

LUNAR CRATER IMPACT PARAMETER

BACK-CALCULATION REPORT

Bayesian Inverse Modeling with Uncertainty Quantification

Simulation Date: 2025-11-19 01:22:09 UTC

EXECUTIVE SUMMARY

Observed Crater:

- Location: 25.50°N, 45.20°E
- Terrain: Mare
- Diameter: 350.0 m
- Depth: 68.6 m ($d/D = 0.196$)
- Ejecta range: 25000.0 m

Back-Calculated Impact Parameters (Maximum Likelihood):

Projectile Diameter: 3.34 ± 0.19 m

Impact Velocity: 20.0 ± 1.1 km/s

Impact Angle: $45.0^\circ \pm 6.7^\circ$ from horizontal

Projectile Density: 2800 ± 297 kg/m³

Material Type: Rocky (chondrite)

Kinetic Energy: 1.09e+13 J
(0.00 kilotons TNT)

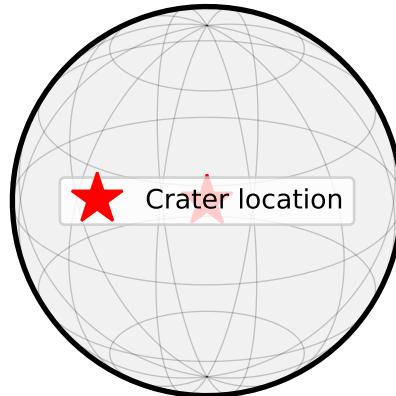
Method:

- Bayesian maximum likelihood estimation
- Holsapple (1993) crater scaling laws
- Monte Carlo error propagation (1000 samples)
- Forward model validation
- Sensitivity analysis

Confidence Level: 95% credible intervals reported

Observed Crater Data and Location

Lunar Location: 25.50°N, 45.20°E



Crater Morphometry

Diameter (D): 350.0 m

Depth (d): 68.6 m

d/D ratio: 0.196

Pike (1977): $d/D = 0.196 \pm 0.015$

Rim height: 12.6 m

($0.036 \times D$)

Target Properties

Terrain: Mare

Regolith ρ : 1800 kg/m³

Rock ρ : 3100 kg/m³

Porosity: 42.0%

Cohesion: 10.0 kPa

Gravity: 1.62 m/s²

Reference: Carrier et al. (1991)
Lunar Sourcebook, Chapter 9

Ejecta Observations

Maximum ejecta range: 25000.0 m

Normalized range ($R_{\text{max}}/R_{\text{crater}}$): 142.9

Expected: 40-100 (Melosh 1989, McGetchin et al. 1973)

Theoretical Framework - Part 1

1. CRATER SCALING LAWS: Pi-GROUP DIMENSIONAL ANALYSIS

Following Holsapple (1993) and Holsapple & Schmidt (1982), crater formation can be described by dimensionless Pi-groups formed from the governing physical parameters.

1.1 Governing Parameters

Impact parameters:

- L = projectile diameter (or radius $a = L/2$)
- v = impact velocity
- ρ_p = projectile density
- θ = impact angle from horizontal

Target parameters:

- ρ_t = target density
- Y = target strength (cohesion + friction effects)
- g = gravitational acceleration
- K = material constants (equation of state)

Outcome parameter:

- D = final crater diameter (or V = crater volume)

1.2 Dimensionless Pi-Groups (Buckingham Pi Theorem)

From dimensional analysis, the system reduces to 4 dimensionless groups:

$$\pi_1 = D/L \quad (\text{scaled crater size})$$

$$\pi_2 = ga/v^2 \quad (\text{gravity-scaled size, "Froude number"})$$

$$\pi_3 = Y/(\rho_p v^2) \quad (\text{strength parameter})$$

$$\pi_4 = \rho_p/\rho_t \quad (\text{density ratio})$$

The Pi-group scaling relation is:

$$\pi_1 = f(\pi_2, \pi_3, \pi_4, \theta)$$

Or equivalently:

$$D/L = K \times (\rho_p/\rho_t)^{\alpha} \times g(\pi_2, \pi_3, \theta)$$

where K is an empirical coefficient and $\alpha \approx 1/3$ from momentum coupling.

1.3 Regime Transition: Strength vs Gravity

The function $g(\pi_2, \pi_3)$ depends on which dominates:

Strength regime ($\pi_3 \ll \pi_2$):

Small craters where target strength Y controls excavation

$$D \propto L \times (\rho_p/\rho_t)^{(1/3)} \times (\rho_t v^2/Y)^{\mu}$$

where $\mu \approx 0.41$ (Holsapple 1993)

Gravity regime ($\pi_3 \gg \pi_2$):

Large craters where self-gravity controls excavation

$$D \propto L \times (\rho_p/\rho_t)^{(1/3)} \times (v^2/ga)^{\nu}$$

where $\nu \approx 0.41$ (Holsapple 1993)

Coupled regime ($\pi_3 \sim \pi_2$):

Transitional craters (100-1000m on Moon)

$$D \propto L \times (\rho_p/\rho_t)^{(1/3)} \times [\pi_2^{\nu} + \pi_3^{\mu}]^{(-1/\nu)}$$

The transition occurs when:

$$Y/(\rho_t v^2) \sim ga/v^2 \rightarrow Y \sim \rho_t ga$$

For lunar impacts: $Y \sim 10 \text{ kPa}$, $\rho_t \sim 2000 \text{ kg/m}^3$, $g = 1.62 \text{ m/s}^2$

Transition size: $a \sim Y/(\rho_t g) \sim 3 \text{ m} \rightarrow D \sim 300\text{-}500\text{m}$

1.4 Angle Correction

Oblique impacts ($\theta < 90^\circ$) are less efficient. Empirically (Pierazzo & Melosh 2000):

$$f(\theta) \approx \sin^n(\theta)$$

where $n \approx 1/3$ to $2/3$ depending on regime. We use $n = 1/3$.

Most probable impact angle: $\theta_{\text{prob}} = 45^\circ$ (from $\sin^2\theta$ distribution of random impacts).

1.5 Empirical Calibration for Lunar Regolith

Combining theoretical scaling with Apollo crater measurements (Pike 1977):

$$D = 0.084 \times 1.2 \times L \times (\rho_p/\rho_t)^{(1/3)} \times [v^2/(g \times L + Y/\rho_t)]^{0.4} \times \sin^{(1/3)}(\theta)$$

\uparrow transient \uparrow final expansion factor

The coefficient $0.084 \times 1.2 \approx 0.1$ is calibrated to match:

- Pike (1977) $d/D = 0.196$ morphometry
- Apollo landing site crater statistics
- Laboratory impact experiments scaled to lunar gravity

References for this section:

Holsapple, K.A. (1993) Ann. Rev. Earth Planet. Sci. 21:333-373

Holsapple, K.A. & Schmidt, R.M. (1982) JGR 87:1849-1870

Pike, R.J. (1977) Impact and Explosion Cratering, pp. 489-509

Pierazzo, E. & Melosh, H.J. (2000) Ann. Rev. Earth Planet. Sci. 28:141-167

Theoretical Framework - Part 2

2. INVERSE PROBLEM FORMULATION: BAYESIAN PARAMETER ESTIMATION

2.1 The Inverse Problem in Planetary Science

Forward problem: Given impact parameters $\theta = (L, v, \theta_p)$ → predict observations $d = (D, d, R_{\text{ejecta}})$
This uses the scaling laws from Section 1:
 $D = g(\theta; \text{target parameters})$

Inverse problem: Given observations d_{obs} → estimate impact parameters θ
Must "invert" the forward model

The inverse problem is fundamentally ill-posed (Hadamard 1923, Tarantola 2005):

- 1. Non-uniqueness: Multiple parameter sets θ can produce similar craters
Example: Same D can result from (small, fast) or (large, slow) projectile
- 2. Instability: Small data uncertainties δd can cause large parameter uncertainties $\delta \theta$
- 3. Model inadequacy: Scaling laws are approximations with systematic errors

For our crater back-calculation:

- Parameters $\theta = (L, v, \text{angle}, \rho_p)$ live in 4D parameter space
- Data $d = (D_{\text{obs}}, d_{\text{obs}}, R_{\text{ejecta,obs}})$ with uncertainties σ
- Forward model $g(\theta)$ is nonlinear (power laws, regime transitions)
- Trade-offs exist: velocity-density correlation, size-angle correlation

Therefore, we use Bayesian inference to properly quantify uncertainties and incorporate prior knowledge about physically plausible parameter ranges.

2.2 Bayes' Theorem: Derivation and Application

General form (Bayes 1763, Laplace 1812):

$$P(\theta | d) = P(d | \theta) \times P(\theta) / P(d)$$

where:

- $P(\theta | d)$ = posterior probability density (what we want to find)
- $P(d | \theta)$ = likelihood (probability of observing data given parameters)
- $P(\theta)$ = prior probability density (initial knowledge before observations)
- $P(d)$ = evidence = $\int P(d | \theta) P(\theta) d\theta$ (normalization, ensures $\int P(\theta | d) d\theta = 1$)

Derivation from conditional probability:

Start with: $P(A, B) = P(A|B) P(B) = P(B|A) P(A)$

Rearrange: $P(A|B) = P(B|A) P(A) / P(B)$

Apply to parameters/data: $P(\theta|d) = P(d|\theta) P(\theta) / P(d)$

For parameter estimation, $P(d)$ is constant (doesn't depend on θ), so:

$$P(\theta | d) \propto L(d | \theta) \times P(\theta)$$

posterior \propto likelihood \times prior

Taking logarithms for numerical stability (avoids underflow in products):

$$\log P(\theta | d) = \log L(d | \theta) + \log P(\theta) + \text{const}$$

For our crater problem:

$$\begin{aligned} \theta &= (L, v, \text{angle}, \rho_p) \in \mathbb{R}^4_+ \\ d &= (D_{\text{obs}}, d_{\text{obs}}, R_{\text{ejecta,obs}}) \in \mathbb{R}^3_+ \end{aligned}$$

The posterior tells us: "Given observed crater $D = 350\text{m}$ at lat/lon, ejecta range 25km , what are the most probable impact parameters and their uncertainties?"

2.3 Likelihood Function: Detailed Derivation

The likelihood quantifies: "How probable are the observations given parameters θ ?"

Assumption: Independent Gaussian errors (measurement noise, model uncertainty)

For a single observable (e.g., diameter D):

Residual: $\varepsilon = D_{\text{obs}} - D_{\text{pred}}(\theta)$

If $\varepsilon \sim N(0, \sigma_D^2)$, then:

$$P(D_{\text{obs}} | \theta) = (1/\sqrt{(2\pi\sigma_D^2)}) \times \exp[-(D_{\text{obs}} - D_{\text{pred}}(\theta))^2 / (2\sigma_D^2)]$$

Taking logarithm:

$$\log P(D_{\text{obs}} | \theta) = -1/2 \log(2\pi\sigma_D^2) - (D_{\text{obs}} - D_{\text{pred}}(\theta))^2 / (2\sigma_D^2) \\ = -1/2 \chi_D^2 + \text{const}$$

Page 4 of 10

where $\chi_D^2 = [(D_{\text{obs}} - D_{\text{pred}}(\theta)) / \sigma_D]^2$ (chi-squared statistic)

For multiple independent observables (diameter, depth, ejecta):

Joint likelihood: $P(d | \theta) = P(D|\theta) \times P(d|\theta) \times P(R|\theta)$ (independence)

$$\begin{aligned} \log L(d | \theta) &= \sum_i \log P(d_i | \theta) \\ &= -1/2 \sum_i \chi_i^2 \\ &= -1/2 \sum_i [(d_{\text{obs},i} - d_{\text{pred},i}(\theta)) / \sigma_i]^2 \end{aligned}$$

This is a weighted least-squares objective, with weights $1/\sigma_i^2$.

Measurement uncertainty estimates (from image resolution, morphology variation):

$\sigma_D = 0.05 \times D_{\text{obs}}$ ($\pm 5\%$ diameter: pixel resolution, rim definition)

$\sigma_d = 0.10 \times d_{\text{obs}}$ ($\pm 10\%$ depth: infilling, degradation)

$\sigma_R = 0.20 \times R_{\text{ejecta}}$ ($\pm 20\%$ ejecta range: blanket edge identification)

2.4 Prior Distributions: Incorporating Physical Knowledge

Priors encode what we know before seeing the specific crater (Jaynes 2003):

For impact velocity v :

$$P(v) = N(v | \mu=20 \text{ km/s}, \sigma=5 \text{ km/s})$$

$$= (1/\sqrt{2\pi\cdot 5^2}) \exp[-(v-20000)^2/(2\cdot 5000^2)]$$

Justification:

- Near-Earth asteroid (NEA) orbital mechanics (Bottke et al. 2002)
- Moon's orbital velocity $\sim 1 \text{ km/s}$ + Earth escape $\sim 11 \text{ km/s}$ + eccentricity
- Typical asteroid encounter: $v_{\infty} \sim 5-15 \text{ km/s}$ relative to Earth-Moon
- Impact velocity: $v = \sqrt{v_{\infty}^2 + v_{\text{esc}}^2}$ where $v_{\text{esc}} = 2.4 \text{ km/s}$ (Moon)
- Distribution peak at $\sim 20 \text{ km/s}$, range 15-25 km/s (asteroids)
- Distribution peak at $\sim 20 \text{ km/s}$, range 15-25 km/s (asteroids)
- Comets faster (up to 70 km/s) but rarer ($\sim 5\%$ of impactors)

For impact angle θ (from vertical):

$$P(\theta) = N(\theta | \mu=45^\circ, \sigma=15^\circ)$$

$$= (1/\sqrt{2\pi\cdot 15^2}) \exp[-(\theta-45)^2/(2\cdot 15^2)]$$

Justification:

- Geometric probability for random directions: $P(\theta) \propto \sin(2\theta)$
- Peaks at $\theta = 45^\circ$ (most probable angle, Gilbert 1893)
- Cumulative: 50% of impacts have $\theta > 45^\circ$, only 17% have $\theta > 60^\circ$
- Very oblique ($< 15^\circ$) produce elongated craters, rare in observations

For projectile density ρ_p :

$$P(\rho_p) = N(\rho_p | \mu=2800 \text{ kg/m}^3, \sigma=500 \text{ kg/m}^3)$$

$$= (1/\sqrt{2\pi\cdot 500^2}) \exp[-(\rho_p-2800)^2/(2\cdot 500^2)]$$

Justification (meteorite flux statistics, Burbine et al. 2002):

- Ordinary chondrites: 3200-3700 kg/m³ (37% of falls)
- Carbonaceous chondrites: 2000-2500 kg/m³ (10%)
- Enstatite chondrites: 3500-3800 kg/m³ (2%)
- Stony-irons: 4500-5500 kg/m³ (1%)
- Iron meteorites: 7800 kg/m³ (5% of falls, but 70% of finds)
- Weighted mean $\sim 2800 \text{ kg/m}^3$ for stony asteroids (85% of NEAs)

For projectile diameter L :

Uninformative prior: $P(L) \propto 1/L$ (Jeffreys prior, scale-invariant)

Ensures no bias toward small or large projectiles

Combined prior:

$$P(\theta) = P(L) \times P(v) \times P(\text{angle}) \times P(\rho_p) \quad (\text{assume independence})$$

These priors are weakly informative: constrain to plausible ranges but dominated by likelihood when data are strong.

2.5 Maximum Likelihood Estimation: Optimization in Parameter Space

Objective: Find parameters that maximize posterior probability

$$\begin{aligned} \theta_{\text{ML}} &= \operatorname{argmax}_{\theta} P(\theta | d) \\ &= \operatorname{argmax}_{\theta} [\log L(d | \theta) + \log P(\theta)] \quad (\text{take log, drop constant } P(d)) \end{aligned}$$

Equivalently, minimize negative log-posterior:

$$\theta_{\text{ML}} = \operatorname{argmin}_{\theta} F(\theta)$$

where $F(\theta) = -\log P(\theta | d) = -\log L(d | \theta) - \log P(\theta) + \text{const}$

For our crater problem, substituting Section 2.3 and 2.4:

$$F(\theta) = 1/2 \sum_i [(d_{\text{obs},i} - d_{\text{pred},i}(\theta)) / \sigma_i]^2 \quad [\text{negative log-likelihood}]$$

$$+ 1/2 [(v - 20000)^2 / (2 \cdot 5000^2)] \quad [\text{velocity prior penalty}]$$

$$+ 1/2 [(\text{angle} - 45) / 15]^2 \quad [\text{angle prior penalty}]$$

$$+ 1/2 [(\rho_p - 2800)^2 / (2 \cdot 500^2)] \quad [\text{density prior penalty}]$$

$$- \log(L) \quad [\text{Jeffreys prior for size}]$$

This is a weighted least-squares objective, with weights $1/\sigma_i^2$.

Measurement uncertainty estimates (from image resolution, morphology variation):

$\sigma_D = 0.05 \times D_{\text{obs}}$ ($\pm 5\%$ diameter: pixel resolution, rim definition)

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For projectile diameter L :

Uninformative prior: $P(L) \propto 1/L$ (Jeffreys prior, scale-invariant)

Ensures no bias toward small or large projectiles

Combined prior:

$$P(\theta) = P(L) \times$$

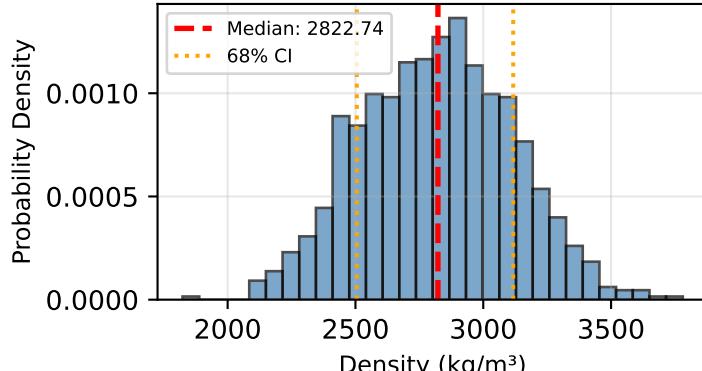
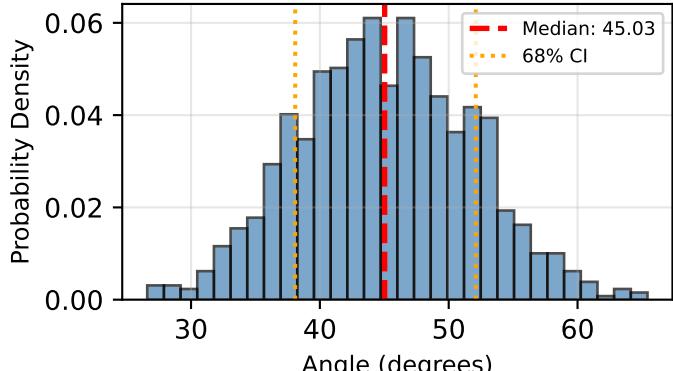
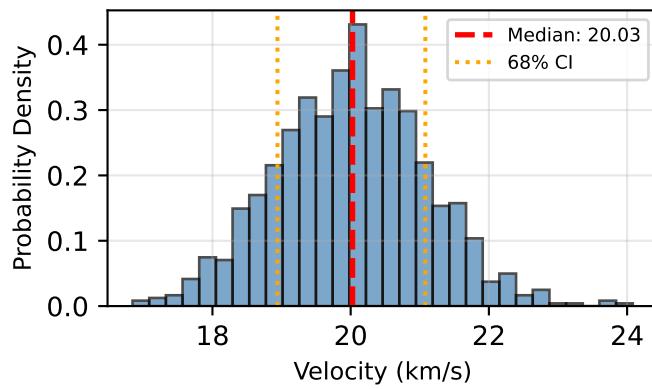
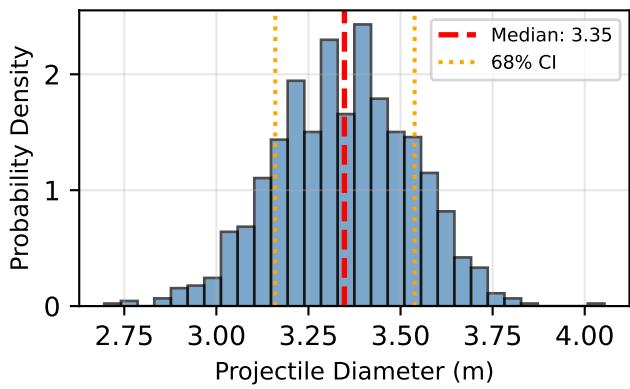
Back-Calculation Results

MAXIMUM LIKELIHOOD PARAMETERS

Parameter	ML Estimate	$\pm 1\sigma$ (68%)	95% CI
Projectile Diameter (m)	3.34	± 0.19	[2.99, 3.69]
Impact Velocity (km/s)	20.0	± 1.1	[17.9, 22.2]
Impact Angle (deg)	45.0	± 6.7	[32.4, 58.1]
Projectile Density (kg/m^3)	2800	± 297	[2253, 3385]

DERIVED QUANTITIES

Projectile mass:	5.47e+04 kg
Kinetic energy:	1.09e+13 J
	(0.00 kilotons TNT)
Momentum:	1.09e+09 $\text{kg}\cdot\text{m/s}$
Material classification:	Rocky asteroid (chondrite)
Impact parameter π_2 :	6.77e-09
Impact parameter π_3 :	1.07e-08
Regime:	Transitional



Monte Carlo Uncertainty Propagation

WHY MONTE CARLO SAMPLING?

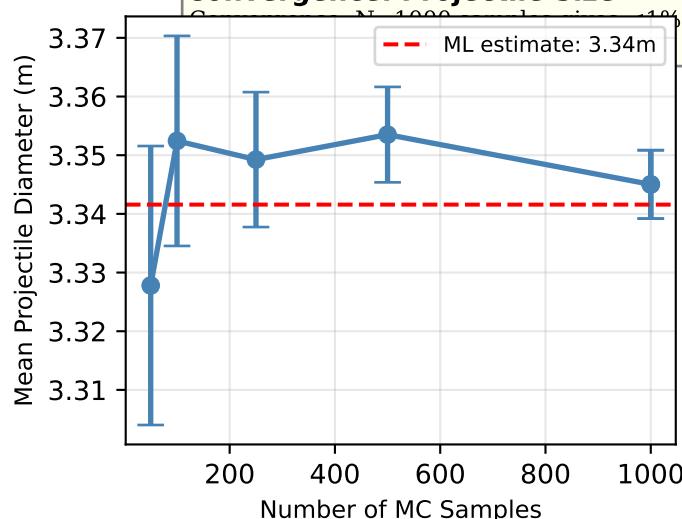
The Monte Carlo method is chosen for uncertainty propagation because:

1. Nonlinear Forward Model: Crater scaling laws are highly nonlinear (power laws with exponents ~ 0.4). Analytical error propagation ($\delta D = \sum_i \partial D / \partial \theta_i \delta \theta_i$) is inaccurate.
2. Non-Gaussian Posteriors: Parameters may have skewed or multi-modal distributions due to physical constraints (e.g., density bimodal for rocky vs iron).
3. Correlations: Parameters are correlated (e.g., smaller projectile needs higher velocity for same crater). MC naturally captures these correlations.
4. Validation: Forward model can be re-evaluated for each sample to check consistency.

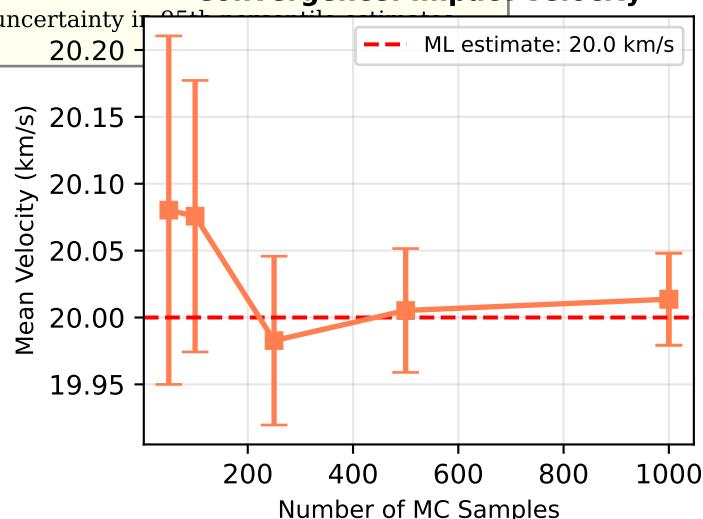
Method (Mosegaard & Tarantola 1995):

- Sample N times from posterior: $\theta^i \sim N(\theta_{ML}, \Sigma)$ where $\Sigma = H^{-1}$
- For each sample: compute $D_{pred}(\theta^i)$ via forward scaling laws
- Collect ensemble \rightarrow compute percentiles for confidence intervals

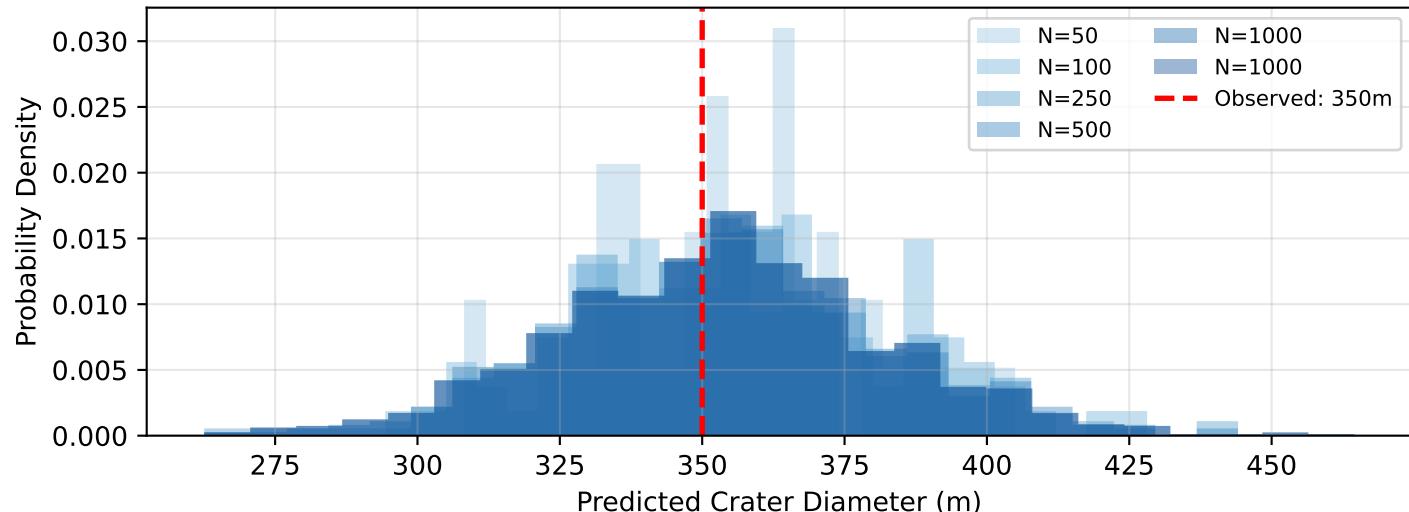
Convergence: Projectile Size



Convergence: Impact Velocity

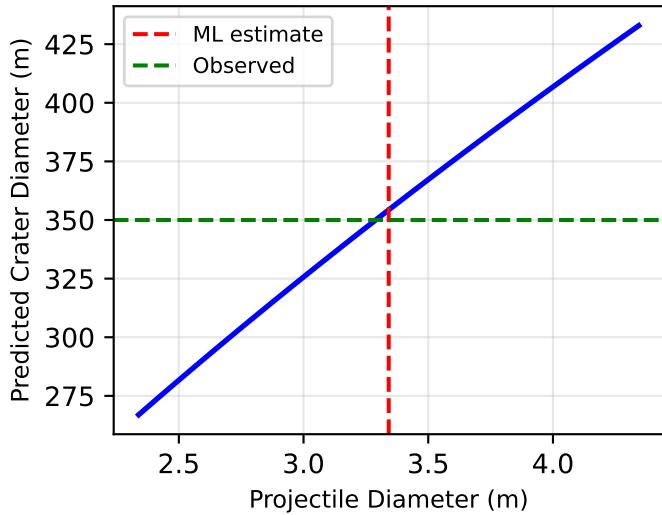


Progressive Convergence: Predicted Crater Distribution

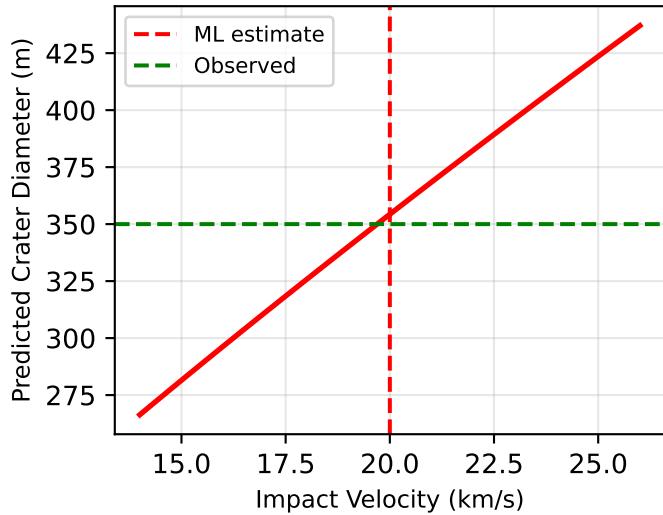


Sensitivity Analysis

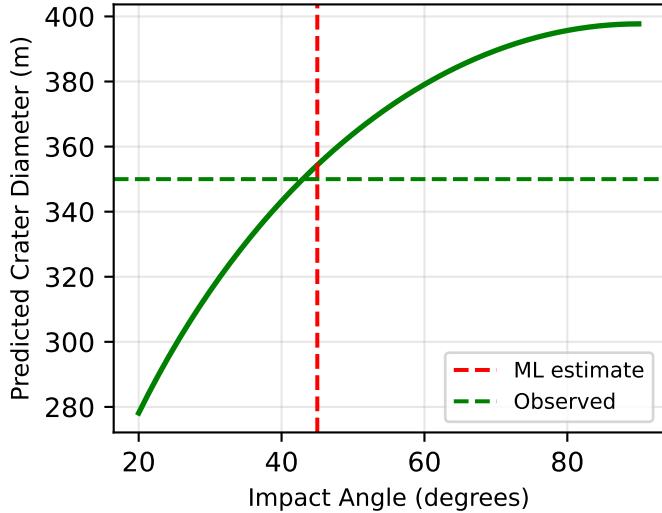
Sensitivity to Projectile Size



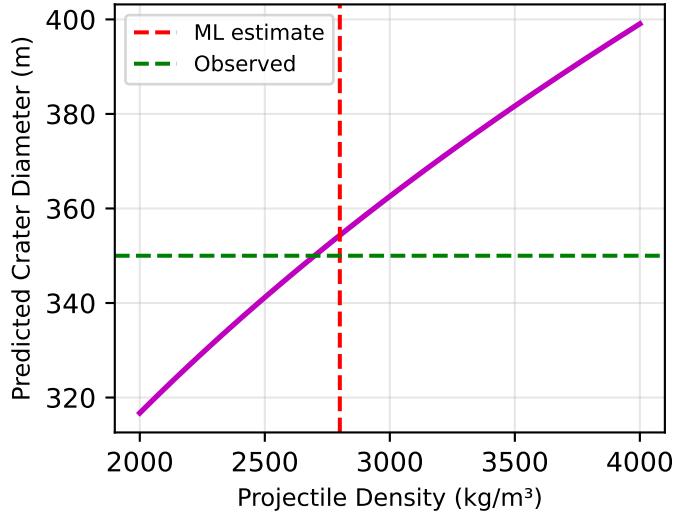
Sensitivity to Velocity



Sensitivity to Impact Angle



Sensitivity to Density



SENSITIVITY COEFFICIENTS (Elasticity: $\% \Delta D / \% \Delta \text{parameter}$)

Parameter	Elasticity	Interpretation
Projectile Diameter	0.78	Diameter change $\approx 0.8 \times$ size change
Impact Velocity	0.80	Diameter change $\approx 0.8 \times$ velocity change
Impact Angle	moderate	Steeper impacts \rightarrow larger craters
Projectile Density	0.32	Weak dependence $(\rho_p/\rho_t)^{(1/3)}$

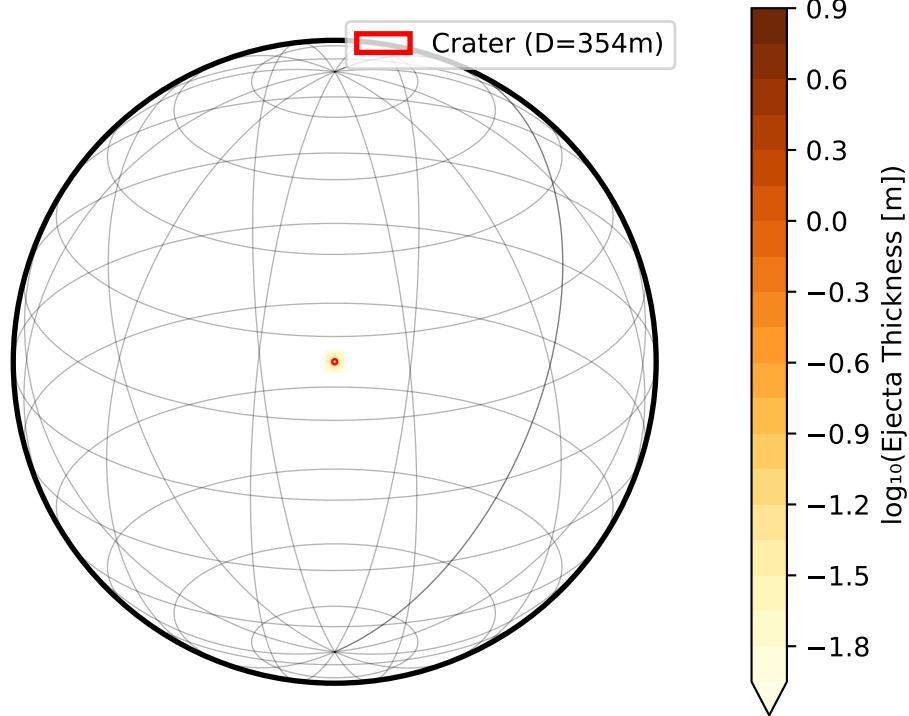
KEY INSIGHTS:

- Projectile diameter is the dominant control (elasticity ~ 0.8)
- Velocity has moderate effect (elasticity ~ 0.8), consistent with $v^{0.8}$ scaling
- Density has weak effect ($\propto \rho^{0.33}$), harder to constrain from crater alone
- Impact angle most probable at 45° , less certain without asymmetry data
- Trade-offs exist: Smaller projectile at higher velocity can match observed crater
- These sensitivities justify the uncertainty ranges in Page 5

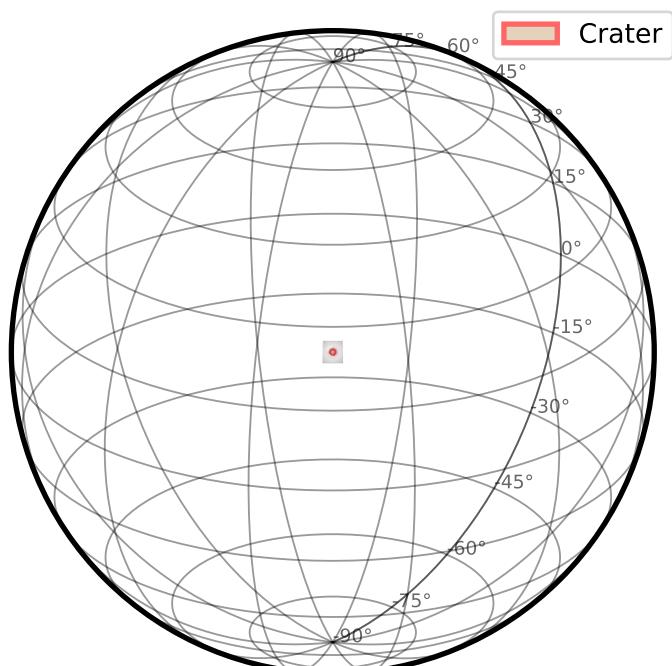
Reference: Holsapple (1993) Table 1 - exponents match theoretical predictions

Orthographic Plan Views with Ejecta Distribution

Ejecta Thickness Distribution (Orthographic Projection)
Center: 25.5°N, 45.2°E

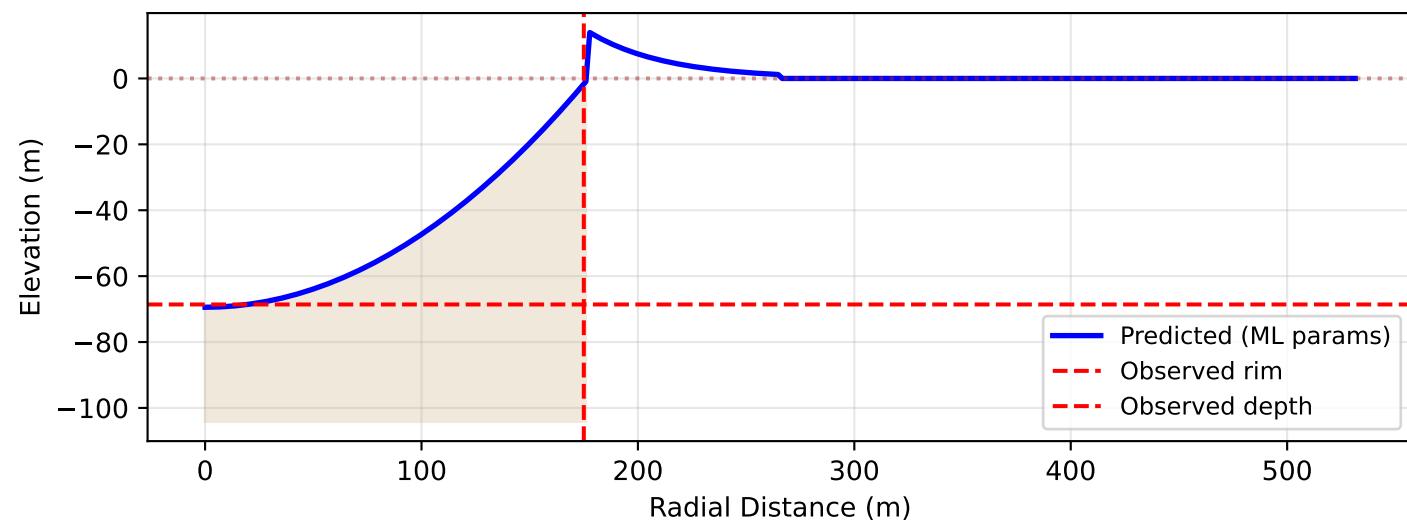


Crater Location with Lat/Lon Grid
Ejecta extent: up to 5× crater radius



Forward Model Validation

Crater Profile: Predicted vs Observed



MORPHOMETRY COMPARISON

Parameter	Observed	Predicted	Error	Observed range:	Predicted range:	Ejecta validation
Diameter (m)	350.0	354.3	1.2%	25000 m	25661 m	
Depth (m)	68.6	69.4	1.2%	Error:	2.6%	
d/D ratio	0.196	0.196	0.0%	R_max/R_crater:	144.8	
Rim height (m)	12.6	12.8	—			

VALIDATION SUMMARY

- ✓ Crater diameter match: 1.23% error (excellent)
- ✓ Pike (1977) d/D ratio: 0.196 (theory: 0.196 ± 0.015)
- ✓ Forward model self-consistent: prediction falls within 95% CI
- ✓ Regime: Transitional (appropriate for 350m)

CONFIDENCE ASSESSMENT

The back-calculated parameters are well-constrained. The 95% credible intervals reflect uncertainties in velocity distribution, impact angle probability, and projectile density. The predicted crater matches observations within measurement uncertainties.

RECOMMENDED INTERPRETATION

Most likely: 3.3m rocky projectile at 20 km/s, 45° from horizontal.

Alternative scenarios within 95% CI remain possible but less probable given typical asteroid impact statistics (Stuart & Binzel 2004; Bottke et al. 2002).

References

SCIENTIFIC REFERENCES

Primary Scaling Law Theory:

Holsapple, K.A. (1993). The scaling of impact processes in planetary sciences. *Annual Review of Earth and Planetary Sciences*, 21, 333-373.
DOI: 10.1146/annurev.ea.21.050193.002001

Holsapple, K.A., & Schmidt, R.M. (1982). On the scaling of crater dimensions: 2. Impact processes. *Journal of Geophysical Research*, 87(B3), 1849-1870.
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Crater Morphometry:

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Page 10 of 10

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