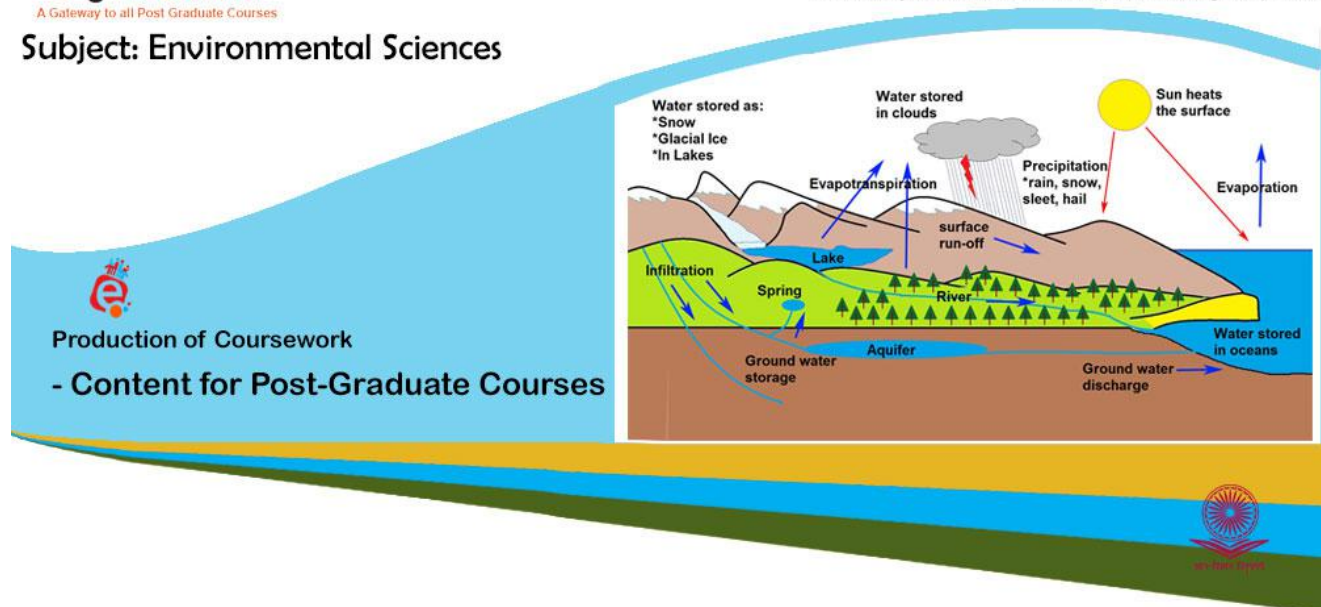


**Subject: Environmental Sciences**



**Paper Name : 5 Water Resources and Management**

**Module : 17 Groundwater Hydrology IV (Coupled Flow and Transport)**



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Description of Module	
<b>Subject Name</b>	<b>Environmental Sciences</b>
<b>Paper Name</b>	Water Resources and Management
<b>Module Name/Title</b>	Groundwater Hydrology IV (Coupled flow and transport)
<b>Module Id</b>	EVS/WRM-V/17
<b>Pre-requisites</b>	
<b>Objectives</b>	We will get introduced to fundamentals of the flow and transport equations or models. We learn about different transport processes and develop a transport model. Finally we will analyze few transport problems.
<b>Keywords</b>	Flow Model and its Derivation, Advection, Dispersion, Diffusion, Transport Model and its Derivation, Analytical Solutions of Few Transport Problems

### Objectives:

This module will present an introductory topics relevant to flow and transport problems in groundwater. Emphasis will be on d processes, coupling them and eventually establishing the mathematical formulation of the problem. The target groups are higher level undergraduate students and the first year PG students. The module, as it is introductory, will only briefly introduce methods for solving transport problems. The specific objectives of the module are:

1. Derivation of the flow model
2. Recognizing factors and processes affecting transport problems
3. Derivation of transport models
4. Developing a systematic approach to solve transport problems.

## 17. Introduction

In this module we will get introduced to fundamentals of the flow and transport equations or models. To distinguish, flow deals with the quantity aspect of groundwater whereas transport deals with the quality. Hence the flow problem is governed by the energy gradient where as mass gradient is significant in the transport problem. The complexities in both problems arises from the properties of the porous media, such as porosity, conductivity etc, which are mostly varying in space (heterogeneity versus homogeneity) and direction (isotropic versus anisotropic). In this module we will first introduce the flow model and subsequently learn about different transport processes and develop a transport model. Emphasising on the transport problem we will introduce methods to solve it. Our approach will remain mostly mathematical as such we will attempt to solve and visualize few examples of transport problems. In this module we will focus on saturated (all voids filled), homogeneous (aquifer properties such as conductivity do not change along the extend) and isotropic (aquifer properties is uniform also in all directions) aquifers.

### 17.1 Flow Model

#### 17.1.1 The Representative Control Volume (RCV)

Aquifers as we know already are extensive, at least in the horizontal and lateral directions, whereas its properties such as conductivity, transmissivity often changes in a very small scale, in many cases also at a pore-scale. As such to understand aquifer properties and underlying processes, we define a Representative Control Volume (RCV, see Fig. 1). Ideally a RCV is very much smaller than the aquifer that is investigated, and at the same it is very much larger than individual grain or pore. The essential condition is that Darcy's law has to be applicable in the RCV. A big advantage of RCV is that in the aquifer processes analysis it can be oriented as per the coordinate system, e.g., rectangular for Cartesian coordinate.



FIG. 1: THE REPRESENTATIVE CONTROL VOLUME (RCV)

### 17.1.2 The derivation of flow model

The groundwater flow problem (and hence the model) can be classified according to: the problem dimensionality (*1D*, *2D*, or *3D*), the time-dependent state (steady and transient), the variation in aquifer properties with respect to space (homogeneous and heterogeneous), the variation in aquifers properties with respect to direction (isotropic and anisotropic), the aquifer condition (confined and unconfined) and based on weather aquifer is with or without source and sink. Each of the above combination corresponds to a certain equation governing groundwater flow. The derivation of the equation requires two fundamental principles: the conservation of volume and the Darcy's law, which is given as:

$$v_f = -k \text{ grad } h = -k \frac{\partial h}{\partial x} \quad (1)$$

With  $q = Q/A$  i.e., the discharge  $Q$  [ $L^3T^{-1}$ ] and Area  $A$  [ $L^2$ ].  $v_f$  is often referred to as Darcy's velocity or specific discharge.  $K$  [ $LT^{-1}$ ] refers to hydraulic conductivity and  $\text{grad}h$  [-] is gradient vector of hydraulic head  $h$  [ $L$ ]. The negative sign is indicator of the flow from larger to smaller head. With Darcy's law already defined, we now develop a groundwater flow model for several scenarios. We begin with the derivation of *3D* groundwater model.

We consider a unit 3D RCV (Fig. 2) in which  $v_{fx}$ ,  $v_{fy}$  and  $v_{fz}$  represents the specific discharge entering the RCV from the left face of the RCV. Since the RCV is very small, the change of the specific discharge across the RCV (at the right face) can be regarded as linear, i.e., we add the product of first-order derivative and the corresponding distance between faces. The volume budget is given by:

$$\frac{\Delta V_w}{\Delta t} = Q_{in} - Q_{out} \quad (2)$$

Where  $\Delta V_w [L^3]$  is change in water volume over the time  $\Delta t [T]$ . The difference  $Q_{in} - Q_{out}$  represents the difference of discharge in the two faces of the RCV. Thus total  $Q_{in}$  in RCV is

$$Q_{in} = v_{fx}\Delta y\Delta z + v_{fy}\Delta x\Delta z + v_{fz}\Delta x\Delta y \quad (3)$$

and the corresponding  $Q_{out}$  is

$$Q_{out} = \left( v_{fx} + \frac{\partial v_{fx}}{\partial x} \right) \Delta y \Delta z + \left( v_{fy} + \frac{\partial v_{fy}}{\partial y} \right) \Delta x \Delta z + \left( v_{fz} + \frac{\partial v_{fz}}{\partial z} \right) \Delta x \Delta y \quad (4)$$

Thus  $Q_{in} - Q_{out}$  becomes

$$Q_{in} - Q_{out} = -\frac{\partial v_{fx}}{\partial x} \Delta x \Delta y \Delta z - \frac{\partial v_{fy}}{\partial y} \Delta x \Delta y \Delta z - \frac{\partial v_{fz}}{\partial z} \Delta x \Delta y \Delta z \quad (5)$$

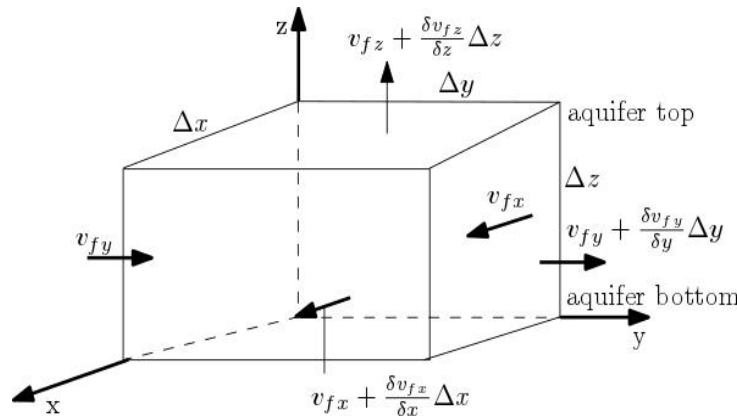


FIG. 2: A 3D RCV UNIT VOLUME FOR THR FLOW MODEL

Next we find the relationship between the change in head and the corresponding change in water volume. This relation is expressed as

$$\frac{\Delta V_w}{\Delta x \Delta y \Delta z} \propto \Delta h \quad (6)$$

The expression equates to

$$\Delta V_w = S_s \Delta h \Delta x \Delta y \Delta z \quad (7)$$

With  $S_s$  [ $L^{-1}$ ] the specific storage coefficient of the volume. Substituting expressions for  $Q_{in} - Q_{out}$  and  $\Delta V_w$  in eq (2), and upon cancelling the common terms from both side of equality we get

$$S_s \frac{\Delta h}{\Delta t} = -\frac{\partial v_{fx}}{\partial x} - \frac{\partial v_{fy}}{\partial y} - \frac{\partial v_{fz}}{\partial z} \quad (8)$$

Letting  $\Delta t \rightarrow 0$ , eq. (8) becomes

$$S_s \frac{\partial h}{\partial t} = -\frac{\partial v_{fx}}{\partial x} - \frac{\partial v_{fy}}{\partial y} - \frac{\partial v_{fz}}{\partial z} \quad (9)$$

Now when we insert Darcy's velocity expression (eq. 1) for all corresponding co-ordinate directions in eq. (9), we get the general groundwater flow model

$$S_s \frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left( K \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial y} \left( K \frac{\partial h}{\partial y} \right) + \frac{\partial}{\partial z} \left( K \frac{\partial h}{\partial z} \right) \quad (10)$$

When sources/sinks are part of the flow system, eq. (10) takes the form:

$$S_s \frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left( K \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial y} \left( K \frac{\partial h}{\partial y} \right) + \frac{\partial}{\partial z} \left( K \frac{\partial h}{\partial z} \right) + q \quad (11)$$

Where  $q$  [ $T^{-1}$ ] represents the volumetric rate of the source sink per unit aquifer volume. Few examples of  $q$  can be water injection/extraction through wells, water transferred to aquifer from/to rivers/streams. For anisotropic case the conductivity ( $K$ ) becomes the function of spatial direction and the eq. (11) takes the following form:

$$S_s \frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left( K_x \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial y} \left( K_y \frac{\partial h}{\partial y} \right) + \frac{\partial}{\partial z} \left( K_z \frac{\partial h}{\partial z} \right) + q \quad (12)$$

For a steady-state, homogeneous, isotropic, confined with source/sink term the groundwater flow equation is

$$\frac{\partial^2 h}{\partial x^2} + \frac{\partial^2 h}{\partial y^2} + \frac{\partial^2 h}{\partial z^2} = -\frac{q}{K} \quad (13)$$

Eq. (13) is referred to as **Poisson equation**. In a case with steady-state flow in homogeneous isotropic confined aquifer without the sources/sink term, we get the well known **Laplace equation** for the groundwater flow, which is of the form

$$\frac{\partial^2 h}{\partial x^2} + \frac{\partial^2 h}{\partial y^2} + \frac{\partial^2 h}{\partial z^2} = 0 \quad (14)$$

For the 2D case the flow equation varies between the confined and the unconfined aquifer. In a confined aquifer the step from three to two dimension conversion requires to sum-up  $K$  values over the entire layer thickness. This is achieved using an aquifer property called transmissivity,  $T$  [ $L^2T^{-1}$ ], which is defined as

$$T_x = K_x \cdot m \quad T_y = K_y \cdot m \quad (15)$$

Where  $K_x$  and  $K_y$  are vertically averaged hydraulic conductivities along  $x$  and  $y$  coordinates, respectively, and  $m$  [L] the thickness of the aquifer. For a confined isotropic aquifer  $T = K m$ . Furthermore, if vertical flow can be neglected we can define storage coefficient,  $S$  [-], as

$$S = S_s \cdot m \quad (16)$$

With  $T$  and  $S$  defined, the 2D groundwater flow equation for a confined anisotropic aquifer with source/sink is

$$S \frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left( T_x \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial y} \left( T_y \frac{\partial h}{\partial y} \right) + q \quad (17)$$

With source/sink term included, eq. (17) is also termed as **Boussinesq equation**. Further, eq. (17) is also valid for the unconfined aquifer, in which the top of the aquifer is the groundwater table, provided



that Dupuit's assumptions, that groundwater flow is horizontal, the flow velocity do not vary along the aquifer thickness, and that Darcy's law (eq. 1) is valid at water table, are observed.

### 17.1.3 Complete formulation of flow model

The complete formulation of the flow model requires that we define the following in addition to an appropriate flow equation (eqs. 10 - 14 and 17):

1. Specify the geometric properties of the region of interest
2. Specify values of aquifer parameters (hydraulic conductivity, storage coefficient) by considering spatial variability and anisotropy, if necessary
3. Specify the initial condition (IC): Head values at time  $t=0$  for the transient problems only.
4. Specify boundary conditions (BCs): Along the complete boundary, and which may be time-dependent. The three type of BCs used in groundwater flow problems are defined next.
  - 4.1 Boundary condition of the first kind or Dirichlet condition: The head value is given.
  - 4.2 Boundary condition of the second kind or Neumann condition: The component of the head gradient, which is perpendicular to the boundary, is provided.
  - 4.3 Boundary condition of the third kind or Cauchy condition or Robin boundary condition: A relationship between the head value and the component of the head gradient, which is perpendicular to the boundary, is given.

## 17.2 Transport Processes

In the transport model we consider mass (solute, chemical) that is flowing along with the groundwater. Before we develop the transport model, we will first briefly learn about the different transport processes that affect the concentration of mass in the aquifer.

### 17.2.1. Advection

Advection (also called convection) is the transport of mass with the movement of a moving medium. The moving medium can be termed carrier, and in the groundwater case it is usually fluid. To



quantify advection process, let us consider a laboratory column of a constant cross-sectional area,  $A$  [ $L^2$ ] and the steady-state discharge,  $Q$  [ $L^3T^{-1}$ ]. Further we will assume  $n_e$  [ ] the effective porosity of the column packing and  $v$  as average linear velocity of water flow. Note that  $v = q/n_e$ , where  $q$  is the Darcy velocity (eq. 1). With these information, we quantify the advective mass flow rate  $J_{adv}[MT^{-1}]$  as

$$J_{adv} = QC \quad (18)$$

We may want to modify eq. (18) by considering the continuity equation in the column

$$Q = n_e Av = \text{constant} \quad (19)$$

Substituting eq. (19) in eq. (18), and defining advective flux,  $j_{adv}[MT^{-1}L^2] = J_{adv}/A$ , we get

$$j_{adv} = \frac{J_{adv}}{A} = n_e v C \quad (20)$$

## 17.2.2 Mechanical Dispersion

The spread of mass causes dispersion. Spread of mass is not entirely a physical processes but physiochemical processes also causes spread. We refer mechanical dispersion to the spread due to physical processes only. Next, we quantify dispersive mass flow rate,  $J_{dis}[MT^{-1}]$ . Standard hydrogeology texts, such as *Domenico and Schwartz*, (1998), Chahar (2014), suggest that dispersion can be explained as a random phenomena. Based on observations, the dispersive mass flow due to mechanical dispersion in porous media is found to be

$$J_{dis} = -\alpha n_e A v \frac{\Delta C}{L} \quad (21)$$

with  $\alpha$  [L] a constant of proportionality called dispersivity,  $n_e$  [ - ] is effective porosity,  $A$  [ $L^2$ ] is surface area of the flow,  $v$  [ $LT^{-1}$ ] is the groundwater flow velocity,  $L$  [L] is the flow length and  $C$  [ $ML^{-3}$ ]. The negative sign is used in equality relation to identify that transport is from a higher concentration towards the lower one. Dispersivity is a property of porous media alone and is not dependent on fluid type or the flow characteristics. In a more formal manner the ratio  $\Delta C/L$  [ $ML^{-4}$ ] is termed as concentration gradient.

As mass transport is influenced by both porous medium and flow properties, a product of dispersivity and linear velocity called the mechanical dispersion coefficient,  $D_{mech}$  [ $L^2T^{-1}$ ], is more often used in analysing dispersive transport. This changes eq. (21) to

$$J_{dis} = -D_{mech} n_e A \frac{\Delta C}{L} \quad (22)$$

with  $D_{mech} = \alpha v$ . Instead of mass flow rate one may use mass flux  $j_{dis}$  [ $MT^{-1}L^2$ ], which is

$$j_{dis} = \frac{J_{dis}}{A} = -D_{mech} n_e \frac{\Delta C}{L} \quad (23)$$

Since  $\alpha$  characterizes spread length, its value is dependent on the flow coordinate direction under consideration. Advection and the mechanical dispersion are two processes that are generally used in analysing groundwater transport problems. Nevertheless, it is important for us to introduce a very important transport mechanism called diffusion that is not flow or mechanically driven but rather it is concentration gradient driven.

### 17.2.3 Diffusion

The quantification of diffusive mass transport,  $J_{dif}$  [ $MT^{-1}$ ] is based on the work in *Fick*, (1855). In that work the mass flow due to diffusion (1D) is described to be

$$J_{dif} = -D_{dif} A \frac{\Delta C}{L} \quad (24)$$

with  $A$  [ $L^2$ ] is surface area of the flow,  $L$  [ $L$ ] is the flow length and  $C$  [ $ML^{-3}$ ] and  $D_{dif}$  [ $L^2T^{-1}$ ], is the constant of proportionality called diffusion coefficient. The negative sign in eq. (24) represents transport of mass from higher concentrations toward the lower one. From eq. (24) we can get the diffusive flux  $j_{dif}$  [ $MT^{-1}L^2$ ] from

$$j_{dif} = \frac{J_{dif}}{A} = -D_{dif} \frac{\Delta C}{L} \quad (25)$$

In the analysis of groundwater transport problems the dispersion coefficient and the diffusion coefficient are summed and a new coefficient called hydrodynamic dispersion coefficient  $D_{hyd}$  is defined as in

$$D_{hyd} = D_{mech} + D_p = \alpha v + D_p \quad (26)$$

This summation is based on the fact that both dispersion and diffusion lead to spread of mass, and can be justified only on practical basis. It is to be noted that dispersion is physically driven process whereas difference is chemically driven process.

### 17.3. Joint Action of advection and Dispersion

The transport of conservative (non-reacting) solute in an aquifer can be understood as a superposition of advection, mechanical dispersion and pore diffusion. To quantify the transported mass over the time, we sum all the mass flow rate,  $J [MT^{-1}]$  that we have considered. i.e., eqs. (18, 22 and 24). Thus we get a combined equation for the mass flow

$$J = J_{adv} + J_{dis} + J_{dif} = J_{adv} + J_{dis,h} \quad (27)$$

Where  $J_{dis,h} = J_{dis} + J_{dif}$  refer to mass flow due to hydrodynamic dispersion. Eq (27) can alternatively written as

$$J = n_e A v C - n_e A v \frac{\Delta C}{L} - A D_p \frac{\Delta C}{L} = n_e A v C - n_e A D_{hyd} \frac{\Delta C}{L} \quad (28)$$

The spreading of mass due to advection and dispersion can be quantified by combining a mass budget and the corresponding laws of motion. This result in a transport equation called advection-dispersion or convection-dispersion equation, which we derive in the next section.

### 17.4. Derivation of transport model

Transport equations are derived by using a representative control volume (or RCV see Fig 1). For the transport problems (2D), the RCV extends from the aquifer bottom to the aquifer top in confined aquifers. Furthermore average concentration along the thickness is used. The fundamental transport equation is based on mass balance (similar to eq 2, volume replaced by mass) relation and laws of

motion used in advection and dispersion. For the derivation we will consider a 2D transport in the horizontal (xy-plane) with advection along the x-axis. The groundwater flow is steady, i.e.,  $v_x$  is constant over time. we will consider 2D dispersion. The mass balance equation in the RCV is then (see Fig. 3)

$$\frac{\Delta M}{\Delta t} = J_{in} - J_{out} \quad (29)$$

Where  $\Delta M$  is change in mass in the RCV over time  $\Delta t$ . For deriving the transport equation, we will have to find an expression for each term of eq. (29). Let  $n_e$  be the porosity of RCV then we have:

- I. The volume of the RCV,  $V = \Delta x \Delta y m$
- II. The dissolved mass,  $= n_e V C$
- III. Thus total mass,  $M$ , in RCV is:  $M = n_e V C$

in which  $n_e$  and  $V$  are fixed quantities, then the change of mass  $\Delta M$  over time  $\Delta t$  is given as

$$\Delta M = n_e V \Delta C = n_e \Delta x \Delta y m \Delta C \quad (30)$$

with  $\Delta C$  being the change in solute concentration. Since the RCV is very small compared to the actual study area, the components of mass flux ( $j = J/A$ ) can be assumed to change linearly across the extension of a RCV. These changes are obtained by multiplying first-order derivatives with corresponding distances, i.e., using Taylor series expansion (see Fig. 3) . We thus get

$$J_{in} = j_x m \Delta y + j_y m \Delta x \quad (31)$$

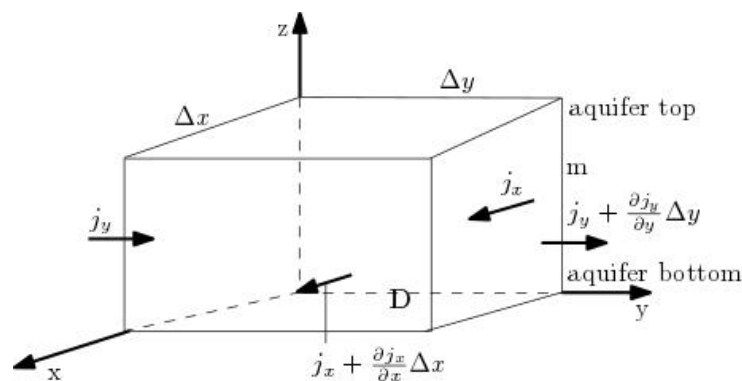


FIG. 3: A RCV WITH DIFFERENT MASS COMPONENTS

and

$$J_{out} = \left( j_x + \frac{\partial j_x}{\partial x} \Delta x \right) m \Delta y + \left( j_y + \frac{\partial j_y}{\partial y} \Delta y \right) m \Delta x \quad (32)$$

and the difference is

$$J_{in} - J_{out} = -\frac{\partial j_x}{\partial x} m \Delta x \Delta y - \frac{\partial j_y}{\partial y} m \Delta x \Delta y \quad (33)$$

Inserting the expressions for  $\Delta M$  (eq. 30) and  $(J_{in} - J_{out})$  (eq. 33) in eq. (29), we get

$$n_e m \Delta x \Delta y \frac{\Delta C}{\Delta t} = -\frac{\partial j_x}{\partial x} m \Delta x \Delta y - \frac{\partial j_y}{\partial y} m \Delta x \Delta y \quad (34)$$

Canceling common factors and letting  $\Delta t \rightarrow 0$  in eq (34) changes it to

$$n_e \frac{\partial C}{\partial t} = -\frac{\partial j_x}{\partial x} - \frac{\partial j_y}{\partial y} \quad (35)$$

Finally to complete our transport equation we will use laws of motion to replace fluxes ( $j_x$  and  $j_y$ ). For this we will use eqs. (20 and 23), which provide the mass flow rate ( $J = jA$ ) due to advection and dispersion. For the 2D case for the flow in parallel to the  $x$ -axis, the fluxes can be represented in matrix form as

$$\begin{pmatrix} j_x \\ j_y \end{pmatrix} = n_e \cdot \begin{pmatrix} v_x \\ 0 \end{pmatrix} \cdot C - n_e \begin{pmatrix} \alpha_L v_x & 0 \\ v_x & \alpha_{Th} v_x \end{pmatrix} \cdot \begin{pmatrix} \frac{\partial C}{\partial x} \\ \frac{\partial C}{\partial y} \end{pmatrix}$$

which represents the following set of equations.

$$j_x = n_e v_x C - n_e \alpha_L v_x \frac{\partial C}{\partial x} \quad (36)$$

and

$$j_y = -n_e \alpha_{Th} v_x \frac{\partial C}{\partial y} \quad (37)$$

Note that we have used  $D_{hyd} = \alpha v$  in the above equations. We will insert eqs. (36 and 37) to complete our transport equation (eq. 37), thus we get

$$n_e \frac{\partial C}{\partial t} = -\frac{\partial}{\partial x} \left( n_e v_x C - n_e \alpha_L v_x \frac{\partial C}{\partial x} \right) - \frac{\partial}{\partial y} \left( -n_e \alpha_{Th} v_x \frac{\partial C}{\partial y} \right) \quad (38)$$

If we let  $n_e$  be constant it can be taken out of differentials, consequently eq. (38) after rearrangement becomes

$$\frac{\partial C}{\partial t} + \frac{\partial}{\partial x} (v_x C) = \frac{\partial}{\partial x} \left( \alpha_L v_x \frac{\partial C}{\partial x} \right) + \frac{\partial}{\partial y} \left( \alpha_{Th} v_x \frac{\partial C}{\partial y} \right) \quad (39)$$

Eq. (39), a  $2D$  transport equation, can be converted to  $1D$  transport equation by removing the  $y$ -components and likewise it can be converted to  $3D$  transport equation by adding the  $z$ -components. Next we learn about other essential requirements for transport problems.

### 17.4.1 Initial and boundary conditions

A complete set of transport problem includes the transport equation, e.g., eq. (39), the initial conditions (time dependent) and the boundary conditions (space dependent). The initial condition describes the distribution of mass or specifies concentration in the entire area of investigation at some starting time, usually  $t=0$  is considered. The initial conditions are required for the transport problems that are time dependent (e.g., eq. 39). These are so called transient problems and they comprise of general transport problems. A mathematical statement of initial condition is

$$C(t=0) = C_0 \quad (40)$$

i.e., at  $t=0$ ,  $C_0$  is the concentration in the entire investigated area.

For time-independent problems, i.e.,  $\partial C / \partial t = 0$ , also called steady-state problems boundary conditions are required to be specified. The boundary conditions quantify the impacts of condition at the surrounding just adjacent to the area under investigation. As the name suggests, the boundary conditions have to be specified for the entire boundary of the investigation area, even for the cases when the investigated area in unlimited or extend are infinite. The boundary conditions may or may not be time dependent. For transport problems three types of boundary conditions have been defined.

A first type of boundary condition, also called Dirichlet type, specifies the value of mass or concentration at the boundary. The mathematical statement of the first type boundary condition is

$$C(x, y, z, t) = C_0 \quad (41)$$

i.e., at any  $t$ , the concentration at specified space  $(x, y, z)$  is  $C_0$ . Tracer injection, effluent concentration from polluted area can be few examples of first type boundary conditions. In a second type boundary condition, also called Neumann type condition, the component of the concentration gradient perpendicular to the boundary is specified. The mathematical statement of the second type boundary condition along  $x$ -axis is

$$\frac{dC(t)}{dx} = C_0 \quad (42)$$

Since the gradient of concentration is proportional to diffusive flux, the second type boundary condition is also called a flux specified boundary. Second type boundary conditions are also used for specifying the no-flow condition, i.e., no concentration gradient exist across the boundary. However it has to be noted that absolute no-flow can only exist when there is no concentration gradient and the velocity vector is zero. Rock formation at the bottom of aquifer will result to a no flow type of boundary condition.

The third type, Cauchy or Robin type, conditions combines the first two conditions, i.e. both a specific value and gradient is specified. This condition generalizes the other two boundary conditions. Mathematical statement of the third type boundary condition (along  $x$ -axis) is

$$\frac{dC(t)}{dx} + C(x, t) = C_0 \quad (43)$$

## 17.5 Solving Transport Problems

In the previous sections we learned about the transport problems and developed the transport equation, also called Advection-Dispersion (AD) equation. In this section we will attempt to solve the AD model, which has been one of our objective from this module. The techniques for solving the transport problems are extremely



varied and therefore we will restrict to simpler problems (1D, and few reaction components etc.). We will very briefly get introduced to complex problems (higher dimensions, multiple reaction).

Direct integrating of the AD equation (e.g. eq. 39) is only possible in very simplified cases (e.g., only advection case). However, transformation methods (Laplace, Fourier etc.) can be applied to directly solve ADR equation under some conditions. Thus obtained solutions are called analytical models, which are exact and often a closed-form. The conditions very generally include that aquifer is of regular geometry, is homogeneous and isotropic and the chemical reactions are linear models based. Numerical methods, e.g. Finite-Difference Method (FDM), Finite-Element Method (FEM), that provide approximate solution of the ADR equation are very commonly used to solve transport problems. In this module we will restrict to an analytical model and very briefly discuss numerical methods.

### 17.5.1 Analytical Solution

The solution provided in the succeeding subsections generally assumes the following:

1. Fluid of constant density and viscosity
2. Flow in x-direction only, and velocity is constant.
3. The longitudinal dispersion coefficient  $D_x$  is constant.

In addition we assume aquifer is homogeneous and isotropic. Additional assumptions is provided when required.

### 17.5.2 The Advection-Dispersion (AD) equation

The simplest ADR equation can be a transient 1D equation without the reaction term, i.e. AD equation, which is

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2} - v \frac{\partial C}{\partial x} \quad (44)$$

we consider the initial condition

$$C(x, t = 0) = C_0$$

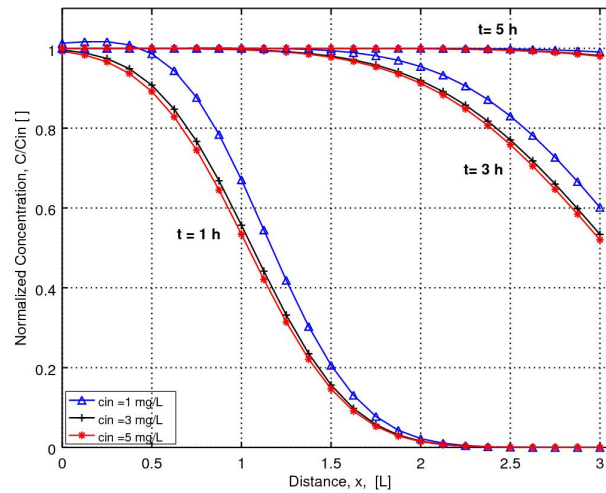


Fig. 4: Visualizing the Ogata and Banks (1961) solution and boundary conditions

$$C(x=0, t) = C_{in} \quad \text{and} \quad C(x=\infty, t) = 0$$

The conditions implies that our investigation area is semi-infinite and contain  $C_0$  concentration initially and  $C_{in}$  concentration is continuously injected to the investigation area from the location  $x=0$ . The following analytical solution of this problem is provided in *Ogata and Banks*, (1961):

$$C(x, t) = C_0 + \frac{C_{in} - C_0}{2} \left( \operatorname{erfc} \left( \frac{x - vt}{2\sqrt{Dt}} \right) + \exp \left( \frac{v}{D} x \right) \operatorname{erfc} \left( \frac{x + vt}{2\sqrt{Dt}} \right) \right) \quad (45)$$

in which  $\operatorname{erfc}$  denotes the complementary error-function, which is defined as

$$\operatorname{erfc}(x) = 1 - \operatorname{erf}(x) \quad \text{with} \quad \operatorname{erf}(x) = \frac{1}{\sqrt{\pi}} \int_{-x}^x e^{-t^2} dt$$

Readers should check the wikipedia site ([https://en.wikipedia.org/wiki/Error\\_function](https://en.wikipedia.org/wiki/Error_function)) to find methods to approximate the value of  $\operatorname{erfc}$ .

In order to visualize eq. (45) we perform a column (3 m long) experiment with  $C_0 = 0$  mg/L,  $v = 1$  m/h and  $D = 1$  m<sup>2</sup>/h. Let our  $C_{in} = 1, 3$  and  $5$  mg/L. We will check the concentration in the column at different times. Fig. 4 provides the results of our experiment. The result suggest that it will require 5 hours for our homogeneous distribution of mass in the column. For further visualisation of the Ogata and Banks solution, readers may check

the GNU-Octave <https://www.gnu.org/software/octave/>) code (ogata1D.m) provided at <https://github.com/prabhasyadav/UGC-Transport> .

### 17.5.3 An introduction to numerical method for transport problems

Numerical modelling is generally used for solving transport problems. The main reason behind this is that analytical solution are possible for only few cases, often limited to regular domain geometry and linear (differential) equations. These limitations are easily overcome by numerical method. This topic is a very extensive. Here, we intend to only introduce the topic very briefly. Among few numerical methods that are available, Finite Difference Method (FDM) and Finite Element Method (FEM) are the two very extensively used in solving transport problems. We will focus on this two methods. Towards the end of the section we will use an example problem to illustrate FDM, the more common of the two methods mentioned above.

In the FDM the governing partial differential equation is replaced by a set of difference equations applicable to the system of nodes, i.e., the area under investigation is discretized into structured meshes. Taylor series expansion or polynomial fitting techniques are used to approximate all space/time derivative in terms of concentration. For example the first and second order derivatives of the transport equation is approximated using

$$\left( \frac{\partial C}{\partial x} \right)_i \approx \frac{C_{i+1} - C_{i-1}}{2\Delta x} \quad (46)$$

and

$$\left( \frac{\partial^2 C}{\partial x^2} \right)_i \approx \frac{C_{i+1} + 2C_i - C_{i-1}}{(\Delta x)^2} \quad (47)$$

where  $i$  refers to the node  $i$  and  $i+1$  and  $i-1$  are two adjacent nodes.  $\Delta x$  represents the width of the node, which may be a constant. A system of algebraic equation results when each node in the domain is considered. The system of equation is then solved using matrix algebra techniques. The most important aspect of FDM is the ease of formulating difference equations. *Anderson and Woessner*, (1992) provides an excellent introduction to the method. Readers are suggested to start with that

literature if interested in learning the FDM for groundwater studies. For modelling, MODFLOW (an open-source) developed and maintained by United States Geological Survey (USGS) can be considered academic, research and industry standard. Further details can be obtained from <http://water.usgs.gov/ogw/modflow/>.

The mathematics of FEM is not as straightforward as that of FDM. In this method the area under investigation is subdivided into elements that are defined by nodes. The element can be of different shapes, although triangular and quadrilateral shapes are the most common one used. Generation of FEM mesh is very tedious, as such FEM codes usually include mesh generation software. The solution of the differential equation using FEM is found as a combination of shape functions. The shape function (also called Lagrangian function) for the triangular element are linear within each element. Thus in Cartesian coordinates the shape function  $f$  is

$$f(x, y) = a_{\alpha 0} + a_{\alpha 1x} + a_{\alpha 2y} \quad (48)$$

within element  $\alpha$ . All coefficients  $a_{\alpha j}$  for all elements are computed, which is derived from the integral equivalent of differential equation, also called weak formulation (*Huyakorn and Pinder, 1983*).

## Summary:

This sub-module provided:

1. A systematic approach to understand the flow and transport problems
2. Introduced several processes that influence the coupled (flow and transport) problems..
3. The derivation of governing equation to solve for the flow and transport problem.
4. Approaches to solve the flow and transport problems.

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