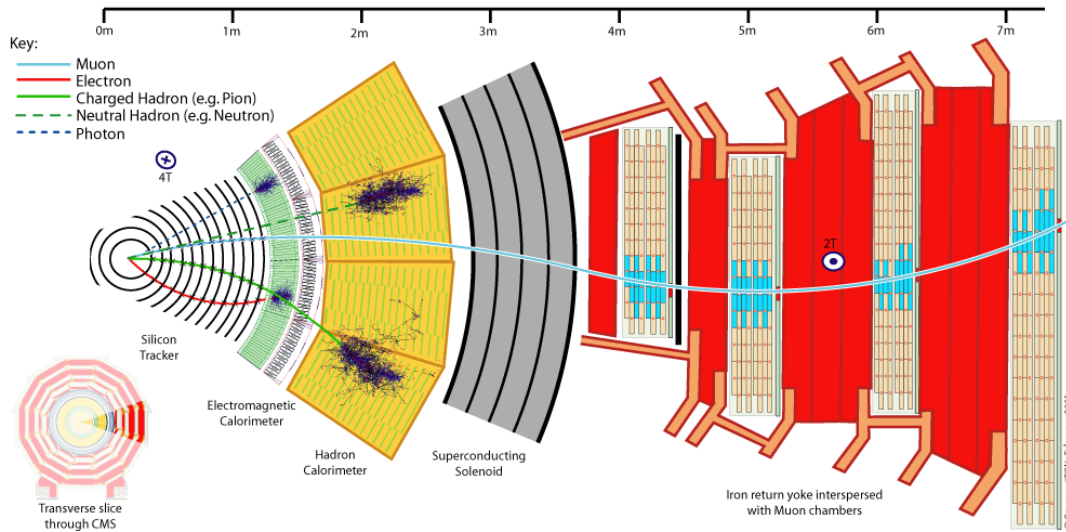
two beams are kept travelling in a circle by 1232 superconducting dipole magnets and steered to four collision points around the LHC. Detectors are situated at these collision points to observe the collisions, the main four being: ALICE [4], ATLAS [5], CMS [6] and LHCb [7]. A schematic of the LHC accelerator chain and the detectors can be seen in Figure 2.1

The number of times any physical process will occur in particle collisions can be expressed as:

$$N = \mathcal{L}\sigma, \quad (2.1)$$

where  $\mathcal{L}$  is the integrated luminosity and depends only on the parameters of the collisions and  $\sigma$  is the cross-section which depends only on the process. In order to observe rare (i.e. low cross-section) processes it is therefore necessary to use very high luminosity datasets. The integrated luminosity is obtained by integrating the instantaneous luminosity over time. For collisions at the LHC the instantaneous luminosity is given by:

$$\mathcal{L} = \frac{k_b N_b^2 f_{rev} \gamma}{4\pi \epsilon_n \beta}, [8] \quad (2.2)$$



**Figure 2.2:** A schematic cross-section of the CMS experiment showing the path taken by several types of particles.

where  $k_b$  is the number of bunches per beam,  $N_b$  the number of protons per bunch,  $f_{rev}$  the revolution frequency,  $\epsilon_n$  the normalised transverse beam emittance,  $\beta^*$  the beta-function at the interaction point and  $\gamma$  the Lorentz factor.

The cross-section for several processes is shown in Figure ?? and it can be seen that the cross-section for VBF Higgs production is approximately 1 pb.

## 2.2 The CMS experiment

### 2.2.1 Tracker

### 2.2.2 Electromagnetic calorimeter

### 2.2.3 Hadronic calorimeter

### 2.2.4 Muon system

### 2.2.5 Trigger system

# Chapter 3

## Physics objects and event reconstruction

### 3.1 Primary vertex

### 3.2 Jets

### 3.3 Missing transverse energy

### 3.4 Electrons

### 3.5 Muons

### 3.6 Taus

### 3.7 Photons

## Chapter 4

### Methods for limit setting

# Chapter 5

## Search for invisibly decaying Higgs bosons in run I prompt data

### 5.1 Trigger and event selection

### 5.2 Background estimation

#### 5.2.1 $W \rightarrow e\nu + \text{jets}$

#### 5.2.2 $W \rightarrow \mu\nu + \text{jets}$

#### 5.2.3 $W \rightarrow \tau\nu + \text{jets}$

#### 5.2.4 $Z \rightarrow \nu\nu + \text{jets}$

#### 5.2.5 QCD

#### 5.2.6 Minor backgrounds

### 5.3 Systematic uncertainties

### 5.4 Results





## Chapter 6

# Search for invisibly decaying Higgs bosons in run I parked data

### 6.1 Trigger

### 6.2 Event selection

### 6.3 Background estimation

#### 6.3.1 $W \rightarrow e\nu + \text{jets}$

#### 6.3.2 $W \rightarrow \mu\nu + \text{jets}$

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#### 6.3.5 QCD

#### 6.3.6 Minor backgrounds

### 6.4 Systematic uncertainties

### 6.5 Results

## Chapter 7

# Combinations and interpretations of run I searches for invisibly decaying Higgs bosons

### 7.1 Searches in other channels

### 7.2 Combination with prompt VBF search

### 7.3 Combination with the parked VBF search

### 7.4 Dark matter interpretations

## Chapter 8

### Search for invisibly decaying Higgs bosons in run II data



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