

Unit : II

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Chapter: 3.

Electromagnetic Waves

Plane Electromagnetic Waves in Free space!

The Maxwell's Equations are

$$\left. \begin{aligned} \text{div } D &= \nabla \cdot D = \rho \\ \text{div } B &= \nabla \cdot B = 0 \\ \text{curl } E &= -\frac{\partial B}{\partial t} \\ \text{curl } H &= J + \frac{\partial D}{\partial t} \end{aligned} \right\} \text{ and } \begin{aligned} B &= \mu H \\ D &= \epsilon E \\ J &= \sigma E \end{aligned} \quad \text{--- (1)}$$

The characteristics of free space are:

$$\boxed{\rho = 0, \sigma = 0, \mu = \mu_0 \text{ and } \epsilon = \epsilon_0} \quad \text{--- (2)}$$

\therefore By using # (2) results, the equation (1) can be reduce to

$$\left. \begin{aligned} \text{div } E &= 0 \quad \text{--- (a)} \\ \text{div } H &= 0 \quad \text{--- (b)} \\ \text{curl } E &= -\mu_0 \frac{\partial H}{\partial t} \quad \text{--- (c)} \\ \text{curl } H &= \epsilon_0 \frac{\partial E}{\partial t} \quad \text{--- (d)} \end{aligned} \right\} \quad \text{--- (3)}$$

Now Taking curl of eqⁿ 3(c), we get

$$\text{curl} \cdot \text{curl } E = -\mu_0 \frac{\partial}{\partial t} (\text{curl } H)$$

Substituting curl H from # 3(d), we get

$$\text{curl curl } E = -\mu_0 \frac{\partial}{\partial t} (\epsilon_0 \frac{\partial E}{\partial t}) \quad (3)$$

$$\text{curl curl } E = -\mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2} \quad (4)$$

Now we know that,

$$\boxed{\text{curl curl } E = \text{grad div } E - \nabla^2 E} \quad [\text{Vector Identity}]$$

$$\therefore \text{curl curl } E = -\nabla^2 E \quad [\because \text{div } E = 0 \text{ from \# 3(a)}]$$

\therefore On substituting this result in # (4), becomes

$$-\nabla^2 E = -\mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2}$$

$$\therefore \boxed{\nabla^2 E - \mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2} = 0} \quad (5)$$

Now taking curl of equation 3(d)

$$\text{curl curl } H = \epsilon_0 \frac{\partial}{\partial t} (\text{curl } E)$$

substituting curl E from equation (3(c)), we get

$$\text{curl curl } H = \epsilon_0 \frac{\partial}{\partial t} \left(-\mu_0 \frac{\partial H}{\partial t} \right) = -\mu_0 \epsilon_0 \frac{\partial^2 H}{\partial t^2} \quad (6)$$

Again Using identity $\boxed{\text{curl curl } H = \text{grad div } H - \nabla^2 H}$ and with

$\text{div } H = 0$ from [3(b)], we get

$$\text{curl curl } H = -\nabla^2 H$$

\therefore On substituting this value in equation (6) becomes

$$\boxed{\nabla^2 H - \mu_0 \epsilon_0 \frac{\partial^2 H}{\partial t^2} = 0} \quad (7)$$

Equation (5) and (7) represents wave equations governing electromagnetic fields E and H in free space.

This equation resembles with the general wave equation

$$\nabla^2 u = \frac{1}{v^2} \frac{\partial^2 u}{\partial t^2} \quad \text{--- (8)} \quad v \rightarrow \text{velocity of wave.}$$

Comparing (7) and (8), the field vectors E and H are propagated in free space as waves at speed equal to

$$v = \frac{1}{\sqrt{\mu_0 \epsilon_0}}, \quad \epsilon_0 = 8.8542 \times 10^{-12} \text{ farad/m.}$$

$$= \frac{1}{\sqrt{4\pi \times 10^{-7} \times 8.8542 \times 10^{-12}}}$$

$$c = 3 \times 10^8 \text{ m/sec.} \quad \text{--- speed of light.}$$

$$\left. \begin{aligned} \nabla^2 E - \frac{1}{c^2} \frac{\partial^2 E}{\partial t^2} &= 0 \\ \nabla^2 H - \frac{1}{c^2} \frac{\partial^2 H}{\partial t^2} &= 0 \\ \nabla^2 u - \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} &= 0 \end{aligned} \right\} \quad \text{--- (9)}$$

Now let us find the solution of above equations for a plane electromagnetic wave.

A plane wave is defined as a wave whose amplitude is same at any point in a plane \perp to a specified direction.

The plane waves solutions of above equations can be written as

$$E(r, t) = E_0 e^{i(k \cdot r - \omega t)} \quad \text{--- (10)}$$

$$H(r, t) = H_0 e^{i(k \cdot r - \omega t)} \quad \text{--- (11)}$$

where E_0, H_0 are complex amplitude which are constant in space and time, while k is a wave propagation vector.

$$k = kn = \frac{2\pi n}{\lambda} = \frac{2\pi \omega}{c} \cdot n = \frac{\omega}{c} \cdot n \quad (12)$$

where n is a unit vector in the direction of wave propagation.

Now in order to apply the conditions $\nabla \cdot E = 0$, and $\nabla \cdot H = 0$, let us find $\nabla \cdot E$ and $\nabla \cdot H$.

$$\begin{aligned} \nabla \cdot E &= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot E_0 e^{i k \cdot r - i \omega t} \\ &= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot \left[(\hat{i} E_{0x} + \hat{j} E_{0y} + \hat{k} E_{0z}) e^{i(k_x x + k_y y + k_z z) - i \omega t} \right] \end{aligned}$$

$$\text{Since } [k \cdot r = \hat{i} k_x + \hat{j} k_y + \hat{k} k_z] = [k_x x + k_y y + k_z z]$$

$$\begin{aligned} \nabla \cdot E &= (E_{0x} i k_x + E_{0y} i k_y + E_{0z} i k_z) e^{i(k_x x + k_y y + k_z z) - i \omega t} \\ &= i(k_x + k_y + k_z) \cdot (\hat{i} E_{0x} + \hat{j} E_{0y} + \hat{k} E_{0z}) e^{i k \cdot r - i \omega t} \\ &= i k \cdot E_0 e^{i k \cdot r - i \omega t} = i k \cdot E \end{aligned}$$

$$\text{Similarly } \nabla \cdot H = i k \cdot H$$

$\therefore \nabla \cdot E = 0$ and $\nabla \cdot H = 0$ demand that

$$k \cdot E = 0 \text{ and } k \cdot H = 0$$

This means electromagnetic vectors E and H are both \perp to the direction of propagation vector k . This shows that electromagnetic waves are Transverse in nature.

Summary For electromagnetic waves in free space:

- ① In free space EM Waves travels with the speed of light.
- ② The electromagnetic field vectors E and H are mutually perpendicular and they are also \perp to the direction of propagation of EM Waves. Hence EM waves are Transverse in nature.
- ③ The field vectors E and H are in same phase.