

# Automatic extraction of figures from PDF documents

Piotr Praczyk (piotr.praczyk@gmail.com)

# Invenio/INSPIRE

- Invenio – digital library software developed at CERN to manage the repository of documents created in the institution
- SPIRES – The digital library of preprints created at SLAC.
- Invenio + SPIRES = INSPIRE

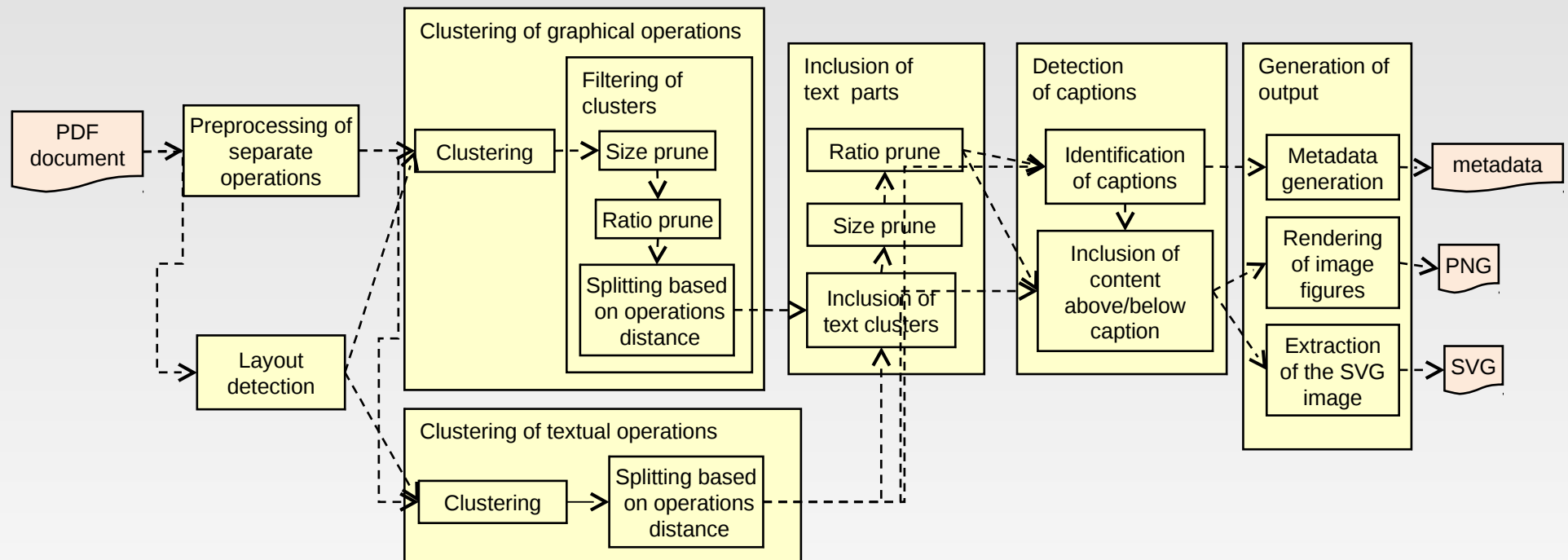
# Possible sources of figures

- Extraction from LaTeX sources
  - Only figures attached via the `\figure` tag, constructed from a single file, without complicated transformations described in the LaTeX source.
- Extraction from PDF files
  - Heuristic estimation, where the figures are but an universal source of data
- Directly from authors (Web interface)
- Others

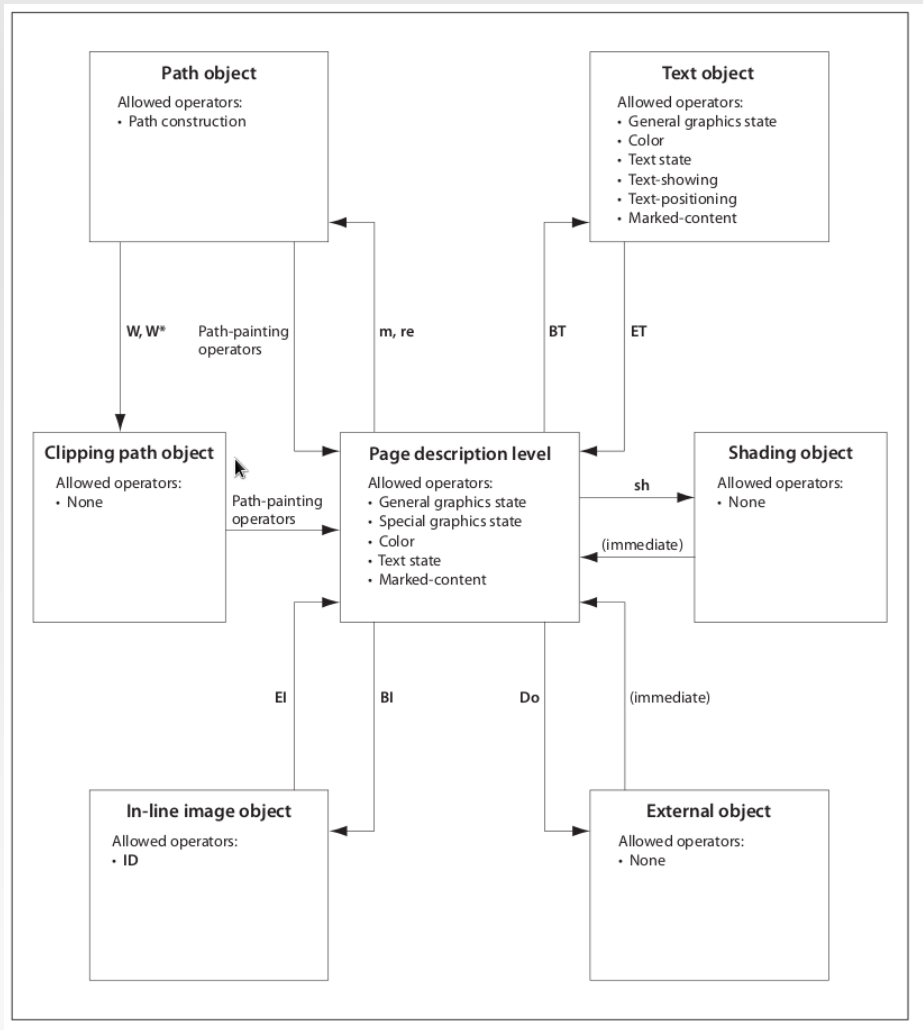
	Foreign code	PDF code
<b>Referenced as external document</b>	External objects	External forms
<b>Included in the content stream</b>	Inline Objects	Standard code

- We can not assume that external objects or forms will always contain figures

# The complete process



# Analysis of the PDF stream content



- Operations modify the state of the interpreter
  - Current transformation matrix, current colour etc...
- 3 types: graphical, text, transformation
- Every operation can be annotated with the boundary of its effect in some canvas

# Pseudo-code of the clustering

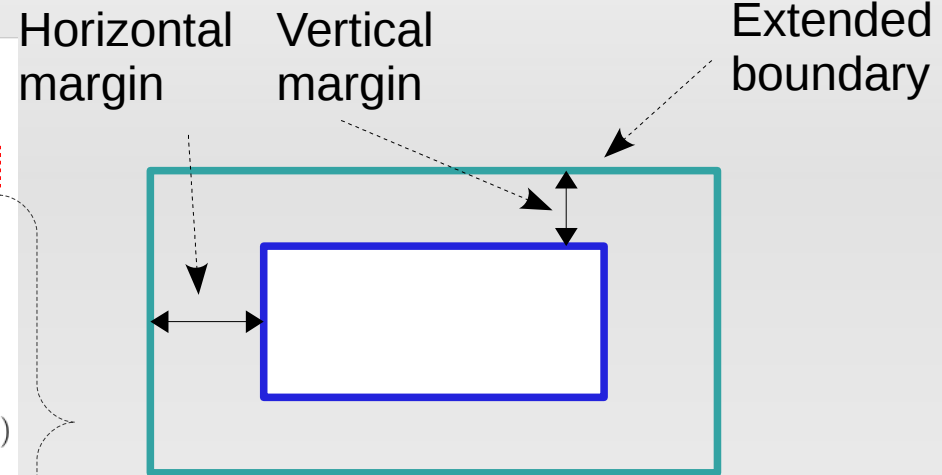
```
1:  $t_x \leftarrow \text{IntervalTree}()$ 
2:  $t_y \leftarrow \text{IntervalTree}()$ 
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4:    $\text{boundary} \leftarrow \text{extend\_by\_margins}(op.\text{boundary})$ 
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```

Original  
boundary



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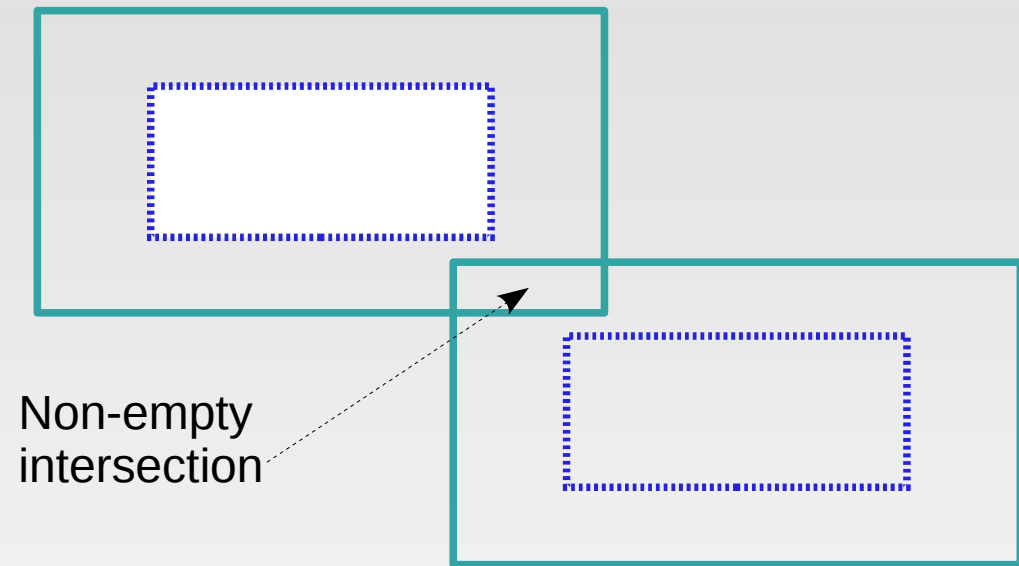
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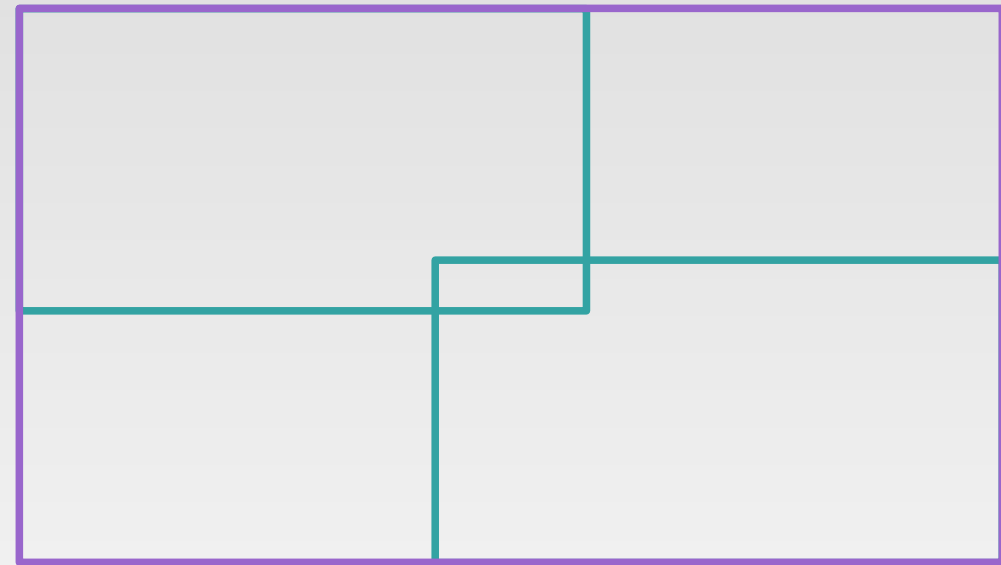
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Intersection of 2 rectangles having edges parallel to axis = intersection of both projections in X and Y direction

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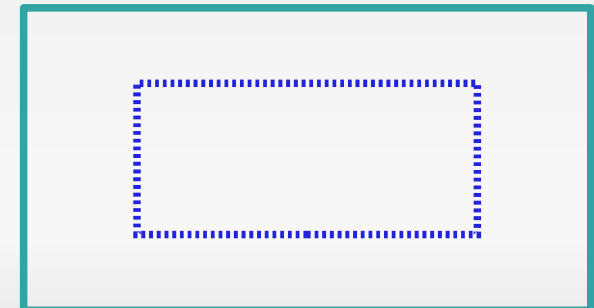
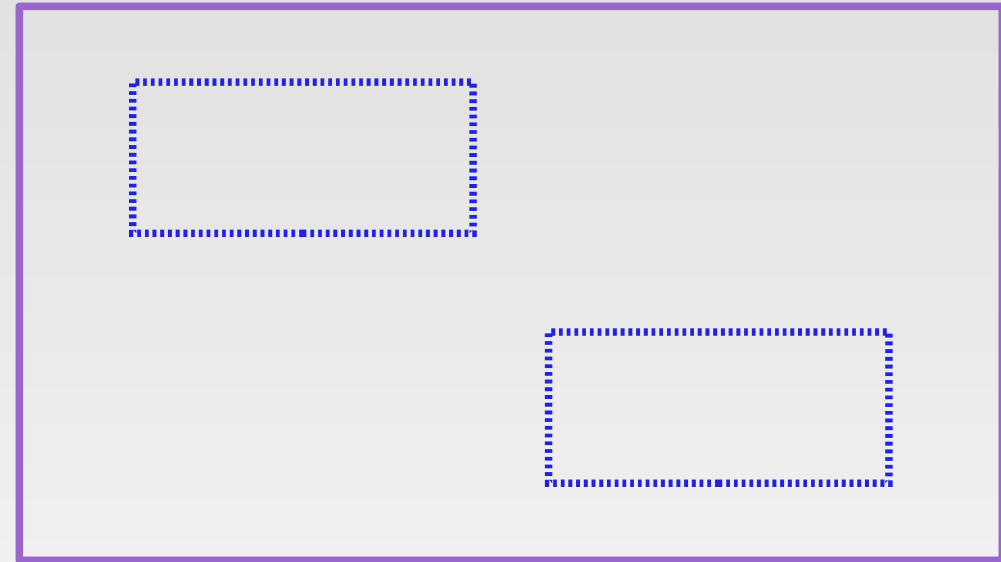


Smallest enclosing  
boundary

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```

Repeat until no  
Intersections  
With existing areas



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Collection of results  
Cluster boundary  $\rightarrow$  list of original operators

# One-dimensional intersections - what data structure ?

- Interface:
  - `add_interval(beginning, end, identifier)`
  - `remove_interval(identifier)`
  - `contains_interval(beginning, end)`
  - `get_intersecting()`



# Interval trees

Interval: (0, 100)  
0 intervals  
contained intervals:

Empty

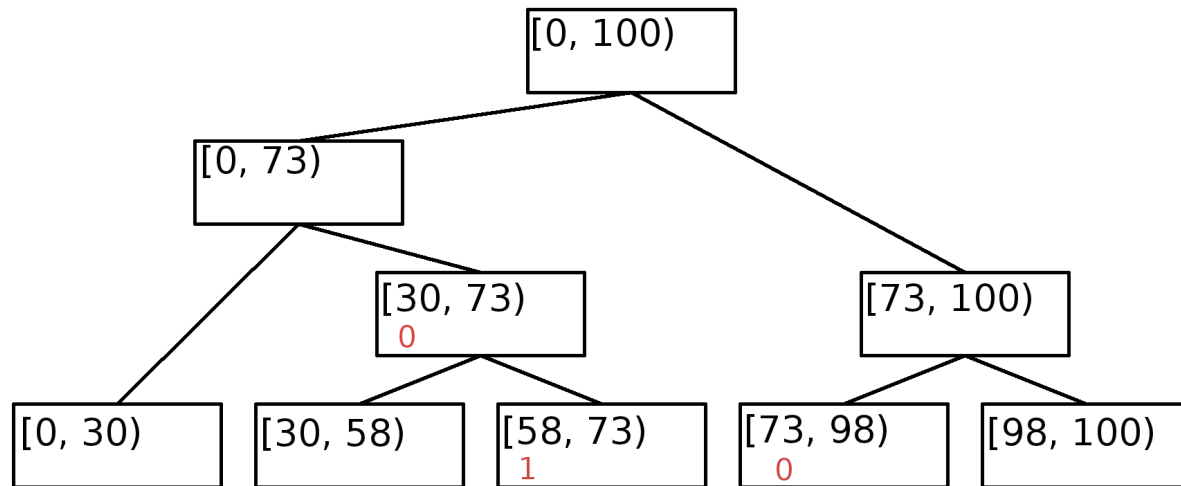
Number: 37  
Interval: (0, 100)  
0 intervals  
contained intervals:

Number: 46  
Interval: (37, 100)  
0 intervals  
contained intervals:

Interval: (0, 37)  
0 intervals  
contained intervals:

Interval: (37, 46)  
1 intervals  
contained intervals: 0

Interval: (46, 100)  
0 intervals  
contained intervals:

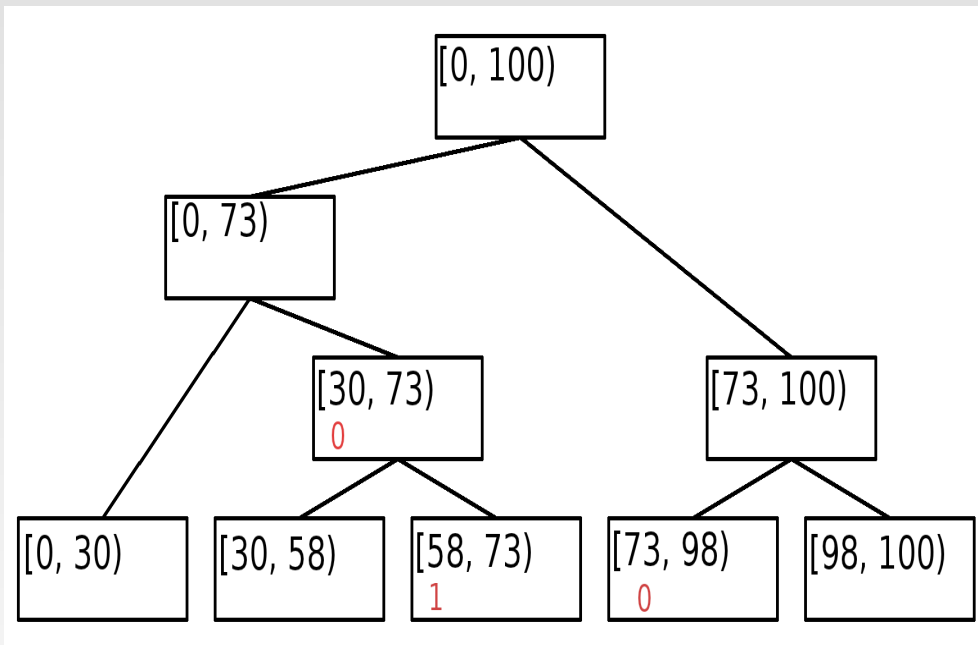


3 intervals

One interval

In the worst case,  
depth  $\log(2*n)$  for  $n$  intervals

# Principles



Balanced binary tree

- Root represents the entire space
- Every node represents an interval
- Interval represented by a node is equal to a sum of intervals from its children
- Every node remembers interval identifiers

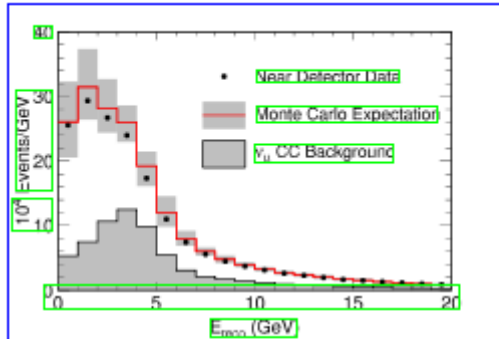
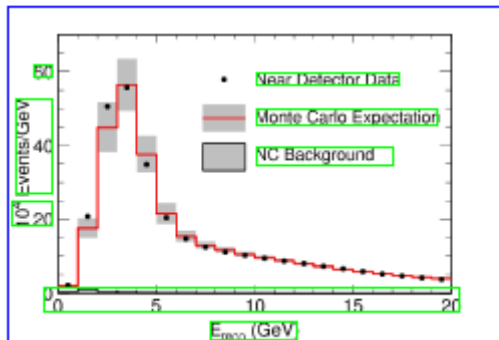


FIG. 7: Distribution of reconstructed visible energy for selected neutral-current events in the near-detector, for the data (solid points) versus the Monte Carlo prediction (open histogram). The systematic errors ( $1\sigma$ ) for the Monte Carlo are shown by the shaded band. Also shown is the Monte Carlo prediction for the background of misidentified charged-current events in the near-detector sample (hatched histogram).



near and far-detectors. Specifically, for the case of the  $\nu_\mu$  charged-current component of the neutral-current and charged-current samples, the F/N extrapolation predicts the number of events at the far-detector for the  $i$ -th bin of reconstructed energy to be

$$F_i^{\text{prediction}} = N_i^{\text{data}} \left( \frac{\sum_x \sum_j F_{ij}^{\text{MC}} P_{\nu_\mu \rightarrow \nu_e}(E_j)}{N_i^{\text{MC}}} \right) \quad (1)$$

where  $N_i^{\text{data}}$  is the number of selected events in the  $i$ -th reconstructed energy bin in the near-detector and  $N_i^{\text{MC}}$  is the number of events expected in that bin from the near-detector Monte Carlo simulation. The  $F_{ij}^{\text{MC}}$  represents the number of events expected from the far-detector Monte Carlo simulation in the  $i$ -th bin of reconstructed energy and  $j$ -th bin of true neutrino energy. In the equation,  $E_j$  is the true neutrino energy and  $P_{\nu_\mu \rightarrow \nu_e}$  the probability of muon-neutrino transition to any other flavor.

In particular, for the neutral-current spectrum, the extrapolation must take neutrino oscillations into account to properly characterize the predominant background arising from misidentified charged-current  $\nu_\mu$ , and it must include the small spectral distortion resulting from misidentified charged-current  $\nu_\tau$  and  $\nu_e$  events. Thus, there are five separate classes of events that must be extrapolated to the far-detector: (i) genuine neutral-current interactions, (ii)  $\nu_\mu$  charged-current interactions, (iii)  $\nu_\tau$  charged-current interactions, (iv) possible  $\nu_e$  charged-current interactions originating from  $\nu_\mu$  oscillations, and (v) charged-current  $\nu_e$  interactions initiated by the intrinsic  $\nu_e$  beam component. The muon neutrinos in the simulation include oscillations and are integrated in bins of reconstructed energy to account for the changing background. Oscillations of the intrinsic beam  $\nu_e$  into  $\nu_\mu$  are not taken into account as those  $\nu_e$  comprise only 1.3% of the neutrinos in the beam and

- Graphical and textual operations clustered separately
- Figures contain graphical and textual areas
- Logical parts of text tend to get separated.



# Is a figure or not?

- Is large enough ?
- Is the aspect ratio in appropriate interval?
- Is the figure candidate consisted in the content string ?
- Does it contain enough graphics compared to text?

# The importance of page layout

4

L.C. Bland, for the STAR collaboration

for 'direct photon' and  $W^\pm$  production processes. The EMC provides critical phase space coverage for both  $\gamma$ +jet and  $W^\pm$  production studies. By detecting large- $p_T$  processes at forward angles, asymmetric initial states, where one parton has a larger  $x$  than the other, are emphasized. For photon production, such collisions effectively select large- $x$  quarks as an *analyzer* of the polarization of small- $x$  gluons. In addition to kinematic selection of quarks with large polarization (the pDIS asymmetry  $A_1^p$  increases with increasing  $x_{quark}$ ), the EMC will detect photons produced at partonic CM angles where the partonic-level spin correlation parameter ( $\hat{a}_{LL}$ ) approaches its maximum value, unlike the situation encountered for midrapidity photon production.

The large acceptance of STAR is critical for the measurement of the away-side jet in coincidence with the produced photon (Fig. 1). Coincident  $\gamma$ +jet detection allows for the reconstruction of the initial-state partonic kinematics. With the momentum fractions  $x_{quark}$  and  $x_{gluon}$  determined by the experiment, a more direct connection between the measured polarization observables and the gluon helicity asymmetry distribution can be made. This is advantageous to in trying understand how the measurement errors will influence the determination of  $\Delta G(x)$ .

### §3. Gluon polarization measurements at STAR

The existing data for scaling violations in pDIS provide only very loose constraints on  $\Delta G(x)$ . Several analyses of these constraints have been made <sup>3), 4)</sup> and have generally concluded that the *integral*  $\Delta G$  is positive. The variation of the gluon helicity asymmetry distribution with the gluon momentum fraction ( $x_{gluon}$ ) has significant differences in these different analyses. There is generally always a positive peak of  $x\Delta G(x)$ , but the  $x_{gluon}$  value of the peak is not well constrained. Consequently, the gluon polarization,  $(\Delta G(x)/G(x))$  can be either large or small, depending on where the peak in  $x\Delta G(x)$  occurs. Many parameterizations of  $\Delta G(x)$  that are consistent with existing measurements result in negatively polarized gluons at some  $x_{gluon}$  values. In leading order perturbative QCD (pQCD), the spin correlation parameter ( $A_{LL}$ ) that will be measured in  $\bar{p} + \bar{p} \rightarrow \gamma + X, \gamma + \text{jet} + X$  reactions at RHIC is proportional to the gluon polarization.

To illustrate the sensitivity of  $\gamma$  + jet coincidence measurements planned for STAR, simulations using the three  $\Delta G(x)$  models in Ref. <sup>4)</sup> (hereafter referred to as GS sets A,B and C) have been performed. In all cases, the input  $\Delta G(x)$  must be

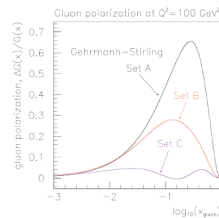


Fig. 2. Gluon polarizations computed from models of  $\Delta G(x)$  consistent with polarized deep inelastic scaling violations<sup>4)</sup>. The structure functions are evolved to the scale that will be probed at RHIC.

- Content from different columns can easily get mixed
- Content columns are detected
- Clustering operates in every column separately

# Extraction captions

- Captions appear as clusters of text content lying close to the figure
- Captions tend to follow a particular grammar
- Captions can be used to improve the quality of graphics detection

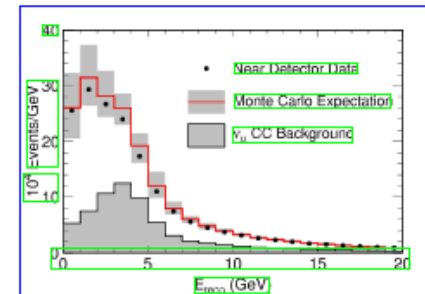
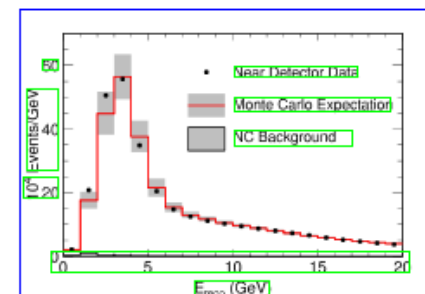


FIG. 7: Distribution of reconstructed visible energy for selected neutral-current events in the near-detector, for the data (solid points) versus the Monte Carlo prediction (open histogram). The systematic errors ( $1\sigma$ ) for the Monte Carlo are shown by the shaded band. Also shown is the Monte Carlo prediction for the background of misidentified charged-current events in the near-detector sample (hatched histogram).



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# Complete extraction process

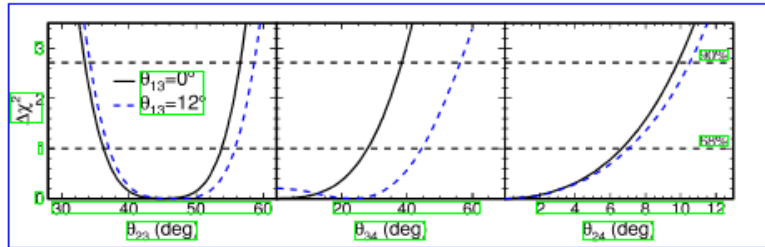


FIG. 14: Projections of  $\Delta\chi^2$  as a function of the mixing angles for the  $m_4 \gg m_3$  model. The solid line is obtained for the case of null  $\nu_e$  appearance whereas the dashed line represents solutions with  $\nu_e$  appearance at the CHOOZ limit. The ranges of values allowed at 68% and 90% confidence levels lie within contours below the horizontal dashed lines.

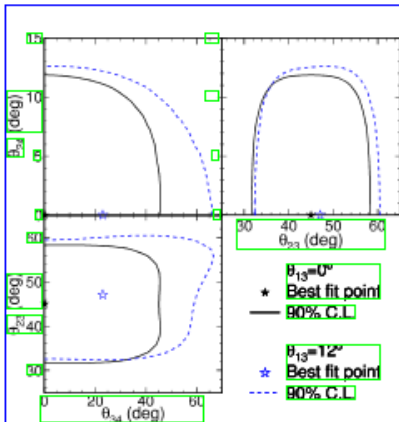


FIG. 15: Contours representing 90% confidence level for the  $m_4 \gg m_3$  model. The solid line and best-fit point (solid symbol) are obtained for the case of null  $\nu_e$  appearance, whereas the dashed line and corresponding best-fit point (open symbol) is obtained with  $\nu_e$  appearance included with  $\theta_{13}$  at the CHOOZ limit.

disappearance probability is a maximum. The determination of the limit follows the procedure described above, but with the addition of selecting a value of  $\theta_{24}$  for each test case as well. At 90% confidence level  $f_d < 0.52$  (0.55) for  $E_\nu = 1.4$  GeV in this model. Thus, in either model, approximately 50% of the disappearing  $\nu_\mu$  can convert to  $\nu_e$  at 90% confidence level as long as the amount of  $\nu_e$  appearance is less than the limit presented by the CHOOZ collaboration.

## IX. OSCILLATIONS WITH DECAY

It was noted more than a decade ago that neutrino decay, as an alternative or companion process to neutrino oscillations, offers some capability for reproducing neutrino disappearance trends [18]. The model investigated here [36] includes neutrino oscillations occurring in parallel with neutrino decay. Normal neutrino-mass ordering is assumed, and the mass eigenstates  $\nu_1, \nu_2$  are approximately degenerate, so that  $m_3 \gg m_2 \approx m_1$ . The heaviest neutrino-mass state  $\nu_3$  is allowed to decay into an invisible final state. With these assumptions, and neglecting the small contributions from  $\nu_e$  mixing, only the two neutrino flavor states  $\nu_\mu$  and  $\nu_\tau$ , and the corresponding mass states  $\nu_2$  and  $\nu_3$ , are considered. The evolution of the neutrino flavor states is given by [36]:

$$\begin{aligned} \frac{d\nu}{dt} &= \frac{\Delta m_{32}^2}{4E} \begin{pmatrix} -\cos 2\theta & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \nu_\tau \end{pmatrix} \\ &- i \frac{m_3}{4\tau_3 E} \begin{pmatrix} 2 \sin^2 \theta & \sin 2\theta \\ \sin 2\theta & 2 \cos^2 \theta \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \nu_\tau \end{pmatrix} \end{aligned} \quad (16)$$

where  $\tau_3$  is the lifetime of the  $\nu_3$  mass state and  $\theta$  is the mixing angle governing oscillations between  $\nu_\mu$  and  $\nu_\tau$ . Solving Eq. (16) one obtains probabilities for  $\nu_\mu$  survival or decay:

$$P_{\mu\mu} = \cos^4 \theta + \sin^4 \theta e^{-\frac{m_3^2 L}{2E\tau_3}} + \frac{2 \cos^2 \theta \sin^2 \theta e^{-\frac{m_3^2 L}{2E\tau_3}} \cos \left( \frac{\Delta m_{32}^2 L}{2E} \right)}{2E} \quad (17)$$

$$P_{\text{decay}} = \left( 1 - e^{-\frac{m_3^2 L}{2E\tau_3}} \right) \sin^2 \theta \quad (18)$$

The limits  $\tau_3 \rightarrow \infty$  and  $\Delta m_{32}^2 \rightarrow 0$  correspond to a pure oscillations or a pure decay scenario, respectively.

In a conventional neutrino oscillations scenario, the ratio of the predicted charged-current spectrum in the far-detector with the null-oscillation expectation displays the characteristic "dip" at the assumed  $\Delta m_{32}^2$  value that is

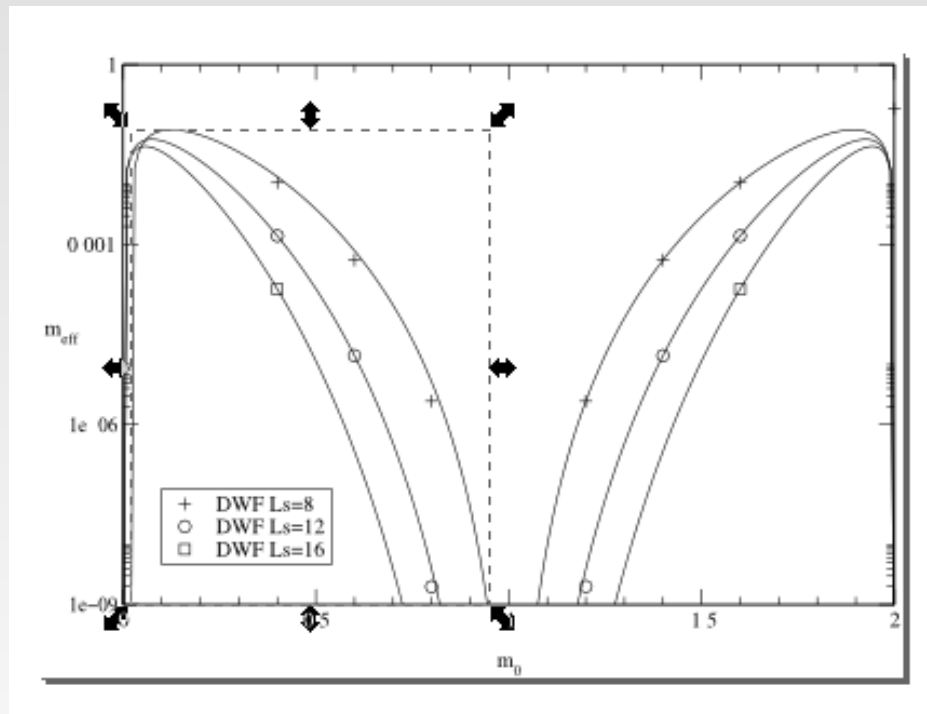
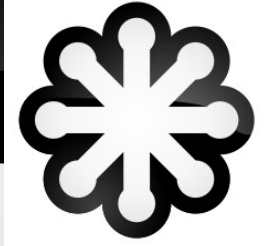
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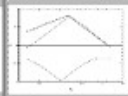
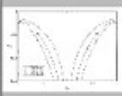
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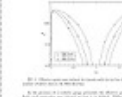
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## Domain-Wall Induced Quark Masses in Topologically-Nontrivial Background

Vakariya Gadiyak, Xiangdong Ji, Chulwoo Jung  
*Department of Physics  
University of Maryland  
College Park, Maryland 20742*

(UMD PP#00-055 DOE/ER/40762-201 February 2000)

### Abstract

In the domain-wall formulation of chiral fermion, the finite separation between domain-walls ( $L_s$ ) induces an effective quark mass ( $m_{\text{eff}}$ ) which complicates the chiral limit. In this work, we study the size of the effective mass as the function of  $L_s$  and the domain-wall height  $m_0$  by calculating the smallest eigenvalue of the hermitian domain-wall Dirac operator in the topologically-nontrivial background fields. We find that, just like in the free case,  $m_{\text{eff}}$  decreases exponentially in  $L_s$  with a rate depending on  $m_0$ . However, quantum fluctuations amplify the wall effects significantly. Our numerical result is consistent with a previous study of the effective mass from the Gell-Mann-Oakes-Renner relation.

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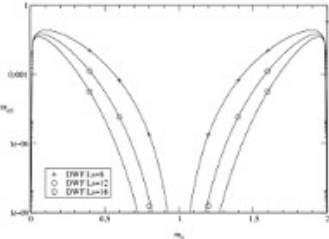


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