

ME/AER 676 Robot Modeling & Control

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Twists and Wrenches

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- ▶ for constant angular velocity $\omega(t) \equiv \omega$, we can solve (1) like a linear system $\dot{x}(t) = Ax(t)$:

$$R(t) = e^{[\omega]t} R(0) \quad (2)$$

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- ▶ Let $\hat{\omega} = \omega / \|\omega\|_2$ (unit norm) and $\theta = \dot{\theta} \cdot 1 = \|\omega\|_2$.
- ▶ Exponential map $\exp: [\hat{\omega}]\theta \in \mathfrak{so}(3) \rightarrow R \in \text{SO}(3)$.
Logarithm map $\log: R \in \text{SO}(3) \rightarrow [\hat{\omega}]\theta \in \mathfrak{so}(3)$.

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- If we view ω as $\hat{\omega}\theta$, where $\hat{\omega}$ is unit norm and $\theta = \dot{\theta} \cdot 1 = \|\omega\|_2$ then

$$R = e^{[\hat{\omega}]\theta} = I + (\sin \theta)[\hat{\omega}] + (1 - \cos \theta)[\hat{\omega}]^2, \quad (4)$$

where $e^{(At)} = I + At + A^2t^2/2! + \dots$ is the matrix exponential

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- ▶ $\hat{\omega}\theta$ are therefore the exponential coordinates of R
- ▶ We may interpret exponential coordinates as coming from a constant angular velocity applied for one second

Angular Velocity Frames

- ▶ The equation $\dot{R}(t) = [\omega(t)]R(t)$ involves terms defined in a fixed reference frame, called the *space frame* $\{s\}$ in MR, so really

$$\dot{R}(t) = [\omega_s]R(t) \quad (R = \underbrace{R_b^s}_{\text{RMC}} = \underbrace{R_{sb}}_{\text{MR}})$$

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- ▶ The equations may therefore be rewritten as

$$\dot{R}(t) = R(t)[\omega_b]$$

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- ▶ Any Lie group has a similar set of manipulations
- ▶ $SE(3)$ (homogenous transformations) are also a Lie group
- ▶ The 'angular velocity' corresponding to $SE(3)$ is a *twist*
- ▶ Twists for SE are not as intuitive as angular velocities for $SO(3)$.

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- ▶ If $\omega \in \mathbb{R}^3$ represent velocities for $SO(3)$, twists \mathcal{V} represent velocities for $SE(3)$
- ▶ Consider a homogenous transformation $T(t) \in SE(3)$ representing a rigid body pose of $\{b\}$ in $\{s\}$:

$$T(t) = \begin{bmatrix} R(t) & p(t) \\ 0 & 1 \end{bmatrix} \quad (T = \underbrace{T_b^s}_{\text{RMC}} = \underbrace{T_{sb}}_{\text{MR}} = \underbrace{H_b^s}_{\text{HP}}) \quad (5)$$

Body Twist

- If the angular velocity in the body frame is ω_b , and the velocity of the origin is v_b , then

$$\dot{R}(t) = R(t)[\omega_b], \quad \dot{p}(t) = R(t)v_b$$

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- ▶ The body twist has simple physical meaning:
instantaneous angular velocity of $\{b\}$ as seen in $\{b\}$, and
instantaneous velocity of origin of $\{b\}$ as seen in $\{b\}$

Spatial Twist

- ▶ We can convert the body twist $\mathcal{V}_b = \begin{bmatrix} \omega_b \\ v_b \end{bmatrix}$ into a spatial twist

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- ▶ v_s a fictitious velocity of the origin of $\{s\}$ as if the space frame $\{s\}$ was rotating about axis ω_s that passes through origin of $\{b\}$ (Fig 3.17 in MR).

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- ▶ The remaining notation is about how to transform between \mathcal{V}_b and \mathcal{V}_s , or between $[\mathcal{V}_b]$ and $[\mathcal{V}_s]$

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- ▶ $\mathcal{V} = \mathcal{S}\dot{\theta}$, just like $\omega = \hat{\omega}\dot{\theta}$
- ▶ The result of following a constant spatial twist \mathcal{V}_s for one second can be interpreted as a screw motion : translation along an axis and rotation about that axis.

Screws

- Given $\mathcal{V} = [\omega \quad v]^T$, we may derive a \mathcal{S} as a normalized twist

$$\mathcal{S} = \begin{cases} \frac{\mathcal{V}}{\|\omega\|} = \begin{bmatrix} \hat{\omega} \\ \frac{v}{\|\omega\|} \end{bmatrix} & \text{if } \omega \neq 0 \\ \frac{\mathcal{V}}{\|v\|} = \begin{bmatrix} 0 \\ \frac{v}{\|v\|} \end{bmatrix} & \text{if } \omega = 0, v \neq 0 \end{cases} \quad (6)$$

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Exponential map $\exp: [\mathcal{S}]\theta \in \mathfrak{se}(3) \rightarrow T \in SE(3)$.
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- ▶ If $\omega = 0$ and $\|v\| = 1$

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- ▶ As long as two wrenches are expressed in the same frame, they belong to the same vector space, and can be summed.
- ▶ As usual, a generic force and moment can be expressed in either $\{b\}$ or $\{s\}$
- ▶ The conversion between expressions is related to the conversion for $[\mathcal{V}_b]$ and $[\mathcal{V}_s]$ (pg. 108 in MR)