

Tutorial 10 Solutions

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Physical properties are same in all directions

5. * A two-dimensional isotropic harmonic oscillator has the Hamiltonian

$$H = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) + \frac{1}{2}k(x^2 + y^2)$$

- (a) Show that the energy levels are given by

$$E_{n_x, n_y} = \hbar\omega(n_x + n_y + 1) \quad \text{where } n_x, n_y \in (0, 1, 2, \dots) \quad \omega = \sqrt{\frac{k}{m}}$$

- (b) What is the degeneracy of each level?

Solⁿ (a) TISE

$$\hat{H}\psi(n, y) = E\psi(n, y)$$

$$-\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial n^2} + \frac{\partial^2}{\partial y^2} \right] \cdot \psi(n, y) + \frac{1}{2}k(n^2 + y^2)\psi(n, y) = E\psi(n, y)$$

$$\text{Expressing } \psi(n, y) = X(n)Y(y)$$

$$-\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial n^2} + \frac{\partial^2}{\partial y^2} \right] X(n)Y(y) + \frac{1}{2}k(n^2 + y^2)X(n)Y(y) = EX(n)Y(y)$$

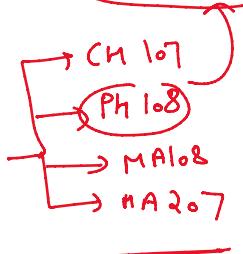
Dividing by $X(n)Y(y)$

$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2 X(n)}{\partial n^2} \cdot \frac{1}{X(n)} + \frac{1}{2}kn^2 \right] + \left[-\frac{\hbar^2}{2m} \frac{\partial^2 Y(y)}{\partial y^2} \cdot \frac{1}{Y(y)} + \frac{1}{2}ky^2 \right] = E$$

E_n

E_y

$$-\frac{\partial^2 V}{\partial n^2} = \frac{\delta}{\epsilon_0}$$



$$-\frac{\hbar^2}{2m} \frac{\partial^2 X(n)}{\partial n^2} + \frac{1}{2}kn^2 X(n) = E_n X(n)$$

$$(k = m\omega^2)$$

Solutions

$$A e^{-\frac{\epsilon_n}{2kT}}$$

$$X_n(n) = \frac{1}{\sqrt{2^n \cdot n!}} \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-m\omega n^2/2\hbar} H_n \left(\sqrt{\frac{m\omega}{\hbar}} n \right)$$

$$E_{n,n} = \hbar\omega(n_n + 1/2)$$

$$E_{y,y} = \hbar\omega(n_y + 1/2)$$

$$E_{\text{Total}} = \hbar\omega(n_n + n_y + 1)$$

$$E_0 = E_{0,0} = \hbar\omega$$

$$E_1 = E_{0,1} = E_{1,0} = 2\hbar\omega$$

$$E_2 = E_{0,2} = E_{1,1} = 3\hbar\omega$$

degeneracy \rightarrow (No. of microstates a particular macrostate has)

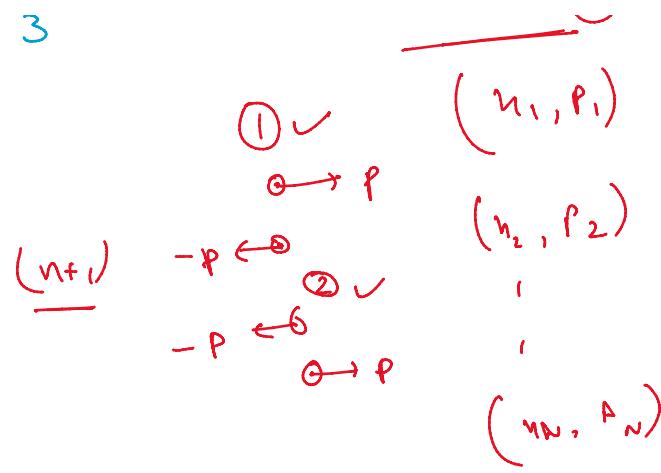
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P, V, T, $\langle E \rangle \rightarrow$ Macrostate
 (n, P)

$$\begin{aligned} \rightarrow E_1 &= E_{01} = \underline{E}_{10} = \underline{\omega \hbar \omega} \\ E_2 &= E_{20} = \underline{E}_{02} = \underline{E}_{11} = 3 \hbar \omega \\ &\vdots \\ E_n &= E_{n0} = \dots E_{0n} = \underline{(n+1) \hbar \omega} \end{aligned}$$



6. * Consider the Hamiltonian of a two-dimensional anisotropic harmonic oscillator ($\omega_1 \neq \omega_2$)

$$H = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + \frac{1}{2}m\omega_1^2 q_1^2 + \frac{1}{2}m\omega_2^2 q_2^2$$

- (a) Exploit the fact that the Schrödinger eigenvalue equation can be solved by separating the variables and find a complete set of eigenfunctions of H and the corresponding eigenvalues.
- (b) Assume that $\frac{\omega_1}{\omega_2} = \frac{3}{4}$. Find the first two degenerate energy levels. What can one say about the degeneracy of energy levels when the ratio between ω_1 and ω_2 is not a rational number.

$q_1, q_2 \rightarrow \underline{\text{generalized coordinates}}$

can be $\underline{u}, \underline{y}$

$\underline{y}, \underline{z}$

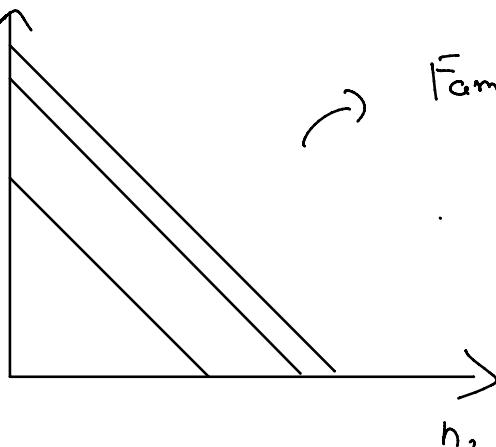
$\underline{u}, \underline{z}$

Soln (a) $\Psi(\underline{q}_1, \underline{q}_2) = \underline{\Theta_1(q_1)} \underline{\Theta_2(q_2)}$

$$E(n_1, n_2) = \hbar \omega_1 \left(n_1 + \frac{1}{2} \right) + \hbar \omega_2 \left(n_2 + \frac{1}{2} \right)$$

(b) $\omega_1 = 3\omega$ $\omega_2 = 4\omega$ $n_1, n_2 \in \mathbb{Z}$
 $E(n_1, n_2) = \hbar \omega \left(\underline{3n_1 + 4n_2 + \frac{1}{2}} \right) = \underline{\text{const.}}$

$$\frac{\omega_1}{\omega_2} = \frac{P}{Q} \quad P, Q \in \mathbb{Z} \quad Q \neq 0$$



$$3n_1 + 4n_2 = C$$

$$n_1, n_2 \in \mathbb{Z} \Rightarrow C \in \mathbb{Z}$$

First two degenerate levels
 $(n_1, n_2) = \underline{(0, 3)}, \underline{(4, 0)}$

$$\frac{31}{2} \hbar \omega$$

Take $\frac{\omega_1}{\omega_2} = P ; P \in \mathbb{I}$

$$E(n_1, n_2) = \hbar \omega_2 \left(\underline{Pn_1 + n_2} + \frac{P+1}{2} \right)$$

For degenerate state $Pn_1 + n_2 = C$ for at least two pairs

$$\begin{aligned} Pn_1 + n_2 &= Pn'_1 + n'_2 \\ P(n_1 - n'_1) &= (n'_2 - n_2) \\ \underbrace{P}_{\in \mathbb{I}} & \quad \underbrace{C}_{\in \mathbb{S}} \\ n'_1 &\neq n_1 \\ n'_2 &\neq n_2 \end{aligned}$$

$\dots + \dots (n'_1, n'_2) \text{ exist}$

$\leftarrow \in I$ $\hookrightarrow \in \mathcal{O}$
 so no two distinct pairs (n_1, n_2) & (n'_1, n'_2) exist
 Hence degeneracy is not possible in this case.

7. A particle of mass m is confined to move in the potential $(m\omega^2x^2)/2$. Its normalized wave function is

$$\psi(x) = \left(\frac{2\beta}{\sqrt{3}}\right) \left(\frac{\beta}{\pi}\right)^{1/4} x^2 e^{-\beta x^2/2}$$

where β is a constant of appropriate dimension.

(a) Obtain a dimensional expression for β in terms of m, ω and \hbar .

(b) It can be shown that the above wave function is the linear combination

$$\psi(x) = a\psi_0(x) + b\psi_2(x)$$

where $\psi_0(x)$ is the normalized ground state wave function and $\psi_2(x)$ is the normalized second excited state wave function of the potential. Evaluate b and hence calculate the expectation value of the energy of the particle in this state $\psi(x)$.

Given: $I_0(\beta) = \int_{-\infty}^{+\infty} e^{-\beta x^2} dx = \sqrt{\frac{\pi}{\beta}}$, $I_n(\beta) = \int_{-\infty}^{+\infty} (x^2)^n e^{-\beta x^2} dx = (-1)^n \frac{\partial^n}{\partial \beta^n} (I_0(\beta))$,
 $\psi_0(x) = \left(\frac{\beta}{\pi}\right)^{1/4} e^{-\beta x^2/2}$

Sol $[Bn^2] = M^a L^b T^c$
 $[\beta] = L^{-2} = [M]^a [\omega]^b [\hbar]^c = [M]^a [T^{-1}]^b [M L^2 T^{-1}]^c$

$$c = -a$$

$$c = -b$$

$$c = -1$$

(b) $\varphi(n) = a\varphi_0(n) + b\varphi_2(n)$

$$\int_{-\infty}^{\infty} (\varphi(n) \cdot \varphi_2^*(n) \cdot dn) = \int_{-\infty}^{\infty} a \varphi_0(n) \cdot \varphi_2(n) \cdot dn + b \int_{-\infty}^{\infty} \varphi_2(n) \cdot \varphi_2^*(n) \cdot dn$$

(orthogonal wave functions)

Normalized wavefunction

$$b = \int_{-\infty}^{\infty} \left(\frac{2\beta}{\sqrt{3}}\right) \left(\frac{\beta}{\pi}\right)^{1/2} n^2 e^{-\beta n^2/2} \cdot \frac{1}{2\sqrt{2}} \left(\frac{\beta}{\pi}\right)^{1/2} e^{-\beta n^2/2} H_2(\sqrt{\beta} n) dn$$

$$\varphi_n(n) = \frac{1}{2}$$

$$H_2(z) = e^{-z^2} \frac{d^2}{dz^2} (e^{-z^2}) = e^{-z^2} \left(\frac{d}{dz} (-2z e^{-z^2}) \right)$$

$$= 4z^2 e^{-2z^2} - 2e^{-2z^2}$$

$$b = \left(\frac{2\beta}{\sqrt{3}}\right) \left(\frac{\beta}{\pi}\right)^{1/2} \cdot \frac{1}{2\sqrt{2}} \int_{-\infty}^{\infty} n^2 e^{-\beta n^2} \left[4\beta n^2 e^{-2\beta n^2} - 2e^{-2\beta n^2} \right] dn \rightarrow 2\sqrt{\frac{\pi}{\beta^3}}$$

$$\approx \frac{1}{\pi} \int_{-\infty}^{\infty} 4\beta n^4 e^{-3\beta n^2} dn - \int_{-\infty}^{\infty} 2n^2 e^{-3\beta n^2} dn$$

$$b = \left(\frac{2\beta}{\sqrt{3}}\right) \left(\frac{\beta}{\pi}\right)^{1/2} \cdot \frac{1}{2\beta} \cdot \frac{\pi}{\beta^3}$$

$$\boxed{b = \sqrt{\frac{2}{3}}}$$

$$\boxed{a = \frac{1}{\sqrt{3}}}$$

$$\varphi(n) = a\varphi_0 + b\varphi_2$$

$$\varphi^*(n) = a\varphi_0^* + b\varphi_2^*$$

$$\boxed{a^2 + b^2 = 1}$$

$$\int_{-\infty}^{\infty} 4B_n^4 e^{-3\beta n^2} dn - \int_{-\infty}^{\infty} 2n^- e^{-3\beta n^2} dn$$

$$\downarrow$$

$$I_1(3\beta)$$

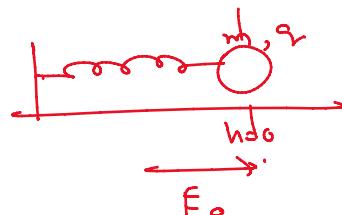
$$= (-1)^1 \left(-\frac{1}{2}\right) \sqrt{\frac{\pi}{\beta^3}}$$

$$\int \varphi(n) \cdot \varphi_0^*(n) = \int a^2 \varphi_0 \cdot \varphi_0^* + b^2 \int \varphi_2 \cdot \varphi_2^*$$

$$1 = a^2 + b^2$$

$$+ ab \int \varphi_0^* \varphi_2 + ab \int \varphi_0 \varphi_2^*$$

9. * A charged particle of mass ' m ' and charge ' q ' is bound in a 1-dimensional simple harmonic oscillator potential of angular frequency ' ω '. An electric field E_0 is turned on.
- What is the total potential $V(x)$ experienced by the charge?
 - Express the total potential in the form of an effective harmonic oscillator potential.
 - Sketch $V(x)$ versus x .
 - What is the ground state energy of the particle in this potential?
 - What is the expectation value of the position (x) if the charge is in its ground state?



$$h = \frac{qE_0}{m}$$

$$\text{Sol C1 } V(n) = -qE_0 n + \frac{1}{2} m \omega^2 n^2$$

; The $\vec{E} = E_0$ is assumed to be in +n dir.

$$E = -\frac{dV}{dn}$$

$$-\frac{qE_0}{m}$$

$$\frac{1}{2} m \omega^2 n^2 + V_0$$

$$(b) V(n) = \frac{1}{2} m \omega^2 n^2 - qE_0 n$$

$$= \frac{1}{2} m \omega^2 \left(n^2 - \frac{2qE_0}{m \omega^2} n \right)$$

$$= \frac{1}{2} m \omega^2 \left(n - \frac{qE_0}{m \omega^2} \right)^2 - \frac{1}{2} m \omega^2 \left(\frac{q^2 E_0^2}{m^2 \omega^4} \right)$$

$$\frac{1}{2} m \omega^2 \left(n^2 - \frac{2qE_0}{m \omega^2} n \right)$$

$$\left(n - \frac{qE_0}{m \omega^2} \right)^2$$

$$V(n) = \frac{m \omega^2}{2} \left(n - \frac{qE_0}{m \omega^2} \right)^2 - \frac{q^2 E_0^2}{2 m \omega^2}$$

$$- \frac{1}{2} m \omega^2 n^2 + V_0$$

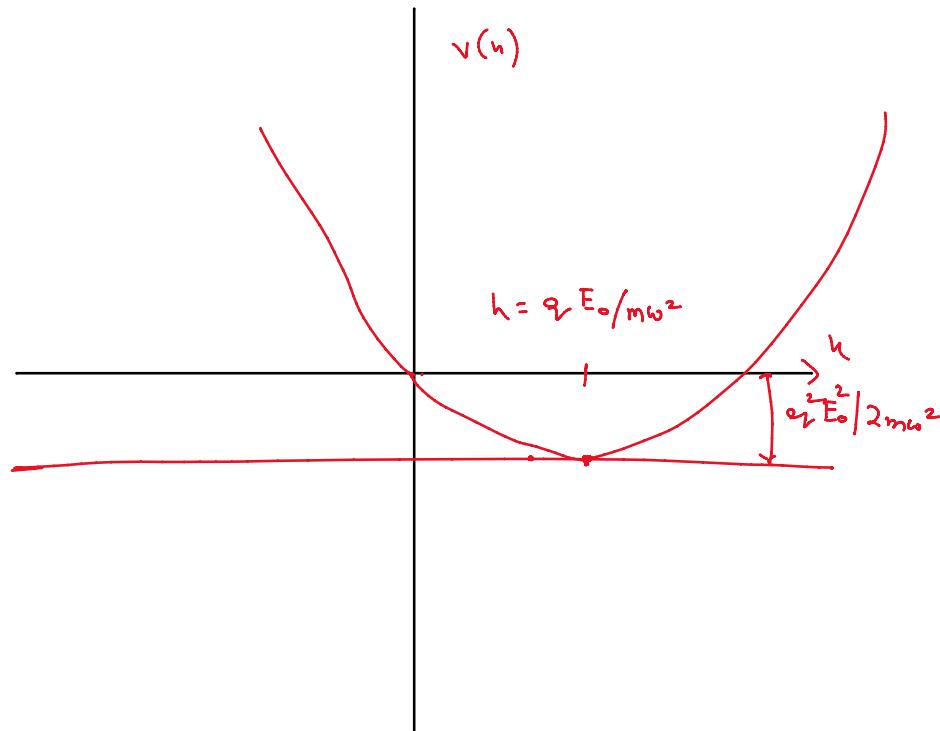
$$\rightarrow n' = n_0 - \frac{qE_0}{m \omega^2} \rightarrow \begin{matrix} \text{mean pos} \\ \text{shifted} \end{matrix}$$

$$= \frac{1}{2} m \omega^2 n'^2 + V_0 \rightarrow n' = n - \frac{q E_0}{m \omega^2} \rightarrow \text{mean shifted}$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial n^2} + \left(\underbrace{\frac{1}{2} m \omega^2 \left(n - \frac{q E_0}{m \omega^2} \right)^2}_{n - \frac{q E_0}{m \omega^2}} + V_0 \right) \psi(n) \rightarrow V_0 = -\frac{q^2 E_0^2}{2 m \omega^2}$$

$$= E \psi(n)$$

(c)

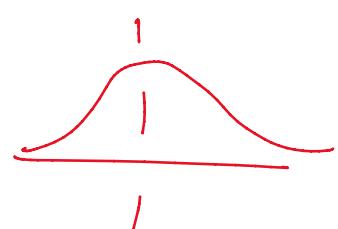


$$(d) E = \hbar \omega \left(n + \frac{1}{2} \right) - \frac{q^2 E_0^2}{2 m \omega^2}$$

(e) Ground state wave function

$$\psi_0(n) = \left(\frac{m \omega}{\pi \hbar} \right)^{1/4} e^{-(m \omega / 2 \hbar) n'^2} \quad \left\{ n' = n - \frac{q E_0}{m \omega^2} \right\}$$

$$\langle n \rangle = \int_{-\infty}^{\infty} \psi_0(n) \cdot n \cdot \psi_0^*(n) \cdot dn$$



$$n' \rightarrow n - \frac{q E_0}{m \omega^2}$$

$$\langle n \rangle = \int_{-\infty}^{\infty} \psi_0(n) n' \cdot \psi_0^*(n) \cdot dn + \int_{-\infty}^{\infty} \psi_0(n) \cdot \psi_0^*(n) \cdot dn \left(\frac{q E_0}{m \omega^2} \right)$$

$$= \frac{q E_0}{m \omega^2}$$

$$= \frac{q^2 \omega}{m \omega^2}$$

$$\textcircled{1} \quad \Delta n = \langle n^2 \rangle - \underline{\underline{\langle n \rangle}} \quad \nabla(n) = \frac{1}{2} m \omega^2 n^2$$

$$= \int \varphi_0(n) \cdot n^2 \cdot \varphi_0^*(n) dn$$

$$= \int \left(\frac{\beta}{\pi} \right)^{1/2} e^{-\frac{\beta n^2}{2}} \cdot n^2 \cdot e^{-\beta n^2/2} \cdot \left(\frac{\beta}{\pi} \right)^{1/2} dn \quad \beta = \frac{m\omega}{\hbar}$$

$$\frac{\Delta n}{\Delta p} = \langle p^2 \rangle - \underline{\underline{\langle p \rangle}}^2$$

$$E \sim \frac{(\Delta p)^2}{2m}$$

$$\int \varphi_0^*(n) - i \frac{\hbar}{m} \frac{\partial}{\partial n} \varphi_0(n)$$

$$\int A e^{-kn^2} - i \hbar (-\alpha A e^{-kn^2}, 2n)$$

$$\textcircled{8} \quad \frac{n_x + n_y + n_z}{n} = \eta$$

$$g_n = \frac{1}{2} (n+1)(n+2)$$