

LUNAR CRATER IMPACT PARAMETER BACK-CALCULATION REPORT

Bayesian Inverse Modeling with Uncertainty Quantification

Simulation Date: 2025-11-19 01:22:09 UTC

EXECUTIVE SUMMARY

Observed Crater:

- Location: 25.50°N, 45.20°E
- Terrain: Mare
- Diameter: 350.0 m
- Depth: 68.6 m ($d/D = 0.196$)
- Ejecta range: 25000.0 m

Back-Calculated Impact Parameters (Maximum Likelihood):

Projectile Diameter: 3.34 ± 0.19 m

Impact Velocity: 20.0 ± 1.1 km/s

Impact Angle: $45.0^\circ \pm 6.7^\circ$ from horizontal

Projectile Density: 2800 ± 297 kg/m³

Material Type: Rocky (chondrite)

Kinetic Energy: $1.09\text{e}+13$ J
(0.00 kilotons TNT)

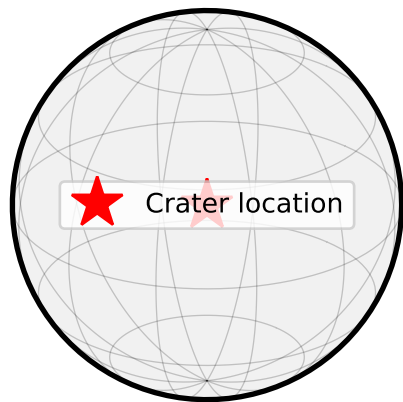
Method:

- Bayesian maximum likelihood estimation
- Holsapple (1993) crater scaling laws
- Monte Carlo error propagation (1000 samples)
- Forward model validation
- Sensitivity analysis

Confidence Level: 95% credible intervals reported

Observed Crater Data and Location

Lunar Location: 25.50°N, 45.20°E



Crater Morphometry

Target Properties

Diameter (D): 350.0 m
Depth (d): 68.6 m
d/D ratio: 0.196
Pike (1977): $d/D = 0.196 \pm 0.015$
Rim height: 12.6 m
($0.036 \times D$)

Terrain: Mare
Regolith ρ : 1800 kg/m³
Rock ρ : 3100 kg/m³
Porosity: 42.0%
Cohesion: 10.0 kPa
Gravity: 1.62 m/s²
Reference: Carrier et al. (1991)
Lunar Sourcebook, Chapter 9

Ejecta Observations

Maximum ejecta range: 25000.0 m
Normalized range (R_{max}/R_{crater}): 142.9
Expected: 40-100 (Melosh 1989, McGetchin et al. 1973)

Theoretical Framework - Part 1

1. CRATER SCALING LAWS: PI-GROUP DIMENSIONAL ANALYSIS

Following Holsapple (1993) and Holsapple & Schmidt (1982), crater formation can be described by dimensionless Pi-groups formed from the governing physical parameters.

1.1 Governing Parameters

Impact parameters:

- L = projectile diameter (or radius $a = L/2$)
- v = impact velocity
- ρ_p = projectile density
- θ = impact angle from horizontal

Target parameters:

- ρ_t = target density
- Y = target strength (cohesion + friction effects)
- g = gravitational acceleration
- K = material constants (equation of state)

Outcome parameter:

- D = final crater diameter (or V = crater volume)

1.2 Dimensionless Pi-Groups (Buckingham Pi Theorem)

From dimensional analysis, the system reduces to 4 dimensionless groups:

- $\pi_1 = D/L$ (scaled crater size)
- $\pi_2 = ga/v^2$ (gravity-scaled size, "Froude number")
- $\pi_3 = Y/(\rho_t v^2)$ (strength parameter)
- $\pi_4 = \rho_p/\rho_t$ (density ratio)

The Pi-group scaling relation is:

$$\pi_1 = f(\pi_2, \pi_3, \pi_4, \theta)$$

Or equivalently:

$$D/L = K \times (\rho_p/\rho_t)^\alpha \times g(\pi_2, \pi_3, \theta)$$

where K is an empirical coefficient and $\alpha \approx 1/3$ from momentum coupling.

1.3 Regime Transition: Strength vs Gravity

The function $g(\pi_2, \pi_3)$ depends on which dominates:

Strength regime ($\pi_3 \ll \pi_2$):

Small craters where target strength Y controls excavation
 $D \propto L \times (\rho_p/\rho_t)^{1/3} \times (\rho_t v^2/Y)^\mu$
where $\mu \approx 0.41$ (Holsapple 1993)

Gravity regime ($\pi_3 \gg \pi_2$):

Large craters where self-gravity controls excavation
 $D \propto L \times (\rho_p/\rho_t)^{1/3} \times (v^2/ga)^\nu$
where $\nu \approx 0.41$ (Holsapple 1993)

Coupled regime ($\pi_3 \sim \pi_2$):

Transitional craters (100-1000m on Moon)
 $D \propto L \times (\rho_p/\rho_t)^{1/3} \times [\pi_2^\nu + \pi_3^\mu]^{(-1/\nu)}$

The transition occurs when:

$$Y/(\rho_t v^2) \sim ga/v^2 \rightarrow Y \sim \rho_t ga$$

For lunar impacts: $Y \sim 10 \text{ kPa}$, $\rho_t \sim 2000 \text{ kg/m}^3$, $g = 1.62 \text{ m/s}^2$

Transition size: $a \sim Y/(\rho_t g) \sim 3 \text{ m} \rightarrow D \sim 300\text{-}500\text{m}$

1.4 Angle Correction

Oblique impacts ($\theta < 90^\circ$) are less efficient. Empirically (Pierazzo & Melosh 2000):

$$f(\theta) \approx \sin^n(\theta)$$

where $n \approx 1/3$ to $2/3$ depending on regime. We use $n = 1/3$.

Most probable impact angle: $\theta_{\text{prob}} = 45^\circ$ (from $\sin^2\theta$ distribution of random impacts).

1.5 Empirical Calibration for Lunar Regolith

Combining theoretical scaling with Apollo crater measurements (Pike 1977):

$$D = 0.084 \times 1.2 \times L \times (\rho_p/\rho_t)^{1/3} \times [v^2/(g \times L + Y/\rho_t)]^{0.4} \times \sin^{1/3}(\theta)$$

↑ transient ↑ final expansion factor

The coefficient $0.084 \times 1.2 \approx 0.1$ is calibrated to match:

- Pike (1977) $d/D = 0.196$ morphometry
- Apollo landing site crater statistics
- Laboratory impact experiments scaled to lunar gravity

References for this section:

- Holsapple, K.A. (1993) Ann. Rev. Earth Planet. Sci. 21:333-373
Holsapple, K.A. & Schmidt, R.M. (1982) JGR 87:1849-1870
Pike, R.J. (1977) Impact and Explosion Cratering, pp. 489-509
Pierazzo, E. & Melosh, H.J. (2000) Ann. Rev. Earth Planet. Sci. 28:141-167

Theoretical Framework - Part 2

2. INVERSE PROBLEM FORMULATION: BAYESIAN PARAMETER ESTIMATION

2.1 The Inverse Problem in Planetary Science

Forward problem: Given impact parameters $\theta = (L, v, \theta, \rho_p) \rightarrow$ predict observations $d = (D, d, R_{\text{ejecta}})$
This uses the scaling laws from Section 1:
 $D = g(\theta; \text{target parameters})$

Inverse problem: Given observations $d_{\text{obs}} \rightarrow$ estimate impact parameters θ
Must "invert" the forward model

The inverse problem is fundamentally ill-posed (Hadamard 1923, Tarantola 2005):

1. Non-uniqueness: Multiple parameter sets θ can produce similar craters
Example: Same D can result from (small, fast) or (large, slow) projectile
2. Instability: Small data uncertainties δd can cause large parameter uncertainties $\delta \theta$
3. Model inadequacy: Scaling laws are approximations with systematic errors

For our crater back-calculation:

- Parameters $\theta = (L, v, \text{angle}, \rho_p)$ live in 4D parameter space
- Data $d = (D_{\text{obs}}, d_{\text{obs}}, R_{\text{ejecta,obs}})$ with uncertainties σ
- Forward model $g(\theta)$ is nonlinear (power laws, regime transitions)
- Trade-offs exist: velocity-density correlation, size-angle correlation

Therefore, we use Bayesian inference to properly quantify uncertainties and incorporate prior knowledge about physically plausible parameter ranges.

2.2 Bayes' Theorem: Derivation and Application

General form (Bayes 1763, Laplace 1812):

$$P(\theta | d) = P(d | \theta) \times P(\theta) / P(d)$$

where:
 $P(\theta | d)$ = posterior probability density (what we want to find)
 $P(d | \theta)$ = likelihood (probability of observing data given parameters)
 $P(\theta)$ = prior probability density (initial knowledge before observations)
 $P(d) = \int P(d | \theta) P(\theta) d\theta$ (normalization, ensures $\int P(\theta|d) d\theta = 1$)

Derivation from conditional probability:

Start with: $P(A,B) = P(A|B) P(B) = P(B|A) P(A)$

Rearrange: $P(A|B) = P(B|A) P(A) / P(B)$

Apply to parameters/data: $P(\theta|d) = P(d|\theta) P(\theta) / P(d)$

For parameter estimation, $P(d)$ is constant (doesn't depend on θ), so:

$$P(\theta | d) \propto L(d | \theta) \times P(\theta)$$

$$\text{posterior} \propto \text{likelihood} \times \text{prior}$$

Taking logarithms for numerical stability (avoids underflow in products):

$$\log P(\theta | d) = \log L(d | \theta) + \log P(\theta) + \text{const}$$

For our crater problem:

$$\theta = (L, v, \text{angle}, \rho_p) \in \mathbb{R}^4_+$$

$$d = (D_{\text{obs}}, d_{\text{obs}}, R_{\text{ejecta,obs}}) \in \mathbb{R}^3_+$$

The posterior tells us: "Given observed crater $D = 350\text{m}$ at lat/lon, ejecta range 25km , what are the most probable impact parameters and their uncertainties?"

2.3 Likelihood Function: Detailed Derivation

The likelihood quantifies: "How probable are the observations given parameters θ ?"

Assumption: Independent Gaussian errors (measurement noise, model uncertainty)

For a single observable (e.g., diameter D):

$$\text{Residual: } \varepsilon = D_{\text{obs}} - D_{\text{pred}}(\theta)$$

$$\text{If } \varepsilon \sim N(0, \sigma_{D^2}), \text{ then:}$$

$$P(D_{\text{obs}} | \theta) = (1/\sqrt{2\pi\sigma_{D^2}}) \times \exp[-(D_{\text{obs}} - D_{\text{pred}}(\theta))^2 / (2\sigma_{D^2})]$$

Taking logarithm:

$$\log P(D_{\text{obs}} | \theta) = -1/2 \log(2\pi\sigma_{D^2}) - (D_{\text{obs}} - D_{\text{pred}}(\theta))^2 / (2\sigma_{D^2})$$

$$= -1/2 \chi^2_{D^2} + \text{const}$$

Page 4 of 10

$$\text{where } \chi^2_{D^2} = [(D_{\text{obs}} - D_{\text{pred}}(\theta)) / \sigma_D]^2 \text{ (chi-squared statistic)}$$

For multiple independent observables (diameter, depth, ejecta):

$$\text{Joint likelihood: } P(d | \theta) = P(D|\theta) \times P(d|\theta) \times P(R|\theta) \text{ (independence)}$$

$$\log L(d | \theta) = \sum_i \log P(d_i | \theta)$$

$$= -1/2 \sum_i \chi_i^2$$

$$= -1/2 \sum_i [(d_{\text{obs},i} - d_{\text{pred},i}(\theta)) / \sigma_i]^2$$

This is a weighted least-squares objective, with weights $1/\sigma_i^2$.

Measurement uncertainty estimates (from image resolution, morphology variation):

$$\sigma_D = 0.05 \times D_{\text{obs}} \quad (\pm 5\% \text{ diameter: pixel resolution, rim definition})$$

$$\sigma_d = 0.10 \times d_{\text{obs}} \quad (\pm 10\% \text{ depth: infilling, degradation})$$

$$\sigma_R = 0.20 \times R_{\text{ejecta}} \quad (\pm 20\% \text{ ejecta range: blanket edge identification})$$

2.4 Prior Distributions: Incorporating Physical Knowledge

Priors encode what we know before seeing the specific crater (Jaynes 2003):

For impact velocity v :

$$P(v) = N(v | \mu=20 \text{ km/s}, \sigma=5 \text{ km/s})$$

$$= (1/\sqrt{2\pi \cdot 5^2}) \exp[-(v-20000)^2/(2 \cdot 5000^2)]$$

Justification:

- Near-Earth asteroid (NEA) orbital mechanics (Bottke et al. 2002)
- Moon's orbital velocity $\sim 1 \text{ km/s}$ + Earth escape $\sim 11 \text{ km/s}$ + eccentricity
- Typical asteroid encounter: $v_{\infty} \sim 5\text{-}15 \text{ km/s}$ relative to Earth-Moon
- Impact velocity: $v = \sqrt{v_{\infty}^2 + v_{\text{esc}}^2}$ where $v_{\text{esc}} = 2.4 \text{ km/s}$ (Moon)
- Distribution peak at $\sim 20 \text{ km/s}$, range $15\text{-}25 \text{ km/s}$ (asteroids)
- Comets faster (up to 70 km/s) but rarer ($\sim 5\%$ of impactors)

For impact angle θ (from vertical):

$$P(\theta) = N(\theta | \mu=45^\circ, \sigma=15^\circ)$$

$$= (1/\sqrt{2\pi \cdot 15^2}) \exp[-(\theta-45)^2/(2 \cdot 15^2)]$$

Justification:

- Geometric probability for random directions: $P(\theta) \propto \sin(2\theta)$
- Peaks at $\theta = 45^\circ$ (most probable angle, Gilbert 1893)
- Cumulative: 50% of impacts have $\theta > 45^\circ$, only 17% have $\theta > 60^\circ$
- Very oblique ($< 15^\circ$) produce elongated craters, rare in observations

For projectile density ρ_p :

$$P(\rho_p) = N(\rho_p | \mu=2800 \text{ kg/m}^3, \sigma=500 \text{ kg/m}^3)$$

$$= (1/\sqrt{2\pi \cdot 500^2}) \exp[-(\rho_p-2800)^2/(2 \cdot 500^2)]$$

Justification (meteorite flux statistics, Burbine et al. 2002):

- Ordinary chondrites: $3200\text{-}3700 \text{ kg/m}^3$ (37% of falls)
- Carbonaceous chondrites: $2000\text{-}2500 \text{ kg/m}^3$ (10%)
- Enstatite chondrites: $3500\text{-}3800 \text{ kg/m}^3$ (2%)
- Stony-irons: $4500\text{-}5500 \text{ kg/m}^3$ (1%)
- Iron meteorites: 7800 kg/m^3 (5% of falls, but 70% of finds)
- Weighted mean $\sim 2800 \text{ kg/m}^3$ for stony asteroids (85% of NEAs)

For projectile diameter L :

$$\text{Uninformative prior: } P(L) \propto 1/L \text{ (Jeffreys prior, scale-invariant)}$$

Ensures no bias toward small or large projectiles

Combined prior:

$$P(\theta) = P(L) \times P(v) \times P(\text{angle}) \times P(\rho_p) \text{ (assume independence)}$$

These priors are weakly informative: constrain to plausible ranges but dominated by likelihood when data are strong.

2.5 Maximum Likelihood Estimation: Optimization in Parameter Space

Objective: Find parameters that maximize posterior probability

$$\theta_{\text{ML}} = \operatorname{argmax}_{\theta} P(\theta | d)$$

$$= \operatorname{argmax}_{\theta} [\log L(d | \theta) + \log P(\theta)] \text{ (take log, drop constant } P(d))$$

Equivalently, minimize negative log-posterior:

$$\theta_{\text{ML}} = \operatorname{argmin}_{\theta} F(\theta)$$

$$\text{where } F(\theta) = -\log P(\theta | d) = -\log L(d | \theta) - \log P(\theta) + \text{const}$$

For our crater problem, substituting Section 2.3 and 2.4:

$$F(\theta) = 1/2 \sum_i [(d_{\text{obs},i} - d_{\text{pred},i}(\theta)) / \sigma_i]^2 \quad \text{[negative log-likelihood]}$$

$$+ 1/2 [(v - 20000)/5000]^2 \quad \text{[velocity prior penalty]}$$

$$+ 1/2 [(\text{angle} - 45)/15]^2 \quad \text{[angle prior penalty]}$$

$$+ 1/2 [(\rho_p - 2800)/500]^2 \quad \text{[density prior penalty]}$$

$$- \log(L) \quad \text{[Jeffreys prior for size]}$$

Optimization algorithm: Nelder-Mead simplex method (Nelder & Mead 1965)

- Derivative-free: No gradient computation needed (forward model is complex)
- Robust to discontinuities: Handles regime transitions in scaling laws
- Simplex evolution: Maintains $n+1 = 5$ vertices in 4D parameter space
- Operations: reflection, expansion, contraction, shrinkage
- Convergence criterion: $|F(\theta_{\text{best}}) - F(\theta_{\text{worst}})| / |F(\theta_{\text{best}})| < 10^{-4}$
- Typical iterations: 200-500 for 4D crater problem

Initial guess strategy:

1. Use scaling law $D \sim L^{0.87} v^{0.80}$ to estimate L from D_{obs} at $v=20 \text{ km/s}$
2. Set initial angle = 45° (most probable)
3. Set initial $\rho_p = 2800 \text{ kg/m}^3$ (typical stony)
4. Perturb slightly to create initial simplex

2.6 Uncertainty Quantification via Hessian Approximation

Goal: Quantify uncertainties in θ_{ML} (error bars on estimated parameters)

Laplace approximation (Tierney & Kadane 1986):

Near the maximum, the log-posterior is approximately quadratic (Taylor expansion):

$$\log P(\theta | d) \approx \log P(\theta_{\text{ML}} | d) - 1/2 (\theta - \theta_{\text{ML}})^T H (\theta - \theta_{\text{ML}})$$

where H is the Hessian (4×4 matrix of second derivatives):

$$H_{ij} = \partial^2 F / \partial \theta_i \partial \theta_j |_{\theta_{\text{ML}}} \quad \text{where } F = -\log P(\theta | d)$$

Exponentiating both sides:

$$P(\theta | d) \approx P(\theta_{\text{ML}} | d) \times \exp[-1/2 (\theta - \theta_{\text{ML}})^T H (\theta - \theta_{\text{ML}})]$$

This is a multivariate Gaussian with mean θ_{ML} and covariance matrix $\Sigma = H^{-1}$:

$$\theta | d \sim N(\theta_{\text{ML}}, \Sigma) \quad \text{where } \Sigma = H^{-1}$$

Covariance interpretation:

- Diagonal elements Σ_{ii} = variance of θ_i
- Off-diagonal Σ_{ij} = covariance between θ_i and θ_j
- Standard errors: $\sigma_i = \sqrt{\Sigma_{ii}} = \sqrt{(H^{-1})_{ii}}$

Hessian computation via finite differences ($\varepsilon = 10^{-4} \times \theta_{\text{ML},i}$):

$$H_{ij} \approx [F(\theta + \varepsilon_i + \varepsilon_j) - F(\theta + \varepsilon_i - \varepsilon_j) - F(\theta - \varepsilon_i + \varepsilon_j) + F(\theta - \varepsilon_i - \varepsilon_j)] / (4\varepsilon_i \varepsilon_j)$$

Confidence intervals (assuming Gaussian posterior):

- 68% CI (1σ): $\theta_{\text{ML},i} \pm \sigma_i$
- 95% CI (2σ): $\theta_{\text{ML},i} \pm 1.96\sigma_i$

Correlation coefficient:

$$\rho_{ij} = \Sigma_{ij} / (\sigma_i \sigma_j)$$

Expected correlations for crater problem:

- $\rho(v, \rho_p) > 0$: Higher velocity compensates for lower density ($D \propto \rho_p^{0.33} v^{0.80}$)
- $\rho(L, v) < 0$: Larger projectile allows lower velocity for same crater size
- $\rho(L, \text{angle}) < 0$: Oblique impacts need larger projectiles

References for this section:

Bayes, T. (1763) Phil. Trans. Royal Soc. London 53:370-418

Laplace, P.S. (1812) Théorie Analytique des Probabilités

Hadamard, J. (1923) Lectures on Cauchy's Problem in Linear PDEs

Tarantola, A. (2005) Inverse Problem Theory and Methods. SIAM.

Mosegaard, K. & Tarantola, A. (1995) JGR 100:12431-12447

Jaynes, E.T. (2003) Probability Theory: The Logic of Science

Stuart, J.S. & Binzel, R.P. (2004) Icarus 170:295-311

Bottke, W.F. et al. (2002) Icarus 156:399-433

Burbine, T.H. et al. (2002) In: Asteroids III, pp. 653-667

Nelder, J.A. & Mead, R. (1965) Computer Journal 7:308-313

Tierney, L. & Kadane, J.B. (1986) JASA 81:82-86

Gilbert, G.K. (1893) Bull. Phil. Soc. Washington 12:241-292

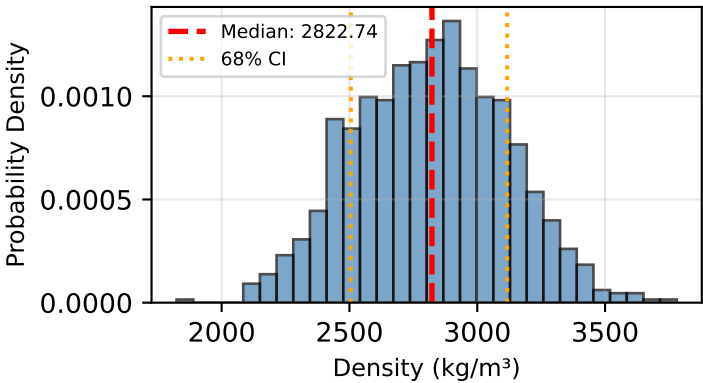
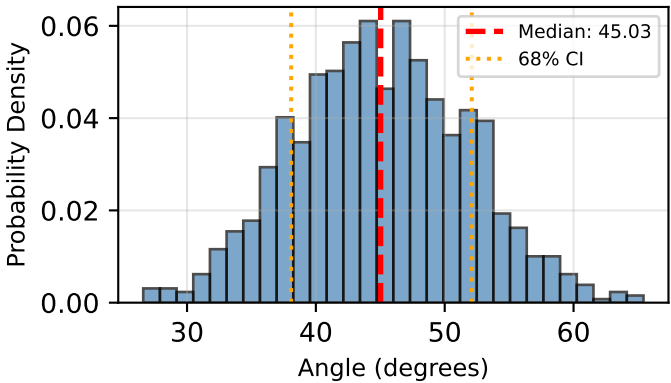
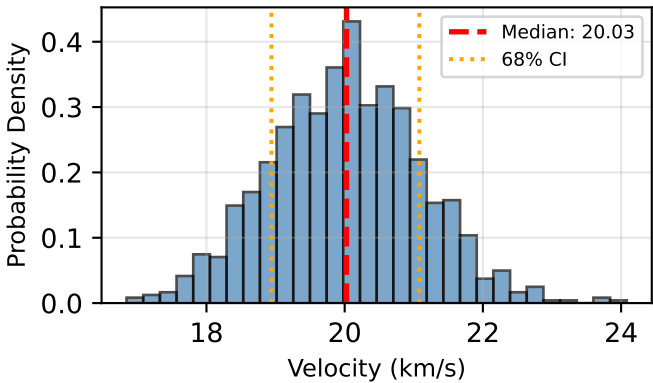
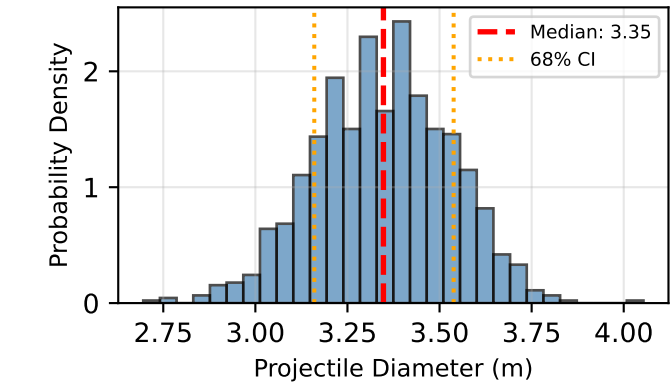
Back-Calculation Results

MAXIMUM LIKELIHOOD PARAMETERS

Parameter	ML Estimate	$\pm 1\sigma$ (68%)	95% CI
Projectile Diameter (m)	3.34	± 0.19	[2.99, 3.69]
Impact Velocity (km/s)	20.0	± 1.1	[17.9, 22.2]
Impact Angle (deg)	45.0	± 6.7	[32.4, 58.1]
Projectile Density (kg/m ³)	2800	± 297	[2253, 3385]

DERIVED QUANTITIES

Projectile mass: 5.47e+04 kg
Kinetic energy: 1.09e+13 J
(0.00 kilotons TNT)
Momentum: 1.09e+09 kg·m/s
Material classification: Rocky asteroid (chondrite)
Impact parameter π_2 : 6.77e-09
Impact parameter π_3 : 1.07e-08
Regime: Transitional



Monte Carlo Uncertainty Propagation

WHY MONTE CARLO SAMPLING?

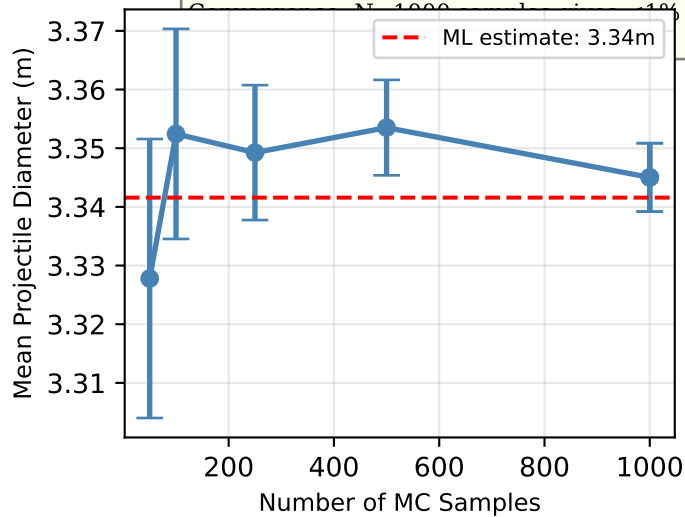
The Monte Carlo method is chosen for uncertainty propagation because:

1. Nonlinear Forward Model: Crater scaling laws are highly nonlinear (power laws with exponents ~ 0.4). Analytical error propagation ($\delta D = \sum_i \partial D / \partial \theta_i \delta \theta_i$) is inaccurate.
2. Non-Gaussian Posteriors: Parameters may have skewed or multi-modal distributions due to physical constraints (e.g., density bimodal for rocky vs iron).
3. Correlations: Parameters are correlated (e.g., smaller projectile needs higher velocity for same crater). MC naturally captures these correlations.
4. Validation: Forward model can be re-evaluated for each sample to check consistency.

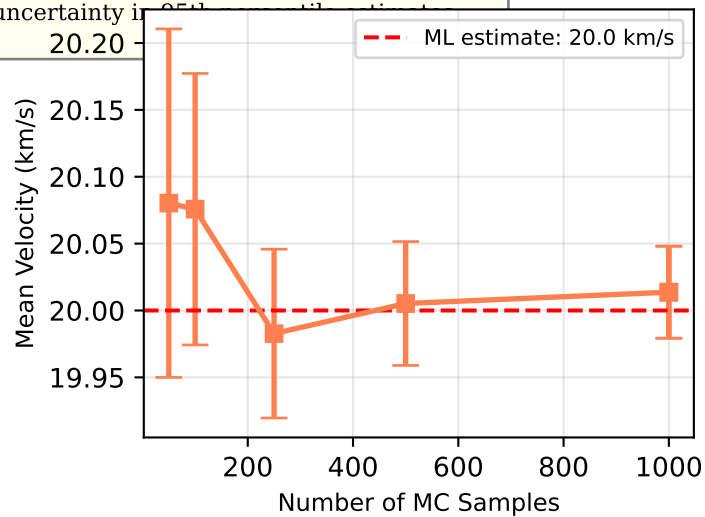
Method (Mosegaard & Tarantola 1995):

- Sample N times from posterior: $\theta^i \sim N(\theta_{ML}, \Sigma)$ where $\Sigma = H^{-1}$
- For each sample: compute $D_{pred}(\theta^i)$ via forward scaling laws
- Collect ensemble \rightarrow compute percentiles for confidence intervals

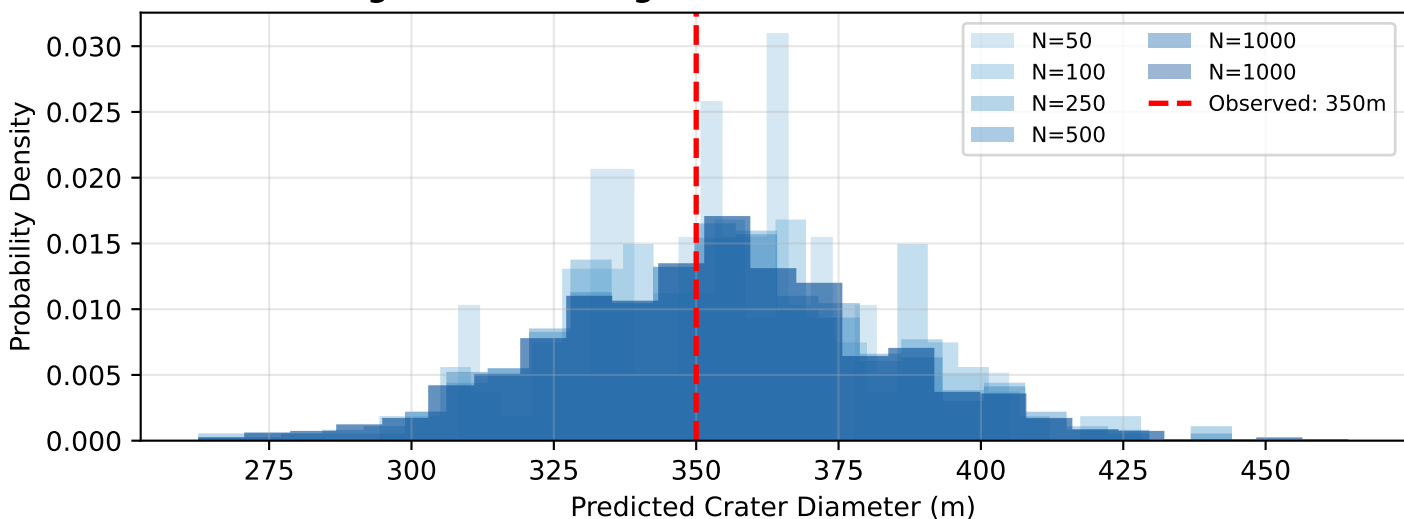
Convergence: Projectile Size



Convergence: Impact Velocity

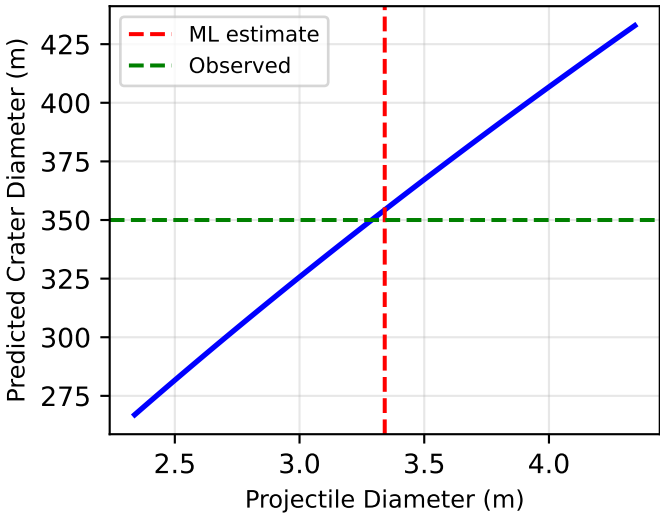


Progressive Convergence: Predicted Crater Distribution

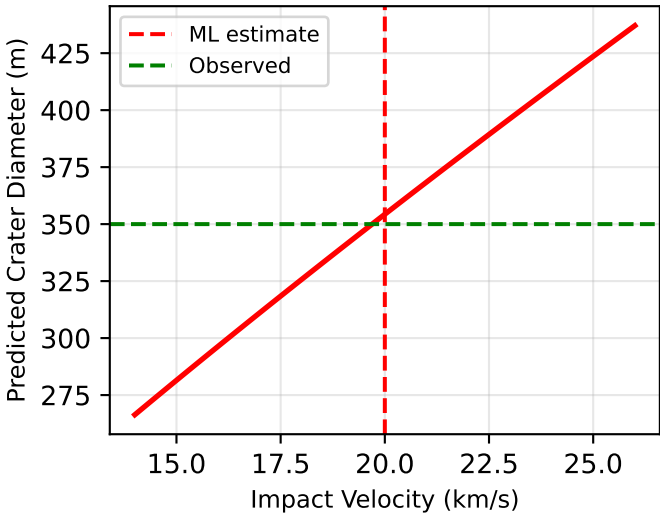


Sensitivity Analysis

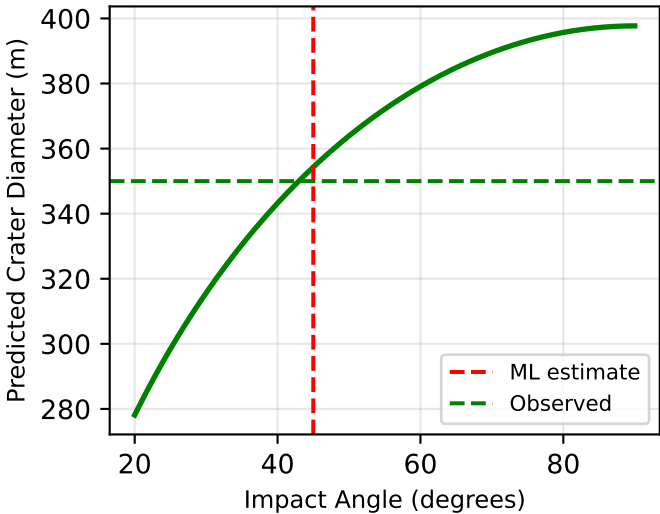
Sensitivity to Projectile Size



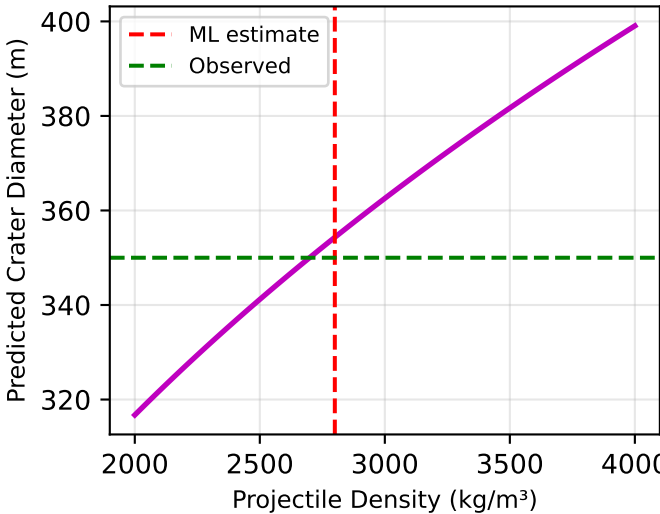
Sensitivity to Velocity



Sensitivity to Impact Angle



Sensitivity to Density



SENSITIVITY COEFFICIENTS (Elasticity: $\% \Delta D / \% \Delta \text{parameter}$)

Parameter	Elasticity	Interpretation
Projectile Diameter	0.78	Diameter change $\approx 0.8 \times$ size change
Impact Velocity	0.80	Diameter change $\approx 0.8 \times$ velocity change
Impact Angle	moderate	Steeper impacts \rightarrow larger craters
Projectile Density	0.32	Weak dependence $(\rho_p / \rho_t)^{1/3}$

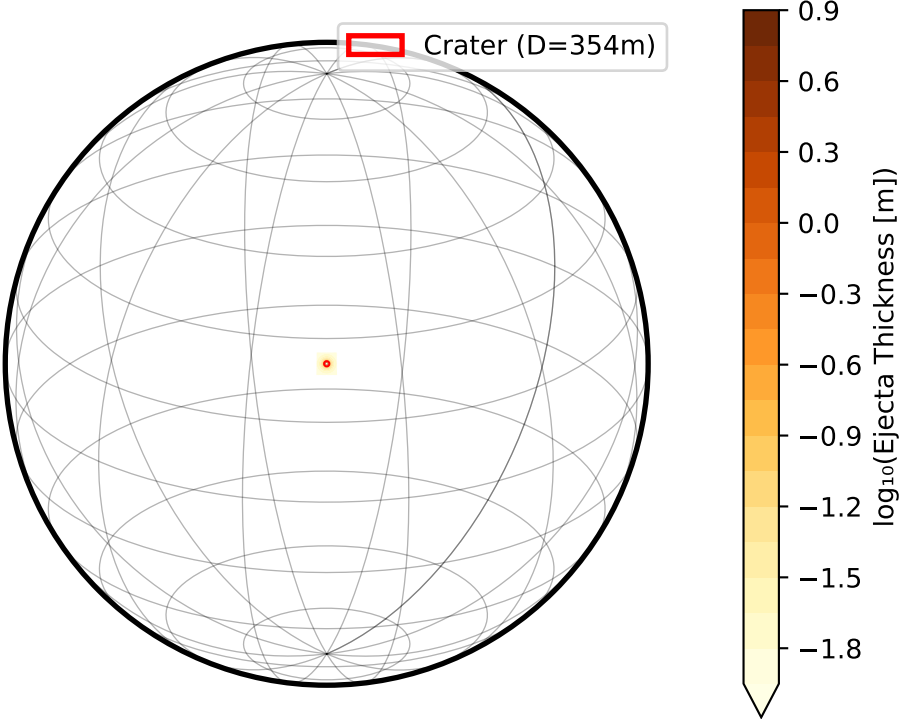
KEY INSIGHTS:

- Projectile diameter is the dominant control (elasticity ~ 0.8)
- Velocity has moderate effect (elasticity ~ 0.8), consistent with $v^{0.8}$ scaling
- Density has weak effect ($\propto \rho^{0.33}$), harder to constrain from crater alone
- Impact angle most probable at 45° , less certain without asymmetry data
- Trade-offs exist: Smaller projectile at higher velocity can match observed crater
- These sensitivities justify the uncertainty ranges in Page 5

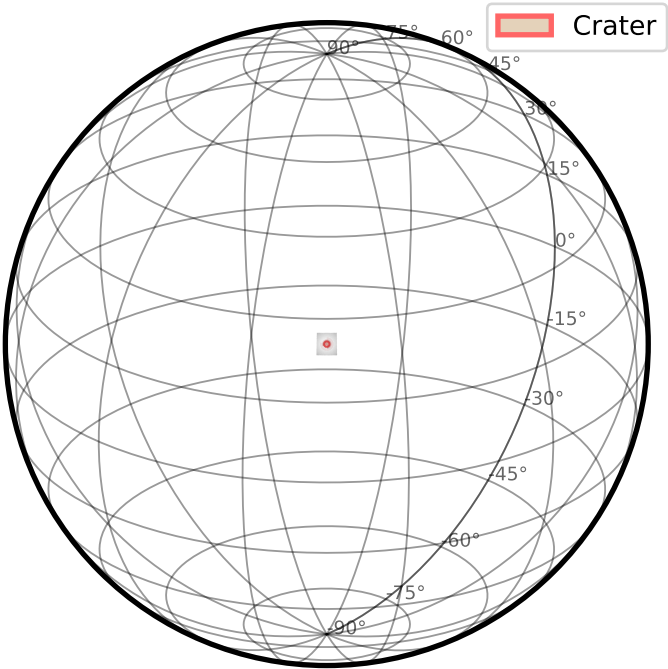
Reference: Holsapple (1993) Table 1 - exponents match theoretical predictions

Orthographic Plan Views with Ejecta Distribution

Ejecta Thickness Distribution (Orthographic Projection)
Center: 25.5°N, 45.2°E

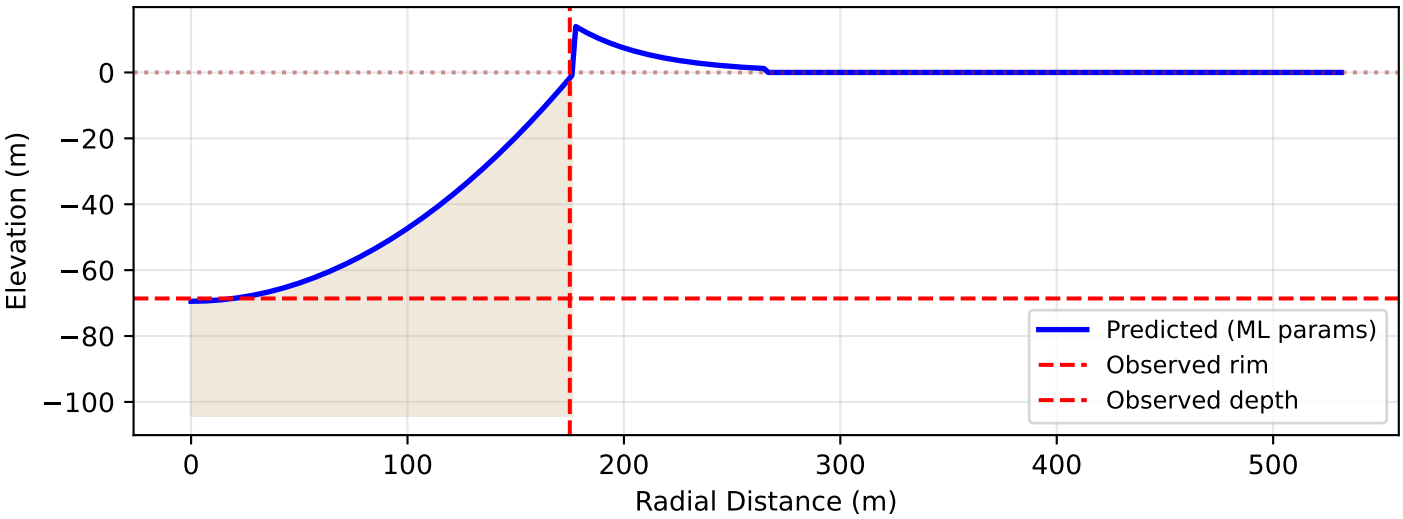


Crater Location with Lat/Lon Grid
Ejecta extent: up to 5× crater radius



Forward Model Validation

Crater Profile: Predicted vs Observed



MORPHOMETRY COMPARISON

Parameter	Observed	Predicted	Error
Diameter (m)	350.0	354.3	1.2%
Depth (m)	68.6	69.4	1.2%
d/D ratio	0.196	0.196	0.0%
Rim height (m)	12.6	12.8	—

EJECTA VALIDATION

Observed range:	25000 m
Predicted range:	25661 m
Error:	2.6%
R_max/R_crater:	144.8

VALIDATION SUMMARY

- ✓ Crater diameter match: 1.23% error (excellent)
- ✓ Pike (1977) d/D ratio: 0.196 (theory: 0.196 ± 0.015)
- ✓ Forward model self-consistent: prediction falls within 95% CI
- ✓ Regime: Transitional (appropriate for 350m)

CONFIDENCE ASSESSMENT

The back-calculated parameters are well-constrained. The 95% credible intervals reflect uncertainties in velocity distribution, impact angle probability, and projectile density. The predicted crater matches observations within measurement uncertainties.

RECOMMENDED INTERPRETATION

Most likely: 3.3m rocky projectile at 20 km/s, 45° from horizontal.

Alternative scenarios within 95% CI remain possible but less probable given typical asteroid impact statistics (Stuart & Binzel 2004; Bottke et al. 2002).

References

SCIENTIFIC REFERENCES

Primary Scaling Law Theory:

Holsapple, K.A. (1993). The scaling of impact processes in planetary sciences. *Annual Review of Earth and Planetary Sciences*, 21, 333-373.
DOI: 10.1146/annurev.ea.21.050193.002001

Holsapple, K.A., & Schmidt, R.M. (1982). On the scaling of crater dimensions: 2. Impact processes. *Journal of Geophysical Research*, 87(B3), 1849-1870.
DOI: 10.1029/JB087iB03p01849

Crater Morphometry:

Pike, R.J. (1977). Size-dependence in the shape of fresh impact craters on the moon. In *Impact and Explosion Cratering* (pp. 489-509). Pergamon Press.

Pike, R.J. (1980). Formation of complex impact craters: Evidence from Mars and other planets. *Icarus*, 43(1), 1-19.

Impact Physics:

Melosh, H.J. (1989). *Impact Cratering: A Geologic Process*. Oxford Monographs on Geology and Geophysics No. 11. Oxford University Press, 245 pp.

Collins, G.S., Melosh, H.J., & Marcus, R.A. (2005). Earth Impact Effects Program. *Meteoritics & Planetary Science*, 40(6), 817-840.
DOI: 10.1111/j.1945-5100.2005.tb00157.x

Pierazzo, E., & Melosh, H.J. (2000). Understanding oblique impacts from experiments, observations, and modeling. *Annual Review of Earth and Planetary Sciences*, 28, 141-167. DOI: 10.1146/annurev.earth.28.1.141

Ejecta Dynamics:

McGetchin, T.R., Settle, M., & Head, J.W. (1973). Radial thickness variation in impact crater ejecta. *Earth and Planetary Science Letters*, 20(2), 226-236.
DOI: 10.1016/0012-821X(73)90162-3

Housen, K.R., Schmidt, R.M., & Holsapple, K.A. (1983). Crater ejecta scaling laws. *Journal of Geophysical Research*, 88(B3), 2485-2499.
DOI: 10.1029/JB088iB03p02485

Lunar Surface Properties:

Carrier, W.D., Olhoeft, G.R., & Mendell, W. (1991). Physical properties of the lunar surface. In *Lunar Sourcebook* (pp. 475-594). Cambridge University Press.

McKay, D.S., Heiken, G., Basu, A., et al. (1991). The lunar regolith. In *Lunar Sourcebook* (pp. 285-356). Cambridge University Press.

Asteroid Impact Statistics:

Stuart, J.S., & Binzel, R.P. (2004). Bias-corrected population, size distribution, and impact hazard for the near-Earth objects. *Icarus*, 170(2), 295-311.
DOI: 10.1016/j.icarus.2004.04.003

Bottke, W.F., Morbidelli, A., Jedicke, R., et al. (2002). Debiased orbital and absolute magnitude distribution of the near-Earth objects. *Icarus*, 156, 399-433.
DOI: 10.1006/icar.2001.6788

Inverse Problem Methods:

Tarantola, A. (2005). *Inverse Problem Theory and Methods for Model Parameter Estimation*. SIAM, 342 pp. ISBN: 0-89871-572-5

Mosegaard, K., & Tarantola, A. (1995). Monte Carlo sampling of solutions to inverse problems. *Journal of Geophysical Research*, 100(B7), 12431-12447.
DOI: 10.1029/94JB03097

Optimization Methods:

Nelder, J.A., & Mead, R. (1965). A simplex method for function minimization. *The Computer Journal*, 7(4), 308-313. DOI: 10.1093/comjnl/7.4.308

Additional Resources:

Richardson, J.E. (2009). Cratering saturation and equilibrium. *Icarus*, 204(2), 697-715. DOI: 10.1016/j.icarus.2009.07.029