# **Robust Optimal Control of Quadrotor UAVs**

Aykut C. Satici, Hasan Poonawala and Mark W. Spong

#### Abstract

This paper provides the design and implementation of an  $L_1$ -optimal control of a quadrotor unmanned aerial vehicle (UAV). The quadrotor UAV is an underactuated rigid body with four propellers that generate forces along the rotor axes. These four forces are used to achieve asymptotic tracking of four outputs, namely the position of the center of mass of the UAV and the heading. Under perfect knowledge of plant parameters and no measurement noise, the magnitude of the errors are shown to exponentially converge to zero. In the case of parametric uncertainty and measurement noise, the controller yields an exponential decrease of the magnitude of the errors in an  $L_1$ -optimal sense. In other words, the controller is designed so that it minimizes the  $L_{\infty}$ -gain of the plant with respect to disturbances. The performance of the controller is evaluated in experiments and compared with that of a related robust nonlinear controller in the literature. The experimental data indicates that the proposed controller rejects persistent disturbances which is quantified by the very small magnitude of the mean error.

#### I. Introduction

Quadrotors have become increasingly popular in recent years due to their ability to hover and maneuver in tight spaces. Stabilizing the quadrotor in its configuration space SE(3) is an interesting and challenging control problem. The challenge arises due to the availability of four control inputs to stabilize the system that lives in a six dimensional configuration space and possesses coupled, nonlinear and unstable open-loop dynamics.

The goal of most quadrotor control schemes is either to stabilize the center of mass (CoM) at a desired location or for the CoM to track a desired trajectory. Practical implementation of these nonlinear control

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schemes for robotic systems requires one to take into consideration the issues of robustness to parameter uncertainty, external disturbances, and sensor noise. It is well-known that the equations of motion of a quadrotor can be partially linearized [1] and decoupled by appropriate nonlinear feedback. known as the method of partial feedback linearization. It is very useful from a control viewpoint, since a part of the complex highly coupled nonlinear dynamics of the quadrotor are replaced by a simple set of second-order linear differential equations. The partial feedback linearization results in a controller that has the same order of complexity as the model of the system being controlled. Good control performance requires a high sampling rate at which the desired control can be computed. A desire to maximize quadrotor payloads and performance results in a constraint on computational power available for autonomous control. It is highly advantageous therefore to consider nonlinear feedback based upon simplified models and to quantify the performance of such controllers.

In this paper the partial feedback linearization method is applied to the nominal model. The linear compensator for the resulting partial linear model is then designed using the method of stable factorizations. The advantages of this control scheme can be summarized as follows.

- Any linear compensator that stabilizes a linear model can be specified using the method of stable factorization. As a consequence, one is able to conclude whether a compensator designed based on the linear model is able to stabilize the "true" nonlinear system as well, provided the deviation between the modeled and the actual plant satisfies certain bounds.
- If the deviation of the model from the plant does not exceed the maximum allowable, it is possible
  to achieve corresponding reductions in the effects on the closed-loop performance of the parameter
  uncertainty, measurement noise, available computational power, etc.

The simplest control method involves linearization of the quadrotor dynamics around the hover state and then using various linear control methods to stabilize the system, such as standard SISO tools in [2], or PID and LQR methods in [3]. In [4] the authors use a sliding mode control method to stabilize the linearized dynamics. The work includes a rate limited PID for control of the height. Controllers designed using such linear methods may be used to track slow motions. However, nonlinear control methods are needed to achieve higher performance and to execute more aggressive motions.

The quadrotor dynamics has a cascade structure. The four motor inputs are mapped to three independent torques acting on the quadrotor frame and one thrust force. The torque inputs directly control the orientation dynamics. The orientation then affects the thrust vector direction which drives the translation dynamics. This cascade structure can be elegantly exploited by using backstepping control techniques [5],

[6] or inner-loop/outer-loop approaches [2], [3], [7]. The work in [5] modifies the backstepping approach by using sliding mode control for one of the cascaded variables, which introduces some robustness to disturbances.

While the position of the CoM is almost always expressed in cartesian coordinates, selection of a representation for the orientation is an important step. All the papers mentioned so far use an Eulerangle parametrization, which suffers from singularities. To address this, the work in [8] represents the orientation using the rotation matrix itself. This leads to an orientation tracking controller that is almost globally attractive. The desired orientation and thrust magnitude are designed such that the translation dynamics are linear. The paper gives a good summary of the various representations. The same method was implemented in the work of [9] in order to achieve aggressive motions of a quadrotor.

Feedback linearization [10] is often employed in order to facilitate the use of linear control methods over a large region of the state space. The choice for the outputs or states to linearize may differ. In [11] the position and heading are taken as the outputs to be linearized whereas in [12] the height and three Euler angles are linearized. The control inputs to the linearized dynamics are selected as PD control.

Most nonlinear control methods are highly dependent on the exact knowledge of the system parameters in order to achieve asymptotic tracking. Various researchers have attempted to address this issue using robust control [13], [14], [15] or adaptive control [16] and even a combination of the two [17]. The performance recovery generally trades off asymptotic stability with uniform ultimate boundedness.

The control method proposed in this paper also uses the idea of controlling the thrust direction to achieve position control. However, the control of orientation is prone to persistent disturbances stemming from uncertain inertial parameters, misalignment of blade axes, discrepancy between the measured quadrotor frame and the actual frame. The rejection of such persistent disturbances is the primary motivation for our control design. Methods addressing this issue were first developed in [18], [19], [20] and [21].

The method in [21] elegantly determines the optimal linear controller that stabilizes a given plant while rejecting persistent disturbances. The method relies on the stable factorization approach [22] which is applicable only to linear plants. This precludes the methods used in [8] to stabilize the quadrotor since the plant remains nonlinear. The key contribution of this work is the parametrization of the orientation and subsequent partial feedback linearization of the dynamics which paves the way for applying the  $L_1$  optimal robust linear compensator design for nonlinear robotic systems similar to [20]. Furthermore, we have implemented this controller on a quadrotor in LARS <sup>1</sup>at the University of Texas at Dallas, and we

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have compared the performance of our controller with that in [15]. The design of the robust controller in [15] is not best suited for rejecting persistent disturbances, whereas our controller is. The results show that our controller can handle the effect of modeling errors such as blade misalignment and shifted center of mass location better than that in [15].

## II. KINEMATICS AND DYNAMICS OF THE QUADROTOR

In the following few paragraphs, we derive the kinematic and dynamic differential equations of motion of the quadrotor, culminating in equations (5). The reader who is familiar with the development may skip to these equations immediately without losing any generality. The quadrotor is a simple mechanical system with configuration space  $Q = SE(3) = \mathbb{R}^3 \ltimes SO(3)$ , the semidirect product of the three-dimensional Euclidean space and the special orthogonal group. The first factor in the semidirect product represents the position of the center of mass of the quadrotor in the world frame. We select the center of mass as the origin of the body-fixed frame  $\Sigma : \{o, x, y, z\}$  and denote the vector from the origin of the inertial frame  $\Sigma_0:\{o_0,x_0,y_0,z_0\}$  to o by  $p\in\mathbb{R}^3$ . Ideally, the vectors  $\{x,y\}$  span the plane parallel to the rotor axes and z is orthogonal to this plane. We denote by  $R \in SO(3)$ , the orientation of the quadrotor bodyfixed frame with respect to the world frame. The relevant axes and forces are illustrated on the Quanser Qball [23] in Figure 1. These four forces, generated by four rotors and propellers, induce a total thrust f, and a net torque  $\tau = \{\tau_x, \tau_y, \tau_z\}$  around the three body axes  $\{x, y, z\}$ . In addition to these definitions, we will denote the linear and angular velocities of the quadrotor by the pair  $(v,\Omega) \in \mathbb{R}^3 \ltimes \mathfrak{so}(3)$ . Under the vector space isomorphism  $\mathfrak{so}(3)\simeq\mathbb{R}^3$ , the skew-symmetric matrix  $\Omega$  will be identified with the 3-vector  $\omega$ . Finally, we denote the (identical) distance from the center of mass of the quadrotor to the rotor axes by L, the total mass by m, and the inertia tensor represented in the body-fixed frame by  $J = \operatorname{diag}(J_p, J_r, J_y).$ 

Since the configuration space of the quadrotor is Q = SE(3), it can be modeled as an underactuated rigid body. The velocity of the quadrotor is an element of the tangent space  $T_{(p,R)}\mathbb{R}^3 \ltimes SO(3)$ . As a Lie group, the tangent space of  $\mathbb{R}^3 \ltimes SO(3)$  can be identified with the tangent space at the identity by the left or right translation to the identity. Since the rotational equations of motion are easier when expressed in the body-fixed frame, we choose to left translate to the identity, which yields the body-fixed angular velocity

$$\Omega = R^T \dot{R} \tag{1}$$

Since the tangent space of  $\mathbb{R}^3$  is again  $\mathbb{R}^3$ , the kinematics of the translational motion is simply  $v = \dot{p}$ .

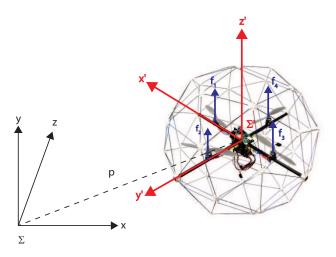


Fig. 1: Quanser Qball Model

The unforced translational dynamics of the quadrotor is the same as the dynamics of a particle (center of mass) under the action of the gravitational potential field. The total thrust  $f = \sum_{i=1}^{n} f_i$  acts on the acceleration of this particle in the direction of the z-axis.

The unforced rotational dynamics of the quadrotor are given by the Lie-Poisson equations of motion of so(3)\* relative to the rigid body bracket. Following [24], we define the angular momentum in the body frame  $\Pi := J\Omega$  whence the kinetic energy function on  $\mathfrak{so}(3)^*$  becomes

$$K(\Pi) = \frac{1}{2} \Pi^T J^{-1} \Pi \tag{2}$$

Define another set of functions  $F_i(\Pi) := \Pi_i$ ,  $i = \{1, 2, 3\}$ . Set F to be the vector-valued function with components  $F_i$  and apply the Lie-Poisson equations,

$$\dot{F} = \{F, K\} \tag{3a}$$

$$\dot{\Pi} = -\Pi \cdot (\nabla F(\Pi) \times \nabla K(\Pi)) \tag{3b}$$

$$= \Pi \times J^{-1}\Pi \tag{3c}$$

Expressing these equations in body angular velocities and including the body torques applied by the rotors,

$$J\dot{\omega} = J\omega \times \omega + \tau \tag{4}$$

In summary, we have the following equations of motion for the quadrotor describing the states at each instant in time.

$$\dot{p} = v \tag{5a}$$

$$\dot{p} = v$$

$$m\dot{v} = (-mg + fR)e_3$$
(5a)
(5b)

$$R = R\Omega \tag{5c}$$

$$\dot{R} = R\Omega$$

$$J\dot{\omega} = J\omega \times \omega + \tau$$
(5c)
(5d)

where  $e_3$  is the third standard unit vector in the Cartesian coordinate system and g is the constant gravitational acceleration. We have already mentioned that the total thrust is the sum of the individual forces of the propellers. Since the front, back and left, right propellers rotate in opposite directions, they generate a net torque around the z axis. While the moment arm L and the difference between the forces  $f_1$  and  $f_3$  generate a net torque around the y-axis, the difference between the forces  $f_2$  and  $f_4$  generate a net torque around the x-axis. This is summarized with the matrix equation given below

$$\begin{bmatrix} f \\ \tau_x \\ \tau_y \\ \tau_z \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 0 & -L & 0 & L \\ L & 0 & -L & 0 \\ K_\tau & K_\tau & -K_\tau & -K_\tau \end{bmatrix} \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \end{bmatrix}$$
(6)

III. CONTROL DESIGN

We parametrize the 12-dimensional state space of the quadrotor by selecting six position variables and their time derivatives. The first three position variables (x,y,z) specify the center of mass, p, of the quadrotor, while the remaining three position variables  $\left(\frac{r_{13}}{r_{33}},\frac{r_{23}}{r_{33}},\arctan_2\left(-r_{12},r_{22}\right)\right)$  specify the orientation of the quadrotor with respect to the inertial frame. The numbers  $r_{ij}$  are the elements of the rotation matrix  $R \in SO(3)$ . The parametrization of the orientation is local and restricted to the connected open subset  $U = \{R \in SO(3) : r_{33} > 0\}$  of SO(3). The first two orientation variables are proportional to the x and y components of the thrust vector in the world frame determining the linear horizontal acceleration of the quadrotor in this frame. The last orientation variable corresponds to the heading of the body x-axis in the world frame, which can also be interpreted as the yaw angle in a Body312 Euler angle representation of the orientation. Unlike the traditional Euler-angle parametrization, the orientation variables in our parametrization above enter linearly in the translation dynamics, as shown below, which facilitates the corresponding controller development.

In this parametrization of the state space, a point is specified by the vectors

$$\begin{split} \eta &= (\eta_1, \dots, \eta_6)^T &= (x, y, z, \dot{x}, \dot{y}, \dot{z})^T \\ \xi &= (\xi_1, \dots, \xi_6)^T &= \left(\frac{r_{13}}{r_{33}}, \frac{r_{13}}{r_{33}}, \arctan_2\left(-r_{12}, r_{22}\right), \frac{\dot{r}_{13}r_{33} - r_{13}\dot{r}_{33}}{r_{33}^2}, \frac{\dot{r}_{23}r_{33} - r_{23}\dot{r}_{33}}{r_{33}^2}, \frac{\dot{r}_{22}r_{12} - r_{22}\dot{r}_{12}}{r_{12}^2 + r_{22}^2}\right)^T \\ \text{whose dynamics can be obtained from the global set of equations (5) as} \end{split}$$

$$\dot{\eta} = \begin{bmatrix} \eta_4 \\ \eta_5 \\ \eta_6 \\ 0 \\ 0 \\ -mg \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ \xi_1 \\ \xi_2 \\ 1 \end{bmatrix} r_{33} f + d_{\eta}$$
(7a)

$$\dot{\xi} = \begin{bmatrix} \xi_4 \\ \xi_5 \\ a_{1x}(R)\tau_x + b_{1x}(R)\tau_y + \sigma_1 \\ a_{2x}(R)\tau_x + b_{2y}(R)\tau_y + \sigma_2 \\ a_{3x}(R)\tau_x + b_{3z}(R)\tau_z + \sigma_3 \end{bmatrix} + d_{\zeta}$$
(7b)

where  $a_{(\cdot)}$ ,  $b_{(\cdot)}$  and  $\sigma_{(\cdot)}$  are given in Appendix A. We have introduced the terms  $d_{\eta}$  and  $d_{\zeta}$  as additional persistent disturbances on the system that may arise, for example, from misalignment of the quadrotor blades, as shown below in section IV.

The control goal is to asymptotically stabilize the center of mass position and the heading of the quadrotor. In other words, if  $(\eta_{1d}, \eta_{2d}, \eta_{3d}, \xi_{3d})$  denote the desired values for  $(\eta_1, \eta_2, \eta_3, \xi_3)$ , we would like to have a real number T>0 such that given any  $\epsilon>0$ 

$$\|\left(\eta_{1}, \eta_{2}, \eta_{3}, \xi_{3}, \dot{\eta}_{1}, \dot{\eta}_{2}, \dot{\eta}_{3}, \dot{\xi}_{3}\right) - \left(\eta_{1d}, \eta_{2d}, \eta_{3d}, \xi_{3d}, \dot{\eta}_{1d}, \dot{\eta}_{2d}, \dot{\eta}_{3d}, \dot{\xi}_{3d}\right)\| < \epsilon \tag{8}$$

for all t > T.

The four available control inputs can be viewed as the thrust along body z-axis and a torque about each body axis. We use feedback linearization to facilitate later design of controllers based on  $L_1$ -optimal control. This is the one of the main reasons the parametrization mentioned at the beginning of this section is selected. The  $L_1$ -optimal controller that is going to be designed assumes that the nominal dynamics can be reduced to a double integrator; however, the global representation of the orientation kinematics (1) cannot be put in this form.

To begin the feedback linearization procedure, we choose the thrust f as

$$f = \frac{1}{r_{33}}h(\eta_3, \eta_6) \tag{9a}$$
 
$$h(\eta_3, \eta_6) := mg + u_z(\eta_3, \eta_6) \tag{9b}$$

$$h(\eta_3, \eta_6) := mg + u_z(\eta_3, \eta_6)$$
 (9b)

where  $u_z$  is an auxiliary controller to be designed in the sequel. The torques  $au_{(\cdot)}$  are chosen as the solution to the equations

$$\begin{bmatrix} a_{1x}(R) & b_{1y}(R) & 0 \\ a_{2x}(R) & b_{2y}(R) & 0 \\ a_{3x}(R) & 0 & b_{3z}(R) \end{bmatrix} \begin{bmatrix} \tau_x \\ \tau_y \\ \tau_z \end{bmatrix} + \begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \end{bmatrix} = \begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix}$$
(10)

Note that a unique solution exists as long as  $r_{33} \neq 0$ . The right-hand side of these equations are auxiliary control variables to be designed as  $L_1$ -optimal controllers that asymptotically stabilize  $\xi$ -dynamics. As a result of these manipulations, the dynamics (7) become

$$\dot{\eta} = f_{\eta}(\eta) + g(\eta)\xi + d_{\eta} \tag{11a}$$

$$\dot{\xi} = f_{\xi}(\xi) + v + d_{\zeta} \tag{11b}$$

where

$$f_{\eta}(\eta) = \begin{bmatrix} \eta_4 & \eta_5 & \eta_6 & 0 & 0 & u_z(\eta) \end{bmatrix}^T$$

$$[12a)$$

$$f_{\xi}(\xi) = \begin{bmatrix} \xi_4 & \xi_5 & \xi_6 & 0 & 0 & 0 \end{bmatrix}^T$$
 (12c)

$$v = \begin{bmatrix} 0 & 0 & 0 & v_x & v_y & v_z \end{bmatrix}^T \tag{12d}$$

We observe that the selected thrust (9) asymptotically stabilizes the height, z, of the quadrotor. In the remainder of this section, we are going to choose v in an  $L_1$ -optimal fashion so that the orientation is stabilized to its desired trajectory, which will, in turn, be constructed to asymptotically stabilize (x, y) and  $\xi_3$ .

# A. Backstepping Control

Consider the resulting system dynamics, given by (11). Concentrating on the first component (11a) of this system of differential equations, we notice that by setting  $\xi = \phi(\eta)$  with  $\phi$  chosen appropriately, we can asymptotically stabilize  $\eta$  to a desired value. With this observation in mind, let

$$\phi(\eta) = \begin{bmatrix} \frac{u_x}{h} & \frac{u_y}{h} & \xi_{3d} & 0 & 0 & \dot{\xi}_{3d} \end{bmatrix}$$
 (13)

where  $u_x$  and  $u_y$  are going to be selected as the output of an  $L_1$ -optimal stabilizing linear controller. Consider the change of variables

$$\zeta = \xi - \phi(\eta) \tag{14}$$

We substitute this relation into (11) and consider the auxiliary control v as functions of  $\zeta$  to obtain

$$\dot{\eta} = A\eta + Bu(\eta) + g(\eta)\zeta + d_{\eta} \tag{15a}$$

$$\dot{\zeta} = A\zeta + Bv(\zeta) - g_2(\eta, \dot{\eta}) + d_{\zeta} \tag{15b}$$

The various other terms are given by

$$A = \begin{bmatrix} 0 & I \\ 0 & 0 \end{bmatrix}, B = \begin{bmatrix} 0 \\ I \end{bmatrix}$$
 (16a)

$$g_2(\eta, \dot{\eta}) = \begin{bmatrix} \frac{\dot{u}_x h - \dot{h} u_x}{h^2} & \frac{\dot{u}_y h - \dot{h} u_y}{h^2} & 0 & 0 & 0 \end{bmatrix}^T$$
 (16b)

The dynamics in (15a) and (15b) consists of the following nominal linear system, perturbed by the terms  $g(\eta)\zeta$  and  $g_2$  respectively.

$$\dot{\eta} = A\eta + Bu(\eta) \tag{17a}$$

$$\dot{\zeta} = A\zeta + Bv(\zeta) \tag{17b}$$

$$\dot{\zeta} = A\zeta + Bv(\zeta) \tag{17b}$$

**Proposition III.1.** Assume that the auxiliary controllers  $u_x, u_y, u_z$  and  $v_x, v_y, v_z$  in Equation 17 are designed so that the origin of (17) is exponentially stable. Then, when the disturbance terms  $d_{\eta}$  and  $d_{\zeta}$ are zero, the origin of the full system 15 is exponentially stable.

Proof: It suffices to show that the perturbation terms in (15) are locally Lipschitz and vanish at the origin  $(\eta,\zeta)\equiv(0,0)$ . This is trivially true for the term  $g(\eta)\zeta$  as  $g(\eta)$  is bounded and is continuously differentiable. On the other hand, the second term  $g_2(\eta, \dot{\eta})$ , is comprised of the time derivative of  $\frac{u_x}{h}$ and  $\frac{u_y}{h}$ . Notice that all of these functions are at least continuously differentiable functions of  $(\eta, \dot{\eta})$  and therefore, are locally Lipschitz. Moreover, since  $u_x, u_y$  are linear functions and h is an affine function of  $\eta$ , the functions  $\frac{u_x}{h}$  and  $\frac{u_y}{h}$  are identically zero whenever  $\eta$  is identically zero. This, in particular, implies that under these circumstances, the time derivative of these functions are also identically zero.

From this analysis, we can conclude by [10, p. 341, Lemma 9.1] that the origin of the perturbed system (15) is exponentially stable.

Remark 1. In purely model inversion techniques, in order to feedback linearize equation (15) one would require knowledge of  $g_2$ . However, the computation of  $g_2$  is not feasible since it requires acceleration level measurements, which are known to be extremely noisy. In contrast, proposition III.1 means that to achieve exponential stability we do not need to use this malevolent signal in our control law.

# B. $L_1$ Optimal Control

In this section, we present our main result, which is the design of the control inputs  $u(\eta)$  and  $u(\zeta)$ for the system (15) using the  $L_1$ -Optimal control design procedure from [20] and [21].

We first use the stable factorization approach, given an exposition in [22], that allows one to parametrize the set of all controllers that stabilize a given (linear) system as well as the set of all stable transfer matrices that are achievable.

The first step of the approach is to factor the plant transfer matrix G into a "ratio" of the form

$$G(s) = N(s) [D(s)]^{-1}$$

$$= \left[\tilde{D}(s)\right]^{-1} \tilde{N}(s)$$
(18)

where  $N, D, \tilde{N}, \tilde{D}$  are stable rational matrices, i.e., every element of each matrix is proper and has all of its poles in the open left half-plane. Next, one solves the so-called Bezout identities:

$$\begin{bmatrix} Y(s) & X(s) \\ -\tilde{N}(s) & \tilde{D}(s) \end{bmatrix} \begin{bmatrix} D(s) & -\tilde{X}(s) \\ N(s) & \tilde{Y}(s) \end{bmatrix} = I$$
 (19)

where  $X, \tilde{X}, Y, \tilde{Y}$  are also stable rational matrices. It can then be shown that the set of all controllers that stabilize G is given by [22]

$$\left\{ -\left( Y-R\tilde{N}\right) ^{-1}\left( X+R\tilde{D}\right) \right\} \tag{20}$$

where R is an arbitrary matrix of appropriate dimensions whose elements are stable rational functions.

We notice that the nominal plant G(s) of equation (18) given by (17) is a system of double integrators. It is then routine to determine the various other matrices in (18) and (19) using the techniques of [25]. This gives

$$G(s) = \frac{1}{s^2} \begin{bmatrix} I \\ sI \end{bmatrix} \tag{21a}$$

$$N(s) = \tilde{N}(s)$$

$$=\frac{1}{(s+1)^2} \begin{bmatrix} I\\ sI \end{bmatrix} \tag{21b}$$

$$D(s) = \frac{s^2}{(s+1)^2} \cdot I \tag{21c}$$

$$\tilde{D}(s) = \frac{1}{(s+1)^2} \begin{bmatrix} (s^2 + 2s)I & -2I \\ -sI & (s^2 + 1)I \end{bmatrix}$$
(21d)

$$X(s) = \tilde{X}(s)$$
  
=  $\frac{1}{(s+1)^2} \begin{bmatrix} 1 + 2s & 2+4s \end{bmatrix} I$  (21e)

$$Y(s) = \frac{s^2 + 4s + 2}{(s+1)^2} \cdot I \tag{21f}$$

$$\tilde{Y}(s) = \frac{1}{(s+1)^2} \begin{bmatrix} (s^2 + 2s + 2)I & 2I\\ sI & (s^2 + 4s + 1)I \end{bmatrix}$$
(21g)

It is shown in [20] how to choose the "free" parameter R in (20) such that the system error rejects "uncertainties" due to system modeling. Thus, the only error that may remain will be due to the measurement noise. The procedure involves producing a sequence  $\{R_k\}$  and examining the resulting controller

$$C_k = -\left(Y - R_k \tilde{N}\right)^{-1} \left(X + R_k \tilde{D}\right) \tag{22}$$

and the error dynamics

$$e = P_{1k}\zeta + P_{2k}w \tag{23}$$

where  $\zeta$  stands for all the persistent disturbances exerted on the system (15), w denotes the measurement error, and the transfer matrices  $P_{1k}$  and  $P_{2k}$  are given by

$$P_{1k} = N(Y - R_k \tilde{N}) \tag{24a}$$

$$P_{2k} = I - D(Y - R_k \tilde{N}) \tag{24b}$$

The sequence  $\{R_k\}$  in the linear compensator expression (22) can be chosen such that as  $k \to \infty$ ,  $\|P_{1k}\|_{\hat{A}} \to 0$  and  $\|P_{2k}\|_{\hat{A}} \to 1$ , with the  $\hat{A}$  norm denoting the  $L_{\infty}$ -gain of the system  $P_i$ . As  $\|P_{1k}\|_{\hat{A}} \to 0$ , the rejection of persistent disturbances gets better and better. Since minimizing the  $L_{\infty}$ -gain of the operator  $P_i$  minimizes the  $L_1$ -gain of the input signal, this procedure amounts to designing an  $L_1$ -optimal controller for the quadrotor. This implies that in the absence of measurement noise, the mean value of the error over time will be minimized.

The linear controllers consist of two different transfer functions that accept position and velocity errors as inputs, respectively. Each yields control forces/torques that are added together to form the final linear controller. The forms of the resulting linear controllers derived using this robust control framework are given below.

$$\frac{U(s)}{E_n(s)} = \frac{2k_1s^2 + k_1k_2s + k_1^2}{s^2}$$
 (25a)

$$\frac{U(s)}{E_v(s)} = \frac{2k_2s^2 + k_2^2s + k_1k_2}{s^2}$$
 (25b)

where  $k_1$  and  $k_2$  positive are gains to be tuned. The tuning process is not so difficult to achieve once it is noted that  $k_1$  acts as the proportional (spring) and  $k_2$  as the derivative (damping) gain of the system. Each of the controllers  $\{u_x, u_y, u_z, v_x, v_y, v_z\}$  mentioned in the section III-A are of this form with appropriately chosen  $k_1$ ,  $k_2$  gains. Particular choices of these gains, suited for our quadrotor, are given in section IV.

#### IV. SIMULATION RESULTS

This section presents simulations of the proposed (LARS) controller on a model of the Quanser Qball quadrotor [23] taking into account the various sources of modeling uncertainties. The controller is compared with that presented in [8], which we refer to as RGTC standing for "Robust Geometric Tracking Controller". The simulations are required in order to demonstrate performance for the case when only model uncertainty is present, whereas in experiment some measurement noise is present. We also present simulations which include measurement noise in the model. The Quanser Qball has the following properties as given in its user manual:

$$J = \text{diag}([0.03, 0.03, 0.03]) \text{ kg} \cdot \text{m}^2$$
 (26a)

$$m = 1.4 \,\mathrm{kg}, \ \omega = 15 \,\frac{\mathrm{rad}}{\mathrm{sec}}$$
 (26b)

where  $\omega$  is the actuator bandwidth of each propeller. That is, the force output of each propeller is achieved by passing the command through a low-pass filter with  $\frac{15}{2\pi}$ Hz. bandwidth. This, of course, is a huge limitation on the range fast of motions achievable. Therefore, we design a simple lead controller that moves the pole at  $\omega = \frac{15}{2\pi}$ Hz to  $\frac{37.5}{2\pi}$ Hz. This mini controller can be expressed in frequency domain as

$$\frac{F(s)}{U(s)} = C(s) = 2.5 \frac{s+15}{s+37.5}$$
 (27)

The quadrotor is subject to several kinds of uncertainties, the most straightforward of which are its mass and its inertia tensor. More subtle uncertainties include misalignment of the rotors with respect to the quadrotor frame. That is, although the rotors of the blades are supposed to be parallel to the frame of the quadrotor, in reality they might be slightly tilted causing the application of unmodeled forces and moments on the quadrotor. Such rotor misalignment has been observed on the quadrotor in our laboratory. Moreover, if the center of mass of the quadrotor is not located exactly at the center, the manner in which the torques enter into the rotational dynamics is no longer given by (6). The visualization of the "actual" quadrotor in Fig. 2 sheds light on the derivation of its "actual" dynamics. In this figure, we see that the plane of the quadrotor frame, P, has a normal vector N. In the ideal model, the forces of each of the propellers are along this normal. The direction of these forces deviate from this normal in reality, as depicted in Fig. 2, giving rise to the dynamical equations of motion presented in (28).

$$J\dot{\omega} + \omega \times J\omega = \sum_{i=1}^{4} f_i \left[ K_{\tau} \sigma(i) I + \hat{p}_i \right] R_i \cdot e_3$$
 (28a)

$$m\ddot{p} = \sum_{i=4}^{4} \left( -mgI + f_i R_i \right) \cdot e_3 \tag{28b}$$

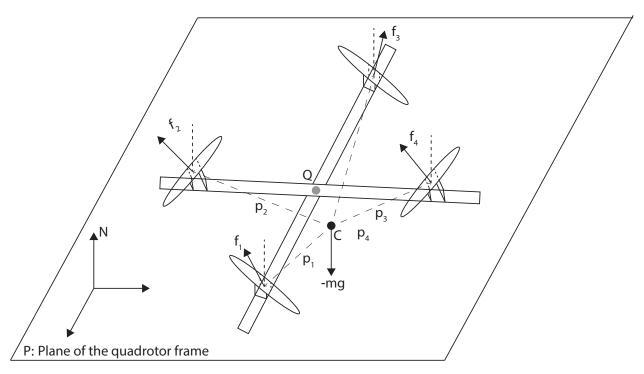


Fig. 2: "Actual" Quadrotor Model: The direction of the thrust forces of each of the propellers deviate from the normal vector to the plane of the quadrotor frame. The CoM is not exactly situated at the center of the quadrotor frame

where  $R_i \in SO(3)$  is the orientation of the  $i^{th}$  propeller with respect to the quadrotor frame,  $e_3$  is the third standard Cartesian unit vector,  $p_i$  is the position vector, in the quadrotor frame, from the center of mass C to to the center of the  $i^{th}$  propeller. The hat operator  $\hat{\ }$  is the usual isomorphism from  $\mathbb{R}^3$  to  $\mathfrak{so}(3).$  Lastly, I is the  $3\times 3$  identity matrix and  $\sigma:\{1,2,3,4\}\to\{-1,1\}$  is defined by

$$\sigma(i) = \begin{cases} 1 & \text{if } i = 1, 2 \\ -1 & \text{if } i = 3, 4 \end{cases}$$

The values of  $R_i e_3$  for the four rotors are given by

$$R_1 e_3 = \begin{bmatrix} 0.03 & 0.03 & 0.9991 \end{bmatrix}^T \tag{29a}$$

$$R_2 e_3 = \begin{bmatrix} 0.06 & 0.03 & 0.9977 \end{bmatrix}^T \tag{29b}$$

$$R_{1}e_{3} = \begin{bmatrix} 0.03 & 0.03 & 0.9991 \end{bmatrix}^{T}$$

$$R_{2}e_{3} = \begin{bmatrix} 0.06 & 0.03 & 0.9977 \end{bmatrix}^{T}$$

$$R_{3}e_{3} = \begin{bmatrix} 0.03 & 0.06 & 0.9977 \end{bmatrix}^{T}$$

$$R_{4}e_{3} = \begin{bmatrix} 0.09 & 0.09 & 0.9919 \end{bmatrix}^{T}$$

$$(29a)$$

$$(29b)$$

$$(29c)$$

$$(29d)$$

$$R_4 e_3 = \begin{bmatrix} 0.09 & 0.09 & 0.9919 \end{bmatrix}^T \tag{29d}$$

Finally, the measurements of the quadrotor pose and velocity are never perfect. In the simulations, they will be modeled as a Gaussian noise acting on both the linear and angular readings additively. In

loop	$k_1$	$k_2$
x	3	3.4293
y	3	2.6944
z	3	2.6944
$\xi_1$	176.2000	26.5481
$\xi_2$	176.2000	26.5481
$\xi_3$	26.4300	10.2820

TABLE I: Gains used in our controller

what follows, we are going to present simulation results of our proposed controller and compare it with one of the most prominent robust controllers in the literature provided by [15].

The value of gain parameters used in our controller is given in Table I The value of parameters defined in [8] are given in Table II where the robust control parameters are unchanged, however the proportional and derivative gains have been scaled to be comparable to the gains in Table I. Finally, the statistics of the additive gaussian white noise used to corrupt the measurements are given in Table III.

parameter	value
$K_R$	5.3394 I <sub>3</sub>
$K_{\Omega}$	0.7
$k_x$	6
$k_v$	5.4
$c_1$	3.6
$c_2$	0.6
$\epsilon_x$	0.04
$\epsilon_R$	0.04

parameter	value
$\sigma_p^2$	$10^{-6}$
$\sigma_{\dot{p}}^2$	0.01
$\sigma_R^2$	$10^{-4}$
$\sigma_{\omega}^2$	0.01

TABLE II: Parameters from [8]

TABLE III: Variances of additive zero mean gaussian white noise

## A. Set-point control

We command a constant desired point in the Cartesian space to which the quadrotor center of mass should converge while potentially suffering from uncertainties in parameters and measurement noise. The results of a simulation for the case when there is no measurement noise but parameter uncertainty exists is given in Fig. 3. The steady state error in position is zero in the case of the LARS controller (Fig. 3c). The orientations errors also have zero steady state values (Fig 3a). For the RGTC controller, the

experiment	noise	controller	$\mu_{e_x}[m]$	$\mu_{e_y}[m]$	$\mu_{e_z}[m]$	$\sigma_{e_x}[m]$	$\sigma_{e_y}[m]$	$\sigma_{e_z}[m]$
set-point	no	LARS	1.4096e-7	9.0064e-8	-1.7917e-7	7.05e-7	4.785e-7	6.509e-7
set-point	no	RGTC	0.0658	0.0288	0.0697	6.404e-4	2.503e-4	1.698e-4
set-point	yes	LARS	-7.4749e-4	-7.2466e-4	-7.4496e-4	0.0111	0.0081	0.0085
set-point	yes	RGTC	0.0704	0.0358	0.0743	0.0082	0.0067	0.0065
circle	no	LARS	2.01e-5	9.4342e-6	-1.33e-6	0.1434	0.1434	0.0001
circle	no	RGTC	0.0791	0.0336	0.0754	0.1291	0.1183	0.0207
circle	yes	LARS	-7.2708e-4	-7.1625e-4	-7.4137e-4	0.1430	0.1451	0.0086
circle	yes	RGTC	0.0811	0.0356	0.0815	0.1249	0.1203	0.0134

TABLE IV: Disturbance rejection of LARS and RGTC controllers

position errors reach a steady non-zero value (Fig 3d). The orientation also does not converge to the desired values (Fig 3b). Table IV gives the errors in steady state for both controllers.

When noise is added, the the results follow a similar trend (Figures 4). As seen in figures 4a and 4c, the average value of the errors in orientation and position are nearly zero for the LARS controller. This is confirmed by the mean values of the errors in position for the experiment in Table IV. In the case of the RGTC controller, we see that the mean values of the errors reach a constant value that is slightly offset from zero (Figures 4b and 4d, the mean values given in Table IV. This shows that the presence of measurement noise does not qualitatively change the relative behaviour of the two controllers.

#### B. Trajectory tracking

The quadrotor is commanded to follow a trajectory

$$\begin{bmatrix} x(t) \\ y(t) \\ z(t) \\ \xi_3(t) \end{bmatrix} = \begin{bmatrix} 0.5\sin(0.5t) \\ 0.5\cos(0.5t) \\ 1.5 \\ 0 \end{bmatrix}$$
(30)

Simulation results are presented for the case when measurement noise is both absent (Figure 5) and present (Figure 6). Paremeter uncertainty is always assumed to exist. First, we look at the case where measurement noise is not included in the simulation. As seen in Figures 5c and 5d, neither controller tracks the desired trajectory perfectly due to the actuator limitations and the prevention of using high gains due to noise. However, the LARS controller rejects the persistent disturbances due to modeling errors, as seen by the mean of the positions in error in Table IV. In the case of the RGTC controller, the mean values of the position errors are not zero. The resulting trajectories can be seen in Figures 5e and

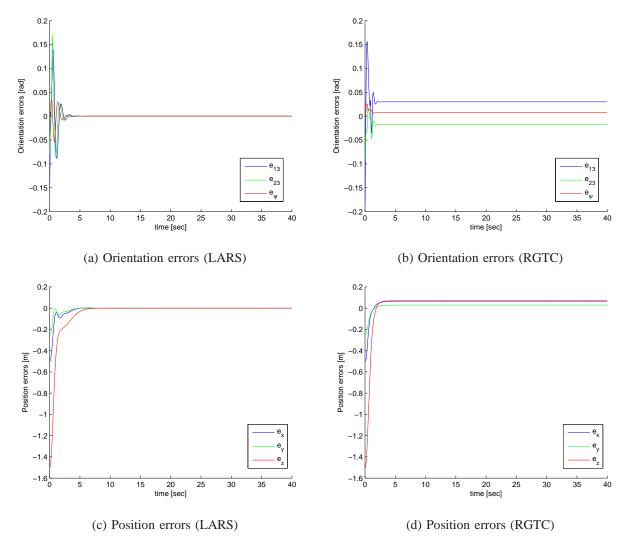


Fig. 3: Set-point regulation with parameter uncertainty and no measurement noise. The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. There is no steady-state error present in the LARS controller (a), (c), while the position errors reach a steady non-zero value in RGTC controller as does the orientation (b), (d).

5f. The LARS controllers results in a trajectory concentric with the desired circle, whereas the trajectory acheived by the RGTC controller has a center different from that of the desired trajectory. The desired orientation is tracked in the case of the LARS controller as seen in Fig 5a, whereas this does not occur for the RGTC controller (Fig 5b). Table IV supports this conclusion, as seen from the mean of the errors. The standard deviations of the RGTC controller are lower. Note that this does not mean that the mean squared error is lower, due to the larger mean in the case of the RGTC controller.

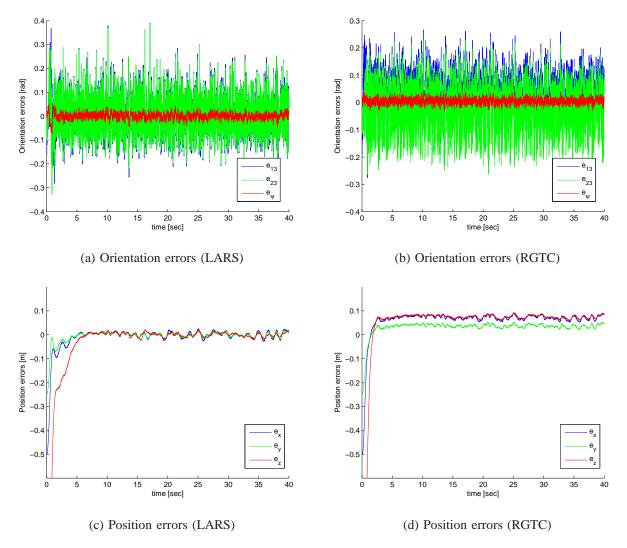


Fig. 4: Set-point regulation with parameter uncertainty and measurement noise. The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. There is no mean steady-state error present in the LARS controller (a), (c), while the average position errors reach a steady non-zero value in RGTC controller as does the orientation (b), (d).

When measurement noise is added, we see that the plots do not change much qualitatively. The effect of measurement noise can easily be seen in Figures 6c and 6d. The statistics of the obtained trajectory in Table IV show that the conclusions from the simulations without noise still hold, that is, the mean value of the trajectory for the LARS controller is still nearly zero, whereas for the RTGC controller it is non zero. Figures 6a and 6b show that the position errors appear to have a zero mean in the case of the LARS controller but not for the RGTC controller. This is confirmed again in Table IV. Similar to

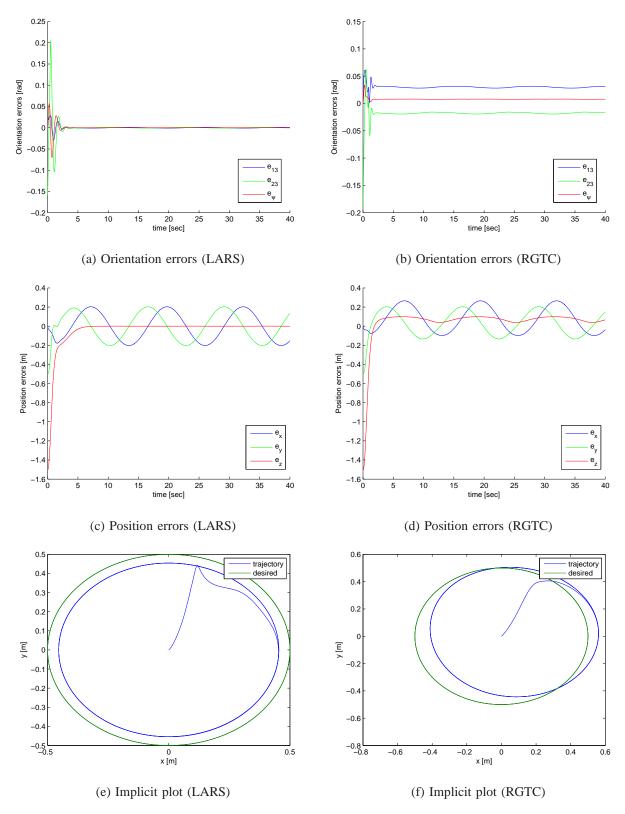


Fig. 5: Trajectory tracking with parameter uncertainty and no measurement noise. The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. The LARS controller rejects persistent distributions (parameter uncertainties) yielding a zero-mean translational steady-state error. In comparison, the RGTC controller keeps a non-zero mean translational steady-state error.

the simulation with no measurement noise, we see that the resulting trajectory for the LARS controller is concentric with the desired one (Fig 6e) whereas that for the RGTC controller is offset (Fig 6f).

### V. EXPERIMENTAL IMPLEMENTATION

The results of the simulation sections are recreated in an experimental environment with an actual Quanser Qball quadrotor [23]. We also continue our foregoing comparison with the RTGC robust controller [8] by implementing this controller and showing the experimental results with that controller. The experimental setup is as follows. The Quanser Qball is assumed to have the same properties as was given at the start of the section IV. The linear position feedback is acquired from an 8-camera Vicon vision system [26] working at a rate of 100 Hz. The linear velocity feedback is derived from this position reading by numerical differentiation coupled with a first-order filter with a cut-off frequency of 40 Hz.

$$H(s) = \frac{40s}{s+40}$$

Since the motions of this particular quadrotor cannot really be faster than 20 Hz. primarily due to the restrictions of the actuators, the aforementioned velocity measurement is justified to work well.

The orientation of the quadrotor is also acquired using vision data from the Vicon system as follows. We start by placing markers on both ends of the rods. The body x-axis of the quadrotor is then identified as the vector  $v_1$  whose tail is the marker at the "back" and whose head is at the "front" of the quadrotor. We then form the vector  $v_2$  from the "back" marker to the "left" marker and the vector  $v_3$  from the "back" marker to the "top" marker (see Figure 7). Using the Gram-Schmidt orthonormalization process [27], we get an orthonormal body-frame for the quadrotor, whose components when expressed in the inertial frame, form the rotation matrix R.

The body angular velocity of the quadrotor, on the other hand, is acquired from the on-board gyroscope. The gyroscope data is downsampled to 100Hz to match the data rate of the Vicon system.

The controller runs on an embedded processor, Gumstix Verdex, stationed on the quadrotor. Its processing power is 600 MHz. The reference commands, control gains and parameters are broadcast to the gumstix via a wireless communication protocol. The software used to prepare the C-code run on the gumstix and to send the various parameters to the gumstix is Matlab/Simulink.

Two types of experiments are conducted using each controller:

- Set-point regulation
- Tracking a circular trajectory at 1 rad/s

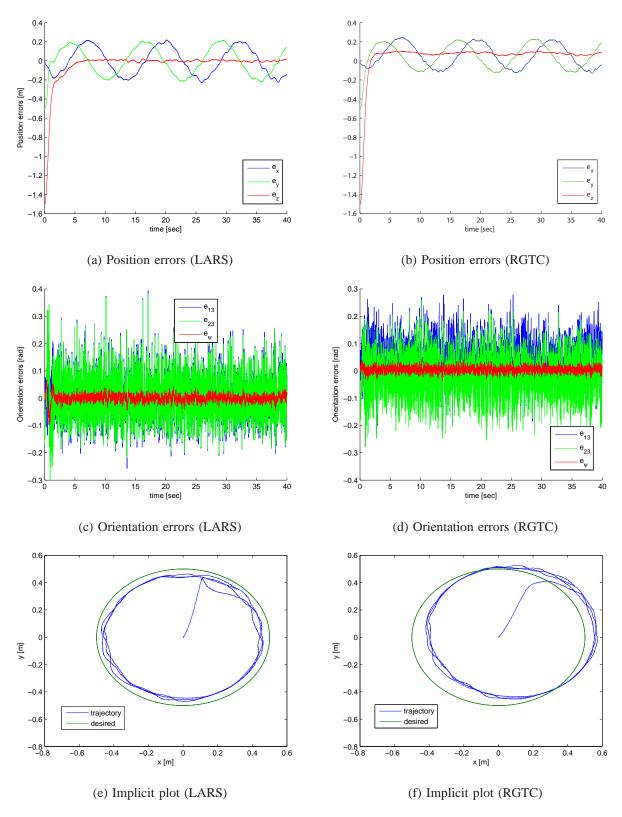


Fig. 6: Trajectory tracking with parameter uncertainty and measurement noise. The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. The LARS controller rejects persistent disturbances (parameter uncertainties) yielding a zero-mean translational steady-state error. In comparison, the RGTC controller keeps a non-zero mean translational steady-state error.

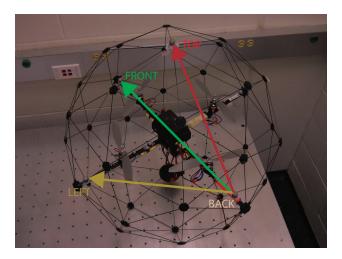


Fig. 7: Rotation matrix construction: The vectors  $v_1$  from BACK to FRONT,  $v_2$  from BACK to LEFT and  $v_3$  from BACK to TOP are orthonormalized and their components expressed in the inertial frame to form the rotation matrix R

Experiment	LARS	RGTC		
Set-point	(0.4,0,0.5)	(0.4,0,0.75)		
Circle 1	$(0.4 + \cos(t), \sin(t), 0.5)$	$(0.4 + \cos(t), \sin(t), 0.75)$		

TABLE V: Desired coordinates in the world frame: (i) Due to the steady state error in height always present in the RGTC controller. This yields sufficient clearance above the ground so as to minimize the ground effect. (ii) The moving trajectory is a circle in the horizontal plane.

The performance of the position controller is demonstrated in two steps. First, the Qball lifts off from the initial position at the desired height  $z_d$ . Once the height has settled, the robot is commanded to the point  $(0.4, 0, z_d)$ .

This is due to the fact that large step changes in all commanded variables will often hit the saturation limits, and in our closed lab space potentially leads to a crash of epic proportions. In the circle experiment, the command is switched from this point to the trajectory given in Table V.

## A. Set-point regulation results

The results of the set-point regulation experiments for both controllers are seen in Fig 8. In fig. 8a we see that the errors decrease from the initial value and tend to a zero-mean behaviour in the case of the LARS controller. There appears to be a disturbance at around t=53s, which is duly rejected. In the case of the RGTC control, after the Qball lifts off, the step position command is applied at around

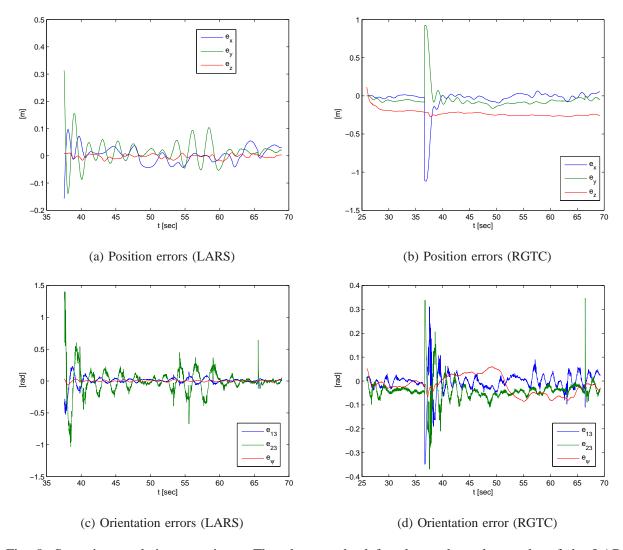


Fig. 8: Set-point regulation experiment The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. The LARS controller tends to zero mean steady-state errors, while the RGTC controller keeps a nonzero mean steady-state error.

t=35s. The errors decrease, however,  $e_y$  and  $e_z$  do not have a zero mean steady state behaviour, as seen in 8b. The x-direction position error,  $e_x$ , appears to have a zero mean behaviour. The orientation error for the LARS controller follows a similar behaviour (Fig 8c): the orientation errors decrease to a mean value close to zero, even after the disturbance. The orientation errors for the RGTC controller are not close to zero, as seen in fig. 8d, except the  $e_{13}$  error. This is logical, since in the dynamics (5) the dynamics in x and y are directly affected by the  $r_{13}$  and  $r_{23}$  components of the rotation matrix R, and hence errors in those values will result in errors in x and y even with ideal thrust generation.

## B. Trajectory tracking results

The Qball quadrotor is desired to track the trajectory in Table V. Figure 9 gives the result of the experiment for both controllers. In figure 9a the quadrotor is commanded to hover at the inital position at (0.4, 0, 0.5) at t = 31s. The command is then switched to the circular trajectory given in Table V at approximately t = 49s. The step command results in the error shooting up, which reduces to the sinusoidal errors with a nearly zero mean, similar to the behaviour seen in fig. 8a. At t = 105s the desired trajectory is switched to a set-point along the tangent to the circle at that instant. At t = 115s the Qball is commanded to return to (0.4, 0, 0.5), resulting in a jump in errors due to the step command. We see that, in spite of the step commands in positions, resulting in step commands in the desired orientations (See fig. 9c) the height error does not increase, as expected due to the LARS control design. The orientation errors are seen to have a nearly zero mean behaviour in fig. 9c. The resulting trajectory is a distorted circle with a center almost identical to that of the desired trajectory (See fig. 9e).

Figure 9a shows the errors in position for a similar experiment using the RGTC controller. The Qball lifts off at t=18s, and hovers at a point slightly offset from the desired point. Due to the steady state error in height always present in the RGTC controller, the command is higher ( $z_d=0.75m$ ). This yields sufficient clearance above the ground so as to minimize the ground effect. The Qball is then commanded to follow the circular trajectory in Table V. The error in x appears to have a zero mean behaviour, whereas y and z have non-zero mean errors. This is reflected in the figure 9f where the resulting trajectory is circular but shifted in the negative y direction with respect to the desired circular trajectory. The orientation plot in fig. 9d shows that there is an error in  $e_{23}$  at least in the set-point tracking part of the experiment.

Videos of experiments conducted with the quadrotor are detailed in Appendix B.

#### VI. CONCLUSION

We have presented an  $L_1$ -optimal robust controller for the quadrotor dynamical system that rejects persistent disturbances. Consequently, the only remaining errors in the steady-state response of the system are due to measurement noise. We have derived an "actual" model for the quadrotor under practical modeling uncertainties. We used this model to conduct simulation studies for both set-point and trajectory tracking applications and compared the performance of our controller with one of the more prominent robust controller in the literature. The same comparison has also been carried out experimentally and the performance of the system, predicted in the simulations, were seen to be very closely replicated.

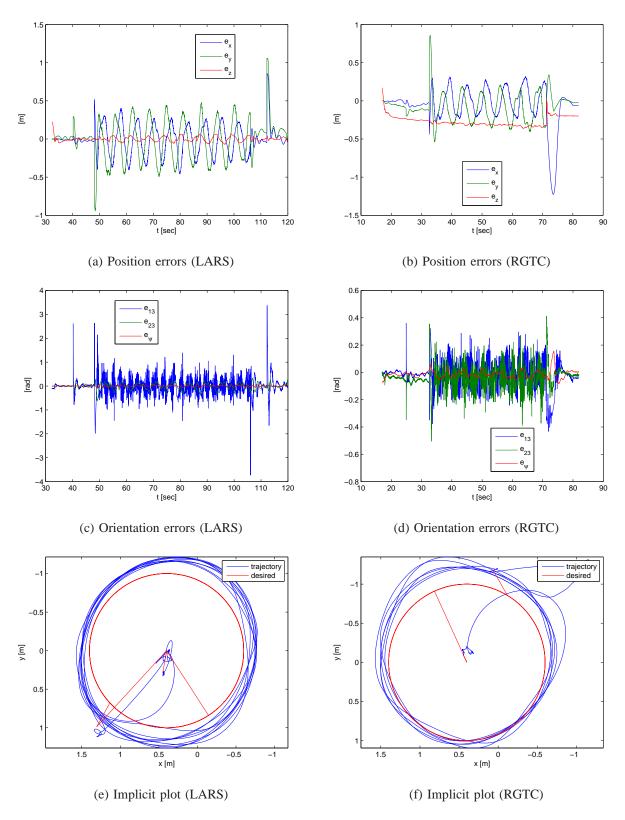


Fig. 9: Trajectory tracking experiment with a circular trajectory. The plots on the left column show the results of the LARS controller while the right column show for the RGTC controller. The LARS controller yields a correctly nearly centered circular trajectory, having a nearly zero mean translational steady state error. The circular trajectory of the RGTC controller has a shifted center, product of nonzero steady state errors.

#### APPENDIX

## A. Equations of Motion under the Selected Parametrization

We use the quadrotor dynamic equations of motion (5) to derive the equations governing the parametrization. We have  $\eta=(x,y,z)$  and  $\xi=\left(\frac{r_{13}}{r_{33}},\frac{r_{23}}{r_{33}},\arctan_2\left(-r_{12},r_{22}\right)\right)$ . We differentiate with respect to time once and substitute from the kinematic equations (1) and differentiate again before substituting from the equations (5). For the translational variables  $\eta$ , this gives

$$\ddot{\eta}_1 = \ddot{x} = \frac{r_{13}}{m}f\tag{31a}$$

$$\ddot{\eta}_2 = \ddot{y} = \frac{\ddot{r}_{23}}{m} f$$
 (31b)

$$\ddot{\eta}_3 = \ddot{z} = -g + \frac{r_{33}}{m}f$$
 (31c)

and for the rotational dynamics,

$$\ddot{\xi}_1 = a_{1x}(R)\tau_x + b_{1y}(R)\tau_y + \sigma_1(R,\omega)$$
(32a)

$$\ddot{\xi}_2 = a_{2x}(R)\tau_x + b_{2y}(R)\tau_y + \sigma_2(R,\omega)$$
(32b)

$$\ddot{\xi}_3 = a_{3x}(R)\tau_x + b_{3z}(R)\tau_z + \sigma_3(R,\omega)$$
(32c)

where the functions  $a_{1x}, a_{2x}, a_{3x}, b_{1y}, b_{2y}, b_{3z}, \sigma_1, \sigma_2, \sigma_3$  are given as follows:

$$\begin{split} a_{1x}(R) &= \frac{r_{13}r_{32} - r_{12}r_{33}}{J_r r_{33}^2} \\ a_{2x}(R) &= \frac{r_{11}r_{33} - r_{13}r_{31}}{J_p r_{33}^2} \\ a_{3x}(R) &= \frac{r_{12}r_{23} - r_{13}r_{22}}{J_r \left(r_{12}^2 + r_{22}^2\right)} \\ b_{1y}(R) &= \frac{r_{13}r_{31} + r_{11}r_{33}}{J_p r_{33}^2} \\ b_{2y}(R) &= \frac{-r_{23}r_{31} + r_{21}r_{33}}{J_p r_{33}^2} \\ b_{3z}(R) &= \frac{-r_{12}r_{21} + r_{11}r_{22}}{J_y \left(r_{12}^2 + r_{22}^2\right)} \\ \sigma_1(R, \omega) &= \frac{1}{J_p J_r r_{33}^3} \left(2J_p J_r r_{13}r_{32}^2 w_x^2 - 4J_p J_r r_{13}r_{31} v_{32} w_x w_y + 2J_p J_r r_{13}r_{31}^2 w_y^2 \\ &\quad + 2J_p J_r r_{12}r_{33} w_x \left(-r_{32}w_x + r_{31}w_y\right) + J_r \left(-J_r + J_y\right) r_{33} \left(-r_{13}r_{31} + r_{11}r_{33}\right) w_x w_z \\ &\quad + J_p \left(J_r + J_y\right) r_{12}r_{33}^2 w_y w_z + J_p^2 r_{33} \left(r_{13}r_{32} - r_{12}r_{33}\right) w_y w_z \\ (33) &\quad - J_p r_{13}r_{33} \left(J_r r_{31}w_x + \left(J_r + J_y\right) r_{32}w_y\right) w_z + J_p J_r r_{11}r_{33} \left(2r_{22}w_x w_y - 2r_{31}w_y^2 + r_{33}w_x w_z\right)\right) \\ \sigma_2(R, \omega) &= \frac{1}{J_p J_r r_{33}^3} \left(2J_p J_r r_{23}r_{32}^2 w_x^2 - 4J_p J_r r_{23}r_{31}r_{32}w_x w_y + 2J_p J_r r_{23}r_{31}^2 w_y^2 \\ &\quad + 2J_p J_r r_{22}r_{33}w_x \left(-r_{32}w_x + r_{31}w_y\right) + J_r \left(-J_r + J_y\right) r_{33} \left(-r_{23}r_{31} + r_{21}r_{33}\right) w_x w_z \\ &\quad + J_p \left(J_r + J_y\right) r_{22}r_{33}^2 w_y w_z + J_p^2 r_{33} \left(r_{23}r_{32} - r_{22}r_{33}\right) w_y w_z \\ &\quad + J_p \left(J_r + J_y\right) r_{22}r_{33}^2 w_y w_z + J_p^2 r_{33} \left(r_{23}r_{32} - r_{22}r_{33}\right) w_y w_z \\ &\quad - J_p r_{23}r_{33} \left(J_r r_{31}w_x + \left(J_r + J_y\right) r_{32}w_y\right) w_z + J_p J_r r_{21}r_{33} \left(2r_{23}w_x w_y - 2r_{31}w_y^2 + r_{33}w_x w_z\right)\right) \end{aligned}$$

$$\sigma_{3}(R,\omega) = \frac{1}{J_{r}J_{y}\left(r_{12}^{2} + r_{22}^{2}\right)^{2}} \left(J_{r}J_{y}r_{12}^{3}r_{21}w_{x}w_{y} + J_{r}J_{y}r_{12}r_{21}r_{22}^{2}w_{x}w_{y} - J_{r}J_{y}r_{12}r_{21}r_{22}^{2}w_{x}w_{y} - J_{r}J_{y}r_{12}^{2}r_{13}r_{22}w_{y}w_{z} - J_{r}J_{y}r_{13}r_{22}^{3}w_{y}w_{z} + J_{r}J_{y}r_{12}^{3}r_{23}w_{y}w_{z} + J_{r}J_{y}r_{12}^{3}r_{23}w_{y}w_{z} + J_{r}J_{y}J_{y}r_{12}^{2}r_{23}v_{y}w_{z} - J_{y}^{2}\left(r_{12}^{2} + r_{22}^{2}\right)\left(-r_{13}r_{22} + r_{12}r_{23}\right)w_{y}w_{z} - 2J_{r}J_{y}r_{12}^{2}r_{13}w_{x}\left(r_{23}w_{x} - r_{21}w_{z}\right) - 2J_{r}J_{y}r_{13}r_{22}^{2}w_{x}\left(-r_{23}w_{x} + r_{21}w_{z}\right) + 2J_{r}J_{y}r_{12}r_{22}\left(r_{13}w_{x} - r_{23}w_{x} - r_{11}w_{z} + r_{21}w_{z}\right)\left(\left(r_{13} + r_{23}\right)w_{x} - \left(r_{11} + r_{21}\right)w_{z}\right) - J_{r}J_{y}r_{11}r_{22}^{2}\left(r_{22}w_{x}w_{y} + 2r_{23}w_{x}w_{z} - 2r_{21}w_{z}^{2}\right) - J_{r}J_{y}r_{11}r_{12}^{2}\left(r_{22}w_{x}w_{y} - 2r_{23}w_{x}w_{z} + 2r_{21}w_{z}^{2}\right)\right)$$

# B. Videos of Experiments

The results of trajectory tracking experiments were recorded for some challenging trajectories: (i) Horizontal circle with sinusoidal height command [28], (ii) Vertical circle [29], and (iii) Lissajous curve [30].

The experiments consist of three phases. In the first phase, the quadrotor is commanded to hover at a height above the starting position, then commanded to move to a starting point at the same height. The second phase consists of trajectory tracking. During the third phase the following landing scheme is implemented. A predicted position is calculated based on the quadrotor position and velocity at the end of the second phase and commanded as a set-point. The desired height is decreased in steps until the quadrotor can be powered down safely.

The first trajectory is a horizontal circle with a varying height command [28]. The varying height creates an increased demand on quadrotor actuators. This is further tested when the frequency of the height command is doubled during the motion of the quadrotor.

The second experiment demonstrates the controller performance when gravity acts as a disturbance non-trivially. This is achieved by requiring the quadrotor to track a vertical circle [29].

The third experiment shows the performance of the controller under a significantly aggressive motion: tracking a three dimensional Lissajous curve. The quadrotor motion is seen to be both stable and smooth throughout the experiment [30].

The reader is invited to view further videos on the UTD Robotics channel [31], not all of which use the  $L_1$ -optimal controller. It can be observed that the tracking performance is poorer and more sluggish when the quadrotor uses naive PID control.

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