TDDFT Routines

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1 RPA matrix equations

We have $\mathbf{A}_{ia,jb}$ and $\mathbf{B}_{ia,jb}$ defined as:

$$\mathbf{A}_{ia,jb} = \delta_{ij}\delta_{ab}(\varepsilon_a - \varepsilon_i) + (\underline{ia}||jb) \tag{1}$$

$$\mathbf{B}_{ia,jb} = (\underline{ia}||bj) \tag{2}$$

1.1 Direct approximation

In the divert approximation, which we are working with because it is simpler to start with, this becomes:

$$\mathbf{A}_{ia,jb} = \delta_{ij}\delta_{ab}(\varepsilon_a - \varepsilon_i) + (\underline{ia}|\underline{jb}) \tag{3}$$

$$\mathbf{B}_{ia,jb} = (\underline{ia}|bj) \tag{4}$$

After the spin integration, we get:

$$\mathbf{A}_{ia,jb} = \delta_{ij}\delta_{ab}(\varepsilon_a - \varepsilon_i) + 2(ia|jb) \tag{5}$$

$$\mathbf{B}_{ia,jb} = 2(ia|bj) \tag{6}$$

We build the eigenvalue equation:

$$\begin{bmatrix} A & B \\ -B & -A \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix} = \omega \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix}$$
 (7)

1.2 Tamm-Dancoff approximation

In the Tamm-Dancoff approximation, we neglect **B**, which means that we have:

$$\mathbf{AX} = \omega \mathbf{X} \tag{8}$$

The eigenvalues ω are the excitation energies, but we need to do some work to find the excitation vectors V^{μ}_{pq} . However, we will use the eigenvectors from this matrix equation to build the excitation vectors.

1.3 Building the excitation doctors

First, we consider within the direct approximation:

$$\mathbf{W}_{p,q,ia} = \sum_{p,q,i,a} (\underline{pq}|\underline{ia}) \tag{9}$$

1.3.1 Spin integration

 $|\Phi_i^a\rangle$ has the CSF

$$|\Phi_{singlet}\rangle = \frac{1}{\sqrt{2}}(|\Phi_{i\alpha}^{a\alpha}\rangle + |\Phi_{i\beta}^{a\beta}\rangle)$$
 (10)

We want to consider something like:

$$\langle \Phi_0 | \frac{1}{4} \sum_{pqrs} V_{pq\underline{r}\underline{s}} a_p^{\dagger} a_q^{\dagger} a_s a_r \frac{1}{\sqrt{2}} (a_a^{\dagger \alpha} a_i^{\alpha} + a_a^{\dagger \beta} a_i^{\beta}) | \Phi_0 \rangle$$
 (11)

Consolidating constants out front and distributing the CSF terms:

$$\frac{1}{4\sqrt{2}}V_{pq\underline{r}\underline{s}}\langle\Phi_{0}|a_{p}^{\dagger}a_{q}^{\dagger}a_{s}a_{r}a_{a}^{\dagger\alpha}a_{i}^{\alpha}|\Phi_{0}\rangle + \frac{1}{4\sqrt{2}}V_{pq\underline{r}\underline{s}}\langle\Phi_{0}|a_{p}^{\dagger}a_{q}^{\dagger}a_{s}a_{r}a_{a}^{\dagger\beta}a_{i}^{\beta}|\Phi_{0}\rangle \quad (12)$$

Let's consider the first term:

$$\frac{1}{4\sqrt{2}} \sum_{pqrs} V_{pq\underline{r}\underline{s}} \langle \Phi_0 | a_p^{\dagger} a_q^{\dagger} a_s a_r a_a^{\dagger \alpha} a_i^{\alpha} | \Phi_0 \rangle \tag{13}$$