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ZD2: Étude de la consistance:
Considérons l'équals de la chaleur:
   θυ -√2 =0, t>0, e e ω:= 191[
                                                                           ej=j∆ne, yje
                                                                         tn= nat
                                                                             θ= ± : KranK-Nicobson
θ-Schéma, θε[0,1]:
                      \frac{\Delta c^2}{\Delta c^2} = 0
     \frac{u_{ij}^{n+4} - u_{ij}^{n}}{\Delta t} = \frac{(4-\theta)\sqrt{\frac{u_{jn+4}^{n} - 2u_{ij}^{n} + u_{j+4}^{n}}{\Lambda c^{2}}}}{\Lambda c^{2}}
       Combinaison
 On se place en (t,ce) = (t, rej) et on applique Taylor pour u solut?
DL à l'ardre 2 en temps:
  u(tn.4,00) = u(tn,00) + 84 At + 2 84 . At + 0(At3)
 = \frac{u(t_{n+1}, \infty_j) - u(t_{n}, \infty_j)}{\Delta t} = \frac{\partial u}{\partial t}(t_{n}, \infty_j) + \frac{1}{2} \frac{\partial^2 u}{\partial t^2} \Delta t + O(\Delta t^3) 
 _ à l'ordre 4 en espace:
     u(t_n, w_{j+1}) = u(t_n, w_j) + \frac{\partial u}{\partial x}(t_n, w_j) \Delta x + \frac{1}{2} \frac{\partial u}{\partial x}(t_n, w_j) + \frac{1}{6} \frac{\partial u}{\partial x}(t_n, w_j) \Delta x^3 + \frac{1}{2} \frac{\partial u}{\partial x}(t_n, w_j) \Delta x^4 + O(\Delta x^5)
 \frac{u(\xi_{n}, \alpha z_{j+1}) - 9u(\xi_{n}, \alpha z_{j-1})}{\Lambda_{m^2}} = \frac{9^2u(\xi_{n}, \alpha z_{j-1})}{9\infty^2} + \frac{1}{40} \frac{9^4u}{9\infty^4} + O(\Delta \alpha^2)  idem pour \xi_{n+1} (pour le schéma implicite)
                                                                                         = \frac{8^{2}u}{200^{2}} \left( \left( \frac{1}{100} + \frac{1}{100} \right) + \frac{1}{100} \frac{8^{1}u}{200} + O(\Delta a^{3}) \right)
        pour tons
                                                                                                  Bir (6,00) + Bir ∆t
                                                                                           Or \frac{9^{4}}{200} = \frac{1}{4} \frac{94}{84} cour u solut
                                                                                          => (*) = \frac{184}{\times \times \text{L}} (\text{L}_{n+1}, \text{ag}) + \frac{1}{40}. \frac{8\text{L}}{\text{Dec}} \text{Ace}^2 + O(\text{Ace}^2)
                                                                                                          = 18 ( (4,0)) + 18 1. At + 1 0 (4,0) 1002 + 1 0 1. Ac. At + 0 (1023)
⇒ Ch remplace par u soluto dans le soléma:
    \frac{u_{j}^{n+1}-u_{j}^{n}}{\Delta t}=(4-\theta)V. \frac{u_{j_{n+1}}^{n}-2u_{j}^{n}+u_{j_{n+1}}^{n}}{\Delta ce^{2}}-\theta.V. \frac{u_{j_{n+1}}^{n+1}-2u_{j}^{n}+u_{j_{n+1}}^{n+1}}{\Delta ce^{2}}=\frac{\theta u}{\theta t}(t_{n}/2e_{j})+\frac{1}{\theta}\frac{\theta^{2}u}{\theta t} \frac{\Delta t}{\Delta t}+O(\Delta t^{2})
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 $\frac{u_{j} - u_{j}}{\Delta t} = \frac{(1-\theta)V}{\Delta t} \frac{u_{j+1} - 2u_{j} + u_{j+1}}{\Delta t} = \frac{\theta u}{\theta t} (t_{n}/2t) + \frac{1}{\alpha} \frac{\theta^{2} u}{\theta t} 2 \frac{\Delta t}{\Delta t} + \frac{1}{\Omega} (\Delta t^{2})$ $- (1-\theta)V \left(\frac{1}{V} \frac{\theta u}{\theta x^{2}} + \frac{1}{42} \frac{\theta^{2} u}{\theta x} \frac{\Delta t^{2}}{\Delta x^{2}} + \frac{1}{\Omega} (\Delta t^{2}) \right)$ $- \theta V \left(\frac{\theta^{2} u}{\theta x^{2}} + \frac{1}{V} \frac{\theta^{2} u}{\theta t^{2}} \frac{\Delta t}{\Delta t} + \frac{1}{42} \frac{\theta^{4} u}{\theta x} \frac{\Delta x^{2}}{\Delta x^{2}} + \dots \right)$ $= 0 + E_{j}^{n}$ $= (\frac{1}{\alpha} - \theta) \frac{\theta^{2} u}{\theta t^{2}} (t_{n}/2t) \Delta t - \frac{V}{6} \frac{\theta^{2} u}{\theta x^{2}} + \dots$ $= \frac{1}{\alpha} \frac{1}$