Aerial Generalized Slung Load Transportation

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***Abstract*—This paper focuses on a pose tracking problem for a system composed of two connected rigid bodies, namely an aerial vehicle – with hovering capabilities, and a rod-like rigid body. The rod-like rigid body has an axis of axial symmetry, and the joint connecting the aerial vehicle and the rod lies somewhere along that axis. The aerial vehicle is meant to transport that object that can swing/slung with respect to the vehicle: we refer to *generalized* slung load transportation because the object being transported is a rod-like object, as opposed to *standard* slung load transportation where the object is a point mass with no moment of inertia. Given this system, we consider two tracking problems. Firstly, we assume that a torque input on the joint is available, and, for this scenario, we formulate a *semi*-pose tracking problem, requiring the rod object to track a desired pose trajectory, apart from rotations around its axis of axial symmetry (and thus the *semi* qualifier). Secondly, we assume that no such torque input is available, and, for this scenario, we formulate a position tracking problem, requiring a specific point along the axis of axial symmetry of the rod object to track a desired position trajectory. Our approach for solving both these problems lies in finding a state and an input transformations, such that the vector field in the new coordinates is of a known form for which controllers are found in the literature, and which we leverage in this paper. Simulations are presented that validate the proposed algorithms.**

I. Introduction

Unmanned aerial vehicles (UAVs) provide a platform for performing and testing a variety of tasks. Among other things, UAVs have been used to map large surface areas, surveil dangerous terrains, and inspect remote infrastructures [1]–[3]. More recently, UAVs with hovering capabilities have been successfully employed in transporting cargoes, which is of interest in dangerous and cluttered environments [4], [5].

Aerial transportation can be roughly categorized as slung-load/tethered-transportation [5]–[16] *vs* manipulator-endowed-transportation [3], [17]–[24]; and as single-UAV transportation *vs* (cooperative) multiple-UAV transportation. Slung-load/tethered transportation, when compared with transportation by manipulators, is mechanically simple, inexpensive, and it does not require any power supply; manipulator-endowed transportation, on the other hand, provides extra degrees of freedom which can be used to perform more complex tasks – such as picking up an object from inside a drawer [18]. In cluttered environments, transportation with a single UAV may be the only feasible option, while transportation with multiple UAVs is primarily necessary when the cargo exceeds the individual UAVs' payload capacity. In this paper, the system at study is composed of a single UAV connected to a rod-like object, and two different cases are considered: in the first case, we assume a torque input at the joint connecting the UAV and the object is available, thus falling into the manipulator-endowed transportation – see Fig. 1a; in the second case, no such torque is made available, thus falling into the slung-load/tethered transportation – see Fig. 1b.

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Different aspects of aerial transportation have been considered in the literature. In slung-load transportation, vision and force sensors have been used to estimate the swing angle of the load, or to autonomously estimate the cargo's pose, which is then used in the feedback loop to avoid/dampen swing excitation [6]–[8], [25]; in manipulator-endowed transportation, vision has been used to correctly place end-effectors with respect to a visual target placed on the object to be transported [8], [26], [27]. Differential flatness has been used for the purposes of motion control [11], [21], [28]–[31], such as planning trajectories that minimize the loads' swing. Motion planning for collision avoidance between the cargo and the UAVs with obstacles in a cluttered environment has also been studied and validated [14], [19], [20], [32]. Adaptive and robust controllers have also been proposed, which estimate and compensate for the static and dynamic effects of the cargo on the UAV [5], [22]–[24].

In aerial tethered transportation, a cable establishes a physical connection between the UAV and the load, and, under tensile forces, it behaves like an un-actuated manipulator (i.e., like a rigid link). To be specific, under tensile forces, the cable imposes a holonomic restriction, namely that the distance between the UAV and the load is equal to the cable length; if that is not the case, no such restriction applies, and the UAV and the load can move independently. In the literature of tethered transportation, some works have focused on the hybrid modeling describing the dynamics when the cable is taut and when it is not, and hybrid controllers have also been proposed that deal with the hybrid dynamics [10], [11], [33]. In this paper, we model the rod-like object as a rigid three-dimensional pendulum that, rather than attached to a fixed pivot-point [34]–[36], is attached to a moving one, namely the UAV – see Fig. 1b. Because the rod is rigid, it can withstand compression forces, and the need for a hybrid model is unnecessary. Nonetheless, this paper's results apply for tethered transportation, provided that the cable remains under tensile forces.

A. Paper's organization

In Section III, we model our system composed of two rigid bodies, namely a UAV and a rod-like object. We consider two systems, one where there is a torque input at the joint connecting the rigid bodies – UAV-manipulator system (Fig. 1a); and one where there is not – UAV-slung-manipulator system (Fig. 1b). In Subection III-A, we formulate the problem statement for the UAV-manipulator system, and in Subection III-B, we formulate the problem statement for the UAV-slung-manipulator system. In Section IV, we explain the control strategy for both problems; and in Section V, we implement that strategy for the UAV-manipulator system,

while in Section VI we implement that same strategy for the UAV-slung-manipulator system. Finally, in Section VII, we provide some illustrative simulations. In [37], one finds mathematica files that corroborate all of the presented results.

B. Contributions

Let us summarize some of this paper's contributions. First of all, we note that the considered system is not fully actuated, and thus inverse dynamics control – where a nonlinear control law linearizes and decouples the system [38] – cannot be implemented.

One of our contributions lies in analyzing two problems under a common framework: namely, we consider a UAV-manipulator system and a UAV-slung-manipulator system, which are the same system, with the exception that the UAV-slung-manipulator system is deprived of a torque input (and thus the *slung* in the naming).

Another main contribution, with respect to the UAV-manipulator system, lies in showing that it is equivalent to two decoupled subsystems: one concerning the position of the center-of-mass of the system, and another concerning the attitude of the manipulator. The first subsystem has the dynamics of a VTOL vehicle, and thus we are able to leverage controllers from the literature to control the position of the center-of-mass [39]–[42]. The second system has the dynamics of a second order system in the unit sphere, and thus we are able to leverage controllers from the literature to control the attitude of the manipulator [43]–[46].

Equivalently, another main contribution, this time with respect to the UAV-slung-manipulator system, lies in showing that it is equivalent to a VTOL vehicle cascaded after a second order system in the unit sphere. For this system, we are also able to leverage controllers for VTOL vehicles from the literature, which must be complemented with two backstepping steps (related to the second order system). We emphasize that this system is a generalization of the slung load system: indeed, rather than a point mass, the aerial vehicle carries a rigid body with some moment of inertia, with the slung load problem being recovered if that moment of inertia vanishes.

This work provides an extension of the results found in [47] and [48], the first related to manipulator-endowed transportation and the latter to slung transportation. This paper unifies those two problems, while generalizing them at the same time. To be specific, in [47] and [48], the load is taken to be a point-mass, whilst in this paper it is taken as a rod-like object with non-zero inertia. In order to analyze these two problems, and because the rod-like object has an axis of axial symmetry, we need to *ignore* the rotations of the rod object around that axis. This is done in a precise manner, by introducing an equivalence relation, and performing the analysis in the quotient space. This also constitutes anther contribution of this paper.

II. Notation

The map $\mathcal{S} : \mathbb{R}^3 \ni x \mapsto \mathcal{S}(x) \in \mathbb{R}^{3\times 3}$ yields a skew-symmetric matrix and it satisfies $\mathcal{S}(a)\, b := a \times b$, for any $a, b \in \mathbb{R}^3$. $\mathbb{S}^n := \{x \in \mathbb{R}^{n+1} : x^T x = 1\}$ denotes the set of unit vectors in $\mathbb{R}^{n+1}$. The map $\Pi : \mathbb{S}^2 \ni x \mapsto \Pi(x) := I_3 - xx^T \in \mathbb{R}^{3\times 3}$ yields a matrix that represents the orthogonal projection onto the subspace perpendicular to $x \in \mathbb{S}^2$. Denote $d_{\mathbb{R}^n} : \mathbb{R}^n \times \mathbb{R}^n \ni (a, b) \mapsto d_{\mathbb{R}^n}(a, b) := \|a - b\| \in \mathbb{R}_{\geq 0}$ as the standard metric in $\mathbb{R}^n$. Denote $A_1 \oplus \cdots \oplus A_n$ as the block diagonal matrix with block diagonal entries A_1 to A_n (square matrices). We denote by $e_1, \cdots, e_n \in \mathbb{R}^n$ the canonical basis vectors in $\mathbb{R}^n$; when clear from the context the Eucledian space the vectors belong to is omitted. For some set A, $\mathrm{id}_A : A \ni x \mapsto \mathrm{id}_A(x) := x \in A$ denotes the identity map on that set. Given some normed spaces A and B, and a function $f : A \ni a \mapsto f(a) \in B$, $Df : A \ni a \mapsto Df(a) \in \mathcal{L}(A, B)$ denotes the derivative of f ($\mathcal{L}(A, B)$ denotes the set of linear maps from A to B). Given a manifold A, $T_a A$ denotes the tangent set of A at a point $a \in A$.

(a) Manipulator that can be modeled as a rod-like rigid body, and with a torque input available at the joint connecting the manipulator to the aerial vehicle.

(b) A rod-like object connected to an aerial vehicle at a point along the axis of axial symmetry: the rod object may be interpreted as an un-actuated manipulator.

Fig. 1: Real system composed of an aerial vehicle (Iris+ from 3D Robotics) and a manipulator/rod-like-object [47].

A. Rigid Body Kinematics and Dynamics

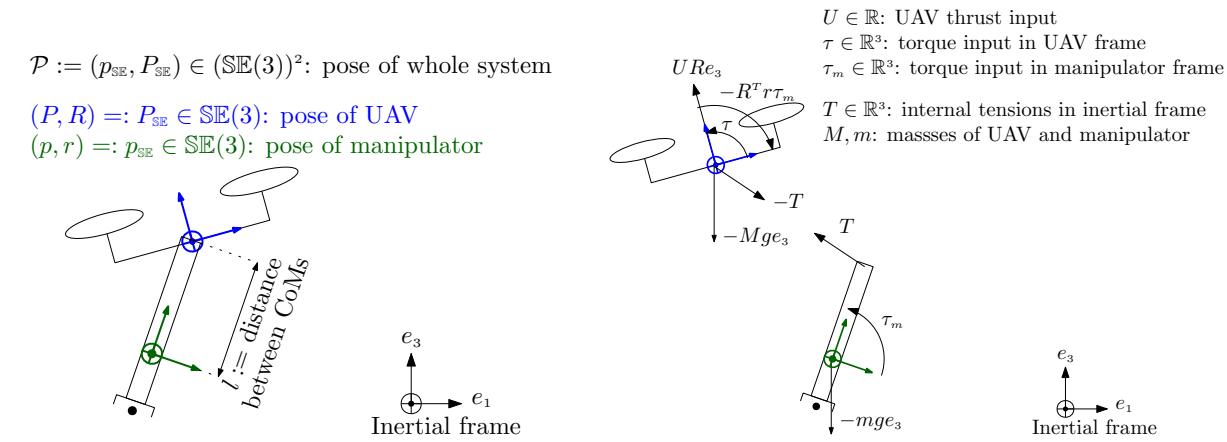
For a brief introduction on rigid body kinematics and dynamics, see Section III in [37]. Consider a rigid body with mass $M > 0$ and moment of inertia $J = J^T \in \mathbb{R}^{3\times 3}$ ($J > 0$). A rigid body is described with a pose $P_{\mathbb{SE}} \in \mathbb{SE}(3) :\Leftrightarrow (P, R) \in \mathbb{R}^3 \times \mathbb{SO}(3)$, as a pair composed of a linear position and angular position (orientation); and a twist $V_{\mathbb{SE}} \in \mathbb{R}^6 :\Leftrightarrow (V, \Omega) \in \mathbb{R}^3 \times \mathbb{R}^3$, as a pair composed of a linear velocity and an angular velocity (the linear velocity V is expressed in the inertial frame, and the angular velocity Ω is expressed in the body frame). A rigid body kinematics are then given by (hereafter, you may read the subindex rb as rigid body)

$$K_{rb} : \mathbb{SE}(3) \times \mathbb{R}^6 \ni P_{\mathbb{SE}} \times V_{\mathbb{SE}} \mapsto K_{rb}(P_{\mathbb{SE}}, V_{\mathbb{SE}}) \in T_{P_{\mathbb{SE}}}\mathbb{SE}(3)$$
$$K_{rb}(P_{\mathbb{SE}}, V_{\mathbb{SE}}) := \begin{bmatrix} V \\ R\mathcal{S}(\Omega) \end{bmatrix} \left(= \begin{bmatrix} \dot{P} \\ \dot{R} \end{bmatrix} = \dot{P}_{\mathbb{SE}} \right). \quad (1)$$

Suppose the rigid body is subjected to a wrench $W_{net} \in \mathbb{R}^6 :\Leftrightarrow (F_{net}, \tau_{net}) \in \mathbb{R}^3 \times \mathbb{R}^3$ with a net force F_{net} expressed in the inertial frame and a net torque τ_{net} expressed in the body frame. The rigid body dynamics are then given by

$$D_{rb} : \mathbb{R}^6 \times \mathbb{R}^6 \ni V_{\mathbb{SE}} \times W_{net} \mapsto D_{rb}(V_{\mathbb{SE}}, W_{net}) \in \mathbb{R}^6 \quad (2)$$
$$D_{rb}(V_{\mathbb{SE}}, W_{net}) := \begin{bmatrix} \frac{1}{M} F_{net} \\ J^{-1}(\tau_{net} - \mathcal{S}(\Omega) J\Omega) \end{bmatrix} \left(= \begin{bmatrix} \dot{V} \\ \dot{\Omega} \end{bmatrix} = \dot{V}_{\mathbb{SE}} \right).$$

Finally, we note that we equip $\mathbb{SO}(3)$ and $\mathbb{SE}(3)$ with the standard metrics ($d_{\mathbb{SO}}$ and $d_{\mathbb{SE}}$), and that the metric $d_{\mathbb{SO}}$ is insensitive to rotations of its arguments [37].



(a) Poses of UAV and manipulator, and constraint between rigid bodies.

(b) Net forces and torques applied on each rigid body, namely the UAV and the manipulator.

Fig. 2: Modeling of coupled UAV and manipulator, the two rigid bodies the system is composed of. Figure 1 provides a screenshot of the real system.

III. Modeling and Problem Statement

Consider the system illustrated in Fig. 2, composed of two rigid bodies, namely one quadrotor (which, for brevity, we refer to as UAV hereafter) and a manipulator (with a load at its end-effector). The two rigid bodies are coupled, with a ball-joint connecting them at the center-of-mass of the UAV. If the ball-joint were absent, the manipulator would behave as a free falling (un-actuated) rigid body, and the UAV would behave as a *standard* UAV. In its presence, the ball-joint imposes a kinematic constraint, specifically, it enforces the UAV's position to be fixed in the manipulator's orientation frame. This kinematic constraint *links* the UAV and the manipulator, and it provides a way to control the manipulator by means of actuation on the UAV.

For brevity, let us list the physical constants that describe the system (we use upper-case symbols for the UAV's constants, and lower-case symbols for the manipulator's):

- distance between rigid bodies centers-of-mass: $l > 0$.
- UAV's mass and moment of inertia: $M > 0$ and $J \in \mathbb{R}^{3\times3}$;
- manipulator's mass and moment of inertia: $m > 0$ and

$$j := j_{xy} \oplus j_{xy} \oplus j_{zz} = \begin{bmatrix} j_{xy} & 0 & 0 \\ 0 & j_{xy} & 0 \\ 0 & 0 & j_{zz} \end{bmatrix} \in \mathbb{R}^{3\times3}. \tag{3}$$

Let us also list the poses (linear and angular positions) and twists (linear and angular velocities) of the two rigid bodies that compose the system:

- UAV's pose – $P_{\mathbb{SE}} \in \mathbb{SE}(3) :\Leftrightarrow (P, R) \in \mathbb{R}^3 \times \mathbb{SO}(3)$;
- manipulator's pose – $p_{\mathbb{SE}} \in \mathbb{SE}(3) :\Leftrightarrow (p, r) \in \mathbb{R}^3 \times \mathbb{SO}(3)$;
- UAV's twist – $V_{\mathbb{SE}} \in \mathbb{R}^6 :\Leftrightarrow (V, \Omega) \in \mathbb{R}^3 \times \mathbb{R}^3$;
- manipulator's twist – $v_{\mathbb{SE}} \in \mathbb{R}^6 :\Leftrightarrow (v, \omega) \in \mathbb{R}^3 \times \mathbb{R}^3$.

Finally, let us list the inputs that act on the system, and for which we propose control laws later:

- UAV's thrust input – $U \in \mathbb{R}$;
- UAV's torque input in the UAV's body frame – $\tau \in \mathbb{R}^3$;
- manipulator's torque input in the manipulator's body frame –

$$\tau_m := \tau_{m,1} e_1 + \tau_{m,2} e_2 \in \mathbb{R}^3. \tag{4}$$

Remark 1: As illustrated in Fig. 2b, the axis of axial symmetry of the manipulator is the third axis, and in (4) one notices that a torque component along that axis has been omitted. In this work, a goal for the "yaw position" of the manipulator has not been prescribed, since the manipulator exhibits axial symmetry. If such a goal exists, a torque component along the third axis may be included, and it will only affect the "yaw motion" of the manipulator, but not the rest of the motion of the system.

For brevity, denote the pose and twist of the system as the pair of poses and twists of the individual rigid bodies, i.e.,

$$\mathcal{P} \in (\mathbb{SE}(3))^2 :\Leftrightarrow (p_{\mathbb{SE}}, P_{\mathbb{SE}}) \in \mathbb{SE}(3) \times \mathbb{SE}(3), \text{and} \tag{5}$$
$$\mathcal{V} \in (\mathbb{R}^6)^2 :\Leftrightarrow (v_{\mathbb{SE}}, V_{\mathbb{SE}}) \in \mathbb{SE}(3) \times \mathbb{SE}(3). \tag{6}$$

We let our state be composed of the system pose and twist as defined in (5) and (6), and define the state space as

$$\mathbb{X} := \big\{(\mathcal{P}, \mathcal{V}) \in (\mathbb{SE}(3))^2 \times (\mathbb{R}^6)^2 : \\ p + lre_3 - P = 0_3, v + lr\mathcal{S}(\omega) e_3 - V = 0_3\big\}, \tag{7}$$

which encapsulates the constraints imposed by the ball-joint, specifically, that the UAV's position is fixed w.r.t. the manipulator's orientation frame, i.e., that $r^T(P-p) = le_3 \Leftrightarrow p + lre_3 - P = 0_3$ (recall that $r \in \mathbb{SO}(3)$ represents the manipulator orientation frame). Hereafter, we always decompose a state and an input in the same way, namely

$$x \in \mathbb{X} :\Leftrightarrow (\mathcal{P}, \mathcal{V}) \in \mathbb{X} :\Leftrightarrow (p_{\mathbb{SE}}, P_{\mathbb{SE}}, v_{\mathbb{SE}}, V_{\mathbb{SE}}) \in \mathbb{X} \\ :\Leftrightarrow (p, r, P, R, v, \omega, V, \Omega) \in \mathbb{X}, \tag{8}$$
$$u \in \mathbb{U} := \mathbb{R}^7 :\Leftrightarrow (U, \tau, \tau_m) \in \mathbb{R} \times \mathbb{R}^3 \times \mathbb{R}^3. \tag{9}$$

Because we are interested in tracking problems for the system just described, we need a metric in the state space $\mathbb{X}$, which is inherited from the metric in $\mathbb{SE}(3)$ [37]; specifically,

$$d_{\mathbb{X}} : \mathbb{X} \times \mathbb{X} \ni (x^1, x^2) \mapsto d_{\mathbb{X}}(x^1, x^2) \in \mathbb{R}_{\geq 0} \tag{10}$$
$$d^2_{\mathbb{X}}(x^1, x^2) := d^2_{\mathbb{SE}}(p^1_{\mathbb{SE}}, p^2_{\mathbb{SE}}) + \|v^1_{\mathbb{SE}} - v^2_{\mathbb{SE}}\|^2 + \\ d^2_{\mathbb{SE}}(P^1_{\mathbb{SE}}, P^2_{\mathbb{SE}}) + \|V^1_{\mathbb{SE}} - V^2_{\mathbb{SE}}\|^2.$$

Given the state space definition in (7), one may compute its tangent set at any point $x \in \mathbb{X}$ (see state decomposition

in (8)), which is given by

$$T_x\mathbb{X} := \big\{(\delta\mathcal{P}, \delta\mathcal{V}) \in T_{\mathcal{P}}(\mathbb{SE}(3))^2 \times (\mathbb{R}^6)^2 : \quad (11)$$
$$\delta p + l\delta r e_3 - \delta P = 0_3,$$
$$\delta v + l\delta r \mathcal{S}(\omega) e_3 + lr\mathcal{S}(\delta\omega) e_3 - \delta V = 0_3\big\},$$

which is critical in computing the internal forces that allow the restrictions in (7) to be satisfied. Given an appropriate input $u : \mathbb{R}_{\geq 0} \ni t \mapsto u(t) \in \mathbb{U}$, a system's trajectory $x : \mathbb{R}_{\geq 0} \ni t \mapsto x(t) \in \mathbb{X}$ evolves according to

$$\dot{x}(t) = X(x(t), u(t)), x(0) \in \mathbb{X}, \quad (12)$$

where the vector field X is given by

$$X : \mathbb{X} \times \mathbb{U} \ni (x, u) \mapsto X(x, u) \in T_x\mathbb{X}$$
$$X(x,u) := \begin{bmatrix} K_X(\mathcal{P}, \mathcal{V}) \\ D_X(x,u) \end{bmatrix} \left(= \begin{bmatrix} \dot{\mathcal{P}} \\ \dot{\mathcal{V}} \end{bmatrix} = \begin{bmatrix} \text{kinematics} \\ \text{dynamics} \end{bmatrix}\right), \quad (13)$$

where the kinematics K_X are given by (recall the rigid body kinematics K_{rb} in (1))

$$K_X(\mathcal{P}, \mathcal{V}) := \left[\begin{array}{c} K_{rb}(p_{\mathbb{SE}}, v_{\mathbb{SE}}) \\ K_{rb}(P_{\mathbb{SE}}, V_{\mathbb{SE}}) \end{array}\right] \left(= \left[\begin{array}{c} \dot{p}_{\mathbb{SE}} \\ \dot{P}_{\mathbb{SE}} \end{array}\right]\right) \quad (14)$$
$$= \left[\begin{array}{c} v \\ r\mathcal{S}(\omega) \\ \hline V \\ R\mathcal{S}(\Omega) \end{array}\right] \left(= \left[\begin{array}{c} \dot{p} \\ \dot{r} \\ \hline \dot{P} \\ \dot{R} \end{array}\right]\right), \quad (15)$$

and the dynamics D_X are given by (recall the rigid body dynamics D_{rb} in (2))

$$D_X(x,u) := \left[\begin{array}{c} D_{rb}(v_{\mathbb{SE}}, W_{net}^{man}) \\ \hline D_{rb}(V_{\mathbb{SE}}, W_{net}^{uav}) \end{array}\right] \left(= \left[\begin{array}{c} \dot{v}_{\mathbb{SE}} \\ \hline \dot{V}_{\mathbb{SE}} \end{array}\right]\right), \quad (16)$$

$$W_{net}^{man} \equiv \begin{bmatrix} F_{net}^{man} \\ \tau_{net}^{man} \end{bmatrix} \equiv \begin{bmatrix} T(x,u) - mge_3 \\ \tau_m + l\mathcal{S}(e_3) r^T T(x,u) \end{bmatrix}, \quad (17)$$

$$W_{net}^{uav} \equiv \begin{bmatrix} F_{net}^{uav} \\ \tau_{net}^{uav} \end{bmatrix} \equiv \begin{bmatrix} URe_3 - T(x,u) - Mge_3 \\ \tau - R^T r \tau_m \end{bmatrix}, \quad (18)$$

with g as the acceleration due to gravity and where the internal tension forces $T(x,u)$ (expressed in the inertial frame) are given by (20). (We emphasize the introduction of the adimensional constant γ in (21); we introduce it for convenience, and, at this point, we only emphasize that $\gamma = 1$ when the manipulator is taken as a point-mass.) The analytic expression of the wrenches (pair net force and net torque) in (17) and (18) applied on the manipulator and the UAV is deduced by inspection of Fig. 2b. The system dynamics in (16) can be written equivalently as

$$D_X(x,u) = \quad (19)$$
$$= \left[\begin{array}{c} \frac{1}{m}(T(x,u) - mge_3) \\ \hline j^{-1}(\tau_m + l\mathcal{S}(e_3) r^T T(x,u) - \mathcal{S}(\omega) j\omega) \\ \frac{1}{M}(URe_3 - T(x,u) - Mge_3) \\ J^{-1}(\tau - R^T r\tau_m - \mathcal{S}(\Omega) J\Omega) \end{array}\right] \left(= \left[\begin{array}{c} \dot{v} \\ \hline \dot{\omega} \\ \dot{V} \\ \dot{\Omega} \end{array}\right]\right)$$
$$\overset{(4)}{\underset{(3)}{=}} \left[\begin{array}{c} \frac{1}{m}(T(x,u) - mge_3) \\ \hline \frac{\tau_m}{j_{xy}} + \frac{l}{j_{xy}}\mathcal{S}(e_3) r^T T(x,u) + \left(\frac{j_{zz}}{j_{xy}} - 1\right) e_3^T \omega \mathcal{S}(e_3)\omega \\ \hline \frac{1}{M}(URe_3 - T(x,u) - Mge_3) \\ J^{-1}(\tau - R^T r\tau_m - \mathcal{S}(\Omega) J\Omega) \end{array}\right].$$

The vector field X in (13) is composed of the system kinematics and dynamics. The system kinematics are given by the kinematics of each individual rigid body, in this particular case of the UAV and of the manipulator. On the other hand, the system dynamics are given by the linear and angular accelerations of each individual rigid body, by considering the net force (expressed in the inertial frame) and the net torque (expressed in the respective body frame) applied on each rigid body. The forces and torques contributing to the net force and net torque of each rigid body are visualized in Fig. 2b. It may be useful to discuss the net wrenches applied to each rigid body. The tension forces ($T(x,u) \in \mathbb{R}^3$ in (20), expressed in the inertial frame) are internal forces, forming an action-reaction pair, and thus explaining its positive contribution to the manipulator's linear acceleration and its negative contribution to the UAV's. The manipulator input torque ($\tau_m \in \mathbb{R}^3$ in (4), expressed in the manipulator frame) also forms an action-reaction pair, thus explaining its positive contribution to the manipulator's angular acceleration and its negative contribution to the UAV's (in the UAV's orientation frame, the manipulator input torque is given by $-R^T r\tau_m$). The reaction forces to the rigid body weights are applied on the earth (an infinite mass rigid body), and are thus ignored in our modeling. The reaction force and reaction torque to the UAV's thrust input (URe_3) and torque input (τ) are applied on the air and on the motor propellers, and are thus ignored in our modeling.

The tension forces ($T(x,u) \in \mathbb{R}^3$) constitute internal forces to the system, and the Newton-Euler's equations of motion do not provide any insight into these forces. However, the constraint that the state must remain in the state space $\mathbb{X}$ in (7), enforces the vector field X in (13) to be in the tangent set of that set. This constraint uniquely defines the tensions, i.e., for any $(x,u) \in \mathbb{X} \times \mathbb{U}$,

$$X(x,u) \in T_x\mathbb{X} \Rightarrow T(x,u) \text{ as in (20).}$$

The Euler-Lagrange formalism provides an alternative, but equivalent, approach for obtaining the vector field in (13).

Proposition 2: The rotation of the manipulator around its axis of axial symmetry (third axis) is constant, i.e., along a solution $t \to x(t) \in \mathbb{X}$ of (12), $t \to e_3^T\omega(t) = e_3^T\omega(0) \in \mathbb{R}$.

Proof 1: For convenience, define the function $f : \mathbb{X} \ni x \mapsto f(x) := e_3^T\omega \in \mathbb{R}$. It follows that, for any $(x,u) \in \mathbb{X} \times \mathbb{U}$, $Df(x)X(x,u) = e_6^T D_X(x,u) = e_3^T j^{-1}(\tau_m + l\mathcal{S}(e_3) T(x,u) - \mathcal{S}(\omega) j\omega) \overset{(3),(4)}{=} 0$, which concludes the proof.

Proposition 2 implies that the rotation of the manipulator around its axis of symmetry, despite uncontrollable, remains constant along solutions. One could, at this point, assume that the rotation around that axis is non-existing at the initial time instant; however, we shall make that assumption only in one of the problems treated in this manuscript (for the UAV-slung-manipulator system, that we describe next).

Proposition 3: If the UAV has axial symmetry around its third axis, then the vector field X in (13) is invariant to rotations of the manipulator around its third axis (the axis of axial symmetry) and to rotations of the UAV around its third axis (the thrust axis).

A proof is found in [37]. Proposition 3 implies that the dynamics do not *see* rotations of the rigid bodies around their third axes. This motivates us to control the space orthogonal

$$T : \mathbb{X} \times \mathbb{U} \ni (x,u) \mapsto T(x,u) \in \mathbb{R}^3 \tag{20}$$

$$\begin{aligned} T(x,u) &:= r\left(\left(\frac{1}{m}+\frac{1}{M}\right) I_3 + \frac{l^2}{j_{xy}}\Pi\left(e_3\right)\right)^{-1}\left(\frac{U}{M} r^T R e_3 + \frac{l}{j_{xy}}\mathcal{S}\left(e_3\right)\tau_m + l\left(\omega^T \Pi\left(e_3\right)\omega e_3 - \frac{j_{zz}}{j_{xy}} e_3^T \omega \Pi\left(e_3\right)\omega\right)\right) \\ &= \frac{1}{\gamma} r\left(\left(\frac{j_{xy}}{ml^2} I_3 + \frac{M}{m+M} e_3 e_3^T\right)\left(\frac{m}{M} U r^T R e_3 + ml\omega^T \Pi\left(e_3\right)\omega e_3\right) + \frac{1}{l}\left(\mathcal{S}\left(e_3\right)\tau_m - j_{zz} e_3^T \omega \Pi\left(e_3\right)\omega\right)\right) \end{aligned}$$

$$\gamma \equiv 1 + \frac{j_{xy}}{ml^2}\frac{m+M}{M} \tag{21}$$

to the manipulator's third axis and the UAV's third axis as separate problems. Loosely speaking, we can treat the yaw position of the manipulator and of the UAV as separate problems: the yaw motion of the UAV can be understood from Re_1, Re_2 and $e_3^T\Omega$; similarly, the yaw motion of the manipulator can be understood from re_1, re_2 and $e_3^T\omega$ (and we shall ignore this motion since $e_3^T\omega$ is uncontrollable – see Proposition 2).

Invariance to rotations around the third axes also motivates us to introduce an equivalence relation in the set $\mathbb{X}$ in (7). To be specific, consider two states in $\mathbb{X}$, namely $x_1 = (\mathcal{P}_1, \mathcal{V}_1) \in \mathbb{X}$ and $x_2 = (\mathcal{P}_2, \mathcal{V}_2) \in \mathbb{X}$; then we define the relation $\sim$ as (below $R_z(\theta) := I_3 + \sin(\theta)\mathcal{S}\left(e_3\right) + (\cos(\theta)-1)\Pi\left(e_3\right)$)

$$x_1 \sim x_2 :\Leftrightarrow \tag{22}$$

$$\begin{bmatrix} p_1 \\ r_1 \\ P_1 \\ R_1 \end{bmatrix} = \begin{bmatrix} p_2 \\ r_2 R_z(\theta_p) \\ P_2 \\ R_2 R_z(\theta_P) \end{bmatrix} \text{ and } \begin{bmatrix} v_1 \\ \omega_1 \\ V_1 \\ \Omega_1 \end{bmatrix} = \begin{bmatrix} v_2 \\ R_z^T(\theta_p)(\omega_2 + be_3) \\ V_2 \\ R_z^T(\theta_P)(\Omega_2 + Be_3) \end{bmatrix},$$

for some real θ_p, θ_P, b and B; i.e., two states are equivalent up to rotations around their third axes and up to the third component of their body-frame angular-velocities. One can verify that the relation (22) is reflexive, symmetric and transitive, which supports the following result.

Proposition 4: The relation defined in (22) is an equivalence relation in $\mathbb{X}$.

Proposition (4) guarantees that one may consider the quotient space induced by the equivalence relation $\sim$ in (22), namely

$$\mathbb{X}\backslash_\sim := \{[x] : x \in \mathbb{X}\}, \tag{23}$$

where $[x] := \{x' \in \mathbb{X} : x' \sim x\}$ denotes the equivalence class associated to some $x \in \mathbb{X}$. We will require the quotient set later, which is the reason why we introduce it at this point. Also, since the standard metric in $\mathbb{SO}(3)$ is insensitive to rotation of its arguments, it follows that the quotient set $\mathbb{X}\backslash_\sim$ can inherit the metric $d_{\mathbb{X}}$ as defined in (10), leading to the definition

$$\begin{aligned} &d_{\mathbb{X}\backslash_\sim} : \mathbb{X}\backslash_\sim \times \mathbb{X}\backslash_\sim \ni ([x^1],[x^2]) \mapsto d_{\mathbb{X}\backslash_\sim}([x^1],[x^2]) \in \mathbb{R}_{\geq 0} \\ &d_{\mathbb{X}\backslash_\sim}([x^1],[x^2]) := d_{\mathbb{X}}(x^1, x^2). \end{aligned} \tag{24}$$

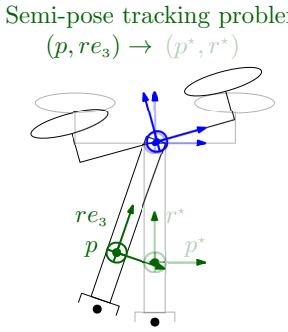
Proposition 5: The metric $d_{\mathbb{X}\backslash_\sim}$ in (24) is well-defined.

A proof is found in [37].

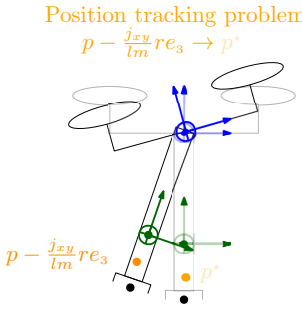
We consider two separate problems – see Fig. 3 – on the system just described: in the first we consider the manipulator torque input in (4) to be available, while in the second we restrict that torque input to be zero.

A. UAV-manipulator system

Hereafter, we refer to *UAV-manipulator system* when considering the manipulator torque input in (4) is available. With



(a) Semi-pose tracking problem for *UAV-manipulator system*: desired semi-pose trajectory $(p^\star, r^\star)$ determines the whole state equilibrium in transparent; i.e., the system is differentially flat w.r.t. (p, re_3).



(b) Position tracking problem for *UAV-slung-manipulator system*: desired position trajectory $p^\star$ determines the whole state equilibrium in transparent; i.e., the system is differentially flat w.r.t. $p - \frac{j_{xy}}{lm} re_3$.

Fig. 3: Two problems: one where the manipulator torque τ_m is available – Fig. 3a; and one where it is not – Fig. 3b.

Fig. 3a in mind, we can then formulate the first problem treated in this paper, which is one of semi-pose tracking: we require the manipulator's position to track a desired position trajectory, and the manipulator axis of axial symmetry (re_3 in our case) to track a desired attitude trajectory. The space orthogonal to the manipulator axis of axial symmetry is ignored, for the reasons discussed after Proposition 2 (and thus the semi-pose tracking rather than (full-)pose tracking).

Problem 1 (UAV-Manipulator): Consider the vector field X in (13). Given a desired position trajectory and a desired orientation trajectory, i.e,

$$p^\star : \mathbb{R}_{\geq 0} \ni t \mapsto p^\star(t) \in \mathbb{R}^3, \tag{25}$$

$$r^\star : \mathbb{R}_{\geq 0} \ni t \mapsto r^\star(t) \in \mathbb{S}^2, \tag{26}$$

design a control law

$$u^{cl} := (U^{cl}, \tau^{cl}, \tau_m^{cl}) : \mathbb{R}_{\geq 0} \times \mathbb{X} \to \mathbb{U}, \tag{27}$$

such that $\lim_{t\to\infty}(p(t) - p^\star(t)) = 0_3$ and $\lim_{t\to\infty}(re_3 - r^\star(t)) = 0_3$ along the solution $\mathbb{R}_{\geq 0} \ni t \mapsto x(t) \in \mathbb{X}$ of

$\dot{x}(t) = X(x(t), u^{cl}(t, x(t)))$ with $x(0) \in \mathbb{X}_0$ for some non-empty $\mathbb{X}_0 \subset \mathbb{X}$.

Remark 6: The problem treated here is a generalization of that treated in [47]. In [47], the manipulator is treated as a point-mass and only the kinematics of the UAV's attitude are considered. The vector field presented there is in fact recovered as follows:

- Since the constraints in (7) must be satisfied, one imposes that $re_3 = \frac{P-p}{l}$ and that $r\omega = \mathcal{S}\left(\frac{P-p}{l}\right)\frac{V-v}{l}$.
- Since the manipulator is taken as a point mass (and thus has no angular position and angular velocity, i.e., r and ω), the manipulator attitude equations ($\dot{r}$ in (15), and $\dot{\omega}$ in (19)) are ignored.
- Finally, the internal forces are found by taking the limit when the moment of inertia of the manipulator vanishes, i.e., when $j \to 0_{3\times 3} \overset{(3)}{\Leftrightarrow} (j_{xy}, j_{zz}) \to (0, 0)$ [37].

B. UAV-slung-manipulator system

Hereafter, we refer to *UAV-slung-manipulator system* when considering that the manipulator torque input in (4) is not available (Remark 7 sheds some light into the name *UAV-slung-manipulator system*). With some abuse of notation, and without hindering comprehension, when referring to the UAV-slung-manipulator system, we reuse and redefine the input u and the vector field X; to be specific, instead of the input in (9), we redefine it as

$$u = (U, \tau) \in \mathbb{R} \times \mathbb{R}^3, \tag{28}$$

and, instead of the vector field in (13), we redefine it as

$$\begin{aligned} &X : \mathbb{X} \times \mathbb{R}^4 \ni (x, u) \mapsto X(x, u) \in T_x\mathbb{X} \\ &X(x, u) := X(x, (U, \tau, 0_3))|_{X \text{ in } (13)}. \end{aligned} \tag{29}$$

With Fig. 3b in mind, we can then formulate the second problem treated in this paper, which is one of position tracking: we require a specific point along the manipulator's axis of axial symmetry to track a desired position trajectory.

Problem 2 (UAV-slung-manipulator): Consider the vector field X in (29). Given a desired position trajectory, i.e,

$$p^\star : \mathbb{R}_{\geq 0} \ni t \mapsto p^\star(t) \in \mathbb{R}^3, \tag{30}$$

design a control law

$$u^{cl} := (U^{cl}, \tau^{cl}) : \mathbb{R}_{\geq 0} \times \mathbb{X} \mapsto \mathbb{R}^4 \tag{31}$$

such that $\lim_{t\to\infty}(p(t) - \frac{j_{xy}}{lm} r(t) e_3 - p^\star(t)) = 0_3$ along a solution $\mathbb{R}_{\geq 0} \ni t \mapsto x(t) \in \mathbb{X}$ of $\dot{x}(t) = X(x(t), u^{cl}(t, x(t)))$ with $x(0) \in \mathbb{X}_0$ for some non-empty $\mathbb{X}_0 \subset \mathbb{X}$.

Remark 7: The problem treated here is a generalization of that treated in [48], which deals with the slung-load problem. In [48], the UAV is tethered to a point mass, which is the system here described if we take the manipulator to be a point-mass. The vector field presented there is in fact recovered following the exact same procedure as in Remark 6.

Remark 8: In Problem 2, rather than requiring the manipulator's position p to track the desired trajectory, we require another position along the manipulator's axis of symmetry ($p - \frac{j_{xy}}{lm} r e_3$) to track the desired trajectory (note however that those positions coincide when the manipulator is taken as a point mass, i.e., when $j_{xy} = 0$). Let us motivate here why. In a UAV-slung-load problem, requiring the UAV to track a desired trajectory is an ill-defined problem, because the system is differentially flat with respect to the load position, but not with respect to the UAV position: in other words, if the UAV tracks a certain trajectory, the trajectory the load tracks depends on the initial condition. Similarly, in the UAV-slung-manipulator problem, requiring the manipulator (center-of-mass) to track a desired trajectory is an ill-defined problem, because the system is not differentially flat with respect to the manipulator's center-of-mass: rather it is differentially flat with respect to another point (to be specific, $p - \frac{j_{xy}}{lm} r e_3$).

Remark 9: Notice that Problem 1 describes a semi-pose tracking problem, while Problem 2 describes a position tracking problem. This difference stems from the fact that the UAV-slung-manipulator system is deprived of the manipulator torque input, which prevents the manipulator pose from tracking an arbitrary pose. In fact, as we verify later, for the UAV-slung-manipulator system, the desired position trajectory in (30) imposes a desired semi-pose trajectory (this is verified by means of differential flatness).

IV. Control strategy summary

In this section, we explain the pursued control strategy, illustrated in Fig. 4. The strategy is composed of three steps, which we explain next. The specifics of these three steps are different when solving Problems 1 and 2, but the overall idea behind them is the same.

Recall then Proposition 3, which states that rotations around the rigid bodies' third axes may be ignored. Our first step is to design a bijective mapping (the specific mapping f and the codomain $\mathbb{Y}$ are specific to each Problem 1 and 2, and are provided later)

$$f : \mathbb{X}\backslash_\sim \ni [x] \mapsto f([x]) =: y \in \mathbb{Y} \tag{32}$$

which reorders the state in a manner that highlights the hierarchical/cascaded structure of the problem, while ignoring the rotations around the rigid bodies third axes: the necessity of the equivalence relation $\sim$ in (22) and of the quotient set $\mathbb{X}\backslash_\sim$ in (23) becomes clear at this point. This step is called state transformation – see Figure 4.

In our second step, we design a control law

$$\bar{u}^{cl} : \mathbb{X} \times \mathbb{U} \times \mathbb{R} \ni (x, v, \tau_\psi) \mapsto \bar{u}^{cl}(x, v, \tau_\psi) \in \mathbb{U} \tag{33}$$

which transforms the input (from v to u – see input transformation in Fig. 4), in such a way that when we compose this input transformation and state transformation with the system dynamics, we obtain the transformed dynamics in $\mathbb{Y}$, i.e.,

$$\begin{aligned} &Y : \mathbb{Y} \times \mathbb{U} \ni (y, v) \mapsto Y(y, v) \in T_y\mathbb{Y} \\ &Y(y, v) := Df([x])X(x, \bar{u}^{cl}(x, v, \tau_\psi))|_{x\in f^{-1}(y)} (= \dot{y}), \end{aligned} \tag{34}$$

where we emphasize that f^{-1} exists since the map f is invertible. Let us emphasize some key points regarding the input transformation $\bar{u}^{cl}$ in (33) and the vector field Y in (34):

- The input transformation $\bar{u}^{cl}$ is designed so as to make make explicit the cascaded structure of the vector field Y (one may think of a chain of integrators);
- The vector field Y is independent of $\tau_\psi \in \mathbb{R}$, which is used to control the space orthogonal to the UAV's third

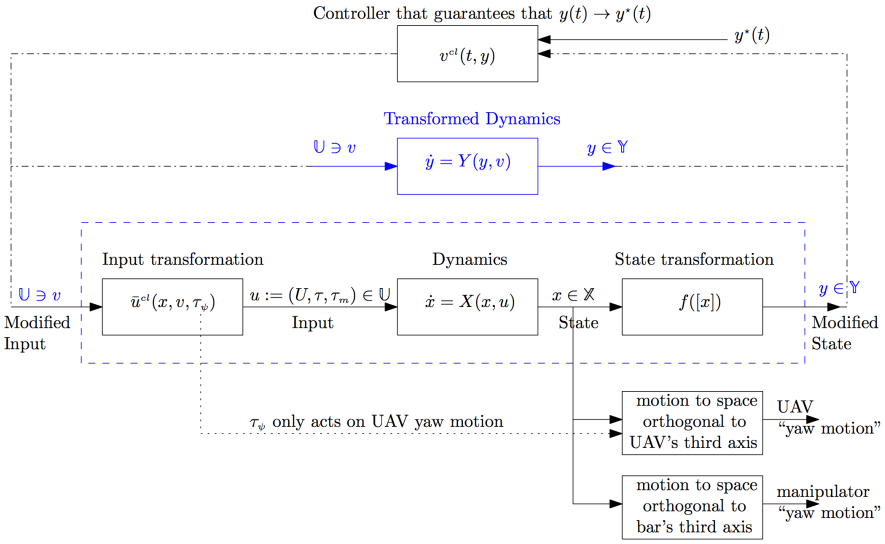


Fig. 4: Control strategy block-diagram.

axis (the yaw motion of the UAV, loosely speaking), but not to control the transformed state y in (32) (the transformed state does not *see* rotations around the rigid bodies third axes since f maps from the quotient set $\mathbb{X}\backslash_\sim$);

- We (must) show that the vector field Y in (34) is well defined, in the sense that (34) is independent of the choice of the representative belonging the equivalence class $f^{-1}(y)$ (notice that for any $y \in \mathbb{Y}$, $f^{-1}(y)$ is an equivalence class, i.e., $f^{-1}(y) \in \mathbb{X}\backslash_\sim$).

In the third and final step, we analyze differential flatness properties of the vector field Y, and based on that determine the equilibrium (more precisely the equilibria)

$$y^\star : \mathbb{R} \ni t \mapsto y^\star(t) \in \mathbb{Y} \tag{35}$$

for which it is guaranteed that Problems 1 and 2 are satisfied. Finally, we design a control law

$$v^{cl} : \mathbb{R} \times \mathbb{Y} \ni (t,y) \mapsto v^{cl}(t,y) \in \mathbb{U} \tag{36}$$

that guarantees that a solution $t \mapsto y(t)$, of $\dot{y}(t) = Y(y(t), v^{cl}(t, y(t)))$, tracks the desired equilibrium $t \mapsto y^\star(t)$. Problems 1 and 2 are then accomplished if one composes the input transformation $\bar{u}^{cl}$ in (33) with the control law v^{cl} in (36), i.e.,

$$\begin{aligned} &u^{cl} : \mathbb{R} \times \mathbb{X} \ni (t,x) \mapsto u^{cl}(t,x) \in \mathbb{U} \\ &u^{cl}(t,x) := \bar{u}^{cl}(x, v^{cl}(t, f([x])), \tau_\psi^{cl}(t,x)), \end{aligned} \tag{37}$$

where $\tau_\psi^{cl} : \mathbb{R} \times \mathbb{X} \ni (t,x) \mapsto \tau_\psi^{cl}(t,x) \in \mathbb{R}$ is some control law designed to control the yaw motion of the UAV (see Fig 4). Let us now look at each Problem individually.

Remark 10: Suppose $f : \mathbb{X} \ni x \to f(x) \in \mathbb{Y}$, instead of that in (32), i.e. that f provides a bijective map from $\mathbb{X}$ to $\mathbb{Y}$, or, in other words, a coordinate change from $\mathbb{X}$ to $\mathbb{Y}$. Then, equation (34) would be instead

$$Y(y,v) := Df(x)X(x, \bar{u}^{cl}(x,v,\tau_\psi))|_{x=f^{-1}(y)} (= \dot{y}),$$

which is the standard equation for computing the vector field in the new coordinate system in $\mathbb{Y}$.

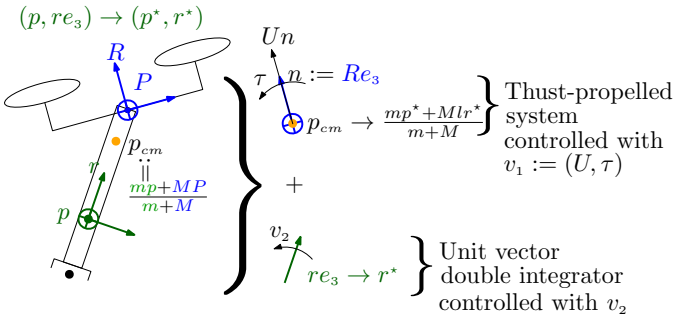


Fig. 5: By means of appropriate state and input transformations, the UAV-manipulator may be understood as two decoupled systems: one thrust-propelled system and one unit-vector double-integrator system.

V. UAV-MANIPULATOR PROBLEM

We verify next that the UAV-manipulator system behaves as two decoupled systems, as illustrated in Fig. 5: one thrust-propelled system by considering the motion of the center-of-mass of the whole system (at this point, the internal forces – $T(x,u)$ – do not play a role); and one unit vector double integrator system corresponding to the attitude dynamics of the manipulator. The three steps we follow next have been described in the previous section.

A. State transformation

Recall the discussion from Section IV Let us then introduce the set $\mathbb{Y}$ as

$$\mathbb{Y} = \mathbb{Y}_1 \times \mathbb{Y}_2 \tag{38a}$$

$$\mathbb{Y}_1 := \{(p_{cm}, v_{cm}, n, \varpi) \in (\mathbb{R}^3)^4 : n^T n = 1, n^T \varpi = 0\}, \tag{38b}$$

$$\mathbb{Y}_2 := \{(r, \omega) \in (\mathbb{R}^3)^2 : r^T r = 1, r^T \omega = 0\}, \tag{38c}$$

with the inherited metrics $d_{\mathbb{Y}_1} := d_{\mathbb{R}^{12}}$, $d_{\mathbb{Y}_2} := d_{\mathbb{R}^6}$ and $d_{\mathbb{Y}} := d_{\mathbb{R}^{18}}$. The set $\mathbb{Y}_1$ will be used for the subsystem that behaves as a thrust-propelled system; while the set $\mathbb{Y}_2$ will be used for the subsystem that behaves as a unit vector double

integrator system (see Fig. 5). Hereafter, and for convenience, we decompose elements in the sets above as

$$y \in \mathbb{Y} :\Leftrightarrow (y_1, y_2) \in \mathbb{Y}_1 \times \mathbb{Y}_2, \tag{39a}$$
$$y_1 \in \mathbb{Y}_1 :\Leftrightarrow (p_{cm}, v_{cm}, n, \varpi) \in \mathbb{Y}_1, \tag{39b}$$
$$y_2 \in \mathbb{Y}_2 :\Leftrightarrow (r, \omega) \in \mathbb{Y}_2. \tag{39c}$$

Hereafter, we denote as well

$$v \in \mathbb{U} :\Leftrightarrow (v_1, v_2) \in \mathbb{R}^4 \times \mathbb{R}^3, \tag{40a}$$
$$v_1 \in \mathbb{R}^4 :\Leftrightarrow (U, \tau) \in \mathbb{R} \times \mathbb{R}^3, \tag{40b}$$
$$v_2 \in \mathbb{R}^3 :\Leftrightarrow \tau_m \in \mathbb{R}^3, \tag{40c}$$

where v_1 will be the input to one subsystem (thrust propelled), and v_2 will be the input to the other subsystem (unit vector double integrator).

Based on Fig. 5, consider then the mapping

$$f : \mathbb{X}\backslash_\sim \ni [x] \mapsto f([x]) \in \mathbb{Y} \tag{41}$$
$$f([x]) := \begin{bmatrix} \frac{mp+MP}{m+M} \\ \frac{mv+MV}{m+M} \\ Re_3 \\ \Pi\left(Re_3\right) R\Omega \\ \hdashline re_3 \\ \Pi\left(re_3\right) r\omega \end{bmatrix} \left(= \begin{bmatrix} p_{cm} \\ v_{cm} \\ n \\ \varpi \\ \hdashline r \\ \omega \end{bmatrix} = \begin{bmatrix} y_1 \\ \\ \hdashline \\ y_2 \end{bmatrix} \right),$$

with $\mathbb{Y}$ as defined in (38a), and consider as well the mapping

$$h : \mathbb{Y} \ni y \mapsto h(y) \in \mathbb{X}\backslash_\sim \tag{42}$$
$$h(y) := \underbrace{\begin{bmatrix} p_{cm} - l\frac{M}{m+M} r \\ \bar{r} \\ p_{cm} + l\frac{m}{m+M} r \\ \bar{R} \\ v_{cm} - l\frac{M}{m+M} \mathcal{S}\left(\omega\right) r \\ \bar{r}^T \omega \\ v_{cm} + l\frac{m}{m+M} \mathcal{S}\left(\omega\right) r \\ \bar{R}^T \varpi \end{bmatrix}}_{\substack{\bar{r}\in\mathbb{SO}(3), \bar{r}e_3 = r \\ \bar{R}\in\mathbb{SO}(3), \bar{R}e_3 = n}} \left(= \begin{bmatrix} p \\ r \\ P \\ R \\ \hdashline v \\ \omega \\ V \\ \Omega \end{bmatrix} = [x] \right),$$

where $[\cdot]$ above denotes the equivalence class as defined immediately after (23).

Proposition 11: The map f in (41) is well defined, and

$$f \circ h = \mathrm{id}_{\mathbb{Y}} \text{ and } h \circ f = \mathrm{id}_{\mathbb{X}\backslash_\sim}, \tag{43}$$

i.e., $h = f^{-1}$ and $f = h^{-1}$.

Proof 2: The map f in (41) is defined with a representative of an equivalence class, so we must show that f is independent of the choice of representative. Notice that (below $R_z(\theta) := I_3 + \sin(\theta)\mathcal{S}\left(e_3\right) + (\cos(\theta) - 1)\Pi\left(e_3\right)$)

$$Re_3 = (RR_z(\theta))e_3, \text{ and that} \tag{44}$$
$$\Pi\left(Re_3\right) R\Omega = \Pi\left(\left(RR_z(\theta)\right)e_3\right)\left(RR_z(\theta)\right)(R_z^T(\theta)\Omega), \tag{45}$$

which suffices to conclude that f is independent of the choice of representative belonging to an equivalence class established by the equivalence relation in (22) (i.e., for any $x \in \mathbb{X}$ and for any $\bar{x}, \tilde{x} \in [x] \in \mathbb{X}\backslash_\sim$, $f([\bar{x}]) = f([\tilde{x}])$). The second part of the Proposition follows from straightforward computations. Let us just illustrate the identities in (43) for the component ϖ of f (denote it $f|_{\varpi}$) and the component Ω of h (denote it $h|_{\Omega}$)

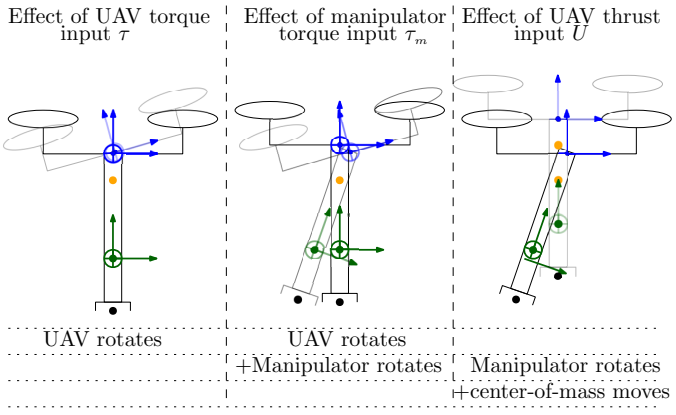


Fig. 6: Effects of original input $u = (U, \tau, \tau_m)$ in UAV-manipulator system.

– i.e., we verify that $f|_{\varpi} \circ h(y) = \varpi$ and that $h|_{\Omega} \circ f([x]) = [\Omega]$ (where, by the equivalence relation in (22), angular velocities are equivalent up to rotations around the third body axis and up to the third-component – i.e., $[\Omega] := [R_z(\theta_P)(\Omega + Be_3)]$ for any real θ_P and B). Then, notice that

$$\begin{aligned} f|_{\varpi} \circ h(y) =& f|_{\varpi}(h(y)) \overset{(41)}{\equiv} \Pi\left(Re_3\right) R\Omega|_{R \text{ and } \Omega \text{ as in (42)}} \\ =& \Pi\left(\bar{R}e_3\right) \bar{R}\bar{R}^T\varpi \\ \overset{\bar{R}\in\mathbb{SO}(3)}{=}& \Pi\left(\bar{R}e_3\right)\varpi \\ \overset{(42):\bar{R}e_3 = n}{=}& \Pi\left(n\right)\varpi = (I_3 - nn^T)\varpi \\ \overset{(39):n^T\varpi = 0}{=}& \varpi. \end{aligned} \tag{46}$$

Finally, notice that

$$\begin{aligned} h|_{\Omega} \circ f([x]) =& h|_{\Omega}(f([x])) \overset{(42)}{\equiv} [\bar{R}^T\varpi]|_{\varpi \text{ as in (41)}} \\ =& [\bar{R}^T\Pi\left(Re_3\right) R\Omega] = [\bar{R}^T R\Pi\left(e_3\right)\Omega] \\ \overset{\substack{Re_3 \overset{(41)}{=} n \overset{(42)}{=} \bar{R}e_3 \Rightarrow \\ \Rightarrow \bar{R}^T R = R_z^T(\theta_P) \text{ for some } \theta_P}}{=}& [R_z^T(\theta_P)\Pi\left(e_3\right)\Omega] \\ \overset{(22)}{=}& [\Pi\left(e_3\right)\Omega] = [\Omega - \Omega^T e_3 e_3] \\ \overset{(22)}{=}& [\Omega]. \end{aligned} \tag{47}$$

The other identities are verified similarly.

B. Input transformation and transformed vector field

Let us construct the input transformation in two steps. For that purpose, consider the transformed dynamics with the original input (not the transformed input), which are given by

$$Df([x])X(x,u) = \begin{bmatrix} \frac{mv+MV}{m+M} \\ \frac{U}{m+M} Re_3 - ge_3 \\ \mathcal{S}\left(\Pi\left(Re_3\right) R\Omega\right) Re_3 \\ \Pi\left(Re_3\right) R(\star_1) \\ \hdashline \mathcal{S}\left(\Pi\left(re_3\right) r\omega\right) re_3 \\ \Pi\left(re_3\right)(\star_2) \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_{cm} \\ \dot{v}_{cm} \\ \dot{n} \\ \dot{\varpi} \\ \hdashline \dot{r} \\ \omega \end{bmatrix} \right),$$
$$(\star_1) = J^{-1}(\tau - R^T r\tau_m - \mathcal{S}\left(\Omega\right) J\Omega) - e_3^T\Omega\mathcal{S}\left(\Omega\right) e_3 \tag{48}$$
$$(\star_2) = \tfrac{m+M}{l^2 mM}\tfrac{1}{\gamma}\left(\tau_m + \mathcal{S}\left(e_3\right)\left(\tfrac{mlU}{m+M} r^T Re_3 + j_{zz} e_3^T \omega\omega\right)\right).$$

If follows from (48) that the original input $u = (U, \tau, \tau_m)$ has the effects illustrated in Fig. 6. (i) The UAV torque input τ only makes the UAV rotate (not the manipulator). (ii) The manipulator torque input τ_m makes the manipulator rotate, but it also makes the UAV rotate (in an opposite direction,

and with a different scaling – factor $\|J^{-1}R^T r\| = \|J^{-1}\|$ in $\star_1$ and factor $\frac{m+M}{l^2 mM}\frac{1}{\gamma}$ in $\star_2$). (iii) The UAV thrust input U makes the whole system (i.e., its center-of-mass) accelerate along the UAV's third body axis; but it also produces a torque on the manipulator third axis, which makes it rotate. The previous discussion provides the basis for the design of the input transformation, whose idea is to design a UAV torque input τ that controls the UAV attitude by canceling the effect of the manipulator torque input τ_m on the UAV attitude; and next, to design a manipulator torque input τ_m that controls the manipulator attitude by canceling the effect of the UAV thrust input U on the manipulator's attitude (notice the design is made in a cascaded fashion). Consider then

$$\bar{u}_1^{cl} : \mathbb{X} \times \mathbb{U} \times \mathbb{R} \ni (x, u, \tau_\psi) \mapsto \bar{u}_1^{cl}(x, u, \tau_\psi) \in \mathbb{U} \tag{49}$$
$$\bar{u}_1^{cl}(x, u, \tau_\psi) := \begin{bmatrix} (m+M)U \\ \left\{ \begin{matrix} J(\Pi(e_3)R^T\tau + \tau_\psi e_3 + e_3^T \Omega \mathcal{S}(\Omega) e_3) + \\ + R^T r \tau_m + \mathcal{S}(\Omega) J \Omega \end{matrix} \right\} \\ \tau_m \end{bmatrix}$$

which when composed in (48) yields $\dot{v}_{cm} = URe_3 - ge_3$ and to $\dot{\varpi} = \Pi(Re_3)\tau$ (for the purpose of τ_ψ in (49), we ask the reader to await until Remark 13). Then consider

$$\bar{u}_2^{cl} : \mathbb{X} \times \mathbb{U} \ni (x, u, \tau_\psi) \mapsto \bar{u}_2^{cl}(x, u) \in \mathbb{U} \tag{50}$$
$$\bar{u}_2^{cl}(x, u) := \begin{bmatrix} U \\ \tau \\ \left\{ \begin{matrix} \frac{l^2 m M}{m+M} \gamma \Pi(e_3) r^T \tau_m - \\ -\mathcal{S}(e_3)\big(mlUr^T Re_3 + j_{zz} e_3^T \omega\omega\big) \end{matrix} \right\} \end{bmatrix}$$

which when composed in (48) (and after applying (49) first) it leads to $\dot{\omega} = \Pi(re_3)\tau_m$.

Consider then the input transformation as a composition of the previous control laws (49) and (50), i.e.,

$$\bar{u}^{cl} : \mathbb{X} \times \mathbb{U} \times \mathbb{R} \ni (x, u, \tau_\psi) \mapsto \bar{u}^{cl}(x, u, \tau_\psi) \in \mathbb{U} \tag{51}$$
$$\bar{u}^{cl}(x, u, \tau_\psi) := \bar{u}_1^{cl}(x, \bar{u}_2^{cl}(x, u), \tau_\psi).$$

Given the input transformation in (51) and the state transformation in (41), one can then compute the transformed dynamics. In fact, it follows that composing (48) with the control law (51) results in

$$Df([x])X(x, \bar{u}^{cl}(x, u, \tau_\psi)) = \begin{bmatrix} \frac{mv+MV}{m+M} \\ URe_3 - ge_3 \\ \mathcal{S}(\Pi(Re_3)R\Omega)Re_3 \\ \Pi(Re_3)\tau \\ \hline \mathcal{S}(\bar{\Pi}(re_3)r\omega)re_3 \\ \Pi(re_3)\tau_m \end{bmatrix} \tag{52}$$

and as such the vector field in the new coordinates is given by

$$Y : \mathbb{Y} \times \mathbb{U} \ni (y, v) \mapsto Y(y, v) \in T_y\mathbb{Y} \tag{53}$$
$$Y(y, v) := Df([x])X(x, \bar{u}^{cl}(x, v, \cdot))|_{x \in h(y)}$$
$$= \begin{bmatrix} Y_1(y_1, v_1) \\ \hline Y_2(y_2, v_2) \end{bmatrix} = \begin{bmatrix} v_{cm} \\ Un - ge_3 \\ \mathcal{S}(\varpi)n \\ \Pi(n)\tau \\ \hline \mathcal{S}(\omega)r \\ \Pi(r)\tau_m \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_{cm} \\ \dot{v}_{cm} \\ \dot{n} \\ \dot{\varpi} \\ \hline \dot{r} \\ \dot{\omega} \end{bmatrix} = \begin{bmatrix} \dot{y}_1 \\ \hline \dot{y}_2 \end{bmatrix} \right),$$

and where we draw attention to the state and input decompositions in (39) and (40). It is now clear that the UAV-manipulator system can be understood as two decoupled systems, with vectors fields Y_1 and Y_2 (and for which controllers are found in the literature). Let us highlight the cascaded structure of those two vector fields (below, and for brevity, $0 \equiv 0_{3\times 3}$ and $I \equiv I_3$), which is better understood if they are rewritten as

$$Y_1(y_1, v_1) = \begin{bmatrix} 0 & I & 0 & 0 \\ 0 & 0 & UI & 0 \\ 0 & 0 & 0 & -\mathcal{S}(n) \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} p_{cm} \\ v_{cm} \\ n \\ \varpi \end{bmatrix} + \begin{bmatrix} 0_3 \\ -ge_3 \\ 0_3 \\ \Pi(n)\tau \end{bmatrix}, \tag{54}$$

$$Y_2(y_2, v_2) = \begin{bmatrix} 0 & -\mathcal{S}(r) \\ 0 & 0 \end{bmatrix} \begin{bmatrix} r \\ \omega \end{bmatrix} + \begin{bmatrix} 0_3 \\ \Pi(r)\tau_m \end{bmatrix}. \tag{55}$$

Recall then Problem 1, and notice that the vector field Y_1 is concerned with the motion of the center-of-mass of the system, while the vector field Y_2 is concerned with the angular position motion of the manipulator. This motivates us to introduce the desired center-of-mass position based on the desired manipulator linear and angular position ($p^\star$ and $r^\star$ in Problem 1), defined as

$$p_{cm}^\star : \mathbb{R} \ni t \mapsto p_{cm}^\star(t) := p^\star(t) + l\frac{M}{m+M} r^\star(t) \in \mathbb{R}^3 \tag{56}$$

and which, for reasons that are made clear next, we assume satisfies

$$p_{cm}^\star \in \mathcal{C}^4, \tag{57a}$$
$$\sup_{t\in\mathbb{R}} \|p_{cm}^{\star(2)}(t) + ge_3\| > 0, \tag{57b}$$
$$\sup_{t\in\mathbb{R}} p_{cm}^{\star(i)}(t) < \infty \text{ for } i \in \{2, 3, 4\}. \tag{57c}$$

C. Differential flatness

It follows from the cascaded structure of the vector field Y_1 in (54) that it is differentially flat with respect to the position of the center-of-mass. In fact, if we require $t \mapsto p_{cm}(t) := p_{cm}^\star(t)$, then we find two equilibria trajectories and two equilibria inputs, namely (below $y_{1\pm}^\star$ stands for the two solutions, namely $y_{1+}^\star$ and $y_{1-}^\star$)

$$y_{1\pm}^\star : \mathbb{R} \ni t \mapsto y_{1\pm}^\star(t) \in \mathbb{Y}_1 \tag{58}$$
$$y_{1\pm}^\star(t) := \begin{bmatrix} p_{cm}^\star(t) \\ v_{cm}^\star(t) \\ n_\pm^\star(t) \\ \varpi^\star(t) \end{bmatrix} := \begin{bmatrix} p_{cm}^{\star(0)}(t) \\ p_{cm}^{\star(1)}(t) \\ \pm\frac{p_{cm}^{\star(2)}(t) + ge_3}{\|p_{cm}^{\star(2)}(t) + ge_3\|} \\ \mathcal{S}\left(\frac{p_{cm}^{\star(2)}(t) + ge_3}{\|p_{cm}^{\star(2)}(t) + ge_3\|}\right) \frac{p_{cm}^{\star(3)}(t)}{\|p_{cm}^{\star(2)}(t) + ge_3\|} \end{bmatrix},$$

and (below we use the definition of $\varpi^\star$ presented above)

$$v_{1\pm}^\star : \mathbb{R} \ni t \mapsto v_{1\pm}^\star(t) := \begin{bmatrix} \pm\|p_{cm}^{\star(2)}(t) + ge_3\| \\ \varpi^{\star(1)}(t) \end{bmatrix} \in \mathbb{R}^4. \tag{59}$$

At this point, it becomes clear the necessity of the constraints (57a), (57b). Since (58) and (59) require $p_{cm}^\star$ to be four times continuously differentiable, (57a) becomes clear. On the other hand, $y_{1\pm}^\star$ is only well defined if (57b) is satisfied. The existence of two equilibria is also clear intuitively: loosely speaking, for $y_{1+}^\star$ the UAV points upwards and its thrust is positive, and for $y_{1-}^\star$ the UAV points downwards and its thrust

is negative (regardless, in both cases, the center-of-mass tracks the desired trajectory).

Similarly, it follows from the cascaded structure of the vector field Y_2 in (55) that it is differentially flat with respect to the angular position of the manipulator. In fact, if we require $t \mapsto r(t) := r^\star(t)$, then we find two equilibria trajectories and one equilibria input, namely

$$\begin{aligned} &\mathbb{R} \ni t \mapsto y^\star_{2\pm}(t) \in \mathbb{Y}_2 \\ &y^\star_{2\pm}(t) := \begin{bmatrix} r^\star_\pm(t) \\ \omega^\star(t) \end{bmatrix} := \begin{bmatrix} \pm r^{\star(0)}(t) \\ \mathcal{S}\left(r^{\star(0)}(t)\right) r^{\star(1)}(t) \end{bmatrix} \end{aligned} \tag{60}$$

and

$$\mathbb{R} \ni t \mapsto v_2^\star(t) := \mathcal{S}\left(r^{\star(0)}(t)\right) r^{\star(2)}(t) \in \mathbb{R}^3. \tag{61}$$

The equilibria $\{y^\star_{1+}, y^\star_{1-}\}$ and $\{y^\star_{2+}, y^\star_{2-}\}$ just defined allows us to compute four equilibria in the original system apart from rotations around the rigid bodies third axis; i.e., it allows us to compute four equilibria equivalence classes (which we denote by $\{[x]^\star_{+,+}, [x]^\star_{-,+}, [x]^\star_{+,-}, [x]^\star_{-,-}\}$), with the help of the map h in (42), namely

$$\begin{aligned} [x]^\star_{\pm,\mp} : \mathbb{R} \ni t \mapsto [x]^\star(t) := h(y^\star_{\pm,\mp}) \in \mathbb{X}\backslash\sim, \\ y^\star_{\pm,\mp} : \mathbb{R} \ni t \mapsto y^\star_{\pm,\mp}(t) := (y^\star_{1\pm}(t), y^\star_{2\mp}(t)) \in \mathbb{Y}. \end{aligned} \tag{62}$$

(Note that the first slot $\pm$ in $y^\star_{\pm,(\cdot)}$ is related to the equilibria $y^\star_{1\pm}$; and the second slot $\pm$ in $y^\star_{(\cdot),\pm}$ is related to the equilibria $y^\star_{2\pm}$.)

Assumption 12: With the equilibria $y^\star_{1,\pm}$ in (58) and $y^\star_{2,\pm}$ in (60) in mind, assume there exist controllers

$$v_1^{cl} : \mathbb{R} \times \mathbb{Y}_1 \to \mathbb{R}^4, \tag{63}$$

$$v_2^{cl} : \mathbb{R} \times \mathbb{Y}_2 \to \mathbb{R}^3, \tag{64}$$

that guarantee that along solutions of

$$\begin{aligned} \dot{y}_1(t) &= Y_1(y_1(t), v_1^{cl}(t, y_1(t))), y_1(0) \in \mathbb{Y}_1, \\ \dot{y}_2(t) &= Y_2(y_2(t), v_2^{cl}(t, y_2(t))), y_2(0) \in \mathbb{Y}_2, \end{aligned}$$

it holds that

$$\begin{aligned} &\lim_{t\to\infty}(y_1(t) - y^\star_{1+}(t)) = 0_{12} \text{ or } \lim_{t\to\infty}(y_1(t) - y^\star_{1-}(t)) = 0_{12}, \\ &\lim_{t\to\infty}(y_2(t) - y^\star_{2+}(t)) = 0_6 \text{ or } \lim_{t\to\infty}(y_2(t) - y^\star_{2-}(t)) = 0_6, \end{aligned}$$

which in turn implies that

$$\begin{aligned} &\lim_{t\to\infty}(p_{cm}(t) - p^\star_{cm}(t)) = 0_3, \\ &\lim_{t\to\infty}(r(t) \pm r^\star(t)) = 0_3. \end{aligned}$$

Assume as well that there exist non-empty $\mathbb{Y}_{1,0} \subset \mathbb{Y}_1$ and $\mathbb{Y}_{2,0} \subset \mathbb{Y}_2$ such that along solutions of

$$\begin{aligned} \dot{y}_1(t) &= Y_1(y_1(t), v_1^{cl}(t, y_1(t))), y_1(0) \in \mathbb{Y}_{1,0}, \\ \dot{y}_2(t) &= Y_2(y_2(t), v_2^{cl}(t, y_2(t))), y_2(0) \in \mathbb{Y}_{2,0}, \end{aligned}$$

it holds that $y^\star_{1+}$ and $y^\star_{2+}$ are stable, and that

$$\lim_{t\to\infty}(y_1(t) - y^\star_{1+}(t)) = 0_{12}, \tag{65a}$$

$$\lim_{t\to\infty}(y_2(t) - y^\star_{2+}(t)) = 0_6. \tag{65b}$$

Finally, assume that these controllers render both equilibria trajectories $y^\star_{1-}$ and $y^\star_{2-}$ unstable.

Note that the vector field Y_1 in (53) is that of a VTOL vehicle, and thus we are able to leverage controllers from the literature to control the position of the center-of-mass [39]–[42]. Similarly, the vector field Y_2 in (53) is that of a second order system in the unit sphere, and thus we are able to leverage controllers from the literature to control the attitude of the manipulator [43]–[46]. As such, Assumption 12 is indeed satisfied (provided that the desired center-of-mass position trajectory satisfies the constraints in (57)), and we can leverage controllers that are found in the literature and which satisfy the requirements described in Assumption 12.

Remark 13: Define $g : \mathbb{X} \ni x \mapsto g := e_3^T\Omega \in \mathbb{R}$, and notice that

$$Dg(x)X(x, \bar{u}^{cl}(x, v, \tau_\psi)) = \tau_\psi$$

for any $(x, u, \tau_\psi) \in \mathbb{X} \times \mathbb{U} \times \mathbb{R}$. As such, given some positive gain $k > 0$, consider the control law

$$\begin{aligned} &\tau_\psi^{cl} : \mathbb{R} \times \mathbb{X} \ni (t, x) \mapsto \tau_\psi^{cl}(t, x) \in \mathbb{R} \\ &\tau_\psi^{cl}(t, x) := -kg(x) = -ke_3^T\Omega. \end{aligned}$$

It then follows that, $t \mapsto e_3^T\dot{\Omega}(t) = -ke_3^T\Omega(t)$, which implies that $t \mapsto e_3^T\Omega(t) = e_3^T\Omega(0)e^{-kt}$ along solutions of the closed loop system. For simplicity, we shall assume that the chosen control law for the UAV yaw motion is always the one shown above, which guarantees only that the UAV does not asymptotically spin around its third axis (loosely speaking, it does not yaw asymptotically). More complex control laws may be chosen for τ^{cl}, but none (including the one above) interferes with the motion of the equivalence class (i.e., with $t \mapsto [x(t)]$).

Let us then present the result that guarantees that Problem 1 is indeed solved. For that purpose, let Assumption 12 be satisfied, and construct the control law in the original system as a composition of the input transformation $\bar{u}^{cl}$ in (51) with the control laws presented in Assumption 12, i.e., (see Remark 13 for τ_ψ^{cl})

$$\begin{aligned} &u^{cl} : \mathbb{R} \times \mathbb{X} \ni (t, x) \mapsto u^{cl}(t, x) \in \mathbb{R}^7 \\ &u^{cl}(t, x) = \bar{u}^{cl}(x, v^{cl}(t, y), \tau_\psi^{cl}(t, x))|_{y=f([x])} \\ &v^{cl}(t, y) := (v_1^{cl}(t, y_1), v_2^{cl}(t, y_2)), \end{aligned} \tag{66}$$

which leads to the closed loop vector field

$$\mathbb{R} \times \mathbb{X} \ni (t, x) \mapsto X^{cl}(t, x) := X(x, u^{cl}(t, x)) \in T_x\mathbb{X}. \tag{67}$$

Theorem 14: Consider the closed loop vector field X^{cl} in (67), constructed under the assumption the controllers in Assumption 12 are provided. Then, along solutions of

$$\dot{x}(t) = X^{cl}(t, x(t)), [x(0)] \in h(\mathbb{Y}_{1,0} \times \mathbb{Y}_{2,0}), \tag{68}$$

where $h(\mathbb{Y}_{1,0} \times \mathbb{Y}_{2,0}) := \{h(y) \in \mathbb{X}\backslash\sim : y \in \mathbb{Y}_{1,0} \times \mathbb{Y}_{2,0}\}$, it follows that

$$\lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}([x(t)], [x]^\star_{+,+}(t)) = 0, \tag{69}$$

where the equilibrium equivalence class $[x]^\star_{+,+}$ is defined in (62) and the metric $d_{\mathbb{X}\backslash\sim}$ is defined in (24); i.e., it follows that Problem 1 is satisfied, where rotations of the rigid bodies around their third axes are ignored.

Proof 3: Since $[x(0)] \in h(\mathbb{Y}_{1,0} \times \mathbb{Y}_{2,0})$ it follows that $(y_1(0), y_2(0)) =: y(0) := f([x(0)]) \in \mathbb{Y}_{1,0} \times \mathbb{Y}_{2,0}$; moreover, if we apply the control law (66), then the conclusion (65) in

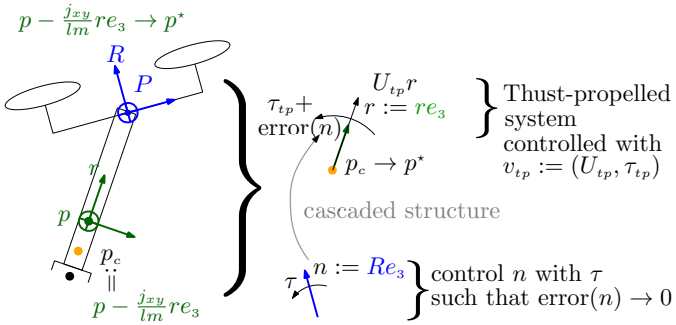


Fig. 7: By means of appropriate state and input transformations, the UAV-slung-manipulator may be understood as two systems in cascade: one thrust-propelled system followed by one unit-vector double-integrator system.

Assumption 12 holds.

Because $t \mapsto y(t) = f([x(t)]) \Rightarrow t \mapsto [x(t)] = h(y(t))$, it follows that

$$\begin{aligned}
&\lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}([x(t)], [x]^\star_{+,+}(t)) = \\
&= \lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}(h(y(t)), h(y^\star_{+,+}(t))) \\
&= \lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}(h(y^\star_{+,+}(t) - y^\star_{+,+}(t) + y(t)), h(y^\star_{+,+}(t))) \\
&\overset{h, d_{\mathbb{X}\backslash\sim} \in C^0}{=} \lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}(h(y^\star_{+,+}(t) + \lim_{t\to\infty}(y(t) - y^\star_{+,+}(t))), h(y^\star_{+,+}(t))) \\
&\overset{(65)}{=} \lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}(h(y^\star_{+,+}(t)), h(y^\star_{+,+}(t))) = 0,
\end{aligned}$$

and where the limits above were derived by invoking continuity of the maps h and $d_{\mathbb{X}\backslash\sim}$.

Theorem 15: Consider the closed loop vector field X^{cl} in (67), constructed assuming that the controllers in Assumption 12 are provided. Then, along solutions of

$$\dot{x}(t) = X^{cl}(t, x(t)), x(0) \in \mathbb{X}, \tag{70}$$

it follows that, for some equilibrium equivalence class $[x]^\star_\star \in \{[x]^\star_{+,+}, [x]^\star_{-,+}, [x]^\star_{+,-}, [x]^\star_{-,-}\}$,

$$\lim_{t\to\infty} d_{\mathbb{X}\backslash\sim}([x(t)], [x]^\star_\star(t)) = 0, \tag{71}$$

where the equilibria equivalence classes are defined in (62) and the metric $d_{\mathbb{X}\backslash\sim}$ is defined in (24); i.e., it follows that the position tracking problem in Problem 1 is satisfied while the orientation tracking is not necessarily satisfied. Moreover, the equilibria equivalence classes $[x]^\star_{-,+}$, $[x]^\star_{+,-}$ and $[x]^\star_{-,-}$ are all unstable.

Proof 4: The first part of the Theorem follows the same steps as in Theorem 14, and its proof is therefore omitted. For the instability of $[x]^\star_{-,+}$, $[x]^\star_{+,-}$ and $[x]^\star_{-,-}$ it suffices to invoke continuity of the map h.

VI. UAV-SLUNG-MANIPULATOR

We shall verify next that the UAV-slung-manipulator system behaves as a thrust-propelled system followed/cascaded after a unit vector double integrator, as illustrated in Fig. 7. The three steps we follow next have been described in the Section IV.

A. State transformation

Recall the discussion from Section IV. Let us then introduce the set $\mathbb{Y}$ defined as (you may read the subindex tp as thrust-propelled)

$$\mathbb{Y}_{tp} := \left\{(p_c, v_c, r, \omega) \in (\mathbb{R}^3)^4 : r^T r = 1, r^T \omega = 0\right\}, \tag{72a}$$

$$\mathbb{Y} := \left\{(y_{tp}, n, \varpi) \in \mathbb{Y}_{tp} \times (\mathbb{R}^3)^2 : n^T n = 1, n^T \varpi = 0\right\}, \tag{72b}$$

with the inherited metrics $d_{\mathbb{Y}} := d_{\mathbb{R}^{18}}$ and $d_{\mathbb{Y}_{tp}} := d_{\mathbb{R}^{12}}$ (the set $\mathbb{Y}_{tp}$ will be used for the part of the system that resembles a thrust-propelled system – see Fig. 7). Hereafter, and for convenience, we decompose elements in the sets above as

$$y_{tp} \in \mathbb{Y}_{tp} :\Leftrightarrow (p_c, v_c, r, \omega) \in \mathbb{Y}_{tp}, \tag{73a}$$

$$y \in \mathbb{Y} :\Leftrightarrow (y_{tp}, n, \varpi) \in \mathbb{Y}. \tag{73b}$$

Hereafter, we denote as well

$$v_{tp} \in \mathbb{R}^4 :\Leftrightarrow (U_{tp}, \tau_{tp}) \in \mathbb{R} \times \mathbb{R}^3, \tag{74a}$$

$$v \in \mathbb{R}^4 :\Leftrightarrow (U_{tp}, \tau) \in \mathbb{R} \times \mathbb{R}^3, \tag{74b}$$

where v_{tp} will be an auxiliary input (to the thrust propelled system), and v will be the actual input to the whole system.

Assumption 16: We must at this point make the assumption that $e_3^T \omega(0) = 0$, i.e., that the manipulator is not initially spinning around itself (an alternative assumption would be to take $j_{zz} = 0$). Based on Proposition 2, it follows that $e_3^T \omega(t) = 0$ for all $t \in \mathbb{R}$, i.e., that the manipulator remains without spinning around itself, and thus, hereafter, and w.l.o.g., we restrict a state of the UAV-slung-manipulator system to be in the set $\{x \in \mathbb{X} : e_3^T \omega = 0\}$.

Recall the definition of the state space $\mathbb{X}$ in (38a), and the state decomposition in (8). Also, with Fig. 7 in mind, consider then the mapping

$$f : \mathbb{X}\backslash\sim \ni [x] \mapsto f([x]) \in \mathbb{Y} \tag{75}$$

$$f([x]) := \begin{bmatrix} p - \frac{j_{xy}}{lm} re_3 \\ v - \frac{j_{xy}}{lm} r\mathcal{S}(\omega) e_3 \\ re_3 \\ \Pi(re_3) r\omega \\ \hline \bar{R}e_3 \\ \Pi(Re_3) R\Omega \end{bmatrix} \left(= \begin{bmatrix} p_c \\ v_c \\ r \\ \omega \\ \hline \bar{n} \\ \varpi \end{bmatrix} = \begin{bmatrix} y_{tp} \\ \\ \hline \bar{n} \\ \varpi \end{bmatrix} = y \right),$$

and consider as well the mapping

$$h : \mathbb{Y} \ni y \mapsto h(y) \in \mathbb{X}\backslash\sim \tag{76}$$

$$h(y) := \underbrace{\begin{bmatrix} p_c + \frac{j_{xy}}{lm} r \\ \tilde{r} \\ p_c + (\frac{j_{xy}}{lm} + l) r \\ \tilde{R} \\ \hline v_c + \frac{j_{xy}}{lm} \mathcal{S}(\omega) r \\ \tilde{r}^T \omega \\ v_c + (\frac{j_{xy}}{lm} + l) \mathcal{S}(\omega) r \\ \tilde{R}^T \varpi \end{bmatrix}}_{\substack{\tilde{r} \in \mathbb{SO}(3), \tilde{r}e_3 = r \\ \tilde{R} \in \mathbb{SO}(3), \tilde{R}e_3 = n}} \left(= \begin{bmatrix} p \\ r \\ P \\ R \\ \hline v \\ \omega \\ V \\ \Omega \end{bmatrix} = [x] \right),$$

where $[\cdot]$ above denotes the equivalence class as defined immediately after (23).

Proposition 17: The map f in (75) is well defined, and

$$f \circ h = \mathrm{id}_{\mathbb{Y}} \text{ and } h \circ f = \mathrm{id}_{\mathbb{X}\backslash\sim}, \tag{77}$$

i.e., $h = f^{-1}$ and $f = h^{-1}$.

The proof follows the same steps as those in Proposition 11.

B. Input transformation

Let us construct the input transformation in two steps. For that purpose, consider the transformed dynamics with the original input (not the transformed input), which are given by

$$Df([x])X(x,u) = \tag{78}$$
$$= \begin{bmatrix} v - \frac{j_{xy}}{lm} r\mathcal{S}(\omega) e_3 \\ re_3(\star_1) - ge_3 \\ \mathcal{S}(\Pi(re_3) r\omega) re_3 \\ \frac{l}{j_{xy}} \mathcal{S}(re_3) rT(x,u) \\ \bar{\mathcal{S}}(\bar{\Pi}(\bar{R}\bar{e}_3)\bar{R}\bar{\Omega})\bar{R}\bar{e}_3 \\ \Pi(Re_3) R(\star_2) \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \bar{\ddot{n}} \\ \dot{\varpi} \end{bmatrix} = \begin{bmatrix} \dot{y}_{tp} \\ \\ \\ \bar{\dot{n}} \\ \end{bmatrix} = \dot{y} \right),$$
$$(\star_1) = \frac{1}{m}(re_3)^T T(x,u) + \frac{j_{xy}}{lm}\omega^T \Pi(e_3)\omega,$$
$$(\star_2) = J^{-1}(\tau - \mathcal{S}(\Omega) J\Omega) - e_3^T \mathcal{S}(\Omega) e_3.$$

Equation (78) sheds some light into how to design the controller. Suppose for now that we had control over the internal tensions – $T(x,u)$ (which we do, by intermediate of the input u; however we do not have control over all of the three components of $T(x,u)$). Then, one can choose the component of those tensions aligned with the manipulator axis of axial symmetry (i.e.,$(re_3)^T T(x,u)$) in a way that behaves as a thrust in the differential equation $\dot{v}_c$ in (78); one can also choose the components of those tensions orthogonal to the manipulator's axis of axial symmetry (i.e., $\Pi(re_3)T(x,u)$) in a way that behaves as a torque on the manipulator axis of symmetry – see the differential equation $\dot{\omega}$ in (78). With the above in mind, and $\dot{v}_c$ and $\dot{\omega}$ in (78) also in mind, we define

$$T^{cl} : \mathbb{Y}_{tp} \times \mathbb{R}^4 \ni (y_{tp}, v_{tp}) \mapsto T^{cl}(y_{tp}, v_{tp}) \in \mathbb{R}^3,$$
$$T^{cl}(y_{tp}, v_{tp}) := \frac{1}{m}\left(U_{tp} - \frac{j_{xy}}{lm}\omega^T\omega\right) r - \frac{j_{xy}}{l}\mathcal{S}(r)\tau_{tp} \tag{79}$$

It then follows that

$$Df([x])X(x,u)|_{T(x,u)=T^{cl}(y_{tp},v_{tp})}|_{x\in h(y)} = \tag{80}$$
$$= \begin{bmatrix} v_c \\ U_{tp} r - ge_3 \\ \mathcal{S}(\omega) r \\ \mathcal{S}(r)\tau_{tp} \\ \star \\ \star \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \dot{n} \\ \dot{\varpi} \end{bmatrix} \right).$$

(The equation above is strictly speaking ill-defined, specifically the last element $\dot{\varpi}$ depends on the choice of representative $x \in h(y) \in \mathbb{X}\backslash_\sim$; the correct formulation is found later in (92)).

However, we do not have full control over the internal tensions, which is simple to verify by inspecting of T in (20): the tensions are three dimensional ($T(x,u) \in \mathbb{R}^3$), and the only input that acts on the tensions is the one-dimensional thrust ($U \in \mathbb{R}$). Moreover, by inspection of T in (20), notice the affine dependence on URe_3, which motivates as to introduce an extra dummy variable U_{3d} and such that $URe_3 = U_{3d} + (URe_3 - U_{3d})$; with this in mind, notice that $T(x,u)$ in (20) is equivalently written as (recall that in the UAV-slung-manipulator system $\tau_m = 0_3$ and $e_3^T\omega = 0$ – see Assumption 16)

$$T(x,u) = \frac{rA_T^{-1}r^T}{M}\left(U_{3d} + lM\omega^T\omega re_3 + (URe_3 - U_{3d})\right), \tag{81}$$

$$A_T \equiv \frac{m+M}{mM} I_3 + \frac{l^2}{j_{xy}}\Pi(e_3), \tag{82}$$

Loosely speaking, since we wish that $T(x,u) = T^{cl}(y_{tp}, v_{tp})$, it naturally leads to the definition of

$$U_{3d}^{cl} : \mathbb{Y}_{tp} \times \mathbb{R}^4 \ni (y_{tp}, v_{tp}) \mapsto U_{3d}^{cl}(y_{tp}, v_{tp}) \in \mathbb{R}^3 \tag{83}$$
$$U_{3d}^{cl}(y_{tp}, v_{tp}) := M\left(rA_T r^T T^{cl}(y_{tp}, v_{tp}) - l\omega^T\omega r\right)$$
$$\overset{(79)}{=} lM\gamma\left(\left(\frac{m+M}{M}\frac{U_{tp}}{l\gamma} - \omega^T\omega\right) r - \mathcal{S}(r)\tau_{tp}\right)$$

and, as such, if we take (81) with $U_{3d} = U_{3d}^{cl}(y_{tp}, v_{tp})$, it follows that the tensions $T(x,u)$ in (20) are equivalently written as

$$T(x,u) = T^{cl}(y_{tp}, v_{tp}) + \frac{rA_T^{-1}r^T}{M}(URe_3 - U_{3d}^{cl}(y_{tp}, v_{tp})), \tag{84}$$

for any $x \in \mathbb{X}$ and any $v_{tp} \in \mathbb{R}^4$, and where $(y_{tp}, \star) = f([x])$ (see (75)). After simple computations, it then follows that for any $v_{tp} = (U_{tp}, \tau_{tp}) \in \mathbb{R}^4$,

$$Df([x])X(x,u)|_{T(x,u)\text{as in (84)}}|_{x\in h(y)} = \tag{85}$$
$$= (80) + \begin{bmatrix} 0_{3\times 3} \\ \frac{1}{m+M} rr^T \\ 0_{3\times 3} \\ \frac{1}{lM\gamma}\mathcal{S}(r) \\ 0_{3\times 3} \\ 0_{3\times 3} \end{bmatrix} (Un - U_{3d}^{cl}(y_{tp}, v_{tp})) \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \dot{n} \\ \dot{\varpi} \end{bmatrix} \right).$$

A common and natural choice for the input U in (85) is one that minimizes the error $Un - U_{3d}^{cl}(y_{tp}, u_{tp})$ in (85), i.e.,

$$\{n^T U_{3d}^{cl}(y_{tp}, v_{tp})\} = \arg\min_{U\in\mathbb{R}} \|Un - U_{3d}^{cl}(y_{tp}, v_{tp})\|. \tag{86}$$

This choice, despite natural, is not ideal, since, if this choice is made, it results in

$$\dot{v}_c = U_{tp} r - ge_3 - \frac{1}{m+M} rr^T \Pi(n) U_{3d}^{cl}(y_{tp}, v_{tp}), \tag{87}$$

which prevents us from using control laws in the literature which rely on the cascaded structure of the thrust-propelled system.

The other option (it comes at a cost we explain next) is one where U is chosen such that the acceleration $\dot{v}_c$ in (85) depends solely on the thrust U_{tp}, i.e., such that $r^T(Un - U_{3d}^{cl}(y_{tp}, v_{tp}))$ vanishes. That leads to the definition of the control law

$$U^{cl} : \tilde{\mathbb{X}} \times \mathbb{U} \ni (x, U_{tp}) \mapsto U^{cl}(x, U_{tp}) \in \mathbb{R} \tag{88}$$
$$U^{cl}(x, U_{tp}) := \frac{r^T U_{3d}^{cl}(y_{tp}, v_{tp})}{r^T n}\Big|_{y=f([x])}$$
$$\overset{(83)}{=} \frac{(m+M)U_{tp} - lM\gamma\omega^T\omega}{(re_3)^T(Re_3)},$$

where

$$\tilde{\mathbb{X}} := \{x \in \mathbb{X} : (re_3)^T(Re_3) \neq 0\}. \tag{89}$$

The cost of this choice is now clear: the control law is not

defined in the whole state space; in fact, it is ill-defined when the manipulator arm is orthogonal to the thrust axis of the UAV. Notice, however, that, since the UAV and manipulator have physical volume, the manipulator should therefore never be orthogonal to the thrust axis of the UAV, since in this case the two rigid bodies would collide.

Again, after simple computations, it follows that for any $v_{tp} = (U_{tp}, \tau_{tp}) \in \mathbb{R}^4$,

$$Df([x])X(x,u)|_{U=U^{cl}(x,U_{tp})}|_{x\in h(y)} = \tag{90}$$
$$= (80) + \tfrac{1}{lM\gamma}\tfrac{\Pi(r)}{r^T n}\begin{bmatrix} 0_{3\times 3} \\ 0_{3\times 3} \\ 0_{3\times 3} \\ I_{3\times 3} \\ 0_{3\times 3} \\ 0_{3\times 3} \end{bmatrix} \mathcal{S}(n)\, U_{3d}^{cl}(y_{tp}, v_{tp}) \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \dot{n} \\ \dot{\varpi} \end{bmatrix} \right).$$

At this point, we can now introduce the input transformation. Consider then (see the input decomposition in (74))

$$\bar{u}^{cl} : \tilde{\mathbb{X}} \times \mathbb{U} \times \mathbb{R} \ni (x, v, \tau_\psi) \mapsto \bar{u}^{cl}(x, v, \tau_\psi) \in \mathbb{U} \tag{91}$$
$$\bar{u}^{cl}(x, v, \tau_\psi) := \begin{bmatrix} U^{cl}(x, U_{tp}) \\ J(\Pi(e_3) R\tau + \tau_\psi e_3 + e_3^T \Omega \mathcal{S}(\Omega) e_3) - \mathcal{S}(\Omega) J\Omega \end{bmatrix}.$$

Given the input transformation $\bar{u}^{cl}$ in (91) and the state transformation f in (75) one can then compute the transformed dynamics. In fact, it follows that

$$Y : \mathbb{Y} \times \mathbb{U} \ni (y, v) \mapsto Y(y, v) \in T_y\mathbb{Y} \tag{92}$$
$$Y(y,v) := Df([x])X(x, \bar{u}^{cl}(x,v,\cdot))|_{x\in h(y)}$$
$$= \begin{bmatrix} v_c \\ U_{tp} r - g e_3 \\ \mathcal{S}(\omega) r \\ \frac{1}{r^T n}\left(\frac{(m+M)}{lM\gamma} U_{tp} - \omega^T\omega\right) \mathcal{S}(r)\, n \\ \mathcal{S}(\varpi)\, n \\ \Pi(n)\, \tau \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \dot{n} \\ \dot{\varpi} \end{bmatrix} = \dot{y} \right)$$

which, for any $\tau_{tp} \in \mathbb{R}^3$, is exactly equivalent to

$$Y(y,v) = \begin{bmatrix} v_c \\ U_{tp} r - g e_3 \\ \mathcal{S}(\omega) r \\ \Pi(r)\, \tau_{tp} + \frac{\mathcal{S}(r)\mathcal{S}(n) U_{3d}^{cl}(y_{tp}, v_{tp})}{\gamma l M r^T n} \\ \mathcal{S}(\varpi)\, n \\ \Pi(n)\, \tau \end{bmatrix} \left(= \begin{bmatrix} \dot{p}_c \\ \dot{v}_c \\ \dot{r} \\ \dot{\omega} \\ \dot{n} \\ \dot{\varpi} \end{bmatrix} \right). \tag{93}$$

The cascaded structure of the this vector fields is now clear, since Y can be equivalently written as (below, and for brevity, $0 \equiv 0_{3\times 3}$ and $I \equiv I_3$)

$$Y(y,v) = \begin{bmatrix} 0 & I & 0 & 0 & 0 & 0 \\ 0 & 0 & U_{tp} I & 0 & 0 & 0 \\ 0 & 0 & 0 & -\mathcal{S}(r) & 0 & 0 \\ 0 & 0 & 0 & 0 & (\star)\mathcal{S}(r) & 0 \\ 0 & 0 & 0 & 0 & 0 & -\mathcal{S}(n) \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} p_c \\ v_c \\ r \\ \omega \\ n \\ \varpi \end{bmatrix} + \begin{bmatrix} 0_3 \\ -g e_3 \\ 0_3 \\ 0_3 \\ 0_3 \\ \Pi(n)\, \tau \end{bmatrix},$$
$$(\star) = \tfrac{1}{r^T n}\left(\tfrac{(m+M)}{\gamma} U_{tp} - \omega^T\omega\right).$$

Recall then Problem 2. For reasons that will be clear next, we require the desired position trajectory to satisfy some constraints, specifically

$$p^\star \in \mathcal{C}^6, \text{ and } \sup_{t\in\mathbb{R}} p^{\star(i)}(t) < \infty \text{ for } i \in \{2,3,4,5,6\}, \tag{94a}$$
$$\inf_{t\in\mathbb{R}} \|p^{\star(2)}(t) + g e_3\| > 0, \text{ and} \tag{94b}$$
$$\inf_{t\in\mathbb{R}} \left\{ \begin{array}{c} \|\Pi(r^\star(t))\ddot{r}^\star(t)\|^2 + \\ \left(\|\mathcal{S}(r^\star(t))\dot{r}^\star(t)\|^2 \pm \frac{1}{\gamma}\frac{(m+M)}{M}\frac{\|\ddot{p}^\star(t) + g e_3\|}{l}\right)^2 \end{array} \right\} > 0, \tag{94c}$$

where $r^\star(t) \equiv \frac{p^{\star(2)}(t) + g e_3}{\|p^{\star(2)}(t) + g e_3\|}$. We say that a trajectory $p^\star$ is *non-aggressive* if it satisfies the constraints in (94). Loosely speaking, non-aggressive trajectories are those whose high-order derivatives (in particular, $p^{\star(2)}, p^{\star(3)}, p^{\star(4)}$) are *small* in magnitude (where *small* is quantified in (94)). In particular, if we take those high order derivatives to be zero (constant speed trajectory), the constraints (94b) and (94c) read as $\sup_{t\in\mathbb{R}} \|g e_3\| > 0$ and as $\sup_{t\in\mathbb{R}} \left(\pm\frac{1}{\gamma}\frac{(m+M)}{M}\frac{\|g e_3\|}{l}\right)^2 > 0$, which are indeed satisfied.

C. Differential flatness

It follows from the cascaded structure of the vector field Y in (92) that it is differentially flat with respect to the position p_c; equivalently, it follows that the vector field X in (29) (if we ignore rotations of the rigid bodies around their third axes) is differentially flat with respect to the position $p - \frac{j_{xy}}{l^2 m} r e_3$ (this justifies the claim in Remark 8). Indeed, if we require $t \mapsto p_c(t) := p^\star(t)$, then we find *four* equilibria state trajectories and *two* equilibria input trajectories, namely (below $y^\star_{\pm,\mp}$ stands for four solutions, namely $y^\star_{+,+}$, $y^\star_{-,+}$, $y^\star_{+,-}$ and $y^\star_{-,-}$)

$$y^\star_{\pm,\mp} : \mathbb{R} \ni t \mapsto y^\star_{\pm,\mp}(t) \in \mathbb{Y} \tag{95}$$
$$y^\star_{\pm,\mp}(t) := \begin{bmatrix} y^\star_{tp,\pm} \\ \hdashline n^\star_{\pm,\mp}(t) \\ \varpi^\star(t) \end{bmatrix} := \begin{bmatrix} p^\star_c(t) \\ v^\star_c(t) \\ r^\star_\pm(t) \\ \omega^\star(t) \\ \hdashline n^\star_{\pm,\mp}(t) \\ \varpi^\star(t) \end{bmatrix} := \begin{bmatrix} p^{\star(0)}(t) \\ p^{\star(1)}(t) \\ \pm\frac{p^{\star(2)}(t) + g e_3}{\|p^{\star(2)}(t) + g e_3\|} \\ \mathcal{S}\left(r^\star_\pm(t)\right) r^{\star(1)}_\pm(t) \\ \hdashline \mp\frac{U^{cl}_{3d}(y^\star_{tp,\pm}(t), u^\star_{tp,\pm}(t))}{\|U^{cl}_{3d}(y^\star_{tp,\pm}(t), u^\star_{tp,\pm}(t))\|} \\ \mathcal{S}\left(n^\star_{\pm,\mp}(t)\right) n^{\star(1)}_{\pm,\mp}(t) \end{bmatrix},$$
$$u^\star_{tp,\pm}(t) := \begin{bmatrix} U^\star_{tp}(t) \\ \tau^\star_{tp}(t) \end{bmatrix} := \begin{bmatrix} \pm\|p^{\star(2)}_c(t) + g e_3\| \\ \omega^{\star(1)}(t) \end{bmatrix},$$

and (below we use the definition of $\varpi^\star$ presented above)

$$v^\star_\pm : \mathbb{R} \ni t \mapsto v^\star_\pm(t) := \begin{bmatrix} \pm\|p^{\star(2)}_c(t) + g e_3\| \\ \varpi^{\star(1)}(t) \end{bmatrix} \in \mathbb{R}^4. \tag{96}$$

(Notice that, despite the presence of $r^\star_\pm$ and $n^\star_{\pm,\mp}$ in the definition of $\omega^\star$ and $\varpi^\star$ (respectively), these are independent of the choice of sign).

It should now be clear from (95), specifically from the definition of $r^\star_\pm$ and $n^\star_{\pm,\mp}$, why the desired position trajectory $p^\star$ was required to satisfy the constraints (94b) and (94c) (in particular, (94c) is equivalently expressed as $\sup_{t\in\mathbb{R}} \|U^{cl}_{3d}(y^\star_{tp,\pm}(t), u^\star_{tp,\pm}(t))\| > 0$). Intuitively, (94b) guarantees that the desired attitude of the manipulator third axis is well defined; and (94c) guarantees that the desired attitude of the UAV third axis is well defined.

The equilibria $\{y^\star_{+,+}, y^\star_{-,+}, y^\star_{+,-}, y^\star_{-,-}\}$ just defined allows us to compute four equilibria in the original system apart from rotations around the rigid bodies' third axis; i.e., it allows us to compute four equilibria equivalence classes (which we denote

by $\{[x]^\star_{+,+}, [x]^\star_{-,+}, [x]^\star_{+,-}, [x]^\star_{-,-}\}$), with the help of the map h in (76), namely

$$[x]^\star_{\pm,\mp} : \mathbb{R} \ni t \mapsto [x]^\star_{\pm,\mp}(t) := h(y^\star_{\pm,\mp}(t)) \in \mathbb{X}\backslash_\sim. \tag{97}$$

See Remark 20 in [37] for why one should specify a desired trajectory for the UAV's (or the manipulator's) position: if that is specified, then there exist infinite equilibria equivalence class trajectories (rather than just four), where the manipulator behaves as a pendulum whose base is the UAV position (or the manipulator's) with a pendulum length $l' = \frac{j_{xy}}{lm} + l$ (or $l' = \frac{j_{xy}}{lm}$).

Assumption 18: With the equilibria $y^\star_{\pm,\mp}$ in (95) in mind, we assume there exists a controller

$$v^{cl} : \mathbb{R} \times \mathbb{Y} \to \mathbb{R}^4, \tag{98}$$

and non-empty $\mathbb{Y}_0 \subset \mathbb{Y}$, such that along solutions of

$$\dot{y}(t) = Y(y(t), v^{cl}(t, y(t))), y(0) \in \mathbb{Y}_0 \tag{99}$$

it holds that (i) $t \mapsto y(t)$ does not approach boundary of $\{y \in \mathbb{Y} : r^T n > 0\}$; that (ii) $y^\star_{+,+}$ is stable; and that (iii)

$$\lim_{t\to\infty}(y(t) - y^\star_{+,+}(t)) = 0_{18}. \tag{100}$$

Finally, we assume that the equilibria trajectories $y^\star_{-,+}$, $y^\star_{+,-}$ and $y^\star_{-,-}$ are unstable.

Note that the vector field Y in (93) is that of a VTOL vehicle cascaded after a second order system in the unit sphere. For this vector field, we are able to leverage controllers for VTOL vehicles from the literature [39]–[42], which must be complemented with two backstepping steps (related to the second order system). As such, Assumption 18 is indeed satisfied if the desired position trajectory satisfies the constraints in (94), i.e., a controller can be found that satisfies the conditions in the assumption.

Let us then present the result that guarantees that Problem 2 is accomplished. For that purpose, let Assumption 18 be satisfied, and construct the control law u^{cl} as a composition of input transformation $\bar{u}^{cl}$ in (91) with the control law presented in Assumption 18, i.e., (see Remark 13 for τ^{cl}_ψ)

$$\begin{aligned} &u^{cl} : \mathbb{R} \times \mathbb{X} \ni (t, x) \mapsto u^{cl}(t, x) \in \mathbb{R}^4 \\ &u^{cl}(t, x) = \bar{u}^{cl}(x, v^{cl}(t, y), \tau^{cl}_\psi(t, x))|_{y=f([x])}, \end{aligned} \tag{101}$$

which leads to the closed loop vector field

$$\mathbb{R} \times \mathbb{X} \ni (t, x) \mapsto X^{cl}(t, x) := X(x, u^{cl}(t, x)) \in T_x\mathbb{X}. \tag{102}$$

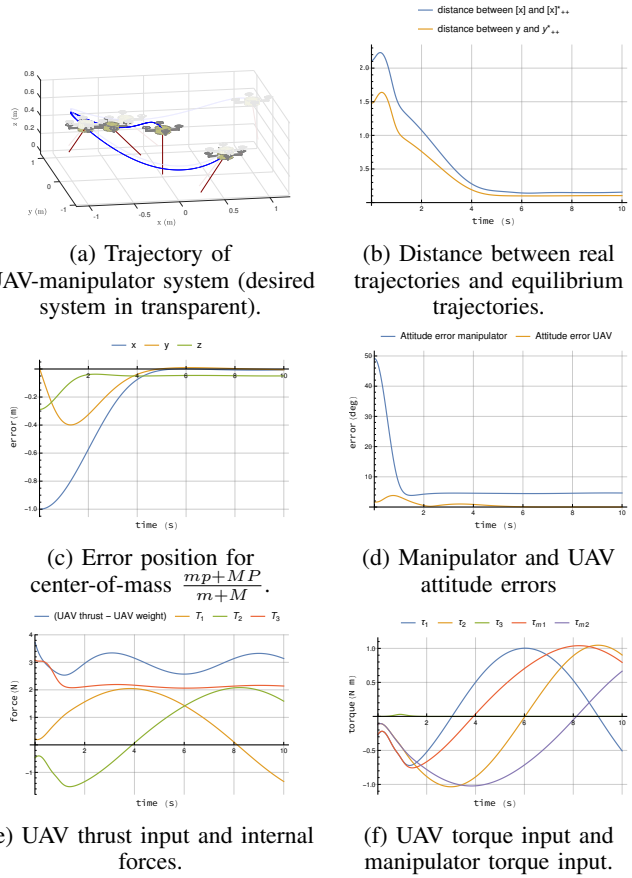
Theorem 19: Consider the closed loop vector field X^{cl} in (102), constructed under the assumption that the controllers in Assumption 18 are provided. Then, along solutions of

$$\dot{x}(t) = X^{cl}(t, x(t)), [x(0)] \in h(\mathbb{Y}_0) \tag{103}$$

where $h(\mathbb{Y}_0) := \{h(y) \in \mathbb{X}\backslash_\sim : y \in \mathbb{Y}_0\}$, it follows that

$$\lim_{t\to\infty} d_{\mathbb{X}\backslash_\sim}([x(t)], [x]^\star_{+,+}(t)) = 0, \tag{104}$$

where the equilibrium equivalence class $[x]^\star_{+,+}$ is defined in (97) and the metric $d_{\mathbb{X}\backslash_\sim}$ is defined in (24). i.e., it follows that Problem 1 is satisfied, where rotations of the rigid bodies around their third axes is ignored. Moreover, the equilibria equivalence classes $[x]^\star_{-,+}$, $[x]^\star_{+,-}$ and $[x]^\star_{-,-}$ are all unstable.



(a) Trajectory of UAV-manipulator system (desired system in transparent).

(b) Distance between real trajectories and equilibrium trajectories.

(c) Error position for center-of-mass $\frac{mp+MP}{m+M}$.

(d) Manipulator and UAV attitude errors

(e) UAV thrust input and internal forces.

(f) UAV torque input and manipulator torque input.

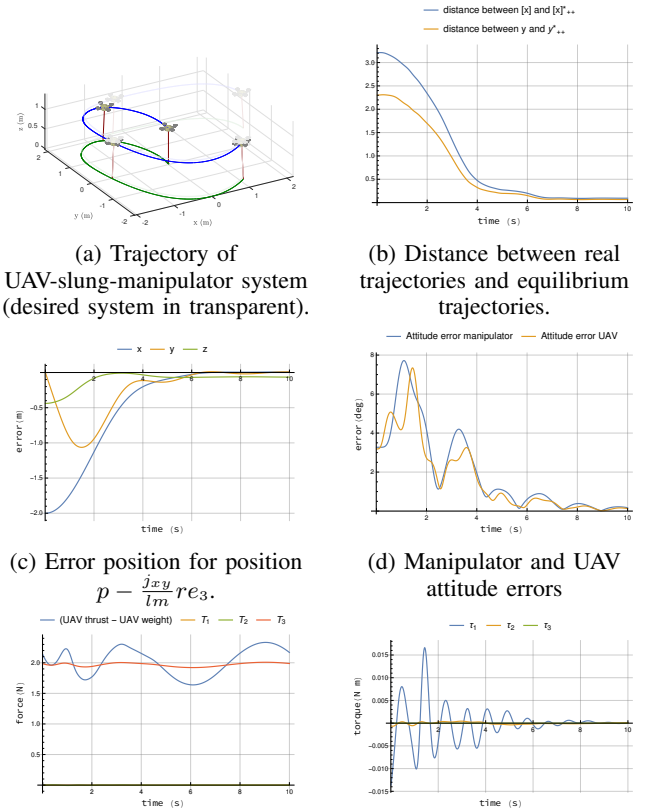
Fig. 8: Simulation for UAV-manipulator system with disturbed parameters (manipulator's mass and moments of inertia): $m_{\text{controller}} = 0.9m_{\text{model}}$ and $j_{\text{controller}} = 0.8j_{\text{model}}$.

Proof 5: Since $[x(0)] \in h(\mathbb{Y}_0)$ it follows that $y(0) := f([x(0)]) \in \mathbb{Y}_0$; moreover, if we apply the control law (101), then the conclusions in Assumption 18 hold. Because $t \mapsto y(t) = f([x(t)]) \Rightarrow t \mapsto [x(t)] = h(y(t))$, it follows that the same logic chain as in the proof of Theorem 14 follows, which completes the proof.

VII. Simulations

Let us provide simulations that illustrate stability, convergence and robustness results. The physical constants were taken as $m = 0.3$ kg, $M = 1.4$ kg, $j = 0.007 \oplus 0.007 \oplus 0.001$ kg m^2 and $J = 0.01 \oplus 0.01 \oplus 0.2$ kg m^2; and with $l = 0.5$ m for the UAV-manipulator system and with $l = 0.9$ m for the UAV-slung-manipulator system. The specific control laws v^{cl}, as presented in Assumptions 12 and 18, are detailed in the mathematica files [37] (including the choice of gains; the chosen controller follows the structure described in [42]). We provide two simulations, in Figs. 8 and 9, one for each system; these are obtained as the solution to $\dot{x}(t) = X^{cl}(t, x(t))$ with $x(0) = (0_3, I_3, le_3, I_3, 0_3, 0_3, 0_3, 0_3) \in \mathbb{X}$ ($x(0)$ in SI units) and where some of the controller parameters have been disturbed (namely, the manipulator's mass and moments of inertia: $m_{\text{controller}} = 0.9m_{\text{model}}$ and $j_{\text{controller}} = 0.8j_{\text{model}}$). For the desired position trajectory, we have chosen (Problems 1 and 2)

$$t \mapsto p^\star(t) := \left(2\cos\left(\tfrac{2\pi}{12}t\right), 2\sin\left(\tfrac{2\pi}{12}t\right), 0.2 + 0.2\cos\left(\tfrac{2\pi}{6}t\right)\right),$$



(a) Trajectory of UAV-slung-manipulator system (desired system in transparent).

(b) Distance between real trajectories and equilibrium trajectories.

(c) Error position for position $p - \frac{j_{xy}}{lm} r e_3$.

(d) Manipulator and UAV attitude errors

(e) UAV thrust input and internal forces.

(f) UAV torque input.

Fig. 9: Simulation for UAV-slung-manipulator system with disturbed parameters (manipulator's mass and moments of inertia): $m_{\text{controller}} = 0.9 m_{\text{model}}$ and $j_{\text{controller}} = 0.8 j_{\text{model}}$.

corresponding to a circle of $2m$ radius in the horizontal plane and with a period of $12s$, superimposed with an oscillatory motion in the vertical direction with an amplitude of $0.2m$ and with a period of $6s$; and for the desired orientation trajectory, we have chosen (Problem 1)

$$t \mapsto r^\star(t) := \frac{2}{\sqrt{7}} \left(-\sin\left(\tfrac{2\pi}{12} t\right), \cos\left(\tfrac{2\pi}{12} t\right), \cos\left(\tfrac{\pi}{6}\right) \right) \in \mathbb{S}^2,$$

corresponding to an orientation trajectory where the desired manipulator is tilted 30° away from the vertical direction, and where the manipulator is tangent (when seen from the top) to the horizontal motion of $p^\star$.

Fig. 8 is related to the UAV-manipulator system, and Fig. 9 is related to the UAV-slung-manipulator system.

In Figs. 8a and 9a, one can visualize the real system pose (pose of both rigid bodies – the UAV's and the manipulator's) in opaque, and the equilibrium system pose in transparent. In Figs. 8b and 9b, one can visualize the distance between the equivalence-class trajectory $t \mapsto [x(t)]$ and the stable equilibrium equivalence-class trajectory $t \mapsto [x(t)]^\star_{+,+}$ as defined in (62) and (97), and by means of the metric $d_{\mathbb{X}\backslash\sim}$ defined in (24); as well as the distance between the transformed trajectory $t \mapsto y(t) := f(x(t))$ and the stable equilibrium transformed trajectory $t \mapsto y^\star_{+,+}(t)$ as defined in (62) and (95) (and by means of the metric $d_{\mathbb{Y}} := d_{\mathbb{R}^{18}}$). Notice that these distances may increase, during some time intervals, but they converge to zero asymptotically in the absence of parameters' mismatch.

In Figs. 8c and 9c, the position error is shown for the positions of the thrust propelled systems, i.e., the center-of-mass position error ($t \mapsto \frac{mp(t)+MP(t)}{m+M} - p^\star_{cm}(t)$) for the UAV-manipulator system, and the position error $t \mapsto p(t) - \frac{j_{xy}}{lm} r(t) e_3 - p^\star(t)$ for the UAV-slung-manipulator system. In Figs. 8d and 9d, the manipulator attitude error ($t \mapsto \arccos((r(t)e_3)^T r^\star_+(t))$), and the UAV attitude error ($t \mapsto \arccos((R(t)e_3)^T n^\star(t))$) are shown ($n^\star \equiv n^\star_+$ in Fig. 8d, and $n^\star \equiv n^\star_{+,+}$ in Fig. 9d). The position errors settle around 5 cm, while the orientation errors settle around 5 degrees; we note, however, that they converge to zero in the absence of parameters' mismatch. As such, these simulations provide some insight into the robustness of the controllers with respect to model parameters mismatch. Moreover, the errors described above can be made smaller than some prescribed upper bound either by choosing bigger gains, or by including some form of integral action.

In Figs. 8e–8f and Figs. 9e–9f, the input as obtained from the proposed control law $t \mapsto u^{cl}(t, x(t))$ is shown, with u^{cl} as defined in (66) and in (101). Finally, in Figs. 8e and 9e, the internal forces $t \mapsto T(x(t), u^{cl}(t, x(t)))$ (with T as defined in (20)) are shown. We emphasize only that the tension along the manipulator third axis is, in general, bigger than the others, since, loosely speaking, this is the component that needs to *cancel* the manipulator's weight.

VIII. Conclusions

This manuscript focuses on pose tracking problems for a system composed of two connected rigid bodies, namely an aerial vehicle – with hovering capabilities, and a rod-like rigid body. When a torque input at the joint connecting the two rigid bodies is available, a semi-pose tracking problem is formulated, requiring the rod object to track a desired pose trajectory. When no such torque input is available, a position tracking problem is formulated, requiring a specific point along the axis of axial symmetry of the rod object to track a desired position trajectory. Our approach for solving both these problems lied in finding a state and an input transformations, such that the vector field in the new coordinates is of a known form for which controllers are found in the literature, and which we leveraged in this manuscript.

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