

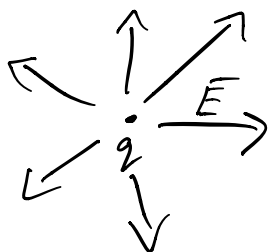
Point charges and the Delta-Function

Divergence theorem assumes that the vector field \vec{A} and its first derivative are continuous (Arfken + Weber)

$$\int_V \vec{\nabla} \cdot \vec{A} \, dv = \oint_S \vec{A} \cdot \hat{n} \, da$$

For a point charge, the electric field is

$$\vec{E} = \frac{q}{r^2} \hat{e}_r$$



What happens as $r \rightarrow 0$?

Is $\frac{\partial \vec{E}}{\partial r}$ continuous at $r=0$?

What is $\nabla \cdot \vec{E}$ for a point charge?

$$\nabla \cdot \vec{E} = \frac{1}{r^2} \frac{d}{dr} (r^2 E_r)$$

$$= \frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{q}{r^2} \right) = \frac{1}{r^2} \frac{d}{dr} (q) = \frac{0}{r^2}$$

For $r \neq 0$, $\nabla \cdot \vec{E} = 0$.

What happens at $r=0$?

From the divergence theorem, we get

$$\int_V \nabla \cdot \vec{E} dV = \oint_S \vec{E} \cdot d\vec{a}$$

↙ Includes origin

Gauss' Law tells us

$$\oint \vec{E} \cdot d\vec{a} = 4\pi q_{enc}$$

$$\Rightarrow \int_V \nabla \cdot \vec{E} dV = \oint \vec{E} \cdot d\vec{a} = 4\pi q_{enc}$$

For a point charge, $q_{enc} = q$ (non-zero)

but $\nabla \cdot \vec{E} = 0$ for $r \neq 0$

→ How does this work?

$\nabla \cdot \vec{E}$ is singular at the origin

and zero everywhere else, and

when integrated it gives a finite

value

→ Answer = Dirac delta function

Dirac delta-function

$$\delta(x-x') = 0 \quad \text{for } x \neq x'$$

$$\int_{-\infty}^{\infty} \delta(x-x') dx = 1$$

$\Rightarrow \delta$ is zero everywhere except
 $x=x'$ where it is infinite

$$\int_{-\infty}^{\infty} f(x) \delta(x-x') dx = f(x')$$

In 3-D:

$$\int_V \delta(\vec{r}-\vec{r}') dV = 1$$

$$\int_V f(\vec{r}) \delta(\vec{r}-\vec{r}') dV = f(\vec{r}')$$

Treating our point charge as a delta-function

$$\rho = q \delta(\vec{r})$$

charge density located at $\vec{r} = \vec{0}$

From Gauss' Law

$$\vec{\nabla} \cdot \vec{E} = 4\pi \rho$$

$$\Rightarrow \int_V \vec{\nabla} \cdot \vec{E} dV = \int_V 4\pi \rho dV$$

$$\Rightarrow \int_V \vec{\nabla} \cdot \frac{q \hat{e}_r}{r^2} dV = \int_V 4\pi q \delta(\vec{r}) dV$$

divide by q

$$\Rightarrow \int_V \vec{\nabla} \cdot \frac{\hat{e}_r}{r^2} dV = \int_V 4\pi \delta(\vec{r}) dV$$

This holds over an arbitrary volume V , so we can say

$$\bar{\nabla} \cdot \frac{\hat{e}_r}{r^2} = 4\pi \delta(\bar{r})$$

Alternatively, in terms of the Laplacian (which comes from the same derivation but starting with $\bar{E} = -\bar{\nabla} \Phi$)

$$\bar{\nabla}^2 \left(\frac{1}{r} \right) = -4\pi \delta(\bar{r})$$

- So, even though $\frac{\partial \bar{E}}{\partial r}$ is discontinuous, we're able to use the Divergence theorem and integrate over the singularity because we get a finite result from the Dirac delta-function.

- There are methods for integrating over discontinuities (see Lebesgue Integration for example) but they are outside the scope of this course.
- For this class, the divergence theorem holds for point charges even though there's a discontinuity