

NOTES ON LADDER OPERATORS

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ABSTRACT. We review ladder operators.

1. COMMUTATION RELATIONS

We have the operators \hat{X} and \hat{N} such that

$$(1.1) \quad [\hat{N}, \hat{X}] = c\hat{X}$$

where c is a scalar value. We have the eigenstates of \hat{N} such that

$$(1.2) \quad \hat{N}|n\rangle = n|n\rangle$$

where we have abused notation letting $|n\rangle$ be a vector and n be a scalar (the eigenvalue of the $|n\rangle$ vector). We see that

$$(1.3a) \quad \hat{N}\hat{X}|n\rangle = (\hat{X}\hat{N} + [\hat{N}, \hat{X}])|n\rangle$$

$$(1.3b) \quad = (\hat{X}\hat{N} + c\hat{X})|n\rangle$$

$$(1.3c) \quad = \hat{X}\hat{N}|n\rangle + c\hat{X}|n\rangle$$

$$(1.3d) \quad = \hat{X}n|n\rangle + c\hat{X}|n\rangle$$

$$(1.3e) \quad = (n + c)\hat{X}|n\rangle.$$

This means that if $|n\rangle$ is an eigenstate of \hat{N} with eigenvalue n , then $(\hat{X}|n\rangle)$ is an eigenstate of \hat{N} with eigenvalue $n + c$. If $c > 0$ then $n + c > n$, so the operator \hat{X} is called the “**Creation Operator**”.

If \hat{N} is a physical observable, it's necessarily self adjoint. This implies that c is real, since n is real and $n + c$ is an eigenvalue of a self-adjoint operator. We should remember from linear algebra the eigenvalues of a self-adjoint operator is always real. The Hermitian adjoint of \hat{X} satisfies

$$(1.4) \quad [\hat{N}, \hat{X}^\dagger] = -c\hat{X}^\dagger$$

We call \hat{X}^\dagger an “**Annihilation Operator**” if \hat{X} is a creation operator.

Proposition 1. Let \hat{A} , \hat{B} , \hat{C} be operators. Then

$$(1.5) \quad [\hat{A}, \hat{B}\hat{C}] = [\hat{A}, \hat{B}]\hat{C} + \hat{B}[\hat{A}, \hat{C}].$$

Proof. We see by direct computation

$$(1.6a) \quad [\hat{A}, \hat{B}\hat{C}] = \hat{A}\hat{B}\hat{C} - \hat{B}\hat{C}\hat{A}$$

$$(1.6b) \quad \hat{B}[\hat{A}, \hat{C}] = \hat{B}(\hat{A}\hat{C} - \hat{C}\hat{A})$$

$$(1.6c) \quad [\hat{A}, \hat{B}]\hat{C} = (\hat{A}\hat{B} - \hat{B}\hat{A})\hat{C}$$

so when we add equation (1.6b) to equation (1.6c) we get

$$(1.7) \quad \hat{A}\hat{B}\hat{C} - \hat{B}\hat{C}\hat{A} = [\hat{A}, \hat{B}\hat{C}]$$

as desired. \square

Proposition 2. Let \hat{A}, \hat{B} be operators. If their commutator vanishes

$$(1.8) \quad [\hat{A}, \hat{B}] = 0$$

then the two operators are equal up to a constant.

Proof. Trivial. \square

Proposition 3. Given these two propositions and the commutations relations, we can find that

$$(1.9) \quad [\hat{N}, \hat{X}\hat{X}^\dagger] = 0.$$

Or equivalently

$$(1.10) \quad \hat{N} = \hat{X}^\dagger \hat{X}$$

up to some constant.

Proof. How to prove this? By direct computation

$$(1.11a) \quad [\hat{N}, \hat{X}^\dagger \hat{X}] = [\hat{N}, \hat{X}^\dagger] \hat{X} + \hat{X}^\dagger [\hat{N}, \hat{X}]$$

$$(1.11b) \quad = (-c\hat{X}^\dagger) \hat{X} + \hat{X}^\dagger (c\hat{X})$$

$$(1.11c) \quad = (c - c) \hat{X}^\dagger \hat{X}$$

$$(1.11d) \quad = 0.$$

This means, up to some constant and ordering,

$$(1.12) \quad \hat{N} = \hat{X}^\dagger \hat{X}$$

just as desired. \square

REFERENCES

- [1] P. A. M. Dirac, *Principles of Quantum Mechanics*. Oxford Science Publications, fourth ed., 2007.
- [2] O. L. de Lange and R. E. Raab, “Ladder Operators for Orbital Angular Momentum,” *American Journal of Physics* **54** (1986) no. 4, 372–375.
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