

Ch/ChE 164 Winter 2024
Homework Problem Set #5

Due Date: Thursday, February 22, 2024 @ 11:59pm PT
Out of 100 Points

- 1.** (15 points) (4.15 from Chandler) Consider an isomerization process $A \rightleftharpoons B$, where A and B refer to the different isomer states of a molecule. Imagine that the process takes place in a dilute gas, and that $\Delta\epsilon$ is the energy difference between state A and state B . According to the Boltzmann distribution law, the equilibrium ratio of A and B populations is given by

$$\frac{\langle N_A \rangle}{\langle N_B \rangle} = \frac{g_A}{g_B} e^{-\beta\Delta\epsilon}, \quad (1)$$

where g_A and g_B are the degeneracies of states A and B , respectively. Show how this same result follows from the condition of chemical equilibria, $\mu_A = \mu_B$.

- 2.** (20 points) (4.16 from Chandler) Consider the system described in problem 1. The canonical partition function is

$$Q = \frac{q^N}{N!}, \quad (2)$$

where N is the total number of molecules, and q is the Boltzmann weighted sum over all single molecule states, both those associated with isomers of type A and those associated with isomers of type B .

- (a) Show that one may partition the sum and write

$$Q = \sum_P \exp \{-\beta F(N_A, N_B)\} \quad (3)$$

with

$$-\beta F(N_A, N_B) = \log[(N_A!N_B!)^{-1} q_A^{N_A} q_B^{N_B}], \quad (4)$$

where \sum_P is over all the partitions of N molecules into N_A molecules of type A and N_B molecules of type B , q_A is the Boltzmann weighted sum over states of isomer A , and q_B is similarly defined.

- (b) Show that the condition of chemical equilibria is identical to finding the partitioning that minimizes the Helmholtz free energy

$$\frac{\partial F}{\partial \langle N_A \rangle} = \frac{\partial F}{\partial \langle N_B \rangle} = 0, \quad (5)$$

subject to the constraint that $\langle N_A \rangle + \langle N_B \rangle = N$ is fixed.

- 3.** (15 points) (4.25 from Chandler) Use the information compiled in Chapter 8 of Hill's *Introduction to Statistical Thermodynamics* to calculate the equilibrium constant, K , for the reaction $I_2 \rightleftharpoons 2I$ when the reaction occurs in the dilute gas phase at $T = 1000^\circ K$. (Note: For the electronic partition function, you need to consider the difference in degeneracies of the ground states of the iodine atom and molecule, $g_{e,I} = 4, g_{e,I_2} = 1$. This is based on the molecular orbital theory. You are not required to show this.)

- 4.** (50 pts.) Consider a one-dimensional monatomic crystal of N atoms with equilibrium nearest neighbor spacing a . In order to minimize boundary effects, assume a periodic boundary condition, i.e., $x_{N+1} = x_1$ where x_i is the position of the i th atom. If the atoms interact only with nearest neighbors via a potential $u(x_{i+1} - x_i)$,

- (i) show that the energy of the crystal can be written in the form

$$H = \frac{m}{2} \sum_{i=1}^N \dot{\xi}_i^2 + \frac{K}{2} \sum_{i=1}^N (\xi_{i+1} - \xi_i)^2 + N u(a) \quad (6)$$

to quadratic order in the displacement $\xi_i \equiv x_i - x_i^{(0)}$, where $x_i^{(0)}$ is the equilibrium position of the i th atom. What is K ?

(ii) Now decompose configurations of atoms into normal modes.

(a) Define normal modes η_k such that ξ_j is a linear superposition of η_k

$$\xi_j = \frac{1}{\sqrt{2N}} \sum_k \eta_k e^{i(jak)}. \quad (7)$$

Show that the periodic boundary condition leads to $k = \frac{2\pi n}{Na}$ where n is any integer.

Show further that adding $\frac{2\pi}{a}$ to k does not change ξ_j . Therefore there are only N independent modes.

We choose $n \in [-\frac{N}{2}, \frac{N}{2} - 1]$ (assuming N even).

(b) Show that the fact that the ξ_j 's are real leads to $\eta_k^* = \eta_{-k}$, where η_k^* is the complex conjugate of η_k .

(iii) It can be shown that the normal mode coordinates η_k diagonalize the Hamiltonian:

$$\sum_j (\xi_{j+1} - \xi_j)^2 = \sum_{k>0} \left[(\eta_k^R)^2 + (\eta_k^I)^2 \right] 4 \sin^2\left(\frac{1}{2}ka\right) \quad (8)$$

and

$$\sum_j \dot{\xi}_j^2 = \sum_{k>0} \left[(\dot{\eta}_k^R)^2 + (\dot{\eta}_k^I)^2 \right], \quad (9)$$

where η_k^R and η_k^I are the real and imaginary parts of η_k . What is the frequency for each normal mode ω_k ? What is the speed of sound for this model? (The speed of sound is defined as $\frac{d\omega_k}{dk}|_{k=0}$).

(iv) Show that in a large 1-D solid, the density (or degeneracy) of normal modes with the frequencies between ω and $\omega + d\omega$ is

$$g(\omega)d\omega = \frac{2N}{\pi\omega_m \sqrt{1 - (\omega/\omega_m)^2}} d\omega, \quad (10)$$

where ω_m is the maximum frequency of normal modes.

(v) Compare the Debye model and the exact results.

(a) If one is to make the Debye model for this 1-D solid, what is its Debye temperature Θ_D ? Why is the Debye frequency ω_D larger than the maximum frequency allowed in the system ω_m ?

(b) Show that one gets the same dependency of the heat capacity on temperature at very low temperature with the Debye approximation and with the exact degeneracy. Why does the Debye model give an accurate result even though unphysical normal modes are considered?

(1) isomerization of A \rightleftharpoons B

- ideal mixture

$$Q = \frac{q_A^{N_A}}{N_A!} \cdot \frac{q_B^{N_B}}{N_B!}$$

$$\mu_A = -\left(\frac{\delta \ln Q}{\delta \beta N_A}\right)_{B,V,N_B}, \quad \mu_B = -\left(\frac{\delta \ln Q}{\delta \beta N_B}\right)_{B,V,N_A}$$

Stirling
approx.

$$\ln Q = \ln [q_A^{N_A} q_B^{N_B}] - \ln [N_A!] - \ln [N_B!]$$

$$= N_A \ln q_A + N_B \ln q_B - N_A \ln N_A + N_A - N_B \ln N_B + N_B$$

$$-\beta \mu_A = \ln q_A - \ln N_A - \cancel{1 + 1}$$

$$-\beta \mu_B = \ln q_B - \ln N_B$$

$$\mu_A = \mu_B$$

$$\ln q_A - \ln N_A = \ln q_B - \ln N_B$$

$$\frac{N_A}{N_B} = \frac{q_A}{q_B} = \frac{q_A}{q_B} e^{-\beta \Delta E}$$

$$\boxed{\frac{N_A}{N_B} = \frac{q_A}{q_B} e^{-\beta \Delta E}}$$

(2) isomerization of $A \rightleftharpoons B$

$$Q = \frac{q^n}{N!} = \frac{q_A^{N_A}}{N_A!} \cdot \frac{q_B^{N_B}}{N_B!}$$

(a) ideal mixture A, B (no interaction)

$$\begin{aligned} Q &= \frac{(q_A + q_B)^N}{N!} \quad N_A + N_B = N \\ &= \frac{1}{N!} \sum_{N_A=0}^N \frac{N!}{N_A! N_B!} q_A^{N_A} q_B^{N_B} \\ &\quad \text{↳ sum over all partitions} \\ &= \sum_P q_A^{N_A} q_B^{N_B} [N_A! N_B!]^{-1} \\ &= \sum_P \exp(\log(q_A^{N_A} q_B^{N_B} [N_A! N_B!]^{-1})) \\ &\quad \text{↳ } -\beta F(N_A, N_B) \end{aligned}$$

$$Q = \sum_P \exp[-\beta F(N_A, N_B)]$$

(b) chemical equilibria + variational principle

$$\mu_A = -\left(\frac{\delta \ln Q}{\delta \beta N_A}\right)_{B, V, N_B}, \quad \mu_B = -\left(\frac{\delta \ln Q}{\delta \beta N_B}\right)_{B, V, N_A}$$

$$\mu_A = \mu_B \rightarrow \frac{\delta F}{\delta N_A} = \frac{\delta F}{\delta N_B}$$

$$\text{b/c } N_A + N_B = N \rightarrow -\frac{\delta F}{\delta N_A} = \frac{\delta F}{\delta N_B}$$

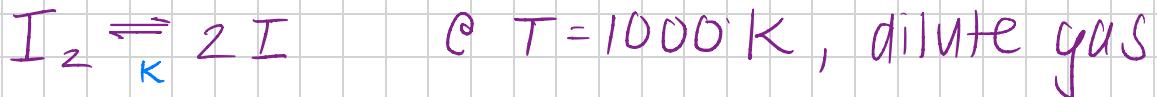
$$+ \frac{\delta F}{\delta N_B} = 0 = \frac{\delta F}{\delta N_A}$$

chemical equilibrium
implies minimization
of F subject to
 $N_A + N_B = N$

(3)

3. (15 points) (4.25 from Chandler) Use the information compiled in Chapter 8 of Hill's *Introduction to Statistical Thermodynamics* to calculate the equilibrium constant, K , for the reaction $I_2 \rightleftharpoons 2I$ when the reaction occurs in the dilute gas phase at $T = 1000\text{ K}$. (Note: For the electronic partition function, you need to consider the difference in degeneracies of the ground states of the iodine atom and molecule, $g_{e,I} = 4, g_{e,I_2} = 1$. This is based on the molecular orbital theory. You are not required to show this.)

calculate equilibrium constant



$$K = \frac{(N_I/V)^2}{(N_{I_2}/V)} = \frac{N_I^2}{N_{I_2} \cdot V}$$

at equil. $\mu_I = \mu_{I_2}$

$$F = -kT \ln Q, \quad dF = -SdT - PdV + \mu dN$$

$$\mu_I = \frac{\delta F}{\delta N_I} = -\frac{\delta \ln Q^I}{\delta \beta N_I}, \quad \mu_{I_2} = \frac{\delta F}{\delta N_{I_2}} = -\frac{\delta \ln Q^{I_2}}{\delta \beta N_{I_2}}$$

$$Q_I = \frac{1}{N_I!} q_I^{N_I}, \quad Q_{I_2} = \frac{1}{N_{I_2}!} q_{I_2}^{N_{I_2}}$$

$$\left(\frac{q_I}{N_I}\right)^2 = \left(\frac{q_{I_2}}{N_{I_2}}\right) \text{ online notes #9}$$

$$q_i = q + r \underbrace{q_{Vi} q_{ro} q_{ne}}_{\text{diatomic only}} \rightarrow \text{no nuclear partition function b/c same}$$

$$q_{ro} = \sum_{j=0}^{\infty} (2j+1) e^{-j(j+1)\theta_{ro}/T}$$

$$\theta_{ro}/T \ll 1 \rightarrow \text{approx w/ integral}$$

$$\rightarrow q_{ro} = T/\sigma \theta_{ro}, \quad \sigma = 2 \text{ (same molecules)}$$

$$\rightarrow q_{Vi} = \frac{1}{1 - e^{-\Theta_{Vi}/T}}$$

$$\rightarrow q_{Tr} = V/\Lambda^3$$

$$q_{n,e} = w_{00} e^{-\beta E_{00}} \quad \begin{cases} q_{ge,I} = 4 \\ q_{ge,I_2} = 1 \end{cases}$$

$$= w_{0,eI} e^{-\beta E_{00}}$$

$$\begin{cases} E_{00,F} = 0 \\ E_{00,I_2} = -D_0 \end{cases}$$

Hill

$$D_0 = 1.54 \text{ eV} \left(\frac{1.6 \cdot 10^{-19} \text{ J}}{\text{eV}} \right)$$

$$q_{n,e} = g_{ei} e^{-\beta E_{00}}$$

$$\begin{cases} q_{n,eI} = g_{eI} \\ q_{n,eI_2} = g_{eI_2} e^{-\beta E_{00}} \end{cases}$$

$$\frac{g_I^2}{N_I^2} = \frac{g_I^2}{KN_{I_2}V} = \frac{g_{I_2}}{N_{I_2}}$$

$$\frac{V^2}{\Lambda_I^6} \cdot 16 \cdot \frac{1}{K} = \frac{V}{\Lambda_{I_2}^3} \cdot 1 \cdot e^{-\beta E_{00}} \cdot \left(T/2\theta_{ro} \right)^2 \cdot \left[\frac{1}{1 - e^{-\Theta_{Vi}/T}} \right]^2$$

$$K = \frac{16}{\Lambda_I^6} e^{+\beta D_0} \left(T/2\theta_{ro} \right)^2 \left[\frac{1}{1 - e^{-\Theta_{Vi}/T}} \right]^2 \cdot \Lambda_{I_2}^3$$

for I_2 : $\Theta_{Vi} = 310 \text{ K}$, $\theta_{ro} = 0.054 \text{ K}$ Hill

$$\Lambda^3 = \left(\frac{\hbar^2}{2mkT} \right)^{3/2} \quad [\text{m}^3] \quad n_I = \frac{(0.1269 \text{ kg/mol})}{(6.023 \cdot 10^{23} \text{ atoms/mol})}$$

$$\Lambda_n^3 \quad [\text{atom/m}^3] \quad \leftarrow \quad n_{I_2} = 2n_I$$

$$K = 2.43 \cdot 10^{22} \quad [\text{atoms/m}^3]$$

$$\Rightarrow K = 0.04 \text{ mol/m}^3$$

(4)

4. (50 pts.) Consider a one-dimensional monatomic crystal of N atoms with equilibrium nearest neighbor spacing a . In order to minimize boundary effects, assume a periodic boundary condition, i.e., $x_{N+1} = x_1$ where x_i is the position of the i th atom. If the atoms interact only with nearest neighbors via a potential $u(x_{i+1} - x_i)$,

(i) show that the energy of the crystal can be written in the form

$$H = \frac{m}{2} \sum_{i=1}^N \dot{\xi}_i^2 + \frac{K}{2} \sum_{i=1}^N (\xi_{i+1} - \xi_i)^2 + N u(a) \quad (6)$$

to quadratic order in the displacement $\xi_i \equiv x_i - x_i^{(0)}$, where $x_i^{(0)}$ is the equilibrium position of the i th atom. What is K ?

i) Hamiltonian needs kinetic, potential energy

$$KE = \frac{1}{2} m v^2 = \frac{m}{2} \sum_i^N \left(\frac{dx_i}{dt} \right)^2 = \frac{m}{2} \sum_i^N \dot{\xi}_i^2$$

$$PE = \sum_i^N u_i = \sum_i^N u(x_{i+1} - x_i)$$

$$= \sum_i^N u(x_{i+1} - x_{i+1}^{(0)} + x_{i+1}^{(0)} - x_i + x_i^{(0)} - x_i^{(0)})$$

$$= \sum_i^N u(\xi_{i+1} - \xi_i + a) \rightarrow PE \text{ in } \xi_{i+1} - \xi_i$$

$$u(\xi_{i+1} - \xi_i + a) = u(a) + (\xi_{i+1} - \xi_i) u'(a) \\ + \frac{1}{2} (\xi_{i+1} - \xi_i)^2 u''(a) + \dots$$

→ periodic nature of displacement, $u'(a)$ vanishes on avg.

$$H = KE + PE$$

$$k = u''(a)$$

$$H = \frac{m}{2} \sum_{i=0}^N \dot{\xi}_i^2 + \frac{k}{2} \sum_{i=0}^N (\xi_{i+1} - \xi_i)^2 + Nu(a)$$

ii)

(ii) Now decompose configurations of atoms into normal modes.

(a) Define normal modes η_k such that ξ_j is a linear superposition of η_k

$$\xi_j = \frac{1}{\sqrt{2N}} \sum_k \eta_k e^{i(jak)}. \quad (7)$$

Show that the periodic boundary condition leads to $k = \frac{2\pi n}{Na}$ where n is any integer.Show further that adding $\frac{2\pi}{a}$ to k does not change ξ_j . Therefore there are only N independent modes. We choose $n \in [-\frac{N}{2}, \frac{N}{2} - 1]$ (assuming N even).

(a)

period BC $\xi_1 = \xi_{N+1}$

$$\xi_1 = \cancel{\frac{1}{\sqrt{2N}}} \sum_k \eta_k e^{iak} = \cancel{\frac{1}{\sqrt{2N}}} \sum_k \eta_k e^{ia k(N+1)} = \xi_{N+1}$$

$$e^{iak} = e^{inak} e^{iak}$$

$$e^{ix} = \cos(x) + i \sin(x)$$

$$1 = e^{inak} = \cos(akN) + i \sin(akN) \quad 0, \text{ can never be real}$$

$$\cos(2\pi n) = 1$$

$$1 = \cos(akN)$$

$$2\pi n = Nak$$

$$k = \frac{2\pi n}{Na}$$

$$\xi_1 = \frac{1}{\sqrt{2N}} \sum_k \eta_k e^{i(ja(k + \frac{2\pi n}{a}))} = \frac{1}{\sqrt{2N}} \sum_k \eta_k e^{ijak} e^{ija \frac{2\pi n}{a}}$$

$$\xi_1 \text{ unchanged: } e^{ij2\pi} = 1 \quad \text{for any int.}$$

$$e^{ij2\pi} = \cos(2\pi j) + i \sin(2\pi j) = 1$$

$\xi_1 \text{ unchanged}$

(b) Show that the fact that the ξ_j 's are real leads to $\eta_k^* = \eta_{-k}$, where η_k^* is the complex conjugate of η_k .

(b) ξ_i is real, $\xi_i = \xi_i^*$

$$\xi_i^* = \frac{1}{\sqrt{2N}} \sum_k \eta_k^* e^{-ijk\alpha k}$$

$$\sum_k \eta_k^* e^{-ijk\alpha k} = \sum_{-k} \eta_{-k} e^{ijk\alpha k}$$

inspection ↪

$$\eta_k^* = \eta_{-k}$$

(iii) It can be shown that the normal mode coordinates η_k diagonalize the Hamiltonian:

$$\sum_j (\xi_{j+1} - \xi_j)^2 = \sum_{k>0} [(\eta_k^R)^2 + (\eta_k^I)^2] 4 \sin^2\left(\frac{1}{2}ka\right) \quad (8)$$

and

$$\sum_j \dot{\xi}_j^2 = \sum_{k>0} [(\dot{\eta}_k^R)^2 + (\dot{\eta}_k^I)^2], \quad (9)$$

where η_k^R and η_k^I are the real and imaginary parts of η_k . What is the frequency for each normal mode ω_k ? What is the speed of sound for this model? (The speed of sound is defined as $\frac{d\omega_k}{dk}|_{k=0}$).

$$\frac{\delta H}{\delta \bar{x}} = -m\bar{x} \rightarrow \frac{\delta H}{\delta \bar{x}_j} = -m\bar{x}_j \rightarrow \frac{\delta H}{\delta \bar{\xi}_j} = -m\bar{\xi}_j$$

$$H = \frac{m}{2} \sum_{i=0}^n \xi_i^2 + \frac{k}{2} \sum_{i=0}^n (\xi_{i+1} - \xi_i)^2 + N u(a)$$

$$\frac{\delta H}{\delta \xi_j} = m \sum_{i=0}^n \xi_i \frac{\delta \xi_i}{\delta \xi_j} \xrightarrow{0} -k(\xi_{j+1} - \xi_j) + k(\xi_j - \xi_{j-1})$$

$$\frac{\delta H}{\delta \bar{\xi}_j} = -m\bar{\xi}_j = k(2\xi_j - \xi_{j+1} - \xi_{j-1})$$

$$-m \sum_{k \in \mathbb{Z}^N} \bar{\eta}_k e^{ijk\alpha k} = \frac{k}{\sqrt{N}} \sum_k \eta_k (2 - e^{i\alpha k} - e^{-i\alpha k}) e^{ijk\alpha k}$$

\downarrow
 \cos

$$\bar{\eta}_k = -\frac{2k}{m} \eta_k (1 - \cos(\alpha k))$$

$$1 - \cos(2x) = 2\sin^2(x)$$

$$\bar{\eta}_k = -\frac{4k}{m} \eta_k \sin^2(\alpha k / 2)$$

$$\bar{\eta}_k = g(k) \eta_k, \quad g(k) = -\frac{4k}{m} \sin^2(\frac{\alpha k}{2})$$

guess: $\eta_k = A e^{i\omega_k t}$

$$\bar{\eta}_k = -A \omega_k^2 e^{i\omega_k t} = -\omega_k^2 \eta_k$$

$$g(k) = -\omega_k^2 \rightarrow \omega_k = 2\sqrt{\frac{k}{m}} \sin(\frac{\alpha k}{2})$$

Speed of sound: $\left. \frac{d\omega_k}{dk} \right|_{k=0} = c$

$$c = \left. (km^2)^{-1/2} \sin(\frac{\alpha k}{2}) + 2\sqrt{\frac{k}{m}} \cdot \frac{q}{2} \cos(\frac{\alpha k}{2}) \right|_{k=0}$$

$$\left. \frac{d\omega_k}{dk} \right|_{k=0} = a\sqrt{\frac{k}{m}}$$

- (iv) Show that in a large 1-D solid, the density (or degeneracy) of normal modes with the frequencies between ω and $\omega + d\omega$ is

$$g(\omega)d\omega = \frac{2N}{\pi\omega_m\sqrt{1-(\omega/\omega_m)^2}}d\omega, \quad (10)$$

where ω_m is the maximum frequency of normal modes.

iv)

$$\omega_k = 2\sqrt{\frac{k}{m}} \sin(\frac{ak}{2}), \quad k = \frac{2\pi n}{Na}$$

$$\omega_k = 2\sqrt{\frac{k}{m}} \sin(\frac{\pi n}{N})$$

$$\omega_{\max} \rightarrow \sin(\frac{\pi n}{N}) = 1 \rightarrow 2\sqrt{\frac{k}{m}} = \omega_{\max}$$

max possible value

$$\frac{dw}{dk} = \frac{dw}{dn} \cdot \frac{dn}{dk} \rightarrow dn \cdot \left[\frac{dw}{dk} \right] = \left[\frac{dn}{dk} \right] dw$$

$$dn = \left[\frac{dn}{dk} \right] \left[\frac{dw}{dk} \right]^{-1} dw \rightarrow \frac{dw}{dn} = f(w)$$

$\underbrace{f^{-1}(w)}$

$$f(w) = \frac{dw}{dn} = \left(2\sqrt{\frac{k}{m}} \right) \cdot \frac{\pi}{N} \cos(\frac{\pi n}{N})$$

$$f(w) = \omega_{\max} \frac{\pi}{N} \cos(\frac{\pi n}{N})$$

$$\omega = \omega_{\max} \sin(\frac{\pi n}{N})$$

positive
negative
 \nearrow
 \searrow
 $k's$

$$\int_0^{\omega_{\max}} f(w) dw = \int_0^{2N} dn = \frac{2N}{2} = N$$

- (i) Show degeneracy of states in 1-D solid
(ii) Define Fermi energy and show that it is a linear approximation of E_F
- $$E_F = \frac{1}{2} \sum_{k=1}^N \omega_k \sin^2(k\pi/N)$$
- Show that the periodic boundary condition leads to $k = \frac{N}{2}$ which is an integer. Thus both the sum above $\sum_{k=1}^N$ does not change. Therefore there are only $N/2$ independent modes. We denote them by $\omega_1, \omega_2, \dots, \omega_{N/2}$.
- (iii) Show that the L is even and looks like $\omega_k = \omega_0 - q_k$, where q_k is the complete conjugate of q_0 .
(iv) Now we show that the density of states is given by
- $$\sum_{k=1}^N \delta(\omega - \omega_k) = \sum_{k=1}^N \left[(\omega_k^2 + \omega_0^2)^{1/2} \sin^2(k\pi/N) \right]$$
- and
- $$\sum_{k=1}^N \delta(\omega - \omega_k) = \sum_{k=1}^N \left[(\omega_k^2 + \omega_0^2)^{1/2} \right]$$
- where ω_k and ω_0 are the real and imaginary parts of ω_k . What is the frequency for each normal mode ω_k ? What is the frequency for each normal mode ω_0 ?
- (v) Show that in a large 1-D solid, the density (or degeneracy) of normal modes with the frequencies between ω and $\omega + d\omega$ is
- $$g(\omega)d\omega = \frac{2N}{\pi\omega_m\sqrt{1-(\omega/\omega_m)^2}}d\omega$$
- where ω_m is the maximum frequency of normal modes.
- (vi) Compute the Debye model for 1-D solid, what is its Debye temperature θ_D ? Why is the Debye model not able to explain the observed behavior of thermal conductivity?
- (vii) Show that one can use the same expression of the heat capacity as important at very low temperature and the same expression of the thermal conductivity as important at high temperature. The Debye model gives an accurate result even though explicit normal modes are considered.

$$g(\omega) = f^{-1}(\omega) = \frac{2N}{\omega_{\max} \pi \cos(\pi n/N)}$$

$$\cos(\pi n/N) = \sqrt{1 - \sin^2(\pi n/N)}$$

\uparrow
 $(\omega/\omega_{\max})^2$

$$g(\omega) \cdot d\omega = \frac{2N}{\omega_{\max} \pi \sqrt{1 - (\omega/\omega_{\max})^2}} d\omega$$

$$\omega_m(k_{\max}), k_{\max}(n_{\max} = \frac{N}{2} - 1)$$

$$k_{\max} = \frac{2\pi}{Na} \left(\frac{N}{2} - 1\right)$$

$$\frac{\pi}{a} - \frac{2\pi}{Na} = \frac{(N-2)\pi}{Na} \rightarrow \frac{\pi}{a}$$

$$\omega_m = \sqrt{\frac{4k}{m}} \sin\left(\frac{(N-2)\pi}{2N}\right) \rightarrow \sqrt{\frac{4k}{m}} \sin\left(\frac{\pi}{2}\right) = \sqrt{\frac{4k}{m}} = \omega_m$$

$$\frac{dw}{dk} = \left(\sqrt{\frac{4k}{m}}\right) \left(\frac{1}{2}\right) k^{-1/2} \sin\left(\frac{ka}{2}\right) + \frac{q}{2} \cos\left(\frac{ka}{2}\right) \left(\frac{4k}{a}\right)^{1/2}$$

$$w = \underbrace{\sqrt{\frac{4k}{m}} \left(\sin \left(\dots \right) \right)}_{w_m}, \quad \frac{w}{w_m} = \sin \left(\dots \right)$$

$$\frac{dw}{dk} = \frac{1}{2} \cdot \frac{1}{k} w + \frac{q}{2} \cos\left(\frac{ka}{2}\right) \left(\frac{4k}{a}\right)^{1/2}$$

$$\cos = \sqrt{\cos^2} = \sqrt{1 - \sin^2} = \sqrt{1 - \sin^2\left(\frac{ka}{2}\right)}$$