
Filament Recycling and Sustained Contractile Flows in a 2D Active Network Model

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Abstract

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Author Summary

In this paper, we develop and analyze a minimal model for 2D active networks based on the cortical cytoskeleton of eukaryotic embryos. Our model introduces a drag-like slip between cross-linked filaments as means to dissipate stored stress, generating a macroscopic effective viscosity. We further introduce an active friction to active stress from microscopic properties. We generate computational simulations based on the model, and demonstrate that active stress is sufficient to drive network contraction only temporarily. By introducing filament recycling, we are able to set up steady state flow profiles such those found in the cortex of developing embryos and migrating cells. The model is used to calculate phenomenological constants measured in prior experiments. Our analysis sheds insight on potential microscopic control parameters governing broad qualitative differences in 2D active networks.

Introduction

Active networks of cross-linked polymers are a class of materials with poorly understood but highly interesting properties. Cross-linked networks of cytoskeletal polymers have been a subject of great interest because of their importance as structural components of cells [2,3].

The fluid-like deformability of actin networks is an important but often over-looked aspect of cell mechanics At long timescales, the purely elastic behavior of cross-linked networks gives way to fluid-like stress relaxation. Additionally, fluid-like flows have been observed in a number of cellular processes [27–32]. In *in vitro*

studies, long timescale creep behaviors are thought to arise predominantly from the transient nature of filament binding for most biologically relevant cross-linkers [33–36]. While the importance of cross-link dynamics in determining the mechanical response of semi-flexible polymer networks has been known for at least 20 years [5], there is still a gap in our understanding of how microscopic cross-link unbinding relates to viscous flows.

Further introductory detail about the components of the actin cortex and how they turn over The dependence of network rheology on cross-link unbinding is an active subject of theoretical research [21].

Several theoretical methods have addressed cross-link binding and unbinding directly [6, 21] in analytical approaches that allowed well-constrained fits for specific cross-linkers. These theories have therefore focused conceptually at the level of the cross-linked filament and were extended analytically to macroscopic networks. In another approach, modelers have taken cross-links as extended springlike structures [37] that are able to bind and unbind in simulated filament networks. Finally, other more ambitious simulations have even sought to interrogate the effects of cross-link unbinding in combination with the more complex mechanics of filament bundles [34, 38].

Ultimately, the complexity of the many theoretical approaches that have been applied to this problem have made it difficult to distinguish what, if any, core physical mechanisms may be sufficient to explain the observed forms of stress relaxation. We believe that serious qualitative understanding can be generated by focusing on some of the common elements exhibited in the aforementioned literature.

Synopsis of mechanical picture and mentions of the few theoretical attempts to describe it Here, we introduce a coarse-grained representation of filament cross-linking in which cross-linked filaments which are able to slide past each other as molecular bonds form and rupture, akin to coarse-grained models of molecular friction [39–41]. This drag-like coupling has been shown to be an adequate approximation in the case of ionic cross-linking of actin [42, 43], and can be found in the theoretical basis of force-velocity curves for myosin bound filaments [44]. We propose that it will form a suitable bulk approximation in the presence of super molecular cross-links as well.

Importantly, this simplification allows us to extend our single polymer models to dynamical systems of larger network models for direct comparison between theory and modeling results. This level of coarse graining will therefore make it easier to understand classes of behavior for varying compositions of cross-linked filament networks. In addition, it allows us to compute a new class of numerical simulations efficiently, which gives us concrete predictions for behaviors in widely different networks with measurable dependencies on molecular details.

Introductory description of how our model is different and what we were able to conclude

Models

Deformable Filaments

We consider individual filaments as chains of springs with relaxed length l_s . The orientations of neighboring springs are linearly coupled. Filaments can therefore be represented as a sequence of nodes with positions \mathbf{x}_i and nearest neighbor interactions of the form

$$|F_{i,i+1}|_{\parallel} = -\mu \cdot \frac{|\mathbf{x}_{i+1} - \mathbf{x}_i| - l_s}{l_s} \quad (1)$$

$$|F_{i,i+2}|_{\perp} = -\frac{\kappa}{l_s^2} \cdot \text{acos} \left(\frac{|\mathbf{x}_{i+2} - \mathbf{x}_{i+1}|}{|\mathbf{x}_{i+1} - \mathbf{x}_i|} \cdot \frac{|\mathbf{x}_{i+1} - \mathbf{x}_i|}{|\mathbf{x}_{i+1} - \mathbf{x}_i|} \right) \quad (2)$$

where, μ represents an extensional modulus of a filament, and κ represents a bending modulus. This is essentially a discretized equivalent to a model of filaments with separable extensional and bending moduli as in [12]. We define the totally elastic force on a node as

$$\nabla \mathcal{H}_i = |F_{i,i+1}|_{\parallel} + |F_{i,i+2}|_{\perp} \quad (3)$$

Here, we take the extensional modulus as a composite quantities related to both filament and cross-linker compliance in a manner similar to a recently proposed effective medium theory [26]. In the limit of highly rigid cross-links and flexible filaments, our model reduces to the pure semi-flexible filament models of [12, 13]. In the opposite regime of nearly rigid filaments and highly flexible cross links, our method is still largely similar to the model of [26] in small strain regimes before any nonlinear cross link stiffening. However, in departure from those models, the magnitude of the force on interior cross-links in our model is still the same as those on the exterior. This is a simplification of the varying levels of strain that would actually be present in these cross-linkers as addressed in [26], but we choose to ignore the slight variation in favor of an approximated, global mean approach.

2D Network Formation

We choose to focus our attention on 2D networks both for their tractability as well as their relevance in the quasi-2D cytoskeletal cortex of many eukaryotic cells [27]. In addition, recent developments in 2D *in vitro* systems [45, 46], make 2D models all the more interesting as a renewed focus of study.

We follow a mikado model approach by initializing a minimal network of connected unstressed linear filaments in a rectangular 2D domain. We generate 2D networks of these semi-flexible filaments by laying down straight lines of length, L , with random position and orientation. We then assume that some fixed fraction of overlapping filaments become cross-linked (defined in) at their point of overlap.

Although real cytoskeletal networks may form with non-negligible anisotropy, we focus on isotropically initialized networks for simplicity. We define the density using the average distance between cross-links along a filament, l_c . A simple geometrical argument can then be used to derive the number of filaments filling a domain as a function of L and l_c [12]. Here, we use the approximation that the number of filaments needed to tile a rectangular domain of size $W \times H$ is $2WH/Ll_c$, and that the length density is therefore $1/l_c$.

In the absence of cross-link slip, we expect the network to form a connected solid with a well defined elastic modulus [12, 13]. These networks are only well-connected when the ratio of filament length to intercross-link spacing, L/l_c , is greater than ~ 6 . Near this percolation threshold, there are only locally connected domains, and discussions of global network properties becomes less reasonable. Additionally, as the filament density is increased beyond this point, there is another transition between non-affine bending and affine stretching of filaments, which changes the dominating term of the network's elastic modulus.

Drag-like Coupling Between Overlapping Filaments

In contrast to previous models, we allow relaxation of the network's stored stress by letting the attachment points slip. We do this by replacing an elastic interaction between pairs of points along filaments with a drag-like coupling between filaments.

$$\mathbf{F}_{\text{drag}} = \xi \cdot \int ds (\mathbf{v}_i(\mathbf{s}) - \mathbf{v}_j(\mathbf{s})) p_{ij}(s) \quad (4)$$

Where $p_{ij}(s)$ represents the locational distribution of cross-link points (equal to 1 at locations of cross-links and 0 elsewhere) and $\mathbf{v}_i(\mathbf{s})$ and $\mathbf{v}_j(\mathbf{s})$ represent the the velocities of the i th and j th filaments. This model assumes a linear relation between applied force and the velocity difference between attached filaments. Obviously, non-linearities can arise in the presence of force dependent detachment kinetics as well as non-linear force extension of cross-links. We address non-linear effects of stress induced unbinding in Appendix ???. Assuming inhomogeneities from non-linear effects are of second order, the motion for the entire network is governed by a dynamical equation of the form

$$\int ds (\zeta \mathbf{v}_i(\mathbf{s}) + \xi \sum_j (\mathbf{v}_i(\mathbf{s}) - \mathbf{v}_j(\mathbf{s})) p_{ij}(s)) = \nabla \mathcal{H}_i \quad (5)$$

Here, the first term in the integral is the filament's intrinsic drag through its embedding fluid, ζ , while the second comes from the drag-like coupling between filaments, ξ .

System of Equations for Applied Stress

We model our full network as a coupled system of differential equations satisfying 5. Although the general mechanical response of this system may be very complex, we focus our attention on low frequency deformations and the steady-state creep response of the system to an applied stress. To do this we introduce a fixed stress, σ along the midline of our domain. This stress points in the direction, $\hat{\mathbf{u}}$, producing either shear ($\hat{\mathbf{u}} = \hat{\mathbf{x}}$) or extensional ($\hat{\mathbf{u}} = \hat{\mathbf{y}}$) stress.

Finally, we add a 0 velocity constraint at the far edges of our domain of interest. We assume that our network is in the "dry," low Reynold's number limit, where inertial effects are so small that we can equate our total force to 0. Therefore, we have a dynamical system of wormlike chain filaments satisfying

$$\int ds (\zeta \mathbf{v}_i(\mathbf{s}) + \xi \sum_j (\mathbf{v}_i(\mathbf{s}) - \mathbf{v}_j(\mathbf{s})) p_{ij}(s)) = \nabla \mathcal{H}_i + \sigma \mathbf{u}(\mathbf{x}) \quad (6)$$

subject to constraints such that $\mathbf{v}_i(\mathbf{x})$ is 0 with $x = 0$. This results in an implicit differential equation for filament segments which can be discretized and integrated in time to produce a solution for the motion of the system.

Computational Simulation Method

We tested our analytical conclusions on a computational model. More technical details of the model can be found in the Appendix, but we summarize the main modeling points here.

We discretize the filaments such that the equations of motion becomes a coupled system of equations for the velocities of filament endpoints, \mathbf{x} . The drag-like force between overlapping filaments results in a coupling of the velocities of endpoints.

$$\mathbf{A} \cdot \dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}) \quad (7)$$

where \mathbf{A} represents a coupling matrix between endpoints of filaments that overlap, and $\mathbf{f}(\mathbf{x})$ is the spring force between pairs of filament segment endpoints. We can then numerically integrate this system of equations to find the time evolution of the positions of all filament endpoints.

We generate a network by laying down filaments with random position and orientation within a domain of size D_x by D_y with periodic boundaries in the y-dimension. The external stress (shear or extensional/compressional) is applied to all filament endpoints falling within a fixed x-distance from the center of the domain. Finally, filament endpoints falling within a fixed x-distance from the edges of the domain are constrained to be nonmoving.

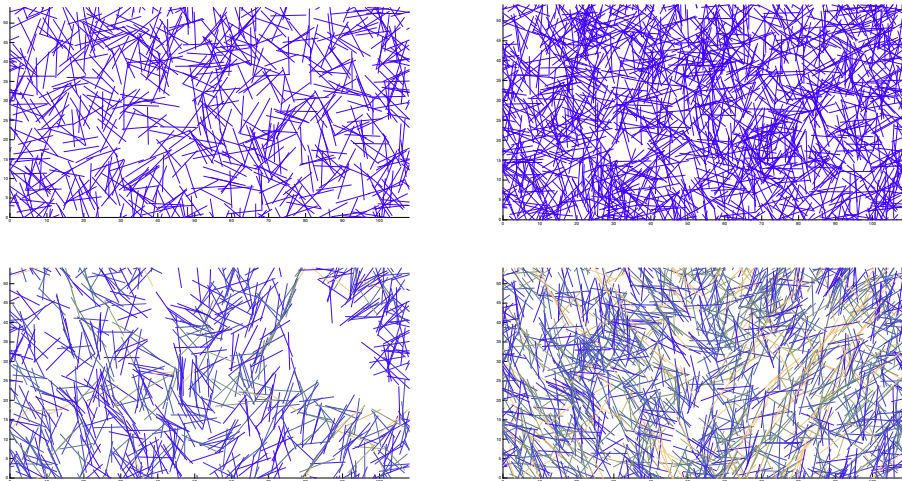


Figure 1. Two Simulation setups with $L = 9\mu m$, $D = 54\mu m$ before (top) and after (bottom) 1000 seconds of applied stress. a) low density $l_c = 2\mu m$, b) moderate density $l_c = 1\mu m$. Scale bar $20\mu m$

The nominal units for length, force, and time are μm , nN, and s, respectively. We explored parameter space around an estimate of biologically relevant parameter values, given in Table 1.

Table 1. Simulation Parameter Values

parameter	symbol	physiological estimate
extensional modulus	μ	$1nN$
bending modulus	κ	$10^{-3}nN \cdot \mu m$
cross-link drag coefficient	ξ	<i>unknown</i>
medium drag coefficient	ζ	$0.0005 \frac{nNs}{\mu m^2}$
filament length	L	$5\mu m$
cross-link spacing	l_c	$0.5\mu m$
domain size	D	$10 - 50\mu m$

Results and Discussion

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Supporting Information

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S1 Video

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S2 Fig

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S1 Table

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