

Boundary Conditions

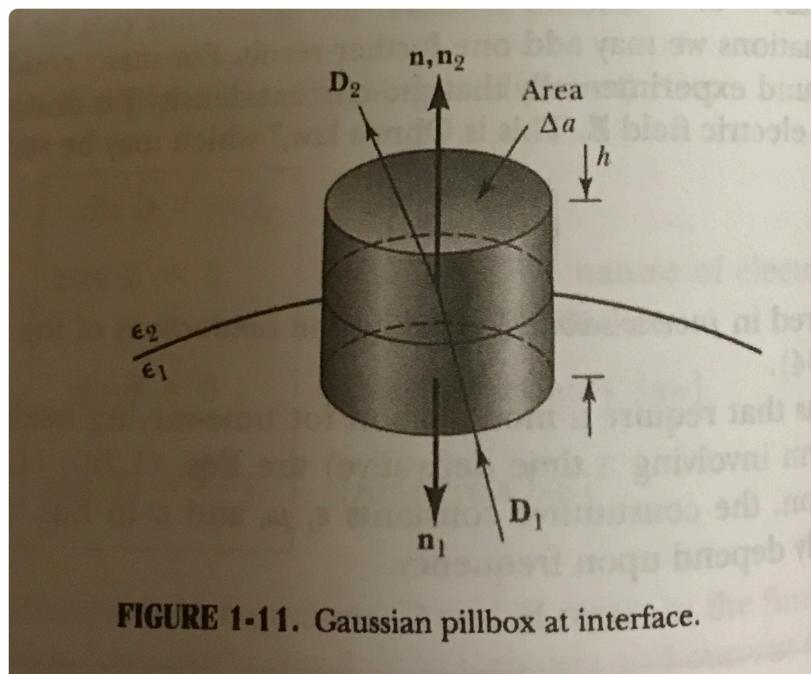
What happens to the field vectors at a boundary?

Starting with the Electric/Displacement Field:

From Gauss' Law

$$\oint_S \vec{D} \cdot \hat{n} d\vec{a} = 4\pi q_f$$

→ The total flux of \vec{D} through a closed surface S depends on the amount of free charge q_f



For a \vec{D} directed from a material with dielectric constant of ϵ_1 to a material with ϵ_2 (figure above), we construct a Gaussian pillbox with area Δa at the top/bottom. $h \rightarrow 0$ so the side flux is ignored.

$$\Rightarrow \oint_S \vec{D} \cdot \hat{n} da = [\vec{D}_2 \cdot \hat{n} + \vec{D}_1 \cdot (-\hat{n})] \Delta a$$

$$= 4\pi (\rho_s)_f \Delta a$$

✓

surface charge density at interface
(charge/area)

$$\Rightarrow (\vec{D}_2 - \vec{D}_1) \cdot \hat{n} = 4\pi (\rho_s)_f$$

Boundary condition on normal component of \vec{D}

In terms of the electric field, this becomes:

$$(\epsilon_2 \bar{E}_2 - \epsilon_1 \bar{E}_1) \cdot \hat{n} = 4\pi (\rho_s)_f$$

If $(\rho_s)_f = 0$ so there is no free charge on the boundary, what does that say about \bar{D}_1 and \bar{D}_2 ? What about \bar{E}_1 and \bar{E}_2 ?

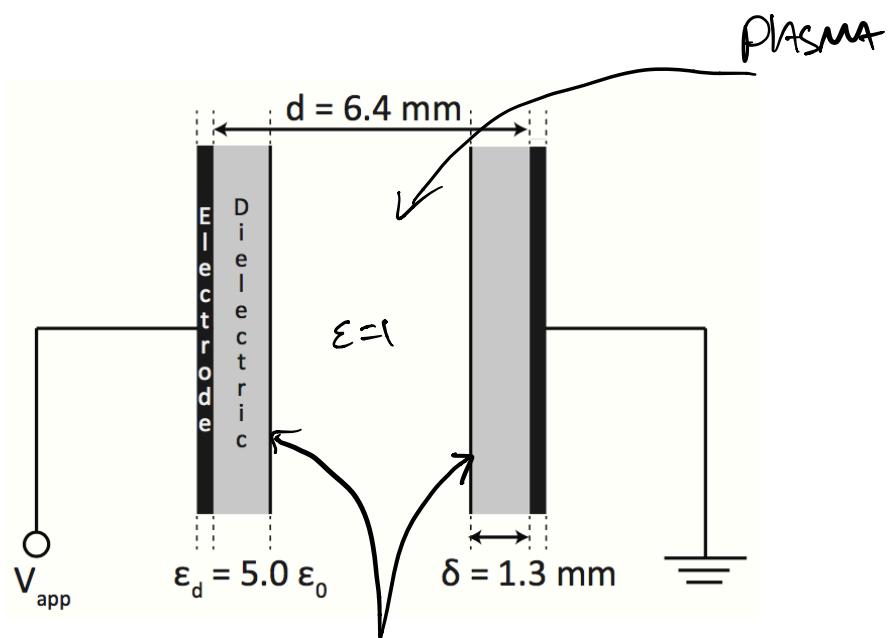
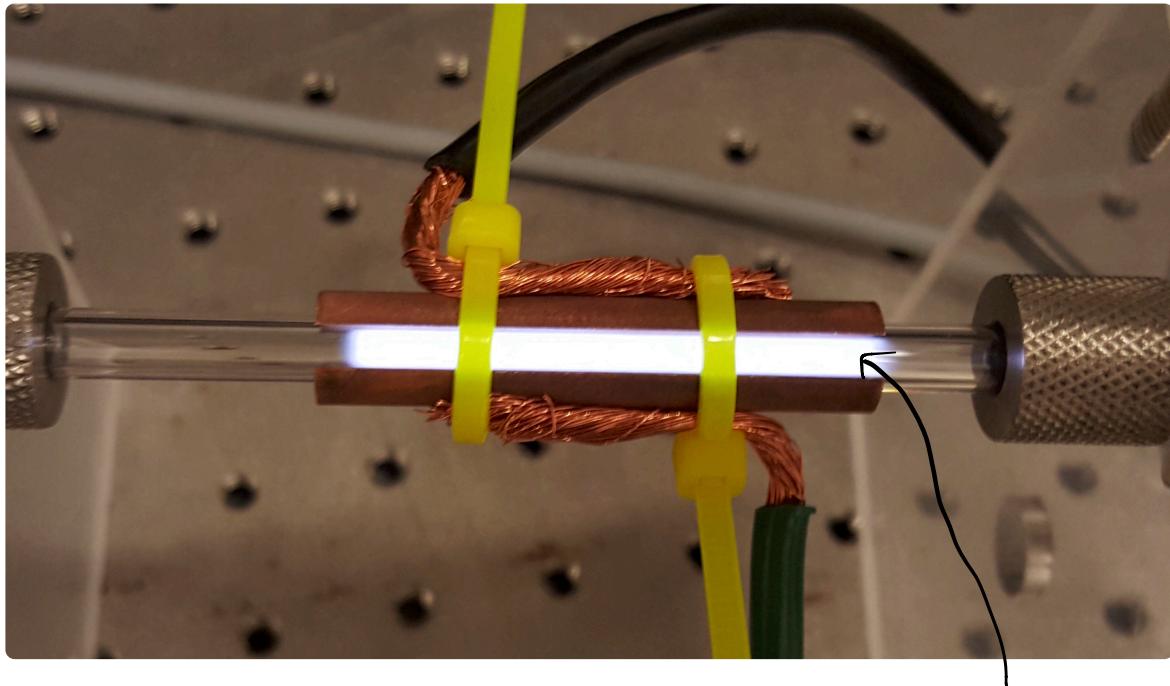
$$\rightarrow \bar{D}_1 = \bar{D}_2 \text{ or } \epsilon_2 \bar{E}_2 = \epsilon_1 \bar{E}_1 \text{ if } (\rho_s)_f = 0$$



continuous across boundary

Example :

Radio Frequency Dielectric Barrier Discharge



What is the electric field at the boundaries?

As the plasma forms, current flows to the dielectric boundaries, creating a surface charge:

$$\hat{n} \cdot (\bar{E}_p - \epsilon_d \bar{E}_d) = 4\pi (\varrho_s)_f$$

✓ field in plasma (vacuum)
✓ dielectric
✓ surface charge density

$$\frac{d(\varrho_s)_f}{dt} = \hat{n} \cdot \bar{\jmath}_e + \hat{n} \cdot \bar{\jmath}_i$$

surface charge density changes over time depending on the electron and ion current densities, $\bar{\jmath}_e$ and $\bar{\jmath}_i$

How does current flow to/froam the electrodes?

Tangential Boundary Condition

From $\vec{\nabla} \times \vec{E} = 0$ we can integrate around a Stokesian loop Γ to obtain:

$$\oint_S \vec{D} \times \vec{E} \cdot \hat{n}_o \, da = \oint_{\Gamma} \vec{E} \cdot d\vec{l} = \sigma$$

normal to surface S Γ

Stoke's theorem

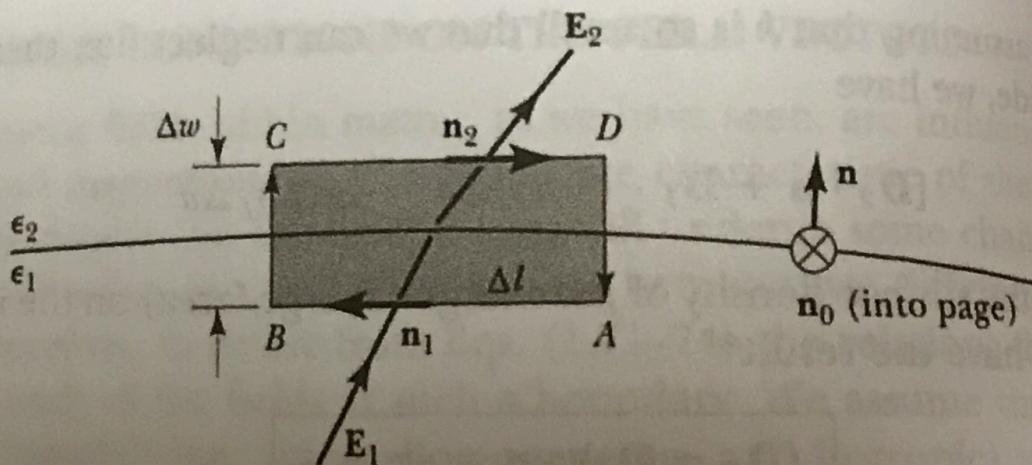


FIGURE 1-12. Stokesian rectangle at interface.

For $\Delta w \rightarrow 0$, we neglect the contribution from the sides

$$\Rightarrow [\bar{E}_1 \cdot \hat{n}_1 + \bar{E}_2 \cdot \hat{n}_2] \Delta l = 0$$

The unit vectors follows

$$\hat{n}_2 = -\hat{n}_1 = \hat{n}_0 \times \hat{n}$$

↓

normal to
surface

normal to
boundary

$$\Rightarrow -\bar{E}_1 \cdot \hat{n}_2 + \bar{E}_2 \cdot \hat{n}_2 = (\bar{E}_2 - \bar{E}_1) \cdot \hat{n}_2 = 0$$

$$\Rightarrow (\bar{E}_2 - \bar{E}_1) \cdot (\hat{n}_0 \times \hat{n}) = 0$$

From the vector identity

$$\bar{C} \cdot (\bar{A} \times \bar{B}) = \bar{A} \cdot (\bar{B} \times \bar{C})$$

$$\Rightarrow \hat{n}_o \cdot [\hat{n} \times (\bar{E}_2 - \bar{E}_1)] = 0$$

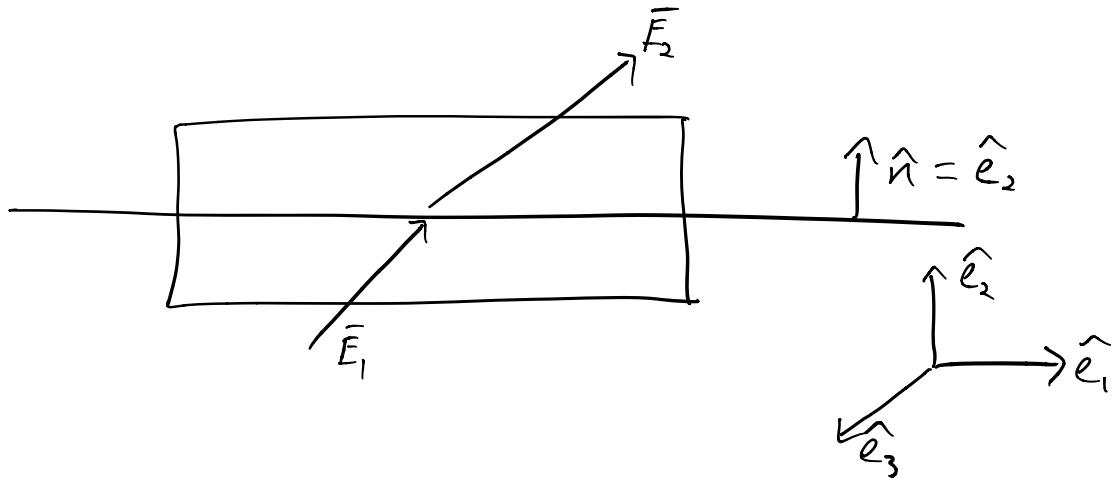
Also, $\bar{A} \times \bar{B} = -\bar{B} \times \bar{A}$

$$\Rightarrow -\hat{n}_o \cdot [(\bar{E}_2 - \bar{E}_1) \times \hat{n}] = 0$$

Because \hat{n}_o is non-zero with arbitrary orientation, we must have

$$(\bar{E}_2 - \bar{E}_1) \times \hat{n} = 0$$

★ Tangential components of \bar{E} are continuous across a boundary



Suppose $\bar{E}_2 - \bar{E}_1$ has a component in the \hat{e}_1 tangential direction with magnitude ΔE

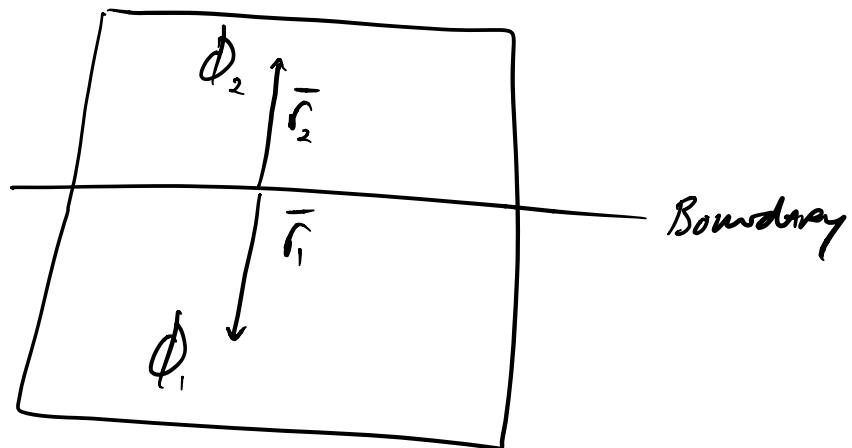
$$\bar{E}_2 - \bar{E}_1 = \Delta E \hat{e}_1$$

$$\Rightarrow (\bar{E}_2 - \bar{E}_1) \times \hat{n} = \Delta E \hat{e}_1 \times \hat{e}_2 = \Delta E \hat{e}_3$$

For $(\bar{E}_2 - \bar{E}_1) \times \hat{n} = 0$ to hold, ΔE must be zero

\Rightarrow tangential component of \bar{E} is continuous

From $\Delta\phi = \phi_2 - \phi_1 = - \int_{r_1}^{r_2} \bar{E} \cdot d\bar{l}$



As $\bar{r}_2 - \bar{r}_1 \rightarrow 0$ and the two points became continuous, $\Delta\phi \rightarrow 0$

$$\lim_{\bar{r}_2 \rightarrow \bar{r}_1} - \int_{r_1}^{r_2} \bar{E} \cdot d\bar{l} = 0 = \Delta\phi$$

$$\Rightarrow \boxed{\phi_2 = \phi_1}$$

Potential is continuous across a boundary

Magnetic Field Boundary Conditions

Following the examples for the electric field:

From $\nabla \cdot \vec{B} = 0$ we have

$$\oint_S \vec{B} \cdot \hat{n} da = 0$$

Using the same Gaussian pillbox as before, the boundary condition on the normal component is

$$(\vec{B}_2 - \vec{B}_1) \cdot \hat{n} = 0$$

Normal components of magnetic field
are continuous across a boundary

For the tangential component we start from Ampère's law:

$$\bar{\nabla} \times \bar{H} = \frac{4\pi}{c} \bar{J}_f$$

$$\Rightarrow \oint_S \bar{\nabla} \times \bar{H} \cdot \hat{n}_0 \, da = \oint_{\Gamma} \bar{H} \cdot d\bar{l} = \int \frac{4\pi \bar{J}_f}{c} \cdot \hat{n}_0 \, da$$

Stokes theorem

$$= \frac{4\pi}{c} (\bar{I}_f)_{\text{link}}$$

↑
free current

Using the same Stokesian Loop as before, and assuming there's a surface current density \bar{K}_f (current/width) in the direction of \hat{n}_0 (normal to surface S enclosed by Stokesian Loop Γ)

$$(\bar{H}_2 - \bar{H}_1) \times \hat{n} = -\frac{4\pi}{c} \bar{K}_A$$

* The tangential component of the magnetic intensity \bar{H} is discontinuous if a surface current is present

$$\text{If } \bar{K}_A = 0 \rightarrow (\bar{H}_2)_+ = (\bar{H}_1)_+$$

↙
tangential to boundary

Using the magnetic field,

$$\left(\frac{\bar{B}_2}{\mu_2} - \frac{\bar{B}_1}{\mu_1} \right) \times \hat{n} = -\frac{4\pi}{c} \bar{K}_A$$

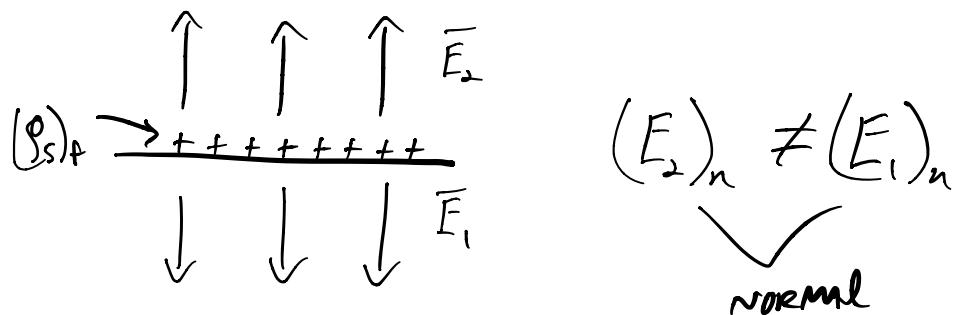
$$\text{If } \bar{K}_f = 0 \rightarrow \frac{(\bar{B}_2)_t}{\mu_2} = \frac{(\bar{B}_1)_t}{\mu_1}$$

$$\Rightarrow (\bar{B}_2)_t \neq (\bar{B}_1)_t \quad \text{unless } \mu_1 = \mu_2$$

To remember the BCs, think

Electric

If charge is on the surface:



Magnetic

If surface current is present:

