

## Chapter 7 可压缩流动：一些预备知识

### 7.1 引言P516

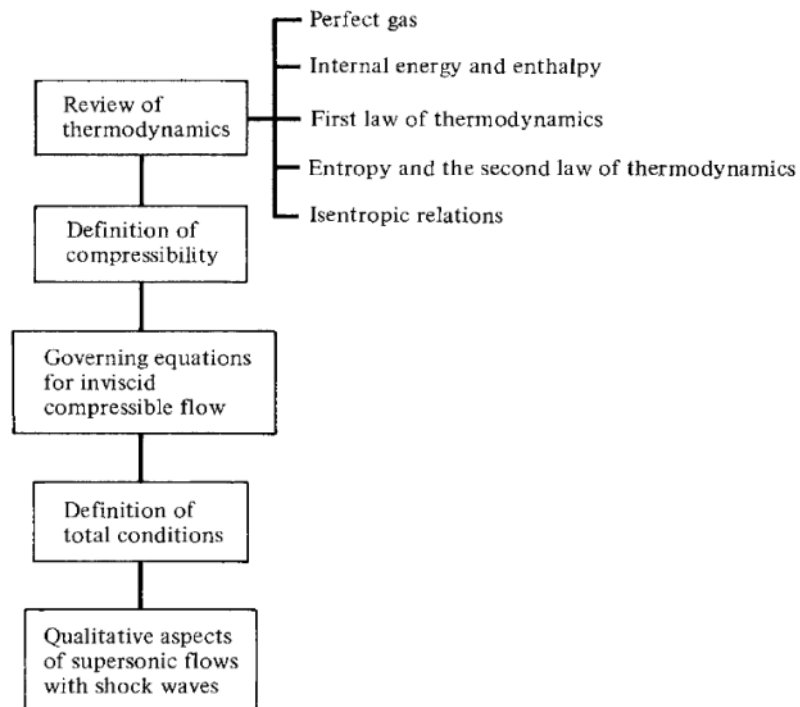


Figure 7.1 Road map for Chapter 7.

### 7.2 热力学的简单回顾P518

#### 7.2.1 完全气体P518

#### 7.2.2 内能和焓P518

### **7.2.3 热力学第一定律P523**

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### **7.2.5 等熵关系式P526**

## **7.3 压缩性的定义P530**

## **7.4 无粘可压缩流动的控制方程P531**

## 7.5 总（滞止）状态的定义P533

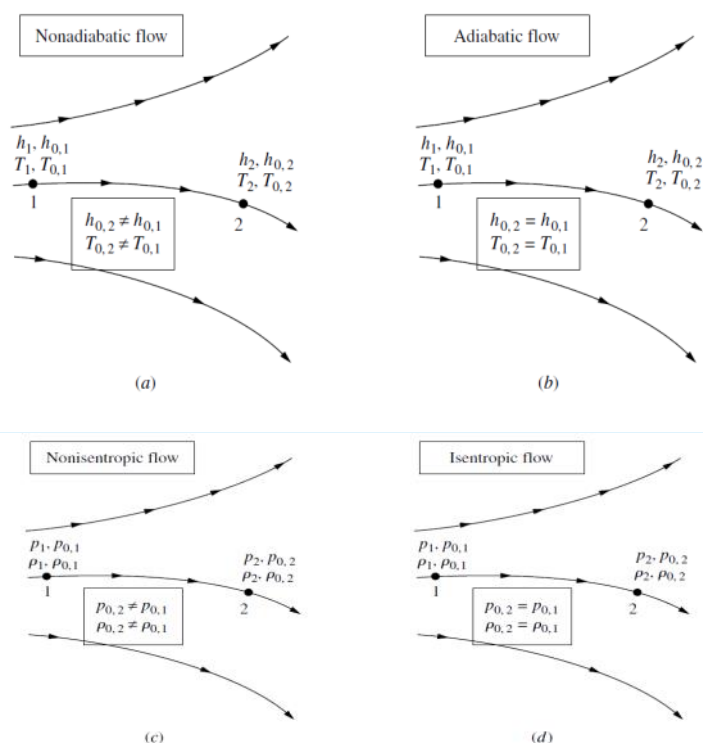


Figure 7.4 Comparisons between (a) nonadiabatic, (b) adiabatic, (c) nonisentropic, and (d) isentropic flows.

## 7.6 超声速流的一些知识：激波P540

## 7.7 总结P544

*Thermodynamic relations:*

Equation of state:  $p = \rho RT$  (7.1)

For a calorically perfect gas,

$$e = c_v T \quad \text{and} \quad h = c_p T \quad (7.6a \text{ and } b)$$

$$c_p = \frac{\gamma R}{\gamma - 1} \quad (7.9)$$

$$c_v = \frac{R}{\gamma - 1} \quad (7.10)$$

Forms of the first law:

$$\delta q + \delta w = de \quad (7.11)$$

$$T ds = de + p dv \quad (7.18)$$

$$T ds = dh - v dp \quad (7.20)$$

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Definition of entropy:

$$ds = \frac{\delta q_{\text{rev}}}{T} \quad (7.13)$$

Also

$$ds = \frac{\delta q}{T} + ds_{\text{irrev}} \quad (7.14)$$

The second law:

$$ds \geq \frac{\delta q}{T} \quad (7.16)$$

or, for an adiabatic process,

$$ds \geq 0 \quad (7.17)$$

Entropy changes can be calculated from (for a calorically perfect gas)

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1} \quad (7.25)$$

and

$$s_2 - s_1 = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1} \quad (7.26)$$

For an isentropic flow,

$$\frac{p_2}{p_1} = \left( \frac{\rho_2}{\rho_1} \right)^\gamma = \left( \frac{T_2}{T_1} \right)^{\gamma/(\gamma-1)} \quad (7.32)$$

General definition of compressibility:

$$\tau = -\frac{1}{v} \frac{dv}{dp} \quad (7.33)$$

For an isothermal process,

$$\tau_T = -\frac{1}{v} \left( \frac{\partial v}{\partial p} \right)_T \quad (7.34)$$

For an isentropic process,

$$\tau_s = -\frac{1}{v} \left( \frac{\partial v}{\partial p} \right)_s \quad (7.35)$$

The governing equations for inviscid, compressible flow are

*Continuity:*

$$\frac{\partial}{\partial t} \iiint_V \rho dV + \iint_S \rho \mathbf{V} \cdot d\mathbf{S} = 0 \quad (7.39)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{V} = 0 \quad (7.40)$$

*Momentum:*

$$\frac{\partial}{\partial t} \rho \mathbf{V} + \nabla \cdot (\rho \mathbf{V} \mathbf{V}) = -\nabla p + \nabla \cdot \boldsymbol{\tau}$$

*Momentum:*

$$\frac{\partial}{\partial t} \iiint_V \rho \mathbf{V} dV + \iint_S (\rho \mathbf{V} \cdot \mathbf{dS}) \mathbf{V} = - \iint_S p \mathbf{dS} + \iiint_V \rho \mathbf{f} dV \quad (7.41)$$

$$\rho \frac{Du}{Dt} = - \frac{\partial p}{\partial x} + \rho f_x \quad (7.42a)$$

$$\rho \frac{Dv}{Dt} = - \frac{\partial p}{\partial y} + \rho f_y \quad (7.42b)$$

$$\rho \frac{Dw}{Dt} = - \frac{\partial p}{\partial z} + \rho f_z \quad (7.42c)$$

*Energy:*

$$\frac{\partial}{\partial t} \iiint_V \rho \left( e + \frac{V^2}{2} \right) dV + \iint_S \rho \left( e + \frac{V^2}{2} \right) \mathbf{V} \cdot \mathbf{dS} \quad (continued)$$

$$= \iiint_V \dot{q} \rho dV - \iint_S p \mathbf{V} \cdot \mathbf{dS} + \iiint_V \rho (\mathbf{f} \cdot \mathbf{V}) dV \quad (7.43)$$

$$\rho \frac{D(e + V^2/2)}{Dt} = \rho \dot{q} - \nabla \cdot p \mathbf{V} + \rho (\mathbf{f} \cdot \mathbf{V}) \quad (7.44)$$

If the flow is steady and adiabatic, Equations (7.43) and (7.44) can be replaced by

$$h_0 = h + \frac{V^2}{2} = \text{const}$$

*Equation of state* (perfect gas):

$$p = \rho R T \quad (7.1)$$

*Internal energy* (calorically perfect gas):

$$e = c_v T \quad (7.6a)$$

Total temperature  $T_0$  and total enthalpy  $h_0$  are defined as the properties that would exist if (in our imagination) we slowed the fluid element at a point in the flow to zero velocity adiabatically. Similarly, total pressure  $p_0$  and total density  $\rho_0$  are defined as the properties that would exist if (in our imagination) we slowed the fluid element at a point in the flow to zero velocity isentropically. If a general flow field is adiabatic,  $h_0$  is constant throughout the flow; in contrast, if the flow field is nonadiabatic,  $h_0$  varies from one point to another. Similarly, if a general flow field is isentropic,  $p_0$  and  $\rho_0$  are constant throughout the flow; in contrast, if the flow field is nonisentropic,  $p_0$  and  $\rho_0$  vary from one point to another.

Shock waves are very thin regions in a supersonic flow across which the pressure, density, temperature, and entropy increase; the Mach number, flow velocity, and total pressure decrease; and the total enthalpy stays the same.

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## 7.8 作业题P547