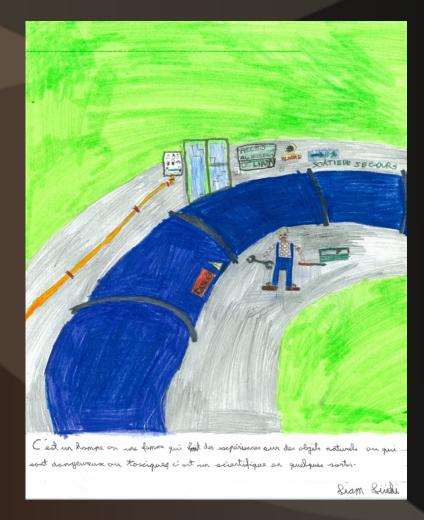
The infrastructure for figures in INSPIRE

Piotr Praczyk,

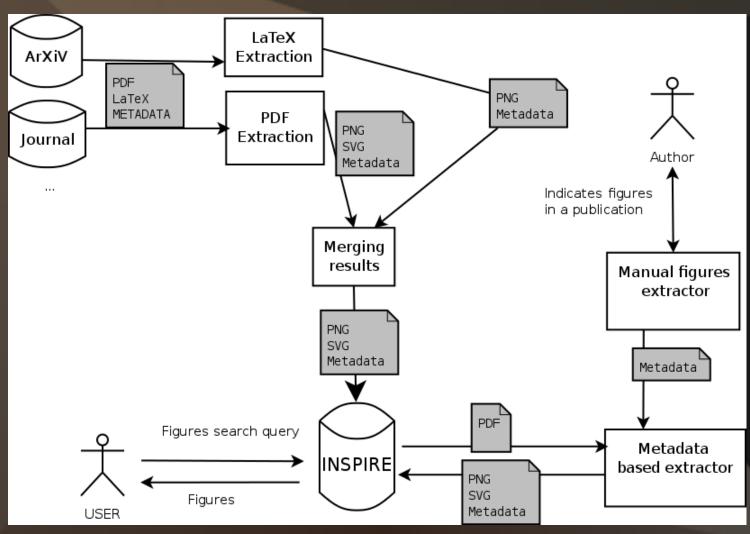
CERN 02/02/2012

Usage of graphics in scholarly communication

- Describe experiments
- Summarise large amounts of data
- Illustrate relations between results
- Present ideas in a schematic manner



Usage of Plots in Inspire



Extracting data from PDF

PDF:

- Stream of instructions
- Embeded objects
 - Fonts
 - External objects
- Meta-description

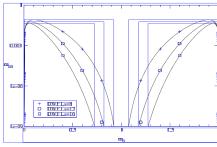


FIG. 1. Effective quark mass induced by domain-walls for the free field configuration D_8 is the

In the presence of a realistic gauge potential, the effective quark mass result from the finite wall separation may depend on how it is defined. Different definitions shall yield results consistent up to a factor of order unity. One approach is to exploit the explicit quark mass dependence in chiral Ward identities such as the Gell-Mann-Oakes-Renner (GMOR) relation as done in Ref. [7]. Here we explore the effective mass in an alternative way. In continuum field theory, the Atiyah-Singer theorem [8] states that the Dirac operator has a zero eigenvalue in the presence of an external background with topological charge [2] \equiv 1.7 The explicit form of the solution was found by 't Hooft in 1976 [9]. On the lattice, however, the notion of topological charge is ill defined: any gauge configuration can be continuously deformed into a null gauge field. Moreover, the discretization of an instanton field can introduce finite lattice-spacing effects lifting any exact zero eigenvalue. Therefore, a test of the Atiyah-Singer theorem on lattice is usually complicated with various lattice artifacts

There exists, however, a definition of lattice topology and fermion zero mode which largely avoids this complication. In the overlap formalism, the Dirac operator is constructed from the overlap of two many-fermion ground states [3]. According to their recipe, one starts from a four-dimensional Wilson-Dirac operator with a negative Wilson mass m_0 and calculates its eigenvalues. For m_0 small and positive, the number of positive eigenvalues is equal to that of negative ones. When m_0 increases, a level might cross from positive to negative or vice versa. When this happens, the gauge field is regarded to have a net topological charge [Q] = 1. Then the overlap determinant is exactly zero by construction. This definition of lattice topology and zero mode do depend on for instance, the Wilson parameters v and m_0 . However, the zero eigenvalue is exact, independent of the lattice spacing v and volume v.

Intermediate steps of the algorithm

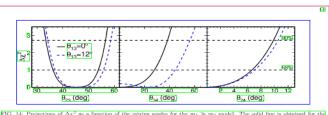


FIG. 14: Projections of $\Delta \chi^2$ as a function of the mixing angles for the $m_4 \gg m_3$ model. The solid line is obtained for the case of null ν_e appearance whereas the dashed line represents solutions with ν_e appearance at the CHOOZ limit. The range of values allowed at 68% and 90% condidence levels lie within contours below the horizontal dashed lines.

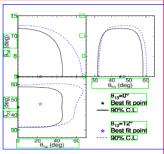


FIG. 15: Contours representing 90% confidence level for the $m_4 \gg m_0$ model. The solid line and best-fit point (solid symbol) are obtained for the case of null ν_e appearance, wherea the dashed line and corresponding best-fit point (open symbol) is obtained with ν_e appearance included with θ_{13} at the

Hisappearance probability is a maximum. The determination of the limit follows the procedure described above but with the addition of selecting a value of θ_{24} for each test case as well. At 10% confidence level $f_s < 0.52$ (0.55 for $E_p = 1.4$ GeV in this model. Thus, in either model approximately 50% of the disappearing ν_{μ} can convert to ν_{μ} at 90% confidence level as long as the amount of ν_{τ} appearance is less than the limit presented by the CHOOZ subhavoration

IX. OSCILLATIONS WITH DECAY

It was noted more than a decade ago that neutrinc locay, as an alternative or companion process to neutrino oscillations, offers some capability for reproducing neutrino disappearance trends [18]. The model investigated here [36] includes neutrino oscillations occurring in parallel with neutrino decay. Norman neutrino-mass ordering is assumed, and the mass eigenstates ν_1 , ν_2 are approximately degenerate, so that $m_2 \gg m_1$. The heaviest neutrino-mass state ν_2 is allowed to decay into an invisible final state. With these assumptions, and neglecting the small contributions from ν_e mixing, only the two neutrino flavor states ν_e , and ν_e , and the corresponding mass states ν_2 and ν_2 , are considered. The evolution of the neutrino flavor states is given by [36]:



where τ_3 is the lifetime of the ν_3 mass state and θ is th mixing angle governing oscillations between ν_μ and ν_τ Solving Eq. (16) one obtains probabilities for ν_μ survivi or decay:

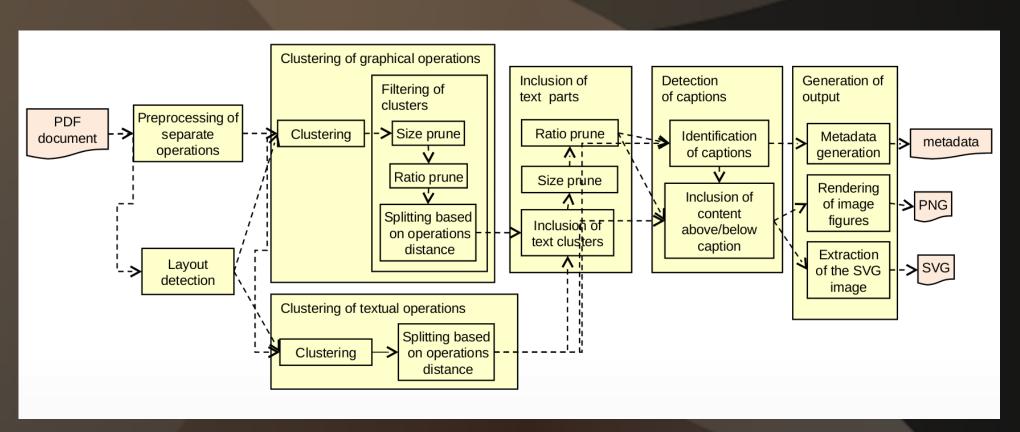
$$\begin{array}{c|c} \mathbb{P}_{10} & = \cos^4 \theta + \sin^4 \theta e^{-\frac{m_{11}}{22L}} \boxplus \\ & 2\cos^2 \theta \sin^2 \theta e^{-\frac{m_{11}}{2r_{10}L}} \cos \left(\frac{\Delta m_{22}^2 L}{2E}\right) \end{array} \quad \boxed{17} \\ \mathbb{P}_{\text{decay}} & = \left(1 - e^{-\frac{m_{11}}{2r_{10}L}}\right) \sin^2 \theta. \quad \boxed{18}$$

The limits $\tau_3 \to \infty$ and $\Delta m_{32}^2 \to 0$ correspond to scillations or a pure decay scenario, respectively.

In a conventional neutrino oscillations scenario, the ratio of the predicted charged-current spectrum in the fardetector with the null-oscillation expectation displays the characteristic "dip" at the assumed Δm_{ϕ}^2 , value that is

- Regions of graphics (blue)
 - Clustered graphic operations
- Regions of text (green)
 - Clustered text operations
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Extraction of figures from PDF



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Automatic extraction of figures



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Distributions

DELPHI Calibration

DELPHI Calibration

Abstract

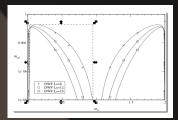
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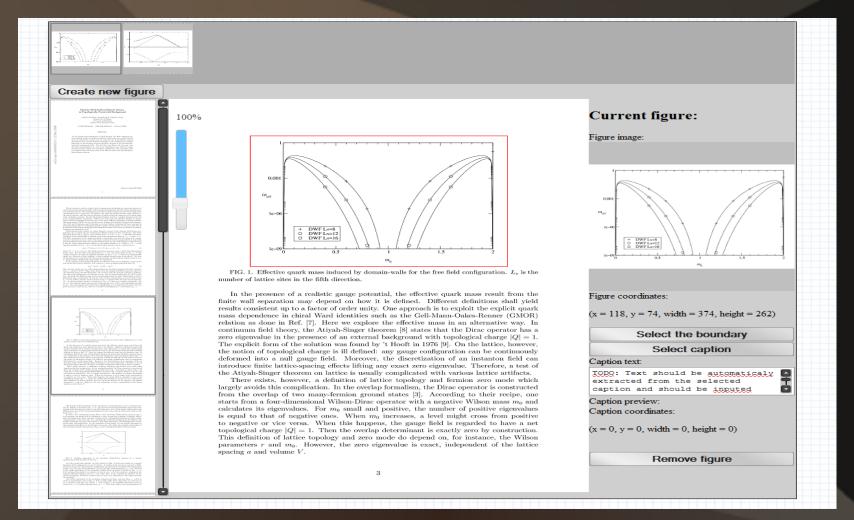


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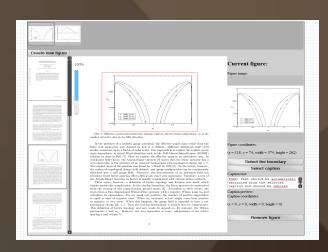
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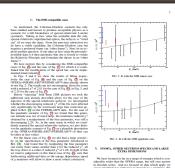


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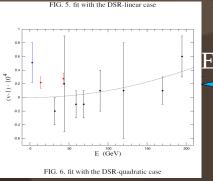
Figures from scientific publications



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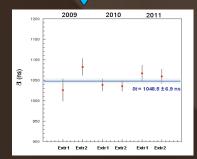
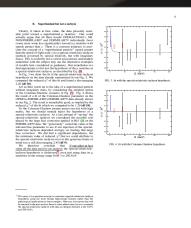


Figure 2 (extracted from different publication)



Questions?