QUANTUM COMPUTATION AND QUANTUM INFORMATION: THE QUANTUM FOURIER TRANSFORM

1

We consider the linear map in \mathbb{C}^N which acts on the computational basis as

$$|j\rangle\mapsto \frac{1}{\sqrt{N}}\sum_{k=0}^{N-1}e^{\frac{2i\pi jk}{N}}\,|k\rangle$$

Let A be the matrix of the transformation in the computational basis.

$$\forall (k,l) \in [0, N-1]^2, \quad a_{kl} = \frac{1}{\sqrt{N}} e^{\frac{2i\pi kl}{N}}$$

The adjoint matrix A^{\dagger} is then

$$\forall (k,l) \in [0, N-1]^2, \quad b_{kl} = a_{lk}^*$$

$$= \frac{1}{\sqrt{N}} e^{-\frac{2i\pi kl}{N}}$$

We compute the coefficient k, l of the product AA^{\dagger} :

$$\begin{split} \forall (k,l) \in [\![0,N-1]\!]^2, \quad c_{kl} &= \sum_{j=0}^{N-1} a_{kj} b_{jl} \\ &= \frac{1}{N} \sum_{j=0}^{N-1} e^{\frac{2i\pi j}{N}(k-l)} \\ &= \frac{1}{N} \sum_{j=0}^{N-1} (e^{\frac{2i\pi}{N}(k-l)})^j \\ &= \begin{cases} \frac{1}{N} \frac{1 - (e^{\frac{2i\pi}{N}(k-l)})^N}{1 - e^{\frac{2i\pi}{N}(k-l)}} = 0 & \text{if } e^{\frac{2i\pi}{N}(k-l)} \neq 1, \\ 1 & \text{if } e^{\frac{2i\pi}{N}(k-l)} = 1. \end{cases} \\ &= \begin{cases} 0 & \text{if } k \neq l, \\ 1 & \text{if } k = l. \end{cases} \\ &= \delta_{kl} \end{split}$$

which shows that $AA^{\dagger} = A^{\dagger}A = I$ i.e. A is unitary.

2

Here the dimension of the state space is $N=2^n$. The Fourier transform of the n qubit state $|00...0\rangle$ is

$$A|0\rangle = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} |k\rangle$$

we can write k in binary $k_{n-1} \dots k_1 k_0$

$$A|0\rangle = \frac{1}{2^{n/2}} \sum_{k_0, k_1, \dots, k_{n-1} = 0}^{1} |k_{n-1} \dots k_1 k_0\rangle$$

or in product representation,

$$= \frac{1}{2^{n/2}} \underbrace{(|0\rangle + |1\rangle)(|0\rangle + |1\rangle) \dots (|0\rangle + |1\rangle)}_{\substack{n \text{ qubits}}}$$

Let $N=2^n$ and $Y=(y_k)_{k\in \llbracket 0,N-1\rrbracket}$ be the classical fourier transform of $X=(x_k)_{k\in \llbracket 0,N-1\rrbracket}$.

$$\forall k \in [0, N-1], \quad y_k = \sum_{j=0}^{N-1} e^{\frac{2i\pi kj}{2^n}} x_j$$

The factor $\frac{1}{\sqrt{N}}$ is omitted for clarity. We can write j in binary $j_{n-1} \dots j_1 j_0$

$$y_{k} = \sum_{j_{0}, j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_{1} + j_{0})}{2^{n}}} x_{j}$$

$$= \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_{1})}{2^{n}}} x_{j_{n-1} \dots j_{1}0} + \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_{1} + 1)}{2^{n}}} x_{j_{n-1} \dots j_{1}1}$$

$$= \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_{1})}{2^{n}}} x_{j_{n-1} \dots j_{1}0} + e^{\frac{2i\pi k}{2^{n}}} \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-1}j_{n-1} + \dots + 2j_{1})}{2^{n}}} x_{j_{n-1} \dots j_{1}1}$$

$$= \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-2}j_{n-1} + \dots + j_{1})}{2^{n-1}}} x_{j_{n-1} \dots j_{1}0} + e^{\frac{2i\pi k}{2^{n}}} \sum_{j_{1}, \dots, j_{n-1} = 0}^{1} e^{\frac{2i\pi k(2^{n-2}j_{n-1} + \dots + j_{1})}{2^{n-1}}} x_{j_{n-1} \dots j_{1}1}$$

We see the first sum is the k^{th} coefficient of the FT of the sequence $(x_{2k})_{k \in [0,N/2-1]}$ and the second is the k^{th} coefficient of the FT of $(x_{2k+1})_{k \in [0,N/2-1]}$. This shows that to compute FT of sequence of length N, we have to compute 2 FT of sequence of length $\frac{N}{2}$ and do 2N complex additions/multiplications. The complexity of the operation T(N) follows the recurrence:

$$T(N) = 2T(\frac{N}{2}) + 2N$$

We can use the Master theorem ¹:

Theorem. Let $a \ge 1$ and b > 1 be constants, let f(n) be a function, and let T(n) be defined on the non negative integers by the recurrence

$$T(n) = aT(\frac{n}{h}) + f(n)$$

where we interpret $\frac{n}{h}$ to mean either $\lfloor \frac{n}{h} \rfloor$ or $\lceil \frac{n}{h} \rceil$. Then T(n) has the following asymptotic bounds:

- (1) If $f(n) = O(n^{\log_b a \epsilon})$ for some constant $\epsilon > 0$, then $T(n) = \Theta(n^{\log_b a})$.
- (2) If $f(n) = \Theta(n^{\log_b a})$, then $T(n) = \Theta(n^{\log_b a} \log n)$.
- (3) If $f(n) = \Omega(n^{\log_b a + \epsilon})$ for some constant $\epsilon > 0$, and if $af(\frac{n}{b}) \leqslant cf(n)$ for some constant c < 1 and n sufficiently large, then $T(n) = \Theta(f(n))$.

Here we are in the second case of the theorem, so $T(N) = \Theta(N \log(N)) = \Theta(n2^n)$.

5

The inverse Fourier Transform

$$|j\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{-\frac{2i\pi jk}{N}} |k\rangle$$

is the adjoint of the Fourier Transform. The quantum circuit of figure 1 is obtained from the FT's circuit, replacing each R_k gate by its adjoint

$$R_k^{\dagger} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-\frac{2i\pi}{2^k}} \end{bmatrix}$$

¹Thomas H. Cormen and Charles E. Leiserson: Introduction to algorithms, MIT Press (2009)

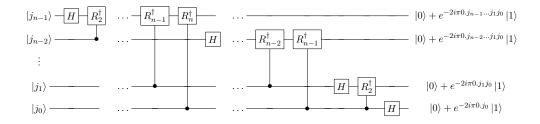


FIGURE 1. Quantum circuit for IFT.

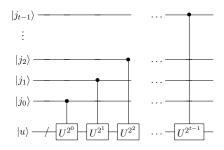


FIGURE 2. Sequence of controlled U.

7

In figure 2, the t qubits of the first register are prepared with $|j\rangle = |j_{t-1}\dots j_1j_0\rangle$, the second register is prepared with some state $|u\rangle$. After the first controlled-U operation, the state is $|j\rangle |U^{j_02^0}u\rangle$. After the second controlled-U, the state is $|j\rangle |U^{j_02^0+j_12^1}u\rangle$ and so on. The final state is $|j\rangle |U^{j_02^0+j_12^1+\dots+j_{t-1}2^{t-1}}u\rangle = |j\rangle |U^ju\rangle$.

8

The phase estimation algorithm takes input $|0\rangle |\Sigma_{u\in A}c_u|u\rangle\rangle$, where A is some orthonormal basis of eigenstates of U, to outut $\sum_{u\in A}c_u|\widetilde{\varphi_u}\rangle |u\rangle\rangle$, where $\widetilde{\varphi_u}$ is an estimation of the phase of the eigenvalue associated with eigenstate u. If we fix $u_0 \in A$, the probability to measure $\widetilde{\varphi_{u_0}}$ when measuring the first register in the computational basis is

The basis is
$$(\sum_{u \in A} c_u^* \langle \widetilde{\varphi_u} | \langle u |) P_{\widetilde{\varphi_{u_0}}} \otimes I(\sum_{u \in A} c_u | \widetilde{\varphi_u} \rangle | u \rangle) = (\sum_{u \in A} c_u^* \langle \widetilde{\varphi_u} | \langle u |) (\sum_{\substack{u \in A \\ \widetilde{\varphi_u} = \widetilde{\varphi_{u_0}}}} c_u | \widetilde{\varphi_u} \rangle | u \rangle)$$

$$= (\sum_{u \in A} c_u^* \langle \widetilde{\varphi_u} | \langle u |) (\sum_{\substack{u \in A \\ \widetilde{\varphi_u} = \widetilde{\varphi_{u_0}}}} c_u | \widetilde{\varphi_{u_0}} \rangle | u \rangle)$$

$$= \sum_{\substack{v \in A \\ u \in A \\ \widetilde{\varphi_u} = \widetilde{\varphi_{u_0}}}} c_v^* c_u \langle \widetilde{\varphi_v} | \widetilde{\varphi_u} \rangle \langle v | u \rangle$$

$$= \sum_{\substack{u \in A \\ \widetilde{\varphi_u} = \widetilde{\varphi_{u_0}}}} |c_u|^2$$

$$\geqslant |c_{u_0}|^2$$

Besides, following the analysis of the book, $\widetilde{\varphi_{u_0}}$ is an approximate to φ_{u_0} to an accuracy 2^{-n} with probability at least $1 - \epsilon$ if we make use of $t = n + \log(2 + \frac{1}{2\epsilon})$ bits in the first register. We conclude that getting the desired approximation at the end of the phase estimation algorithm is at least $|c_{u_0}|^2(1 - \epsilon)$.