

Boundary Value Problems for Partial Differential Equations

Xiao Ma

Winter 2024

1 Introduction and preliminaries

1.1 Definitions of partial differential equations

Definition 1.1 (Notations of partial derivatives). *For $f(x)$ with one variable x , we know $f'(x) = \frac{df}{dx}$. For $u(x, y)$, we introduce partial derivatives as*

$$\frac{\partial u}{\partial x} = \left. \frac{du}{dx} \right|_{y \text{ is fixed}} = \partial_x u = u_x. \quad (1.1)$$

Similarly,

$$\frac{\partial^2 u}{\partial x^2} = \partial_x^2 u = \partial_{xx} u = u_{xx} \quad (1.2)$$

Example 1. For $u(x, y) = xy^2$, we have

$$\begin{aligned} \frac{\partial u}{\partial x} &= \partial_x u = u_x = y^2, & \frac{\partial u}{\partial y} &= \partial_y u = u_y = 2xy, \\ \frac{\partial^2 u}{\partial x^2} &= \partial_x^2 u = \partial_{xx} u = u_{xx} = 0, & \frac{\partial^2 u}{\partial y^2} &= \partial_y^2 u = \partial_{yy} u = u_{yy} = 2x \end{aligned}$$

Definition 1.2 (Definition of general PDEs). *Given a function $u = u(x, y)$ of two variables, (similarly $u = u(x_1, \dots, x_n)$ of n variables) and an expression $F(u_{xx}, u_{xy}, u_{yy}, u_x, u_y, u, x, y)$ of partial derivatives of u , the following equation is a partial differential equation, abbreviated as PDE.*

$$F(u_{xx}, u_{xy}, u_{yy}, u_x, u_y, u, x, y) = 0 \quad (1.3)$$

In the future, we may also use the notation $F[u]$ to represent $F(u_{xx}, u_{xy}, u_{yy}, u_x, u_y, u, x, y)$. And (1.3) can be rewritten as

$$F[u] = 0. \quad (1.4)$$

Remark 1.3. TODO: higher order is possible but we do not discuss it in this course

Example 2 (Examples of PDEs). Here are some examples of PDEs.

$$\begin{aligned}
 u_{xx} - u_y &= 0 && \text{(the heat equation)} \\
 u_{xx} - u_{yy} &= 0 && \text{(the wave equation)} \\
 u_{xx} + u_{yy} &= 0 && \text{(Laplace's equation)} \\
 u_x + u_y &= 0 && \text{(the transport equation)}
 \end{aligned} \tag{1.5}$$

Definition 1.4 (Order of PDEs). *The order of a PDE is the order of the highest-order derivative in the equation. In (1.5), the first three PDEs are second order, and the last one is first order.*

Definition 1.5 (Linear PDEs). **TODO: def of linear pdes**

Proposition 1.6. TODO: general form of second order linear pdes

We have the following proposition,

Proposition 1.7. *The second-order PDEs can always be written as*

$$a(x, y)u_{xx} + b(x, y)u_{xy} + c(x, y)u_{yy} + d(x, y)u_x + e(x, y)u_y + f(x, y)u = g(x, y), \tag{1.6}$$

We assume $a^2 + b^2 + c^2 \neq 0$ for any x, y (at least one of a, b, c is nonzero).

Definition 1.8. *We call a, b, c, d, e, f coefficients and g source term.*

1.2 Classification of second-order PDEs

In this course, we will mainly consider second-order linear PDEs.

$$a(x, y)u_{xx} + b(x, y)u_{xy} + c(x, y)u_{yy} + d(x, y)u_x + e(x, y)u_y + f(x, y)u = g(x, y), \tag{1.7}$$

These equations are classified as follows by the coefficients a, b, c .

Definition 1.9 (Classification of PDEs). *The second-order linear PDEs (1.7) are classified as elliptic, parabolic and hyperbolic by the following,*

$$\begin{cases}
 4ac - b^2 > 0 & \text{elliptic} \\
 4ac - b^2 = 0 & \text{parabolic} \\
 4ac - b^2 < 0 & \text{hyperbolic}
 \end{cases} \tag{1.8}$$

where we note that a, b, c are functions of x, y and the inequalities in (1.8) is required to be true for any x, y .

1.3 Review of ODEs

Definition 1.10 (Separable ODEs). *The following ODE is the separable ODE*

$$y' + p(x)y = 0. \quad (1.9)$$

Theorem 1.11. *Separable ODE can be solved in the following way*

$$y' + p(x)y = 0 \Rightarrow \frac{dy}{dx} + p(x)y = 0 \Rightarrow \frac{dy}{y} = -p(x)dx \Rightarrow \int \frac{dy}{y} = -\int p(x)dx$$

and the solution is

$$y(x) = Ce^{-\int p(x)dx} \quad (1.10)$$

Example 3. Let us solve the following ODEs

$$y' = -6xy$$

Apply the procedure in Theorem 1.11

$$y' = -6xy \Rightarrow \frac{dy}{y} = -6xdx \Rightarrow \int \frac{dy}{y} = -\int 6xdx \Rightarrow \ln |y| = -3x^2 + C'$$

Therefore, the solution is

$$y = Ce^{-3x^2}.$$

where $C(= \pm e^{C'})$ is an arbitrary constant.

Definition 1.12 (Linear ODEs). *The following ODE is the linear ODE*

$$y' + p(x)y = q(x). \quad (1.11)$$

Theorem 1.13. *Linear ODE can be solved by the following procedure*

1. Solve the corresponding separable equation $y' - p(x)y = 0$ to obtain a solution $\hat{y} = e^{\int p(x)dx}$.

2. Multiply the linear ODE by \hat{y} and rewrite the ODE

$$y' + p(x)y = q(x) \Rightarrow \hat{y}(y' + p(x)y) = \hat{y}q(x) \Rightarrow (\hat{y}y)' = \hat{y}q(x)$$

3. Integrate the above equation

$$\begin{aligned} \hat{y}y &= \int \hat{y}q(x)dx + C \Rightarrow y = \frac{1}{\hat{y}} \left(\int \hat{y}q(x)dx + C \right) \\ &\Rightarrow y = e^{-\int p(x)dx} \left(\int q(x)e^{\int p(x)dx}dx + C \right) \end{aligned}$$

Remark 1.14. In (2), we applied the following equation

$$(\hat{y}y)' = \hat{y}(y' + p(x)y)$$

which is a corollary of the Leibniz rule.

$$(\hat{y}y)' = \hat{y}'y + \hat{y}y' = \hat{y}y' + p(x)\hat{y}y = \hat{y}(y' + p(x)y)$$

Example 4. $(x^2 + 1)y' + 3xy = 6x, y(0) = 3$ is solved as $y(x) = 2 + (x^2 + 1)^{-3/2}$ by the procedure in Theorem 1.13.

1. Divide both sides by $(x^2 + 1)$.

$$(x^2 + 1)y' + 3xy = 6x \quad \Rightarrow \quad y' + \frac{3x}{x^2 + 1}y = \frac{6x}{x^2 + 1}$$

2. Solve the corresponding separable equation $y' - \frac{3x}{x^2 + 1}y = 0$ to obtain a solution $(x^2 + 1)^{\frac{3}{2}}$.
3. Multiply the linear ODE by $(x^2 + 1)^{\frac{3}{2}}$ and rewrite the ODE

$$y' + \frac{3x}{x^2 + 1}y = \frac{6x}{x^2 + 1} \quad \Rightarrow \quad ((x^2 + 1)^{\frac{3}{2}}y)' = 6x(x^2 + 1)^{\frac{1}{2}}$$

4. Integrate the above equation

$$\begin{aligned} ((x^2 + 1)^{\frac{3}{2}}y)' = 6x(x^2 + 1)^{\frac{1}{2}} &\quad \Rightarrow \quad y = (x^2 + 1)^{-\frac{3}{2}} \left(\int 6x(x^2 + 1)^{\frac{1}{2}} + C \right) \\ &\quad \Rightarrow \quad y = 2 + C(x^2 + 1)^{-\frac{3}{2}} \end{aligned}$$

Definition 1.15 (Second order ODEs). *The constant coefficient second order ODEs are the following equations*

$$ay'' + by' + cy = 0 \quad (a \neq 0). \quad (1.12)$$

Theorem 1.16. *Constant coefficient second order ODEs can be solved by the following procedure*

1. Solve the characteristic equation $a\lambda^2 + b\lambda + c = 0$ to get two solutions λ_1 and λ_2 .
2. If $\lambda_1 \neq \lambda_2$, the general solution is

$$y(x) = C_1 e^{\lambda_1 x} + C_2 e^{\lambda_2 x} \quad (1.13)$$

3. If $\lambda_1 = \lambda_2 = \lambda$, the general solution is

$$y(x) = (C_1 + C_2 x) e^{\lambda x} \quad (1.14)$$

4. If λ_1, λ_2 are complex roots $\alpha \pm i\beta$, apply the Euler's formula to rewrite (1.13)

$$y(x) = C_1 e^{(\alpha+i\beta)x} + C_2 e^{(\alpha-i\beta)x} = e^{\alpha x} (c_1 \cos \beta x + c_2 \sin \beta x)$$

where $c_1 = C_1 + C_2$, $c_2 = i(C_1 - C_2)$.

Example 5. $y'' - 3y' + 2y = 0$ is solved as $y(x) = C_1 e^x + C_2 e^{2x}$ by the procedure in Theorem 1.16.

TODO:

Example 6. $y'' + y = 0$ is solved as $y(x) = C_1 \cos(x) + C_2 \sin(x)$ by the procedure in Theorem 1.16.

TODO:

Example 7. $y'' + 2y' + y = 0$ is solved as $y(x) = (C_1 + C_2 x)e^x$ by the procedure in Theorem 1.16.

TODO:

1.4 Separation of variables

Many linear PDEs can be reduced to linear ODEs with the method of separation of variables, described below.

We take the Laplace's equation

$$u_{xx} + u_{yy} = 0 \tag{1.15}$$

as an example.

TODO: Boundary conditions

We are looking for a separated solution. Substitute into (1.15), then we get

$$X''Y + XY'' = 0 \quad \Rightarrow \quad \frac{X''}{X} = -\frac{Y''}{Y}$$

The following lemma implies that X''/X and Y''/Y are constants.

Lemma 1.17. TODO: $f(x) = g(y)$ implies that $f(x) = g(y) = \text{const}$,

Proof. $f(x) = g(y) \Rightarrow f'(x) = 0 \Rightarrow f(x) = \text{const}$. □

Let λ be a constant and we write

$$X'' + \lambda X = 0, \quad Y'' - \lambda Y = 0.$$

We call λ the separation constant. At this moment λ is arbitrary. Thus the PDE was reduced to two ODEs.

1.4.1 Solving separated solutions

If $\lambda = 0$, then two ODEs have the following linearly independent solutions.

$$X = 1, x, \quad Y = 1, y. \quad (1.16)$$

If $\lambda \neq 0$, then two ODEs have the following linearly independent solutions.

$$X = e^{\sqrt{-\lambda}x}, e^{-\sqrt{-\lambda}x}, \quad Y = e^{\sqrt{\lambda}y}, e^{-\sqrt{\lambda}y}. \quad (1.17)$$

In either case, the solution is given by superpositions:

$$u = \begin{cases} (A_1x + A_2)(B_1y + B_2), & \lambda = 0 \\ (A_1e^{\sqrt{-\lambda}x} + A_2e^{-\sqrt{-\lambda}x})(B_1e^{\sqrt{\lambda}y} + B_2e^{-\sqrt{\lambda}y}), & \lambda \neq 0 \end{cases} \quad (1.18)$$

where A_1, A_2, B_1, B_2 are constants.

For $\lambda > 0$, by writing $\lambda = k^2 (k > 0)$ we have

$$u(x, y) = (A_1e^{ikx} + A_2e^{-ikx})(B_1e^{ky} + B_2e^{-ky}), \quad (1.19)$$

and for $\lambda < 0$, by writing $\lambda = -l^2 (l > 0)$ we have

$$u(x, y) = (A_1e^{lx} + A_2e^{-lx})(B_1e^{ily} + B_2e^{-ily}). \quad (1.20)$$

Instead of (1.17) we can also choose

$$X = \cos(\sqrt{\lambda}x), \sin(\sqrt{\lambda}x), \quad Y = \cosh(\sqrt{\lambda}y), \sinh(\sqrt{\lambda}y). \quad (1.21)$$

In this case, we have

$$u(x, y) = (A_1 \cos(kx) + A_2 \sin(kx))(B_1 \cosh(ky) + B_2 \sinh(ky)) \quad (1.22)$$

$$u(x, y) = (A_1 \cosh(lx) + A_2 \sinh(lx))(B_1 \cos(ly) + B_2 \sin(ly)) \quad (1.23)$$

Note that (1.22) becomes (1.19) and (1.23) becomes (1.20) by redefining the coefficients. We call solutions such as (1.18) through (1.23) separated solutions because they are given in the form $u(x, y) = X(x)Y(y)$.

1.4.2 Solving the boundary value problem

TODO: revise the section The separation constant λ and coefficients A_1, A_2, B_1, B_2 are partially determined by boundary conditions. Suppose that our Laplace's equation is considered in the region $0 < x < L, 0 < y < \infty$ with boundary conditions

$$u(0, y) = 0, \quad u(L, y) = 0, \quad u(x, 0) = 0.$$

We find that $u = (A_1x + A_2)(B_1y + B_2)$ in (1.3) satisfies the boundary conditions only when $u = 0$.

We then find that (1.5) and (1.7) satisfy the conditions $u(0, y) = u(L, y) = 0$ only when $A_1 = A_2 = 0$. That is, only the solution $u = 0$ satisfies the boundary conditions.

Finally (1.4) and (1.6) satisfy $u(0, y) = u(L, y) = 0$ when $A_1 = 0$ and $k = n\pi/L$, where n is an integer. Furthermore we find $B_1 = 0$ by the condition $u(x, 0) = 0$. That is, the solution $A_2B_2 \sin(kx) \sinh(ky)$ with $k = n\pi/L (n = 0, \pm 1, \pm 2, \dots)$ satisfies the boundary conditions.

Therefore we obtain the following separated solutions of Laplace's equation satisfying the boundary conditions.

$$u(x, y) = A \sin \frac{n\pi x}{L} \sinh \frac{n\pi y}{L}, \quad n = 1, 2, \dots,$$

TODO: linear combination and Fourier series