

# ALGEBRAIC CURVES AND INTEGRABLE HIERARCHIES

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**ABSTRACT.** In the 1990s, a remarkable correspondence was discovered between the geometry of algebraic curves and infinite-dimensional systems of differential equations. The correspondence has its origin in the study of two-dimensional quantum field theories and is related to many different areas of mathematics. After introducing the relevant objects, I will then state the first result in this story, which was conjectured by Witten and proved by Kontsevich.

## 1. ALGEBRAIC CURVES AND THEIR MODULI

For the purposes of this lecture, we will only define smooth curves. Also, space will really mean algebraic variety/scheme/stack, whichever is needed to make the statements of the results true.

**Definition 1.1.** A *smooth curve* is a compact 1-dimensional complex manifold; i.e. a Riemann surface.

Note that curves will be 1-dimensional over  $\mathbb{C}$  and 2-dimensional over  $\mathbb{R}$ . We will also need our curves to develop singularities, but fortunately we only need one kind of singularity:

**Definition 1.2.** A *nodal singularity* is one that locally looks like  $\{xy = 0\} \subseteq \mathbb{C}^2$ .

**1.1. Moduli spaces.** As the simplest algebraic varieties, much has been written on the geometry of curves. For this lecture, we will focus on the study of families of curves – a very productive technique in algebraic geometry is to study how properties vary in families. We will begin with a very natural question.

**Question 1.3.** *Can we classify all algebraic curves?*

If we first restrict to smooth curves, we know that these are topologically surfaces, and the topological classification of surfaces is very simple: they are controlled by a discrete parameter, the genus. However, for a surface of genus  $g \geq 1$ , there will be many ways to give it the structure of a complex manifold or algebraic variety.

**Example 1.4.** For any  $\lambda \neq 0, 1$ , there is a genus 1 smooth curve given by the equation

$$y^2 = x(x-1)(x-\lambda),$$

and different  $\lambda$  give non-isomorphic curves.

It is natural to wonder if for a given genus  $g$ , there is a space  $\mathcal{M}_g$  parameterizing curves of genus  $g$ , known as a moduli space. More precisely, for any space  $X$ , we want the set  $\text{Maps}(X, \mathcal{M}_g)$  of maps from  $X$  to  $\mathcal{M}_g$  to correspond to families of smooth curves of genus  $g$  over  $X$ . In fact, such a space exists, and the following is true:

**Proposition 1.5** (Riemann, Deligne–Mumford). *For any  $g, n$  such that  $2g - 2 + n > 0$ , there is a space  $\mathcal{M}_{g,n}$  of dimension  $3g - 3 + n$  parameterizing smooth curves of genus  $g$  with  $n$  distinct marked points.*

Unfortunately,  $\mathcal{M}_{g,n}$  is not compact, which geometrically means that there are families of smooth curves over  $\mathbb{C} \setminus 0$  which cannot be filled in (meaning adding a fiber over 0) with a smooth curve. In order to compactify  $\mathcal{M}_{g,n}$ , we will add nodal curves.

**Proposition 1.6** (Deligne–Mumford). *For any  $g, n$  such that  $2g - 2 + n > 0$ , there exists a compact space  $\overline{\mathcal{M}}_{g,n}$  of dimension  $3g - 3 + n$  parameterizing families of possibly nodal curves  $C$  of genus  $g$  with  $n$  distinct marked points satisfying the following conditions:*

- (1) *The marked points are away from the nodes; i.e. they lie on smooth points of  $C$ .*
- (2) *For any irreducible component  $C'$  of  $C$  with genus  $g'$ , if  $n'$  denotes the number of nodes and marked points on  $C'$ , then  $2g' - 2 + n' > 0$ .*

Now, the identity map  $\overline{\mathcal{M}}_{g,n} \rightarrow \overline{\mathcal{M}}_{g,n}$  corresponds to some family of curves over  $\overline{\mathcal{M}}_{g,n}$ , and this is the *universal family*, which we will denote  $\mathcal{C}_{g,n}$ . This means that for any map  $f: X \rightarrow \overline{\mathcal{M}}_{g,n}$ , the corresponding family is given by  $f^*\mathcal{C}_{g,n}$ .

**1.2. Integrals on the moduli space of curves.** We will define line bundles  $L_i$  for  $i = 1, \dots, n$  on  $\overline{\mathcal{M}}_{g,n}$  as follows. For a point  $[C, x_1, \dots, x_n] \in \overline{\mathcal{M}}_{g,n}$ , we will define

$$(L_i)_{[C, x_1, \dots, x_n]} = T_{x_i} C.$$

More precisely, if  $x_i: \overline{\mathcal{M}}_{g,n} \rightarrow \mathcal{C}_{g,n}$  is the map corresponding to the  $i$ -th marked point, then  $L_i = x_i^* \Omega^1_{\mathcal{C}_{g,n}/\overline{\mathcal{M}}_{g,n}}$ , where  $\Omega^1_{\mathcal{C}_{g,n}/\overline{\mathcal{M}}_{g,n}}$  is the relative cotangent bundle.

Now, we will define  $\psi_i = c_1(L_i) \in H^2(\overline{\mathcal{M}}_{g,n})$  to be the first Chern class of  $L_i$ . These  $\psi$  classes are known as *gravitational descendents* in the physics literature, where the integrals below are related to 2D topological gravity.

**Definition 1.7.** For any integers  $a_1, \dots, a_n \geq 0$ , define

$$\langle \psi_1^{a_1} \dots \psi_n^{a_n} \rangle_{g,n} := \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{a_1} \psi_2^{a_2} \dots \psi_n^{a_n}.$$

## 2. THE KDV HIERARCHY

The behavior of waves in shallow water is described by the *Korteweg-de Vries equation*:

$$\frac{\partial u}{\partial t_1} = u \frac{\partial u}{\partial t_0} + \frac{1}{12} \frac{\partial^3 u}{\partial t_0^3}.$$

The coefficients can be scaled to any nonzero real numbers by scaling the variables. The KdV equation has a solution given by

$$u_1(t_0, t_1) = \frac{c}{2} \operatorname{sech}^2 \left( \frac{\sqrt{c}}{2} (t_0 - ct_1 + \delta) \right),$$

where  $c$  and  $\delta$  are constants and we impose  $u(x, t) = f(t_0 - ct_1)$  to be a travelling wave. In fact, there are infinitely many exact solutions  $u_n(t_0, t_1)$  to the KdV equation, which behave like independent waves for large time (for both  $t_1 \ll 0$  and  $t_1 \gg 0$ ) but can collide in small time. Waves which exhibit this type of behavior are called *solitons*. The infinitely many exact solutions make the KdV equation an *integrable system* with infinitely many degrees of freedom.

The KdV equation generates an infinite-dimensional system of differential equations using the following procedure. Define the *Schrodinger operator* by

$$L := (\partial_0^2 + u).$$

Its square root can be written as

$$(\partial_0^2 + u)^{\frac{1}{2}} = \partial_0 + \frac{1}{2}u\partial_0^{-1} - \frac{1}{4}u_0\partial_0^{-2} + \left( \frac{u_{00}}{8} - \frac{u^2}{8} \right) \partial_0^{-3} + \dots$$

Now for an expression of the form  $M = \sum_{\ell=0}^{\infty} g_{\ell} \partial^{n-\ell}$ , define

$$M_+ := \sum_{\ell=0}^n g_{\ell} \partial^{n-\ell}, \quad M_- := M - M_+.$$

We can now define the operators

$$K_a(u) := -[L, L_+^{\frac{2a+1}{2}}]$$

for nonnegative integers  $a$  and obtain the *KdV hierarchy*

$$\frac{\partial u}{\partial t_a} = K_a(u).$$

The first two equations are in fact

$$\frac{\partial u}{\partial t_0} = \frac{\partial u}{\partial t_0}, \quad \frac{\partial u}{\partial t_1} = \frac{3}{2}u \frac{\partial u}{\partial t_0} + \frac{1}{4} \frac{\partial^3 u}{\partial t_0^3},$$

and we see the first is a tautology and the second is the KdV equation up to scaling. What justifies calling this an integrable hierarchy is the commutation relation

$$\frac{\partial}{\partial x_b} K_a(u) = \frac{\partial}{\partial x_a} K_b(u)$$

for all  $a, b$ . This states that the evolution of the system in the  $t_b$  direction and the evolution of the system in the  $t_a$  direction commute.

### 3. THE KONTSEVICH-WITTEN THEOREM

Define the partition function (precisely, this is the all-genus Gromov-Witten potential of a point)

$$\mathcal{Z} = \exp \left( \sum_{\substack{g=0 \\ n=1 \\ 2g-2+n>1}}^{\infty} \frac{\hbar^{2g-2+n}}{n!} \sum_{a_1, \dots, a_n} \langle \psi_1^{a_1} \dots \psi_n^{a_n} \rangle_{g,n} t_{a_1} \dots t_{a_n} \right).$$

**Theorem 3.1** (Kontsevich 1992). *The function*

$$u(t_0, t_1, \dots) = \frac{\partial^2}{\partial t_0^2} \log \mathcal{Z}$$

*is the unique solution to the KdV hierarchy with initial condition  $u(t_0, 0, 0, \dots) = t_0$ .*

This was proven by Kontsevich using a combinatorial description of  $\overline{\mathcal{M}}_{g,n}$  in terms of thickened graphs and computing various matrix integrals. This theorem has several other proofs using various techniques. We will briefly outline three of them.

- Okounkov and Pandharipande (2001) gave a proof which proceeds first by relating the integrals of  $\psi$ -classes on  $\overline{\mathcal{M}}_{g,n}$  to counts of permutations known as *Hurwitz numbers*, continues by constructing a different matrix model than the one considered by Kontsevich, and concludes by relating their matrix model to that of Kontsevich using techniques from probability theory.
- Mirzakhani (2007) gave a proof which proceeds first by defining a metric (the so-called *Weil-Petersson metric*) on  $\overline{\mathcal{M}}_{g,n}$  and computing a recursive formula for the volume of  $\overline{\mathcal{M}}_{g,n}$  and concludes by relating the volumes to the integrals we defined in the first section.
- Kazarian and Lando (2007) gave a proof which relates the Hurwitz numbers considered by Okounkov and Pandharipande to a different integrable hierarchy known as the KP hierarchy, and then concludes by reducing the KP hierarchy to the KdV hierarchy.