Propagators

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Compiling with pdflatex. Can use split screen; it seems like it compiles the document everytime I save.

The propagator is simply the inverse of the quadratic term in the Lagrangian. Let's do a few examples.

1 Abelian case (photon propagator)

After applying Fadeev-Popov to fix the gauge, we have the gauge-fixed Lagrangian:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \overline{\psi}(i\gamma^{\mu}D_{\mu} - m)\psi - \frac{1}{2\xi}(\partial_{\mu}A^{\mu})^{2}$$
(1)

We consider the parts of \mathcal{L} which are pure gauge, and we expand them out, integrating by parts to make this more compact:

$$\mathcal{L}_{gauge} = -\frac{1}{4} (\partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}) (\partial^{\mu} A^{\nu} - \partial^{\nu} A^{\mu}) - \frac{1}{2\xi} (\partial_{\mu} A^{\mu}) (\partial_{\nu} A^{\nu})
= \frac{1}{2} A^{\mu} \left(g_{\mu\nu} \partial^{2} - \left(1 - \frac{1}{\xi} \right) \partial_{\mu} \partial_{\nu} \right) A^{\nu}$$
(2)

$$=\frac{1}{2}A^{\mu}D_{\mu\nu}A^{\nu}\tag{3}$$

where we have defined the operator $D_{\mu\nu}$ by:

$$D_{\mu\nu} := g_{\mu\nu}\partial^2 - \left(1 - \frac{1}{\xi}\right)\partial_\mu\partial_\nu \tag{4}$$

The differential operator $D_{\mu\nu}$ is the term which we will need to invert (i.e. we need to find a Green's function for $D_{\mu\nu}$) by solving the equation:

$$D_{\mu\nu}\Pi^{\nu\alpha}(x) = \delta^{\alpha}_{\mu}\delta^{4}(x) \tag{5}$$

The Green's function $\Pi^{\mu\nu}(x)$ is the **propagator**. We solve this by taking it to momentum space:

$$\left(g_{\mu\nu}k^2 - \left(1 - \frac{1}{\xi}\right)k_{\mu}k_{\nu}\right)\tilde{\Pi}^{\nu\alpha} = \delta^{\alpha}_{\mu} \tag{6}$$

One can then verify the result in any QFT textbook works for $\tilde{\Pi}^{\mu\nu}$. The easiest way to invert this equation is to write it as a 4×4 matrix in Lorentz space, then find the inverse of the matrix.

2 Non-abelian case