

#### Equação da Quantidade de Movimento



Sabemos a equação da quantidade de movimento na sua forma integral.

$$\vec{F} = \vec{F}_s + \vec{F}_B = \frac{\partial}{\partial t} \int_{VC} \vec{V} \rho d \nabla + \int_{SC} \vec{V} \rho \vec{V} d \vec{A}$$

Na forma diferencial expressamos as equações para um sistema infinitesimal de massa *dm*, para a qual a segunda lei de Newton pode ser expressa como:

$$d\vec{F} = dm \frac{d\vec{V}}{dt}$$

o termo *dm* é facilmente determinado pelo produto entre a massa específica do fluido dentro do V.C. e o volume diferencial.

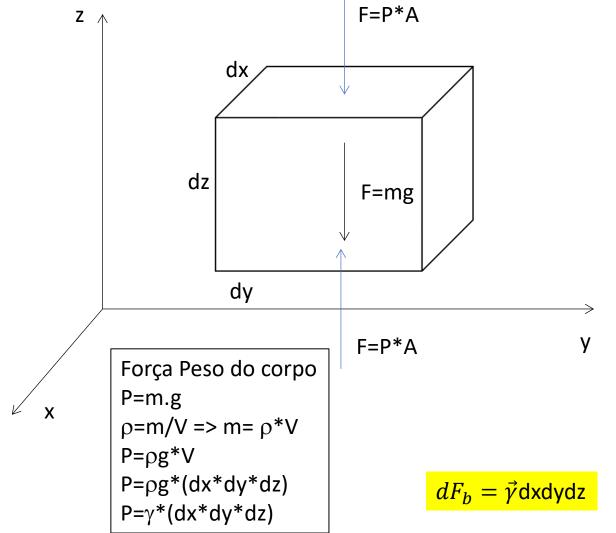
$$\left(\frac{d\vec{V}}{dt}\right)_{\text{sistema}} = \frac{D\vec{V}}{Dt}$$



#### Equação da Quantidade de Movimento



#### Força Campo (peso do Corpo)



$$\rho = \frac{m}{V} \implies m = \rho * V$$

$$F = Força Peso = P$$

$$m\frac{d\vec{V}}{dt} = \rho*(\mathsf{dx}*\mathsf{dy}*\mathsf{dz})*\mathsf{g}$$

$$\rho dx dy dz \frac{d\vec{v}}{dt} = \rho * (dx * dy * dz) * g$$

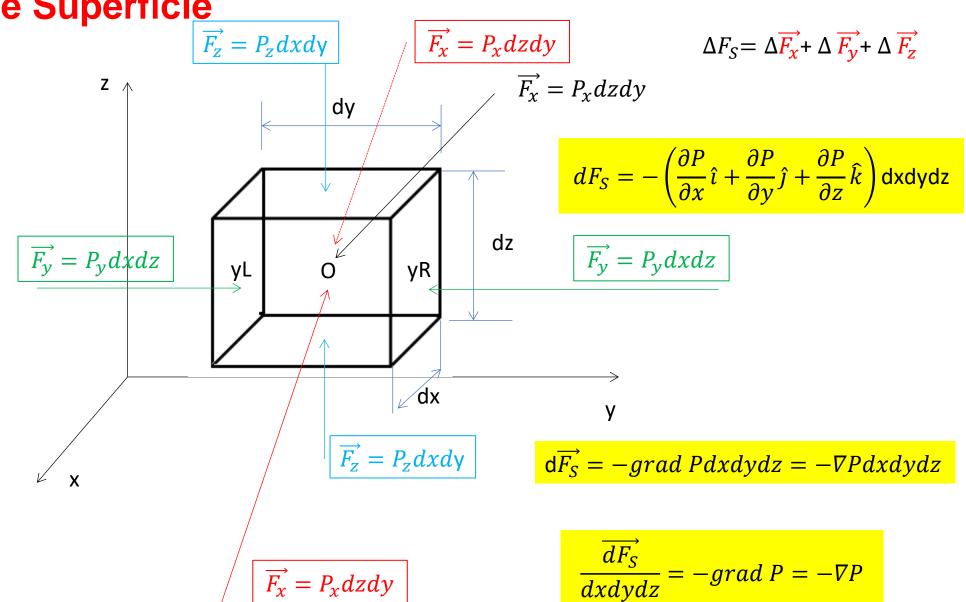
$$\frac{d\vec{V}}{dt} = \vec{g}$$



#### Equação da Quantidade de Movimento







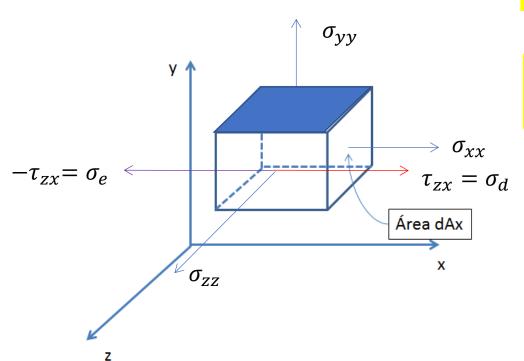


# FÍSICA DOS FLUÍDOS



#### Tensões normais e tangenciais num elemento de fluido

#### Tensões de Superfície num Elemento de Fluido



$$\Delta F_{SX} = \left(\frac{\partial \sigma_{XX}}{\partial x} + \frac{\partial \tau_{YX}}{\partial y} + \frac{\partial \tau_{ZX}}{\partial z}\right) \Delta x \Delta y \Delta z$$

$$\Delta F_{sy} = \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right) \Delta x \Delta y \Delta z$$

$$\tau_{zx} = \sigma_d$$

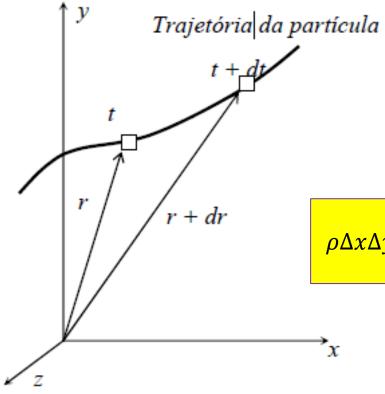
$$\Delta F_{sy} = \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right) \Delta x \Delta y \Delta z$$





$$\rho = \frac{m}{V} \implies m = \rho * V$$

$$\Delta \vec{F} = \Delta m \vec{a} = \rho \Delta x \Delta y \Delta z \vec{a}$$



$$\Delta \vec{F} = \Delta m \vec{a} = \rho \Delta x \Delta y \Delta z \frac{d\vec{V}}{dt}$$

$$\vec{a}_p = \frac{D\vec{V}}{Dt} = u\frac{\partial\vec{V}}{\partial x} + v\frac{\partial\vec{V}}{\partial y} + w\frac{\partial\vec{V}}{\partial z} + \frac{\partial\vec{V}}{\partial t}$$

$$\rho \Delta x \Delta y \Delta z \frac{D\vec{V}}{Dt} = \left( u \frac{\partial \vec{V}}{\partial x} + v \frac{\partial \vec{V}}{\partial y} + w \frac{\partial \vec{V}}{\partial z} + \frac{\partial \vec{V}}{\partial t} \right) \rho \Delta x \Delta y \Delta z$$





$$\vec{F}_{total} = \vec{F}_{press\~ao} + \vec{F}_{gravidade} + \vec{F}_{superficie} + \dots + \vec{F}_{n}$$

$$\vec{F}_{press\~ao} = -grad \ Pdxdydz = -\vec{\nabla}P\Delta x\Delta y\Delta z$$

$$\vec{F}_{gravidade} = 
ho dx dy dz \frac{d\vec{V}}{dt} = 
ho \vec{g} * (\Delta x * \Delta y * \Delta z)$$

$$\vec{F}_{superficie} = \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right) \Delta x \Delta y \Delta z$$

$$+ \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right) \Delta x \Delta y \Delta z + \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right) \Delta x \Delta y \Delta z$$

$$\vec{F}_{total} = \left( u \frac{\partial \vec{V}}{\partial x} + v \frac{\partial \vec{V}}{\partial y} + w \frac{\partial \vec{V}}{\partial z} + \frac{\partial \vec{V}}{\partial t} \right) \rho \Delta x \Delta y \Delta z$$

$$\rho = \frac{m}{V} \implies m = \rho * V$$

$$\vec{F}_{total} = \vec{F}_{press\~ao} + \vec{F}_{gravidade} + \vec{F}_{superficie} + \dots + \vec{F}_{n}$$





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$$\left( u \frac{\partial \vec{V}}{\partial x} + v \frac{\partial \vec{V}}{\partial y} + w \frac{\partial \vec{V}}{\partial z} + \frac{\partial \vec{V}}{\partial t} \right) \rho \Delta x \Delta y \Delta z$$

$$= -\vec{\nabla} P \Delta x \Delta y \Delta z + \rho \vec{g} * (\Delta x * \Delta y * \Delta z) + \left( \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \right) \Delta x \Delta y \Delta z$$

$$+ \left( \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \right) \Delta x \Delta y \Delta z + \left( \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right) \Delta x \Delta y \Delta z$$

$$\left(u\frac{\partial \vec{V}}{\partial x} + v\frac{\partial \vec{V}}{\partial y} + w\frac{\partial \vec{V}}{\partial z} + \frac{\partial \vec{V}}{\partial t}\right)\rho$$

$$= -\vec{\nabla}P + \rho\vec{g} + \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right) + \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right) + \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right)$$





#### Separando nas componentes x,y,z

$$\left(u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z} + \frac{\partial u}{\partial t}\right)\rho = -\frac{\partial P}{\partial x} + \rho \overrightarrow{g_x} + \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right)$$

$$\left(u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z} + \frac{\partial v}{\partial t}\right)\rho = -\frac{\partial P}{\partial y} + \rho \overrightarrow{g_y} + \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right)$$

$$\left(u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z} + \frac{\partial w}{\partial t}\right)\rho = -\frac{\partial P}{\partial z} + \rho \overrightarrow{g_z} + \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right)$$





$$\left(\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z}\right)\rho = -\frac{\partial P}{\partial x} + \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right)$$

$$\left(\frac{\partial v}{\partial t} + u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z}\right)\rho = -\frac{\partial P}{\partial y} + \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right)$$

$$\left(\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}\right)\rho = -\frac{\partial P}{\partial z} + \rho \overrightarrow{g_z} + \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right)$$

Considerando que as variações na densidade é pequena, a densidade  $\rho$  pode ser substituído em todos os lugares pela constante (referência) densidade  $\rho_0$ , com uma grande exceção.

Esta exceção é o termo gravitacional, através do qual, mesmo uma pequena variação da densidade pode causar um efeito de flutuação importante na maioria das aplicações ambientais.

A simplificação resultante é conhecida como a aproximação Boussinesq.





$$\left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z}\right) \rho_0 = -\frac{\partial P}{\partial x} + \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right)$$

$$\left(\frac{\partial v}{\partial t} + u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z}\right)\rho_0 = -\frac{\partial P}{\partial y} + \left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right)$$

$$\left(\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}\right)\rho_0 = -\frac{\partial P}{\partial z} + \rho \overrightarrow{g}_z + \left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right)$$

Dividindo por  $\rho_0$ 

$$\left(\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial x} + \frac{1}{\rho_0}\left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}\right)$$

$$\left(\frac{\partial v}{\partial t} + u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial y} + \frac{1}{\rho_0}\left(\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z}\right)$$

$$\left(\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial z} + \frac{\rho}{\rho_0}\overrightarrow{g_z} + \frac{1}{\rho_0}\left(\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}\right)$$





$$\left(\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} + w\frac{\partial u}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial x} + \vec{\tau}_x$$

$$\left(\frac{\partial v}{\partial t} + u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} + w\frac{\partial v}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial y} + \vec{\tau}_y$$

$$\left(\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + v\frac{\partial w}{\partial y} + w\frac{\partial w}{\partial z}\right) = -\frac{1}{\rho_0}\frac{\partial P}{\partial z} + \frac{\rho}{\rho_0}\overrightarrow{g_z} + \overrightarrow{\tau}_z$$

Na Forma Vetorial

$$\frac{\partial \vec{V}}{\partial t} + (\vec{V}.\nabla)\vec{V} = -\frac{1}{\rho_0}\vec{\nabla}P + \frac{\rho}{\rho_0}\vec{g} + \vec{F}_{tens\~ao}$$



## CINEMATION BESUM FILE SITURES FLUIDO



Para *fluidos newtonianos*, as tensões podem ser expressas em termos de **gradientes de velocidades** e propriedades dos fluidos:

Sabemos que p/ um fluido newtoniano a tensão viscosa é proporcional a taxa de deformação por cisalhamento (taxa de deformação angular). As tensões podem ser expressas em termos dos gradientes de velocidade e das propriedades dos fluidos como:

$$\begin{split} \tau_{xy} &= \tau_{yx} = \mu \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right) & \sigma_{xx} = -p - \frac{2}{3} \mu \nabla \vec{V} + 2\mu \frac{\partial u}{\partial x} \\ \tau_{yz} &= \tau_{zy} = \mu \left( \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \right) & \sigma_{yy} = -p - \frac{2}{3} \mu \nabla \vec{V} + 2\mu \frac{\partial v}{\partial y} \\ \tau_{zx} &= \tau_{xz} = \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) & \sigma_{zz} = -p - \frac{2}{3} \mu \nabla \vec{V} + 2\mu \frac{\partial w}{\partial z} \end{split}$$

Onde *p* é a pressão termodinâmica local.



#### Equações de Navier Stokes



No caso de *fluido incompressível*,  $\nabla \cdot V = 0$ , e a equação acima pode ser simplificada. Fazendo desprezíveis as forças de campo (B = 0) se obtém **as** *Equações de Navier Stokes*.

$$\begin{split} \rho \frac{Du}{Dt} &= -\frac{\partial p}{\partial x} + \frac{\partial}{\partial x} \left( 2\mu \frac{\partial u}{\partial x} - \frac{2}{3}\mu \nabla \vec{V} \right) + \frac{\partial}{\partial y} \left[ \mu \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial z} \left[ \mu \left( \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right) \right] \\ \rho \frac{Dv}{Dt} &= -\frac{\partial p}{\partial y} + \frac{\partial}{\partial x} \left[ \mu \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial y} \left( 2\mu \frac{\partial v}{\partial y} - \frac{2}{3}\mu \nabla \vec{V} \right) + \frac{\partial}{\partial z} \left[ \mu \left( \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \right) \right] \\ \rho \frac{Dw}{Dt} &= -\frac{\partial p}{\partial z} + \frac{\partial}{\partial x} \left[ \mu \left( \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right) \right] + \frac{\partial}{\partial y} \left[ \mu \left( \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \right] + \frac{\partial}{\partial z} \left( 2\mu \frac{\partial w}{\partial z} - \frac{2}{3}\mu \nabla \vec{V} \right) \end{split}$$





no caso de **escoamento incompressível permanente** com viscosidade constante e incluindo as forças de campo

$$\rho \frac{Du}{Dt} = \rho g_x - \frac{\partial p}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right)$$

$$\rho \frac{Dv}{Dt} = \rho g_y - \frac{\partial p}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right)$$

$$\rho \frac{Dw}{Dt} = \rho g_z - \frac{\partial p}{\partial z} + \mu \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right)$$

Em forma vetorial pode ser representada como

$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \nabla p + \mu \nabla^2 \vec{V}$$



#### Equações de Euler



Quando os termos viscosos são pequenos e podem ser desprezíveis ( $\mu$ =0) as equações resultantes são conhecidas como **Equações de Euler**, que podem ser representas na forma vetorial como:

$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \nabla p$$