



# High-pressure oxidation of hydrogen diluted in N<sub>2</sub> with added H<sub>2</sub>O or CO<sub>2</sub> at 100 atm in a supercritical-pressure jet-stirred reactor

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## ABSTRACT

The oxidation of H<sub>2</sub> diluted in N<sub>2</sub> with and without 10 % H<sub>2</sub>O or 20 % CO<sub>2</sub> additions are studied at fuel-lean conditions at 100 atm and 500–1000 K in a supercritical-pressure jet-stirred reactor. The mole fractions of H<sub>2</sub> and O<sub>2</sub> are quantified by using micro-gas chromatography ( $\mu$ -GC). Experiment shows that H<sub>2</sub> oxidation is inhibited at lower temperatures (850–950 K) while it is promoted at higher temperatures (950–1050 K) with 10 % H<sub>2</sub>O additions or 20 % CO<sub>2</sub> additions. In addition, the effect of H<sub>2</sub>O is more significant than that of CO<sub>2</sub>. Five models are employed in simulations of the observables. Unfortunately, all of these models fail to capture the effect of H<sub>2</sub>O and CO<sub>2</sub> additions on H<sub>2</sub> oxidation. Pathway and sensitivity analyses of H<sub>2</sub> show that the reactions of H + O<sub>2</sub> + (M) = HO<sub>2</sub> + (M) and H<sub>2</sub>O<sub>2</sub> + (M) = 2OH + (M) dominate the radical production (HO<sub>2</sub> and OH) and H<sub>2</sub> oxidation at 100 atm. A further perturbation of pre-exponential coefficients and collisional factors of these reactions indicates that collisional factors of H<sub>2</sub>O and CO<sub>2</sub> have small effect under the experimental conditions, while a smaller reaction rate for H<sub>2</sub>O<sub>2</sub> + (M) = 2OH + (M) may explain the inhibiting effect of H<sub>2</sub>O and CO<sub>2</sub> additions at lower temperatures. Real-fluid corrections on intermolecular interactions and mixing rules should be further investigated to explain the effect of H<sub>2</sub>O and CO<sub>2</sub> additions.

## 1. Introduction

As around 80 % of global energy is obtained from combustion of fossil fuels, it is critical to adopt innovative combustion strategies of carbon-free renewable fuels, such as H<sub>2</sub>, for a low carbon society [1]. Among the advanced combustion strategies [2,3], such as Homogeneous Charge Compression Ignition and Reactivity Controlled Compression Ignition, an “ideal” burning condition of fuel lean, low temperature, and high pressure are optimal for high thermodynamic efficiencies (50–60 %) and low emissions. The “ideal” condition above generally falls into supercritical combustion regime, where most reactants are burned at the supercritical state. Moreover, with H<sub>2</sub> as the simplest fuel, its kinetics plays a foundational role in the hierarchical development of models for larger hydrocarbons and oxygenated fuels. Therefore, a better understanding of supercritical combustion of H<sub>2</sub> is vital for both kinetic studies and advanced engine designs.

In the supercritical state, intermolecular attractions become significant, and non-ideal fluid behavior can have substantial effects on the

reaction kinetics, as well as the thermodynamic and transport properties. For instance, at ultra-high pressures, multiple-body collisions can cause significant deviations from the isolated binary collision assumptions of most gas-phase kinetic models [4,5]. This results in unconventional pressure dependencies, and the collisional cross-section area may differ from the potential of interactions in real fluids. Notably, thermodynamic properties like the heat capacity for CO<sub>2</sub> and H<sub>2</sub>O can exhibit a 50 % variance from ideal gas values above 200 atm [6]. Consequently, there is a critical need for kinetic experiments and theoretical computations at pressures exceeding 100 atm to explore the chemistry of supercritical combustion.

Unfortunately, only a few experimental apparatuses, such as a shock tube [7,8], a high-pressure laminar flow reactor [9–12], and a supercritical-pressure jet-stirred reactor (SP-JSR) [4,5,13,14], can be used to study supercritical reaction chemistry under combustion conditions. The SP-JSR recently developed by Zhao et al. [4] provides a new experimental platform for studying supercritical fuel chemistry over an especially wide range of pressure and temperature (10–200 atm,

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295–1200 K), with excellent temperature uniformity ( $\pm 3$  K), and with a residence time (0.1–1.0 s) directly relevant to practical internal combustion engines. We recently studied the low and intermediate temperature chemistries of propane [14], n-butane [4], dimethyl ether (DME) [5], and methanol [13] at 100 atm using the SP-JSR, and developed high-pressure kinetic models, respectively.

As to the high-pressure studies of  $H_2$  oxidation chemistry, Hong et al. [15] updated the  $H_2/O_2$  reaction mechanism by incorporating their recent reaction rate determinations in shock tubes, and they emphasized significant uncertainties in the reaction rate of  $H_2O_2$  decomposition at high pressure and in the collisional efficiencies for colliders other than Ar and  $N_2$ . Therefore, these third-body reactions involved in  $H_2O$  and  $CO_2$  as colliders could have large uncertainties at high pressures. Meanwhile, Burke et al. [16] developed a high-pressure hydrogen oxidation mechanism based on their flame speed measurements [17,18] and other combustion targets. Lei et al. [19] also studied the nonlinear multi-component collision of  $HO_2 + M = H + O_2 + M$  under high pressures, highlighting an error of several times in traditional collisional factor computations. Shao et al. [8] updated the high-pressure kinetic model of  $H_2$  by measuring the ignition delay time of  $H_2 / O_2 / CO_2$  at 250 atm. Mueller [20], Li [21], Davis [22], Hong [15], Burke [16], Keromnes [23], and others have respectively developed detailed chemical reaction mechanisms for  $H_2$  either based on quantum chemistry computations or elementary reaction rate measurements. However, key rate coefficients in different kinetic models retain significant uncertainties. For example, errors in the  $H_2$  adiabatic flame speed predictions for existing models reach up to a factor of 4 between different models at elevated pressures [17].

Clearly high uncertainties in the  $H_2$  chemistry remain at high pressures, especially at supercritical conditions. In this work, the supercritical oxidation chemistry of  $H_2$  is studied with the SP-JSR under fuel lean conditions at 100 atm and 500–1000 K, which extends the available data to engine conditions. The mole fractions of  $H_2$  and  $O_2$  are quantified with a micro-gas chromatograph ( $\mu$ -GC). Moreover, the effect of real-fluid  $H_2O$  and  $CO_2$  on the  $H_2$  oxidation is investigated by adding 10 %  $H_2O$  or 20 %  $CO_2$  to the reactant mixture. Comparisons with simulations based on five representative models are made in the exploration of the  $H_2$  oxidation chemistry under these conditions. The key high-pressure reactions are identified through reaction pathway and sensitivity analyses for  $H_2$ . Updates are made to the rate parameters for some of these key reactions on the basis of ab initio kinetics.

## 2. Experimental methods and kinetic models

$H_2$  oxidation experiments are performed in the SP-JSR system, as is shown in Fig. 1. The SP-JSR is a spherical bulb with an internal volume of  $0.5\text{ cm}^3$ . It contains 4 fingers with 2 jets (0.2 mm I. D.) on each finger at the center of the bulb, which enables optimized jet orientations for intense turbulence mixing to improve the mixing and homogeneity of reactant and temperature distributions. The quartz reactor is placed inside a stainless-steel pressure-resistant jacket. The inside and outside pressures of the reactor are balanced by the bath gas flow to reach the high-pressure conditions. The gases issuing from the SP-JSR exit are sampled by a quartz sonic nozzle, and then equilibrate their pressure with vacuum around  $10^{-2}$  Torr generated by a dry pump. The experimental system is designed for experiments over pressures of 10–200 atm and temperatures of 298–1200 K. Further details of the SP-JSR and its jet and heating arrangements are provided in the literature [4]. The gas flow rates are controlled by high-pressure mass flow controllers (Brooks, SLA5800). Liquid water was delivered by a high performance liquid chromatography (HPLC) pump to a home-made high-pressure pre-vaporizer. The vaporized water vapor was carried by  $N_2$  to the oxidizer line. Gas samples are quantified by using a micro-gas chromatography-thermal conductivity detector (GC-TCD, Inficon 3000) within 5 % errors [24]. The axial temperature profiles under the experimental flow conditions are measured in 1 mm steps along the JSR bulb, where

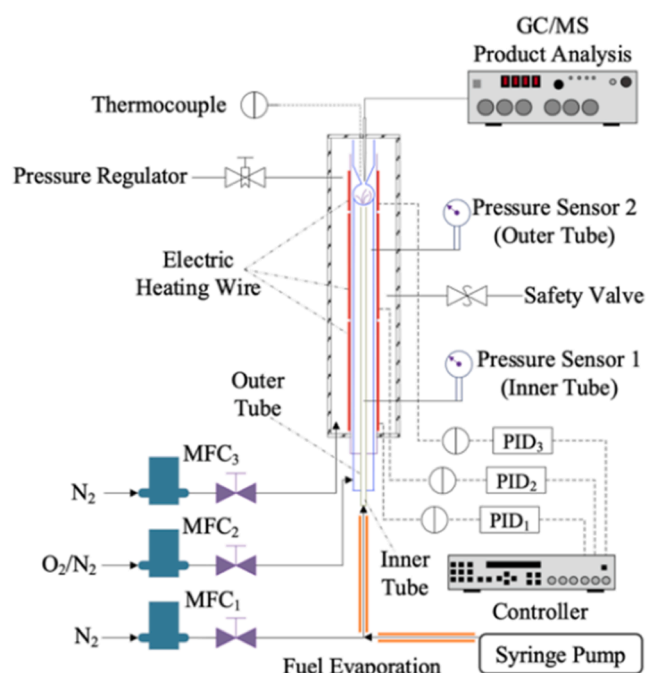


Fig. 1. Schematic of the SP-JSR setup.

the temperature variation is within  $\pm 3$  K between 500 and 1000 K. The details of the temperature profile measurement are provided in [5]. The repeated samples fall within the 5 % error for all conditions, which mainly comes from fuel purity and GC-TCD quantifications, while the temperature variations contribute to the measurement errors.

The experiment is performed at fuel lean conditions with and without  $H_2O$  and  $CO_2$  additions at 100 atm between 500–1000 K, as is shown in Table 1. The residence time in the reactor varies with reactor temperature ( $T$ ) to keep a fixed inlet volume flow rate at 4 L/min at 293 K through  $\tau = (V/v) * T_0 / T$ , where  $V$  is the volume of the JSR,  $v$  is the inlet volume flow rate at room temperature,  $T_0$  is the room temperature (293 K). Simulations are performed in the PSR module of CHEMKIN software [25] by using five  $H_2$  models, Aramco Mech 3.0 [26], Dagaut Mech [27], FFCM Mech [28], Hashemi Mech [10], and the updated HP Mech based on HP Mech [29] with including recently updated key elementary reaction rates for  $H_2$  chemistry as summarized in Table 2. The present theoretical kinetic predictions for the reaction of  $H_2O_2$  with H derive from earlier unpublished work of Wu et al. [30] through high-level quantum chemistry computations, which has been described in detail in the supplementary material. The third-body collision efficiencies of  $H_2O_2 + (M) = 2OH (+M)$  were taken from previous calculations involving a combination of collisional energy transfer trajectories and master equation analyses [31,32]. In particular, an average value of 8.71 was picked for the third-body efficiency of  $H_2O$ .

## 3. Results and discussion

The experimental and modeling results for oxidation of  $H_2$  diluted in  $N_2$  at 100 atm are plotted in Fig. 2. Most models, such as Aramco Mech 3.0, Dagaut Mech, FFCM Mech, Hashemi Mech, and updated HP Mech predict  $H_2$  and  $O_2$  mole fractions that agree reasonably well with the experimental measurements, exhibiting accurate predictions of the  $H_2$  oxidation in  $N_2$  dilutions at 100 atm.

Fig. 3 (a) and (b) depict the  $H_2$  evolution against temperature with 10 %  $H_2O$  additions and 20 %  $CO_2$  additions, respectively, at 100 atm. In Fig. 3 (a), the onset temperature of the  $H_2$  oxidation with 10 %  $H_2O$  additions is delayed, resulting in a right shift of the oxidation scatter between 850 and 950 K. However, at higher temperatures, the  $H_2$

**Table 1**

Experimental conditions.

Case	Equivalence ratio	Pressure (atm)	H <sub>2</sub> (%)	O <sub>2</sub> (%)	N <sub>2</sub> (%)	CO <sub>2</sub> (%)	H <sub>2</sub> O (%)	Residence time (s)	Temperature (K)
1	0.343	100	3.248	4.738	92.014	0	0	0.352-0.195	500-1000
2	0.343	100	3.248	4.738	82.014	0	10.000	0.352-0.195	500-1000
3	0.343	100	3.248	4.738	72.090	19.924	0	0.352-0.195	500-1000

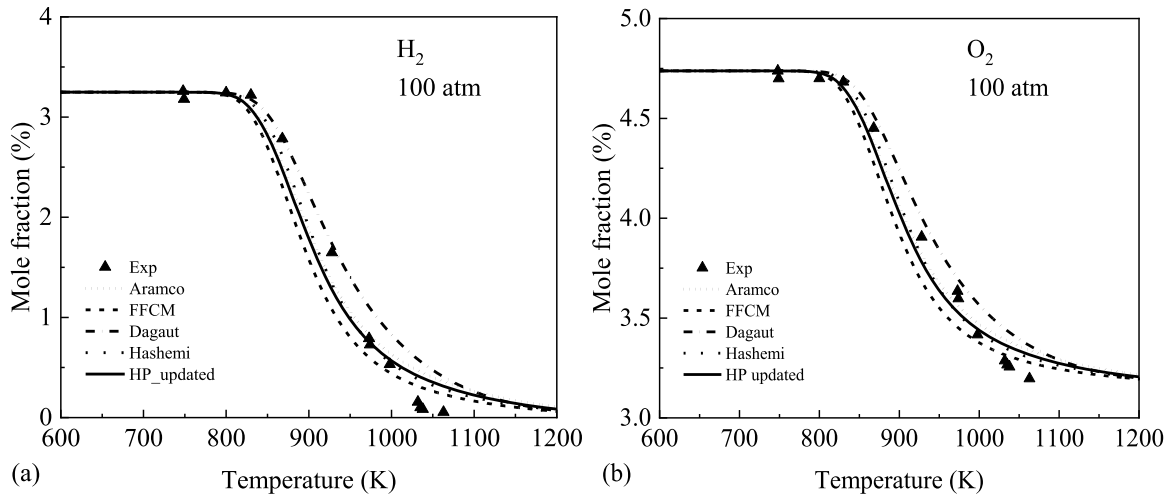
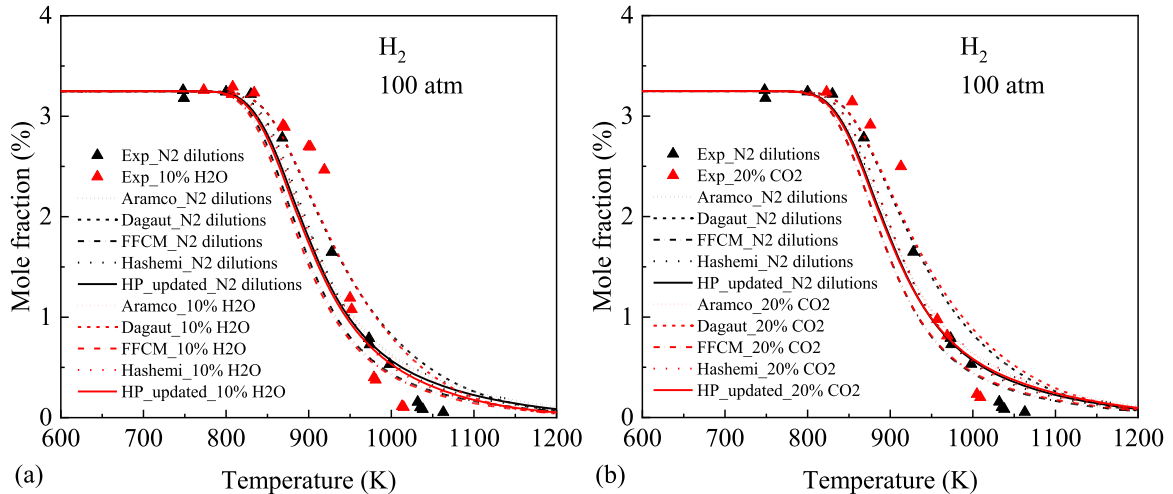
**Table 2**

Updated reaction rates in updated HP Mech from the HP base Mech.

Reactions	A ( $\text{cm}^3/\text{mol.s}$ )	n	Ea ( $\text{cal/mol}$ )	Refs.
$\text{H}_2\text{O}_2 + \text{H} = \text{H}_2\text{O} + \text{OH}$	3.35E+07	1.91	3654	Present
$\text{H}_2\text{O}_2 + \text{H} = \text{H}_2 + \text{HO}_2$	44.00E+00	3.45	712	Present
$\text{H}_2\text{O}_2 (+\text{M}) = 2\text{OH} (+\text{M})$	2.00E+12	0.9	48749	[31,
LOW	2.49E+24	-2.3	48749	32]
TROE	0.43	1.0E-	1.0E+30	
$\text{H}_2\text{O}/8.71/\text{CO}_2/4.79/\text{N}_2/$ $2.07/\text{O}_2/1.5/\text{He}/1.00/$ $\text{H}_2\text{O}_2/8.07/\text{H}_2/3.7/\text{CO}/2.8/$		30		
$\text{HO}_2 + \text{HO}_2 = \text{H}_2\text{O}_2 + \text{O}_2$	1.93E-02	4.12	-4960	[33]
$\text{HO}_2 + \text{HO}_2 = \text{OH} + \text{OH} + \text{O}_2$	6.41E+17	-1.54	8540	[33]

oxidation with H<sub>2</sub>O additions is promoted, leading to a left shift of the scatter between 950 and 1050 K. Therefore, H<sub>2</sub>O additions “twist” the H<sub>2</sub> oxidation more steeply. In Fig. 3 (b), 20 % CO<sub>2</sub> additions show a similar effect on H<sub>2</sub> oxidation as the 10 % H<sub>2</sub>O addition of case 2 did, although its impact is somewhat smaller than for the only 10 % H<sub>2</sub>O additions. That reduction is likely because H<sub>2</sub>O is a highly polar molecule and exhibits a stronger intermolecular interaction than CO<sub>2</sub> at high pressures. As to model predictions, unfortunately, all of the five models exhibit little difference from the purely N<sub>2</sub> diluted case for both H<sub>2</sub>O or CO<sub>2</sub> additions. They clearly fail to capture the experimentally observed effect of H<sub>2</sub>O and CO<sub>2</sub> additions on H<sub>2</sub> oxidation.

The mole fraction of O<sub>2</sub> at 100 atm for 10 % H<sub>2</sub>O addition and 20 % CO<sub>2</sub> addition are plotted in Fig. 4 (a) and (b), respectively. In the experiments, H<sub>2</sub>O and CO<sub>2</sub> additions delay the consumption of O<sub>2</sub> at lower temperatures, but accelerate it at higher temperature, which agree with

**Fig. 2.** Mole fractions of H<sub>2</sub> and O<sub>2</sub> in the oxidation of H<sub>2</sub> diluted in N<sub>2</sub> at 100 atm.**Fig. 3.** Mole fraction of H<sub>2</sub> in H<sub>2</sub> oxidation with 10 % H<sub>2</sub>O additions (a) and 20 % CO<sub>2</sub> additions (b) at 100 atm.

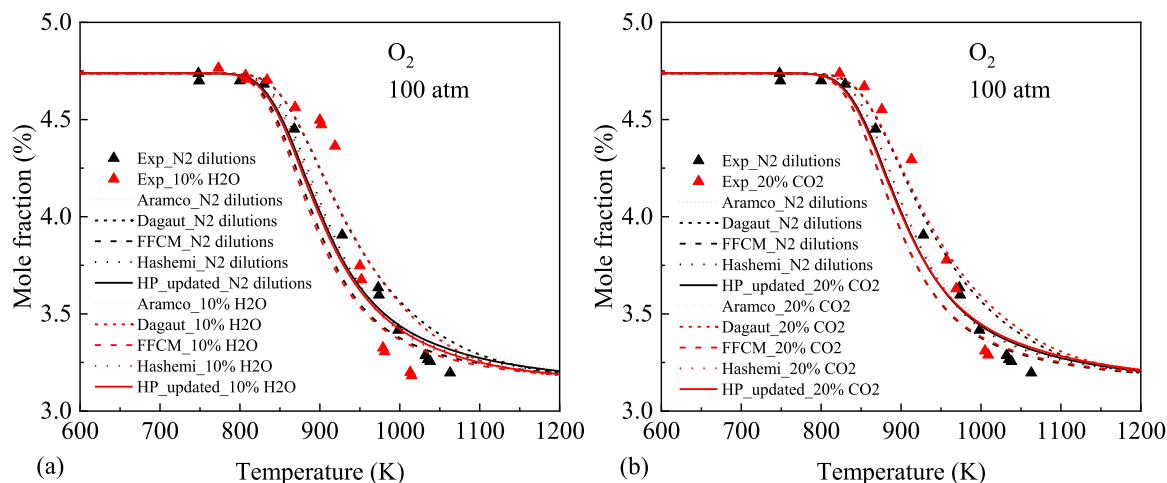


Fig. 4. Mole fraction of  $O_2$  in  $H_2$  oxidation in 10 %  $H_2O$  additions (a) and 20 %  $CO_2$  additions (b) at 100 atm.

the observations in Fig. 3. Similarly,  $H_2O$  and  $CO_2$  additions in the simulations show little impact on  $O_2$  oxidation compared to the case for pure  $N_2$  dilution.

To explain the key characteristics in  $H_2$  oxidation, the pathway analyses of  $H_2$  in  $N_2$  dilution, with 10 %  $H_2O$ , and with 20 %  $CO_2$  are plotted in Fig. 5 (a–c), respectively, at 875 K and 100 atm by using updated HP Mech. In Fig. 5 (a), the H abstraction reaction of  $OH + H_2$  dominates the  $H_2$  oxidation pathway and produces  $H_2O$  and H. The H radical combines with  $O_2$  and forms  $HO_2$  very rapidly through the third-body reaction  $H + O_2 + (M) = HO_2 + (M)$ . Then,  $H_2O_2$ , formed from  $HO_2$ , produces two OH radicals through  $H_2O_2 + (M) = 2OH + (M)$ . Moreover, it is well known that  $H + O_2 = OH + O$  is a crucial branching reaction competing with  $H + O_2 + (M) = HO_2 + (M)$  at low pressures or high temperatures. However, the radical production from  $H + O_2 = OH + O$  almost disappears at 100 atm. Thanks to the high pressure, the third-body reaction  $H + O_2 + (M) = HO_2 + (M)$  is enhanced dramatically and dominates the radical production and  $O_2$  consumption pathways. The discussion above agrees with the typical 3<sup>rd</sup> explosion limit of  $H_2$ . The kinetic effect of  $H_2O$  or  $CO_2$  additions on  $H_2$  oxidation is mainly through altering the reaction rates of the third-body reactions, such as  $H + O_2 + (M) = HO_2 + (M)$  and  $H_2O_2 + (M) = 2OH + (M)$ . Unfortunately, the production fluxes of  $H + O_2 + (M) = HO_2 + (M)$  and  $H_2O_2 + (M) = 2OH + (M)$  in Fig. 5(b) and (c) show little differences from that in Fig. 5 (a). In addition,  $H_2O_2$  formation in these three cases is mainly from  $HO_2 + HO_2$  reactions, which accounts for 73.3 % of the  $H_2O_2$  production flux, while a minor reaction path is  $HO_2 + H_2$ .

To identify key reactions in  $H_2$  oxidation, the sensitivity analyses of  $H_2$  diluted in  $N_2$ , with 10 %  $H_2O$ , and with 20 %  $CO_2$  are plotted in Fig. 6 (a–c), respectively, at 875 K and 100 atm by using updated HP Mech. In Fig. 6(a), it is seen that reactions involved in  $H_2O_2$  pathways are most sensitive, such as  $HO_2 + HO_2 = H_2O_2 + O_2$  (duplicated reactions in

updated HP Mech),  $H_2O_2 + (M) = 2OH + (M)$ , and  $H_2O_2 + OH = HO_2 + H_2O$ . It indicates that  $H_2$  oxidation at high pressure is controlled by  $HO_2$  and  $H_2O_2$  chemistry, which explains the broadly enhanced warm flame chemistry in our previous oxidation studies of propane, dimethyl ether, and n-butane at 100 atm [4,5,14]. Similarly, Fig. 6 (b) and (c) shows little difference in the sensitive reactions from Fig. 6(a).

It is noted that the reactions  $H + O_2 + (M) = HO_2 + (M)$  and  $H + O_2 = O + OH$  are less sensitive than the reactions above. Their reasons are different. The former is because its reaction rate becomes dramatically fast at high pressures and is close to the partial equilibrium assumption, while the latter is relatively much slower than the former, leading to a “frozen” reaction even at high temperatures. This argument also agrees with the pathway analysis in Fig. 5. Therefore, reactions involved with OH become very important as OH is mainly produced from  $H_2O_2 + (M) = 2OH + (M)$  at high pressures, such as  $H_2 + OH = H + H_2O$ ; while reactions with O are frozen as O is mainly produced from  $H + O_2 = O + OH$ , such as  $O + H_2 = OH + H$ . In addition, compared with Fig. 6(a), reactions  $H_2 + OH = H + H_2O$  and  $HO_2 + OH = O_2 + H_2O$  become more sensitive with  $H_2O$  or  $CO_2$  additions in Fig. 6(b) and (c). That is because the average collisional factor of the mixture increases with  $H_2O$  or  $CO_2$  additions and enhances the OH concentrations from  $H_2O_2 + (M) = 2OH + (M)$ .  $H_2O_2 + H = H_2 + HO_2$  also becomes sensitive with  $CO_2$  addition as the reaction  $CO + OH = H + CO_2$  is suppressed with  $CO_2$  addition.

Several possible reasons are proposed here to explain the effects of  $H_2O$  and  $CO_2$  addition on  $H_2$  oxidations: (i) pre-exponential coefficients of  $H + O_2 + (M) = HO_2 + (M)$  and  $H_2O_2 + (M) = 2OH + (M)$ ; (ii) collisional factor (M factor) of the reactions above; (iii) real-fluid corrections on reaction rate calculations.

Therefore, the perturbations of pre-exponential coefficients and collisional factors of  $H + O_2 + (M) = HO_2 + (M)$  and  $H_2O_2 + (M) = 2OH + (M)$  on  $H_2$  oxidation are investigated in Fig. 7(a) and (b), respectively,

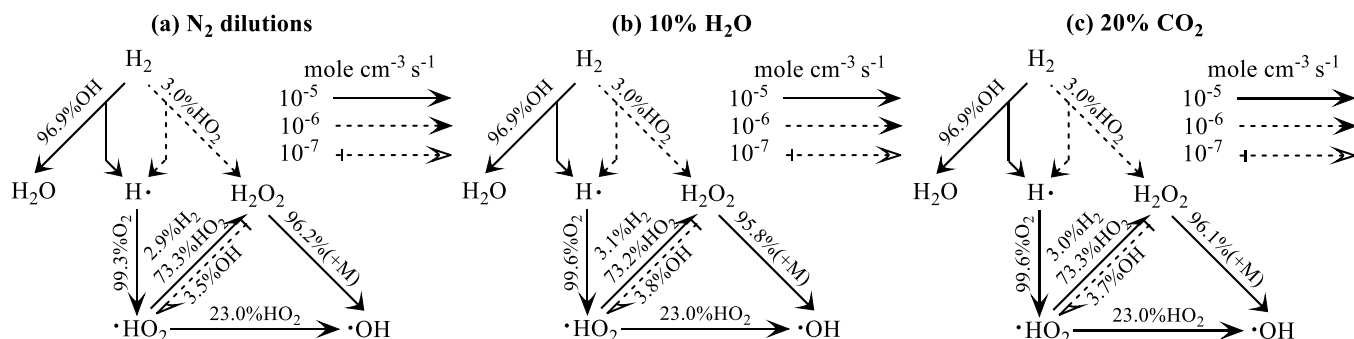


Fig. 5. Reaction pathways for  $H_2$  diluted in  $N_2$  (a), with 10 %  $H_2O$  addition (b), and with 20 %  $CO_2$  addition (c) at 875 K and 100 atm by using updated HP Mech.

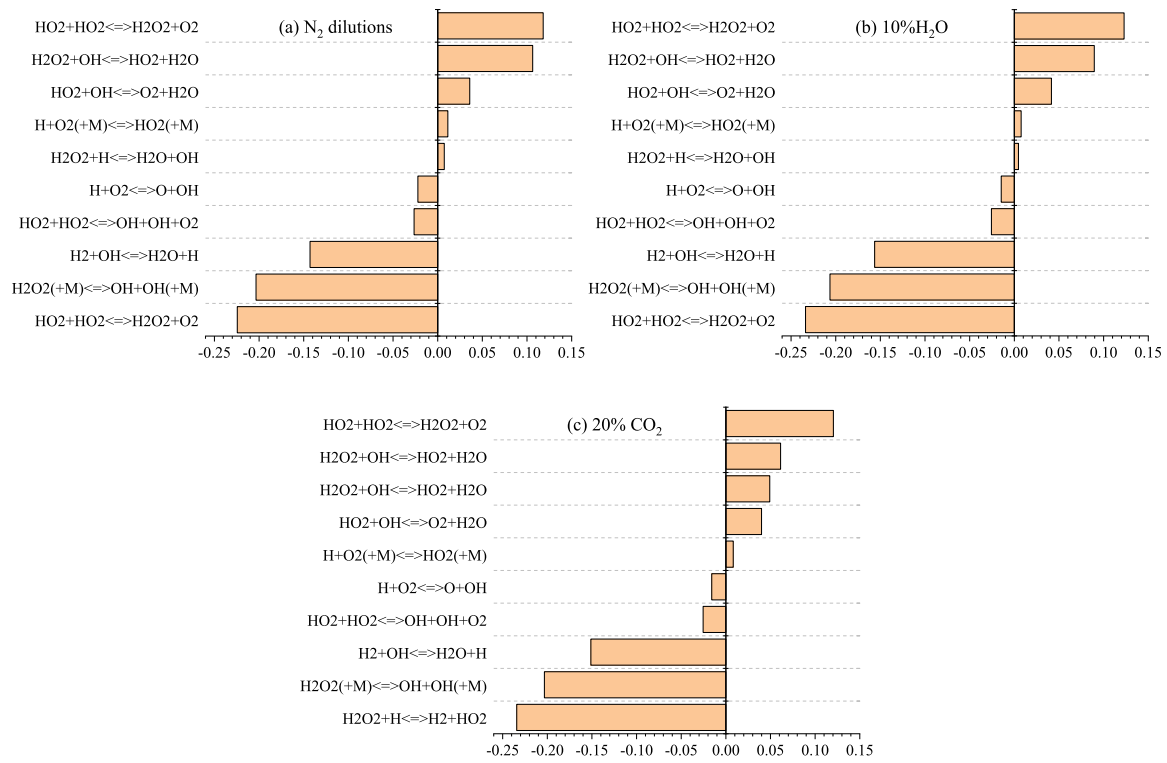


Fig. 6. Sensitivity analyses for  $H_2$  diluted in  $N_2$  (a), with 10 %  $H_2O$  addition (b), and with 20 %  $CO_2$  addition (c) at 875 K and 100 atm by using updated HP Mech.

at 100 atm by using updated HP Mech. In Fig. 7(a), compared with oxidation of  $H_2$  diluted in  $N_2$ , addition of  $H_2O$  exhibits little impact on the oxidation in the modelling. By reducing the pre-exponential coefficient of  $H_2O_2 + (M) = 2OH + (M)$  in the high-pressure limit expression of a Troe format by a factor of 0.1, the oxidation is significantly delayed, while increasing the reaction rate of  $H + O_2 + (M) = HO_2 + (M)$  by a factor of 100 has little impact. This finding agrees with the discussion of the two reactions above in the sensitivity analysis. Therefore, it looks like a smaller rate of  $H_2O_2 + (M) = 2OH + (M)$  in simulations reduces the discrepancy with experiments.

In Fig. 7(b), we further perturb  $H_2$  oxidation by altering the  $H_2O$  addition and the collisional factors of  $N_2$  and  $H_2O$  for  $H_2O_2 + (M) = 2OH + (M)$ . It is seen that a 90 % addition of  $H_2O$  accelerates the oxidation

significantly as the collisional factor of  $H_2O$  is  $\sim 4$  times larger than  $N_2$  as a major bath gas, while it fails to explain the inhibiting effect of  $H_2O$  addition on  $H_2$  oxidation at lower temperatures in experiments. In the case of pure  $N_2$  dilutions or with 10 %  $H_2O$  addition, the change of the collisional factor of  $N_2$  plays a major role in affecting the rate of  $H_2O_2 + (M)$  as  $N_2$  is the major bath gas. However, the change of the collisional factor of  $H_2O$  has no impact, even though the collisional factor of  $H_2O$  is reduced by 100 times. Therefore, we think the update of the collisional factor of  $H_2O$  in  $H_2O_2 + (M) = 2OH + (M)$  from the current model could not explain the present experimental results.

According to the discussion above, a smaller reaction rate of  $H_2O_2 + (M) = 2OH + (M)$  may explain the inhibiting effect of  $H_2O$  addition at lower temperatures. Other possible reasons, such as real-fluid

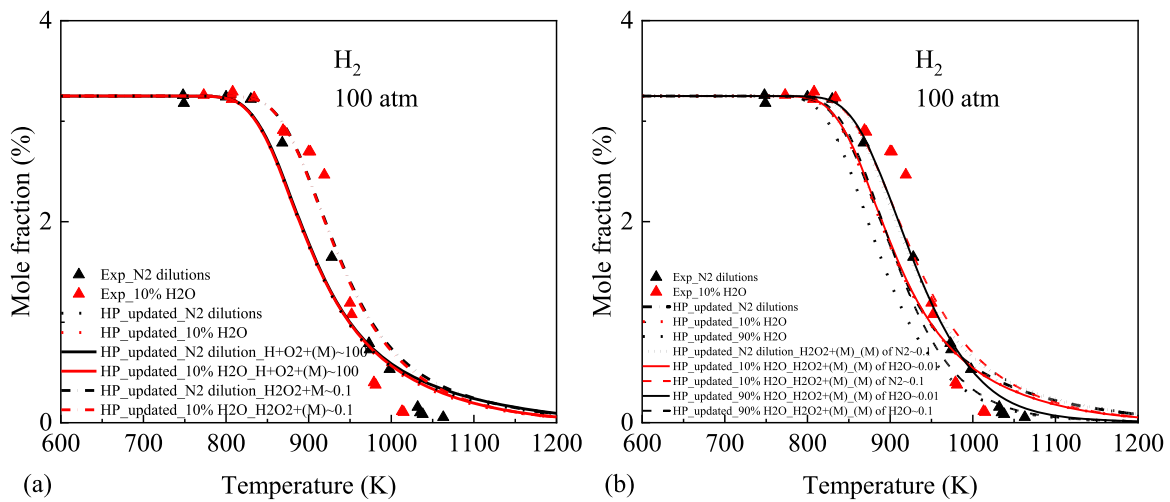


Fig. 7. (a) Effect of pre-exponential coefficients of  $H + O_2 + (M) = HO_2 + (M)$  (increase by a factor of 100) and  $H_2O_2 + (M) = 2OH + (M)$  (decrease by a factor of 0.1) on  $H_2$  oxidation at 100 atm by using updated HP Mech; (b) Effect of different collisional factors (M factor) of  $N_2$  and  $H_2O$  for  $H_2O_2 + (M) = 2OH + (M)$  on  $H_2$  oxidation at 100 atm by using updated HP Mech.



corrections to the intermolecular force field, collisional cross-sections, mixing rules, and multi-body collisions should be further investigated to explain the effect of H<sub>2</sub>O and CO<sub>2</sub> additions.

#### 4. Conclusion

The supercritical pressure jet stirred reactor provides a valuable platform for performing kinetic studies at low and intermediate temperatures at extreme pressures with a uniform temperature distribution and a short flow residence time. The oxidations of H<sub>2</sub> diluted in N<sub>2</sub>, with 10 % H<sub>2</sub>O, or 20 % CO<sub>2</sub> addition are studied at 100 atm and 500–1000 K in SP-JSR experiments. The mole fractions of H<sub>2</sub> and O<sub>2</sub> are quantified with a micro-gas chromatograph ( $\mu$ -GC). Experiment shows that H<sub>2</sub> oxidation is inhibited at lower temperatures (850–950 K), while it is promoted at higher temperatures (950–1050 K) with 10 % H<sub>2</sub>O addition or 20 % CO<sub>2</sub> addition. Simply put, H<sub>2</sub>O or CO<sub>2</sub> additions “twist” the H<sub>2</sub> oxidation more steeply. In addition, the effect of H<sub>2</sub>O is more significant than that of CO<sub>2</sub>, while the reaction H<sub>2</sub>O<sub>2</sub> + H = H<sub>2</sub> + HO<sub>2</sub> becomes more sensitive with CO<sub>2</sub> addition.

Five models (Aramco Mech 3.0, Dagaut Mech, FFCM Mech, Hashemi Mech, and updated HP Mech) are employed in the present simulations. Unfortunately, all of these models fail to capture the effect of H<sub>2</sub>O and CO<sub>2</sub> additions on H<sub>2</sub> oxidation. Through the pathway analyses and sensitivity analyses of H<sub>2</sub>, it is seen that the reactions of H + O<sub>2</sub> + (M) = HO<sub>2</sub> + (M) and H<sub>2</sub>O<sub>2</sub> + (M) = 2OH + (M) dominate the radical production (HO<sub>2</sub> and OH) and H<sub>2</sub> oxidation at 100 atm.

A further perturbation of pre-exponential coefficients and collisional factors for the above reactions indicates that the change of their collisional factors has little effect on H<sub>2</sub> oxidation, while a smaller reaction rate of H<sub>2</sub>O<sub>2</sub> + (M) = 2OH + (M) may explain the inhibiting effect of H<sub>2</sub>O and CO<sub>2</sub> additions at lower temperatures. Other possible reasons, such as real-fluid corrections on intermolecular force field, collisional cross-sections, mixing rules, and multi-body collisions should be further investigated to explain the effect of H<sub>2</sub>O and CO<sub>2</sub> additions.

#### Novelty and significance statement

Supercritical combustion gains increasing attentions for its high thermodynamic efficiency and low emissions, and it is widely employed in advanced combustion engines and rockets. However, non-ideal fluid behaviors under supercritical conditions can have substantial effects on the reaction kinetics, and high uncertainties in the H<sub>2</sub> chemistry remain at high pressures with limited experimental data as well. The present work provides important speciation data in H<sub>2</sub> oxidation at 100 atm by using the first supercritical-pressure jet-stirred reactor, and develops a new H<sub>2</sub> model for supercritical combustion by updating several key elementary reactions. A special real-fluid impact of H<sub>2</sub>O or CO<sub>2</sub> additions on H<sub>2</sub> oxidation has been revealed in experiments, while is not predicted in existing models. This work emphasizes the real-fluid corrections on intermolecular interactions and mixing rules for the community of supercritical combustion as future work.

#### CRediT authorship contribution statement

**Hao Zhao:** Writing – review & editing, Writing – original draft, Software, Methodology, Data curation, Conceptualization. **Chao Yan:** Writing – review & editing, Methodology, Data curation. **Guohui Song:** Formal analysis, Data curation. **Ziyu Wang:** Writing – review & editing, Software, Data curation. **Ahren W. Jasper:** Writing – review & editing, Software, Data curation. **Stephen J. Klippenstein:** Methodology, Writing – review & editing. **Yiguang Ju:** Writing – review & editing, Supervision, Funding acquisition.

#### Declaration of competing interest

The authors declare that they have no known competing financial

interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Supplementary materials

Supplementary material is submitted along with the manuscript. The updated HP Mech for H<sub>2</sub> is also included in the supplementary material.

Supplementary material associated with this article can be found, in the online version, at [doi:10.1016/j.combustflame.2024.113543](https://doi.org/10.1016/j.combustflame.2024.113543).

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