### Two-wire transmission lines

#### Outline:

- What is the purpose of this study?
- What and how was done in the study?
- What are possible future directions for the study?



Gianfranco Scordino Ponomarenko-Timofeev Aleksei

# Why study it?

- Two-wire transmission lines are most common today
- Have a simple structure
- Therefore there are not so many dependencies for the parameters
- However, due to the first point, this topic is of great interest and is not extremely complex as a result

# What we get from this study

- Understanding of how parameters of the line affect it's characteristics
- Can we use analytical formula in case parameters of the line are equal to certain value?
- Basic skills of using simulation software
   (COMSOL, for instance), although that was
   not the main point of the study

# Parallel plate transmission line, theory

Assume  $e^{-j\beta z}$  variation, no variation with x

TM modes: 
$$\nabla_t^2 e_z + k_c^2 e_z = 0$$
,  $k_c^2 = \begin{cases} k_0^2 - \beta^2 & \text{in the air region} \\ k^2 - \beta^2 & \text{in the dielectric region} \end{cases}$ 

Let 
$$k_c = \begin{cases} p & \text{for air region} \\ \ell & \text{for the dielectric region} \end{cases}$$
  
 $\ell^2 - p^2 = (\varepsilon_r - 1)k_0^2$ 

$$h_{x}(y) = \begin{cases} \frac{j\varepsilon_{r} k_{0} Y_{0}}{\ell^{2}} \frac{\partial e_{z}}{\partial y} \\ \frac{jk_{0} Y}{p^{2}} \frac{\partial e_{z}}{\partial y} \end{cases}$$

For dielectric region

for air region

$$e_z(y) = C_1 \sin \ell y$$

$$e_z(y) = C_2 \sin p(b-y)$$

$$C_1 \sin \ell a = C_2 \sin pc$$

$$0 \le y \le a$$

$$a \le y \le b$$

$$\frac{\varepsilon_r}{\ell} C_1 \cos \ell a = -\frac{1}{p} C_2 \cos pc$$

 $\ell \tan \ell a = -\varepsilon_r p \tan pc$  Transcedental equation must be solved simultaneously with  $\ell^2 - p^2 = (\varepsilon_r - 1)k_0^2 \implies \ell$ , p

$$\beta = \sqrt{k_0^2 - p^2} = \sqrt{k^2 - l^2}$$

Most of the modes will be nonpropagating if  $\beta$  is imaginary. The variation is  $e^{-|\beta|z}$  and the field decays exponentially.

The value of  $\beta$  between  $k_0$  and k can occur if  $p = jp_0$ Let  $\ell_0$  to be the corresponding value of  $\ell$  then:

$$\ell_0 \tan \ell_0 a = \varepsilon_r p_0 \tan p_0 c$$

$$\ell_0^2 + p_0^2 = (\varepsilon_r - 1)k_0^2$$

### Low Frequency Solution

#### When the frequency is low,

 $k_0^2$  is very small number,  $\ell_0$  and  $p_0$  are very small  $\Rightarrow$ 

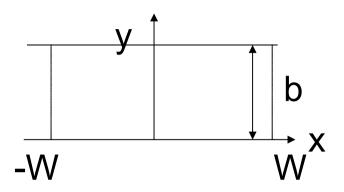
$$\ell_0^2 a = \varepsilon_r p_0^2 c$$

$$(\varepsilon_r - 1)\mathbf{k}_0^2 - p_0^2 = \frac{\varepsilon_r p_0^2 c}{a}$$

or 
$$p_0^2 = \frac{(\varepsilon_r - 1)k_0^2 a}{a + \varepsilon_r c}$$
 The solution for  $\beta$  is

$$\beta = \sqrt{\mathbf{k}_0^2 + p_0^2} = \sqrt{\frac{\varepsilon_r b}{a + \varepsilon_r c}} \ k_0 = \sqrt{\varepsilon_e} k_0$$

 $\varepsilon_e$  is the effective dielectric constant



 $\varepsilon_e = \omega \sqrt{LC}$  *L*, *C* are the static inductance and capacitance per meter.

The time average stored magnetic energy is  $W_m$ 

$$W_{m} = \frac{\mu_{0}}{4} \int_{0}^{b} \int_{-w}^{w} H_{x}^{2} dx dy = \frac{\mu_{0}}{2} WbJ_{z}^{2} = \frac{1}{4} LI_{z}^{2}$$

$$I_{z} = 2 WJ_{z} \qquad \Rightarrow L = \frac{\mu_{0}b}{2W}$$

$$\begin{split} C_2 &= C_1 \sin \ell_0 a I \, j \sinh p_0 c \approx -j C_1 \, \beta \ell_0 \, a I \, p_0 \, c \\ 0 &\leq y \leq a \\ e_z &= C_1 \ell_0 \, y \\ e_y &= -\frac{j \beta}{\ell_0} \, C_1 = -j C_1 \sqrt{\frac{b}{(\varepsilon_r - 1) c}} \\ h_x &= \frac{j \varepsilon_r k_0 Y_0}{\ell_0} \, C_1 = j Y_0 C_1 \sqrt{\frac{(\varepsilon_r c + a) \varepsilon_r}{(\varepsilon_r - 1) c}} \\ \text{The capacitance } C &= \frac{C_a \, C_d}{C_a + C_d} \quad , \quad C_a = \frac{\varepsilon_0 \, 2W}{c} \, , C_d = \frac{\varepsilon_0 \, \varepsilon_r \, 2W}{a} \\ C &= \frac{\varepsilon_0 \varepsilon_r \, 2W}{\varepsilon_r \, c + a} \\ LC &= \frac{\varepsilon_0 \varepsilon_r \, \mu_0 \, b}{\varepsilon_r \, c + a} \qquad \beta = \omega \, \sqrt{LC} \end{split}$$

The capacitance 
$$C = \frac{C_a C_d}{C_a + C_d}$$
,  $C_a = \frac{\varepsilon_0 2W}{c}$ ,  $C_d = \frac{\varepsilon_0 \varepsilon_r 2W}{a}$ 

$$C = \frac{\varepsilon_0 \varepsilon_r 2W}{\varepsilon_r c + a}$$

$$LC = \frac{\varepsilon_0 \varepsilon_r \mu_0 b}{\varepsilon_r c + a} \qquad \beta = \omega \sqrt{LC}$$

Field expressions:

$$C_2 = C_1 \sin \ell_0 a / j \sinh p_0 c \approx -j C_1 \beta \ell_0 a / p_0 c$$

$$0 \le y \le a$$

$$e_z = C_1 \ell_0 y$$

$$e_{y} = -\frac{j\beta}{\ell_{0}} C_{1} = -jC_{I} \sqrt{\frac{b}{(\varepsilon_{r}-1)c}}$$

$$h_{x} = \frac{j\varepsilon_{r}k_{0}Y_{0}}{\ell_{0}}C_{1} = jY_{0}C_{1}\sqrt{\frac{(\varepsilon_{r}c+a)\varepsilon_{r}}{(\varepsilon_{r}-1)c}}$$

In the air region:

In the air region:  

$$e_z = C_1 \ell_0 a \frac{(b-y)}{c}$$

$$e_y = -\frac{j\beta \ell_0 a}{p_{0^2} c} C_1 = -jC_1 \varepsilon_r \sqrt{\frac{b}{(\varepsilon_r - 1)c}}$$

$$h_x = \frac{jk_0 Y_0}{\ell_0} C_1 = jY_0 C_1 \sqrt{\frac{(\varepsilon_r c + a)\varepsilon_r}{(\varepsilon_r - 1)c}}$$

$$V = -\int_0^b e_y dy = jC_1 \sqrt{\frac{b}{(\varepsilon_r - 1)c}} (a + \varepsilon_r c)$$

$$I_z = 2WJ_z = 2WH_z$$

The characteristic impedance is:

$$Z_{c} = \frac{V}{I} = \frac{Z_{0}}{2W} \sqrt{\frac{(a + \varepsilon_{r}c)b}{\varepsilon_{r}}} = \sqrt{\frac{L}{C}}$$

In the low frequency limit, the dominant mode of propagation becomed a TEM mode (quasi-TEM mode) At hight frequency the mode of propagation is an E mod

# Example: Two plate transmission line

#### Problems:

- Plates are infinitely thin
- Simulation mesh might get very big

#### Solutions:

- Define boundaries of inner square (or rectangle) as conductors
- Pick appropriate size by trial and error, so that the simualtion results do not change drastically, compared with previous values

# Calculating impedance

- We need L and C of the line
  - They are the part of the energy equations
  - Calculate them using these equations:

$$L = \frac{2W_m}{I^2}$$

$$U = \frac{2W_e}{U^2}$$

 Optimization: set 1 for current and voltage in the simulation parameters

# Calculating impedance...

The only step that is left now is:

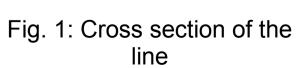
$$Z = \sqrt{\frac{L}{C}} = \sqrt{\frac{2W_m}{2W_e}} = \sqrt{\frac{W_m}{W_e}}$$

- Then we can calculate the analytical function and compare the plots of two solutions
- In the formula above we consider I = 1A and U = 1V

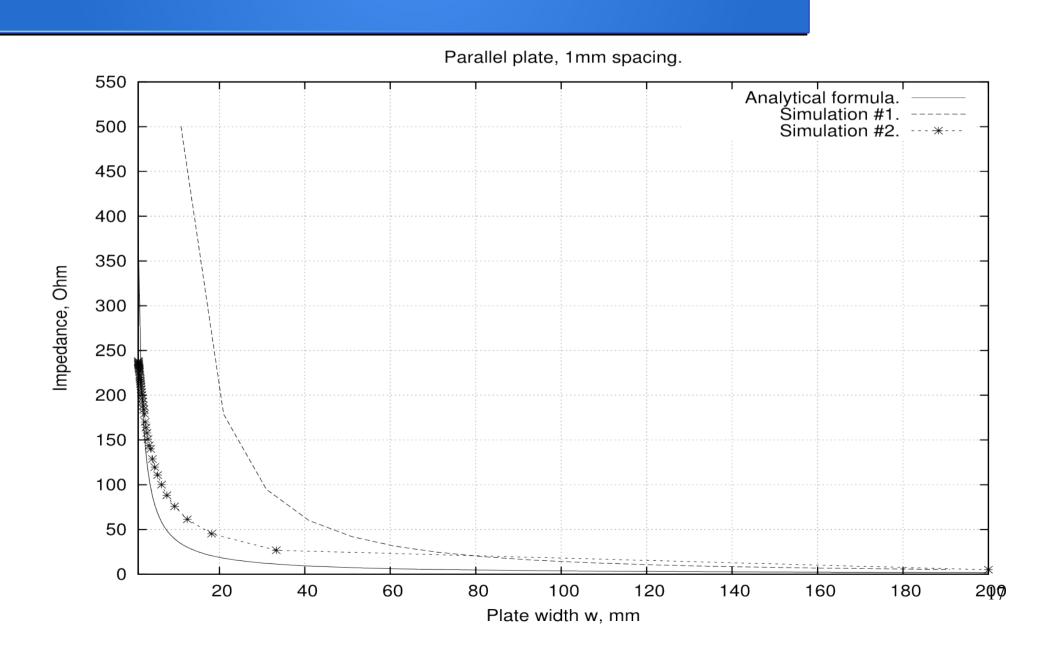
# Two parallel plates

- Consider following:
  - We have two infinitely thin plates of width w
  - They are separated by b meters of insulator
  - Let insulator be air
  - How impedance of such structure can be calculated?
  - Analytic formula:

$$Z \approx \eta \frac{b}{w}$$



### Calculation results



## Simulation screens

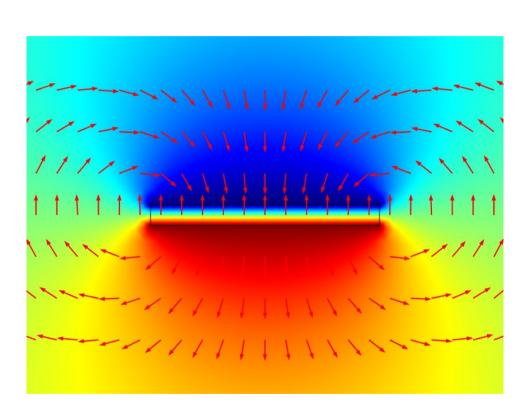


Fig. 2: Electric potential and its density in the two plate transmission line.

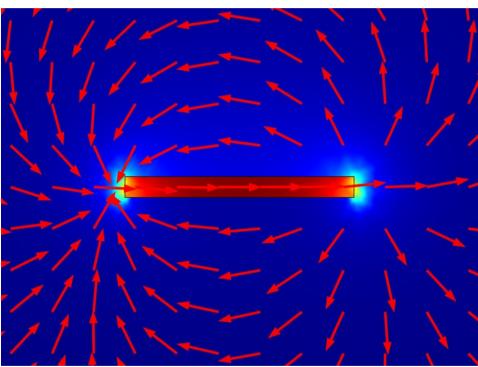


Fig. 3: Magnetic flux and its density in the two plate transmission line.

### Some conclusions on case

- It can be seen from plot that the analytical formula is not always correct
- The bigger the plate width gets, the more accurate our estimation becomes
- The analytical formula approaches values of simulation, meaning that the error is decreasing and consequently getting closer to actual value

### **Future directions**

- Fix simulations that were done improperly
- Make a proper report
- Reviev other cases (Shielded pair, for instance)
- Review different dielectrics (polyethylene, rubber)

### **COMSOL TIME**

- Show some parts of the models in COMSOL
- Questions?