

# QUANTUM COMPUTATION AND QUANTUM INFORMATION: THE QUANTUM FOURIER TRANSFORM

## 1

We consider the linear map in  $\mathbb{C}^N$  which acts on the computational basis as

$$|j\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{\frac{2i\pi jk}{N}} |k\rangle$$

Let  $A$  be the matrix of the transformation in the computational basis.

$$\forall (k, l) \in \llbracket 0, N-1 \rrbracket^2, \quad a_{kl} = \frac{1}{\sqrt{N}} e^{\frac{2i\pi kl}{N}}$$

The adjoint matrix  $A^\dagger$  is then

$$\begin{aligned} \forall (k, l) \in \llbracket 0, N-1 \rrbracket^2, \quad b_{kl} &= a_{lk}^* \\ &= \frac{1}{\sqrt{N}} e^{-\frac{2i\pi kl}{N}} \end{aligned}$$

We compute the coefficient  $k, l$  of the product  $AA^\dagger$ :

$$\begin{aligned} \forall (k, l) \in \llbracket 0, N-1 \rrbracket^2, \quad c_{kl} &= \sum_{j=0}^{N-1} a_{kj} b_{jl} \\ &= \frac{1}{N} \sum_{j=0}^{N-1} e^{\frac{2i\pi j}{N} (k-l)} \\ &= \frac{1}{N} \sum_{j=0}^{N-1} (e^{\frac{2i\pi}{N} (k-l)})^j \\ &= \begin{cases} \frac{1}{N} \frac{1 - (e^{\frac{2i\pi}{N} (k-l)})^N}{1 - e^{\frac{2i\pi}{N} (k-l)}} = 0 & \text{if } e^{\frac{2i\pi}{N} (k-l)} \neq 1, \\ 1 & \text{if } e^{\frac{2i\pi}{N} (k-l)} = 1. \end{cases} \\ &= \begin{cases} 0 & \text{if } k \neq l, \\ 1 & \text{if } k = l. \end{cases} \\ &= \delta_{kl} \end{aligned}$$

which shows that  $AA^\dagger = A^\dagger A = I$  i.e.  $A$  is unitary.

## 2

Here the dimension of the state space is  $N = 2^n$ . The Fourier transform of the  $n$  qubit state  $|00 \dots 0\rangle$  is

$$A|0\rangle = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} |k\rangle$$

we can write  $k$  in binary  $k_{n-1} \dots k_1 k_0$

$$A|0\rangle = \frac{1}{2^{n/2}} \sum_{k_0, k_1, \dots, k_{n-1}=0}^1 |k_{n-1} \dots k_1 k_0\rangle$$

or in product representation,

$$= \frac{1}{2^{n/2}} \underbrace{(|0\rangle + |1\rangle)(|0\rangle + |1\rangle) \dots (|0\rangle + |1\rangle)}_{n \text{ qubits}}$$

## 3

Let  $N = 2^n$  and  $Y = (y_k)_{k \in \llbracket 0, N-1 \rrbracket}$  be the classical fourier transform of  $X = (x_k)_{k \in \llbracket 0, N-1 \rrbracket}$ .

$$\forall k \in \llbracket 0, N-1 \rrbracket, \quad y_k = \sum_{j=0}^{N-1} e^{\frac{2i\pi k j}{2^n}} x_j$$

The factor  $\frac{1}{\sqrt{N}}$  is omitted for clarity. We can write  $j$  in binary  $j_{n-1} \dots j_1 j_0$

$$\begin{aligned} y_k &= \sum_{j_0, j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_1 + j_0)}{2^n}} x_j \\ &= \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_1)}{2^n}} x_{j_{n-1} \dots j_1 0} + \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_1 + 1)}{2^n}} x_{j_{n-1} \dots j_1 1} \\ &= \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_1)}{2^n}} x_{j_{n-1} \dots j_1 0} + e^{\frac{2i\pi k}{2^n}} \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_1)}{2^n}} x_{j_{n-1} \dots j_1 1} \\ &= \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-2}j_{n-1} + \dots + j_1)}{2^{n-1}}} x_{j_{n-1} \dots j_1 0} + e^{\frac{2i\pi k}{2^n}} \sum_{j_1, \dots, j_{n-1}=0}^1 e^{\frac{2i\pi k(2^{n-2}j_{n-1} + \dots + j_1)}{2^{n-1}}} x_{j_{n-1} \dots j_1 1} \end{aligned}$$

We see the first sum is the  $k^{th}$  coefficient of the FT of the sequence  $(x_{2k})_{k \in \llbracket 0, N/2-1 \rrbracket}$  and the second is the  $k^{th}$  coefficient of the FT of  $(x_{2k+1})_{k \in \llbracket 0, N/2-1 \rrbracket}$ . This shows that to compute FT of sequence of length  $N$ , we have to compute 2 FT of sequence of length  $\frac{N}{2}$  and do  $2N$  complex additions/multiplications. The complexity of the operation  $T(N)$  follows the recurrence:

$$T(N) = 2T\left(\frac{N}{2}\right) + 2N$$

We can use the Master theorem <sup>1</sup>:

**Theorem.** Let  $a \geq 1$  and  $b > 1$  be constants, let  $f(n)$  be a function, and let  $T(n)$  be defined on the non negative integers by the recurrence

$$T(n) = aT\left(\frac{n}{b}\right) + f(n)$$

where we interpret  $\frac{n}{b}$  to mean either  $\lfloor \frac{n}{b} \rfloor$  or  $\lceil \frac{n}{b} \rceil$ . Then  $T(n)$  has the following asymptotic bounds:

- (1) If  $f(n) = O(n^{\log_b a - \epsilon})$  for some constant  $\epsilon > 0$ , then  $T(n) = \Theta(n^{\log_b a})$ .
- (2) If  $f(n) = \Theta(n^{\log_b a})$ , then  $T(n) = \Theta(n^{\log_b a} \log n)$ .
- (3) If  $f(n) = \Omega(n^{\log_b a + \epsilon})$  for some constant  $\epsilon > 0$ , and if  $af(\frac{n}{b}) \leq cf(n)$  for some constant  $c < 1$  and  $n$  sufficiently large, then  $T(n) = \Theta(f(n))$ .

Here we are in the second case of the theorem, so  $T(N) = \Theta(N \log(N)) = \Theta(n2^n)$ .

## 5

The inverse Fourier Transform

$$|j\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{-\frac{2i\pi j k}{N}} |k\rangle$$

is the adjoint of the Fourier Transform. The quantum circuit of figure 1 is obtained from the FT's circuit, replacing each  $R_k$  gate by its adjoint

$$R_k^\dagger = \begin{bmatrix} 1 & 0 \\ 0 & e^{-\frac{2i\pi}{2^k}} \end{bmatrix}$$

<sup>1</sup>Thomas H. Cormen and Charles E. Leiserson : *Introduction to algorithms*, MIT Press (2009)



FIGURE 1. Quantum circuit for IFT.



FIGURE 2. Sequence of controlled U.

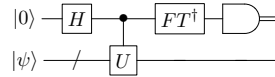
## 7

In figure 2, the  $t$  qubits of the first register are prepared with  $|j\rangle = |j_{t-1} \dots j_1 j_0\rangle$ , the second register is prepared with some state  $|u\rangle$ . After the first controlled-U operation, the state is  $|j\rangle |U^{j_0 2^0} u\rangle$ . After the second controlled-U, the state is  $|j\rangle |U^{j_1 2^1} U^{j_0 2^0} u\rangle = |j\rangle |U^{j_0 2^0 + j_1 2^1} u\rangle$  and so on. The final state is  $|j\rangle |U^{j_0 2^0 + j_1 2^1 + \dots + j_{t-1} 2^{t-1}} u\rangle = |j\rangle |U^j u\rangle$ .

## 8

By linearity, the phase estimation algorithm takes input  $|0\rangle |\sum_{u \in A} c_u |u\rangle\rangle$ , where  $A$  is some orthonormal basis of eigenstates of  $U$ , to output  $\sum_{u \in A} c_u |\widetilde{\varphi}_u\rangle |u\rangle$ , where  $\widetilde{\varphi}_u$  is an estimation of the phase of the eigenvalue associated with eigenstate  $u$ . If we fix  $u_0 \in A$  beforehand, the probability to measure  $\widetilde{\varphi}_{u_0}$  when measuring the first register in the computational basis is

$$\begin{aligned}
 & \left( \sum_{u \in A} c_u^* \langle \widetilde{\varphi}_u | \langle u | \right) P_{\widetilde{\varphi}_{u_0}} \otimes I \left( \sum_{u \in A} c_u |\widetilde{\varphi}_u\rangle |u\rangle \right) = \left( \sum_{u \in A} c_u^* \langle \widetilde{\varphi}_u | \langle u | \right) \left( \sum_{\substack{u \in A \\ \widetilde{\varphi}_u = \widetilde{\varphi}_{u_0}}} c_u |\widetilde{\varphi}_u\rangle |u\rangle \right) \\
 & = \left( \sum_{u \in A} c_u^* \langle \widetilde{\varphi}_u | \langle u | \right) \left( \sum_{\substack{u \in A \\ \widetilde{\varphi}_u = \widetilde{\varphi}_{u_0}}} c_u |\widetilde{\varphi}_{u_0}\rangle |u\rangle \right) \\
 & = \sum_{\substack{v \in A \\ u \in A \\ \widetilde{\varphi}_u = \widetilde{\varphi}_{u_0}}} c_v^* c_u \langle \widetilde{\varphi}_v | \widetilde{\varphi}_u \rangle \langle v | u \rangle \\
 & = \sum_{\substack{v \in A \\ u \in A \\ \widetilde{\varphi}_u = \widetilde{\varphi}_{u_0}}} c_v^* c_u \langle \widetilde{\varphi}_v | \widetilde{\varphi}_u \rangle \delta_{vu} \\
 & = \sum_{\substack{u \in A \\ \widetilde{\varphi}_u = \widetilde{\varphi}_{u_0}}} |c_u|^2 \\
 & \geq |c_{u_0}|^2
 \end{aligned}$$

FIGURE 3. Phase estimation circuit with  $t = 1$ .

$I$  is the identity operator of whatever state space  $U$  operates on, while  $P_{\widetilde{\varphi_{u_0}}}$  is the orthonormal projector onto the space generated by the vector  $|\widetilde{\varphi_{u_0}}\rangle$  of the computational basis. Besides, following the analysis of the book,  $\widetilde{\varphi_{u_0}}$  is an approximation to  $\varphi_{u_0}$  to an accuracy  $2^{-n}$  with probability at least  $1 - \epsilon$  if we make use of  $t = n + \lceil \log(2 + \frac{1}{2\epsilon}) \rceil$  bits in the first register. We conclude we get the desired approximation of  $\varphi_{u_0}$  at the end of the phase estimation algorithm with probability at least  $|c_{u_0}|^2(1 - \epsilon)$ .

## 9

$U$  being unitary with eigenvalues  $-1$  and  $+1$ , the state space is the direct sum of the two eigenspaces  $E_{-1} \oplus E_1$ . Thus we can uniquely decompose any  $|\psi\rangle = |\psi_{-1}\rangle + |\psi_{+1}\rangle$ , with  $|\psi_{-1}\rangle \in E_{-1}$  and  $|\psi_{+1}\rangle \in E_1$ . Then  $-1 = e^{i\pi} = e^{2i\pi 0.1}$  and  $-1 = e^0 = e^{2i\pi 0.0}$  shows that it is sufficient to make use of  $t = 1$  wire in the first register in the phase estimation procedure to read directly the phase of any eigenvector. If we use  $|0\rangle |\psi\rangle$  as input in the circuit of figure 3, the output before the final measurement will be  $|0\rangle |\psi_{+1}\rangle + |1\rangle |\psi_{-1}\rangle$ . When we measure the first register, the state will collapse into either  $\frac{1}{\langle \psi_{+1} | \psi_{+1} \rangle} |0\rangle |\psi_{+1}\rangle$  or  $\frac{1}{\langle \psi_{-1} | \psi_{-1} \rangle} |1\rangle |\psi_{-1}\rangle$ , with probability respectively  $\langle \psi_{+1} | \psi_{+1} \rangle$  and  $\langle \psi_{-1} | \psi_{-1} \rangle$ . Thus if we read 0 in the first register, that means that we have an eigenvector associated to eigenvalue  $+1$  in the second register, and if we read 1 in the first register, that means that we have an eigenvector associated to eigenvalue  $-1$  in the second register.

Once we have noticed that the FT in dimension  $N = 2^1$  is just the Hadamard operator, we conclude the phase estimation circuit in this particular case is just the same as the circuit of exercise 4.34.