

Lindemann ("disorder") parameter  
for a nucleus

3/01/14 (3)

Making use of the harmonic oscillator approximation for the single-particle potential (cf. Fig. 2-22 Bohr and Mottelson, 1969), one can write (cf. Eq. (2-130) of the above reference),

$$\sum_{k=1}^A \langle r_k^2 \rangle = \frac{\hbar}{M\omega_0} \sum_{k=1}^A (N_k + \frac{1}{2}) = \frac{3}{5} AR^2,$$

where  $A = N + Z$  is the nuclear mass number, while the nuclear radius  $R = r_0 A^{1/3}$ , with  $r_0 = 1.2 \text{ fm}$ . It is of notice that  $N_k$  is the oscillator principal quantum number associated with the state  $k$  (cf. Fig. 2-23 Bohr and Mottelson, 1969).

The average internucleon distance can be determined from the relation (Brink and Broglia, 2005, App. C)

$$a' = \left( \frac{3V}{A} \right)^{1/3} = \left( \frac{\frac{4\pi}{3} R^3}{A} \right)^{1/3} = \left( \frac{4\pi}{3} \right)^{1/3} \times 1.2 \text{ fm} \\ \approx 2 \text{ fm}$$

Thus,

$$\Delta_L = \frac{\sqrt{\frac{3}{5}} R}{2 \text{ fm}} \approx 2.3, \quad (A \approx 120).$$

While it is difficult to compare crystal, aperiodic finite crystal and atomic nuclei, arguably, the above value indicates that a nucleus is liquid-like. More precisely, it is made out of a non-Newtonian fluid, which reacts elastically to sudden solicitations, and plastically to strain. In any case, one expects from  $\Delta_L \approx 2.3$  that the <sup>the nucleus</sup> mean free path is long, larger than nuclear dimensions.