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515-518

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Resolving photon number states in a superconducting circuit

Ming Lyu, Elena de la Hoz Lopez-Collado

Final projects for ELE456 at Princeton

May 11, 2017



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Outline

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- System sensitive to number of photons
- System: superconducting qubit + microwave transmission line
- Strong dispersive regime
- Spectroscopic measurements: Qubit's spectral lines different for each photon number state



Cavity QED

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- Cavity QED (cQED) → interaction electromagnetic field modes with atoms (or qubits)

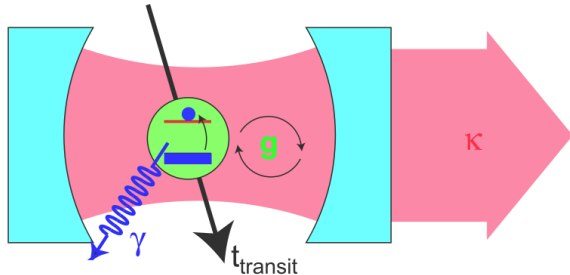


Image from Blais, Alexandre, et al. "Cavity quantum electrodynamics for superconducting electrical circuits: An architecture for quantum computation." Physical Review A 69.6 (2004): 062320.[1]



Cavity QED: Superconducting qubit

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- Cavity QED (cQED) → interaction electromagnetic field modes with superconducting qubit

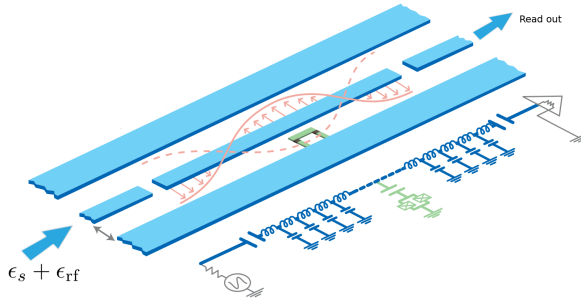


Image from Blais, Alexandre, et al. "Cavity quantum electrodynamics for superconducting electrical circuits: An architecture for quantum computation." Physical Review A 69.6 (2004): 062320.[1]



Cavity QED: the Hamiltonian

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Hamiltonian

$$H = \omega_r \left(a^\dagger a + \frac{1}{2} \right) + \omega_a \frac{\sigma^z}{2} + g \left(a^\dagger \sigma^- + a \sigma^+ \right)$$

- ω_r : cavity resonance frequency
- ω_a : qubit transition frequency
- g : strength qubit-photon coupling
- $\Delta = \omega_r - \omega_a$: detuning between qubit and cavity



Strong Dispersive Regime

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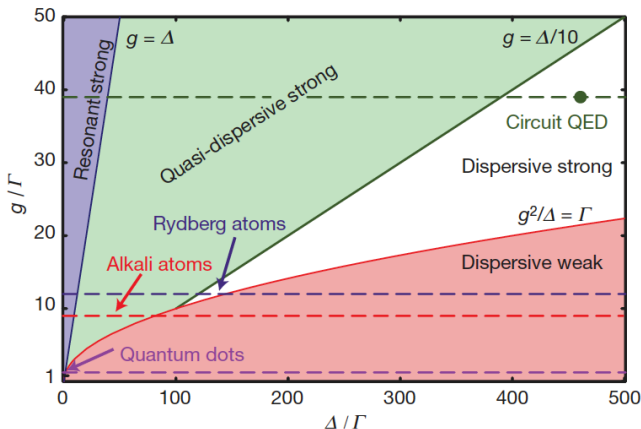


Image from Schuster, D. I., et al. "Resolving photon number states in a superconducting circuit." Nature 445.7127 (2007): 515-518.[3]



Strong dispersive Regime: Diagonalization

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- Transformation:

$$U = \exp \left(\frac{g}{\Delta} \left(a \sigma^+ - a^\dagger \sigma^- \right) \right)$$

- Hamiltonian to first order in $\frac{g}{\Delta}$ (dispersive regime):

$$\begin{aligned} H_0 &= U H U^\dagger \\ &\simeq \omega_r \left(a^\dagger a + \frac{1}{2} \right) + \omega_a \frac{\sigma^z}{2} + \chi \left(a^\dagger a + \frac{1}{2} \right) \frac{\sigma^z}{2} \end{aligned}$$

where $\chi = g/\Delta^2$



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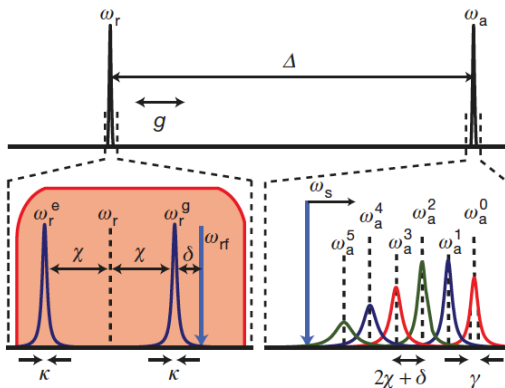


Image from Schuster, D. I., et al. "Resolving photon number states in a superconducting circuit." Nature 445.7127 (2007): 515-518.[3]



Driving terms

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- To conduct a measurement we first drive the cavity:

$$H_{\text{rf}} = \epsilon_{\text{rf}} \left(a^\dagger e^{-i\omega_{\text{rf}}t} + a e^{i\omega_{\text{rf}}t} \right)$$

with ω_{rf} near ω_r

- The frequency shift of the qubit measured with a sweeping signal

$$H_s = \epsilon_s \left(a^\dagger e^{-i\omega_s t} + a e^{i\omega_s t} \right)$$

with ω_s near ω_a

- Note that relative strength of ϵ_s is not mentioned. We treat it as a perturbation.



Rotating frame and Rotating wave approximation

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- Applying the transformation

$$U = \exp \left[\frac{g}{\Delta} (a\sigma^+ - a^\dagger\sigma^-) \right]$$

- And moving to the rotating frame:

$$U_I = \exp \left[i \left(\omega_{\text{rf}} a^\dagger a + \omega_s \sigma^z / 2 \right) t \right]$$

Under rotating frame, H_{rf} and H_s are (with RWA):

$$H_{\text{rf}} = \epsilon_{\text{rf}} (a^\dagger + a)$$

$$H_s = \left(\frac{g}{\Delta} \right) \epsilon_s (\sigma^+ + \sigma^-)$$



Final Hamiltonian and collapse operators

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- Full Hamiltonian:

$$H = \omega_r \left(a^\dagger a + \frac{1}{2} \right) + \omega_a \frac{\sigma^z}{2} + \chi \left(a^\dagger a + \frac{1}{2} \right) \frac{\sigma^z}{2} \\ - \left(\omega_{rf} a^\dagger a + \omega_s \frac{\sigma^z}{2} \right) + \epsilon_{rf} (a^\dagger + a) + \epsilon_s \frac{g}{\Delta} (\sigma^+ + \sigma^-)$$

- Collapse operator:

- Collapse operators cavity: $\sqrt{\kappa(1+n_{th})}a$, $\sqrt{\kappa n_{th}}a^\dagger$
- Collapse operator qubit: $\sqrt{\gamma}\sigma^-$
- Dephasing: $\sqrt{\gamma_\phi}\sigma^z$



- In the experiment, the transmitted amplitude at frequency ω_{rf} is the main observable under steady state.

Steady state

$$\dot{\rho}_s = 0 = -i[H, \rho_s] + \sum_n \left(2C_n \rho_s C_n^\dagger - \{\rho_s, C_n^\dagger C_n\} \right)$$

- What they really measure is the expectation of the electrical field $E \propto \langle a + a^\dagger \rangle$ [2] on a given frequency

$$E \propto \langle a + a^\dagger \rangle = \text{Tr}[\rho_s(a + a^\dagger)]$$



Property of the cavity: Analytical

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- Without the qubit, the cavity state is equivalently a damped harmonic oscillator with driving

$$H = \delta a^\dagger a + \epsilon(a + a^\dagger)$$

Collapse operators: $\sqrt{\kappa(n_{\text{th}} + 1)}a$ and $\sqrt{\kappa n_{\text{th}}}a^\dagger$

- When it's off resonant, its steady state is not but approximately a coherent state
- Analytically the photon number expectation value is

$$\bar{n} = \frac{\epsilon^2}{\delta^2 + \kappa^2/4} + n_{\text{th}}$$



Property of the cavity: Numerical

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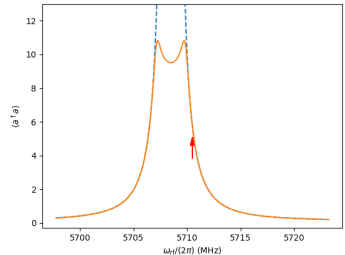
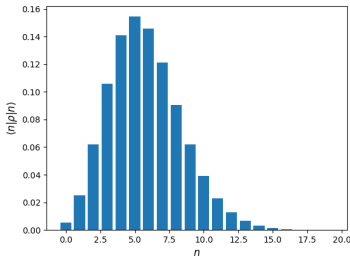
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- Numerically, a truncate on Fock space is needed
- To check the validity of the truncate, we plot the photon distribution and frequency response of the cavity.





Direct spectroscopic observation of quantized cavity photon number

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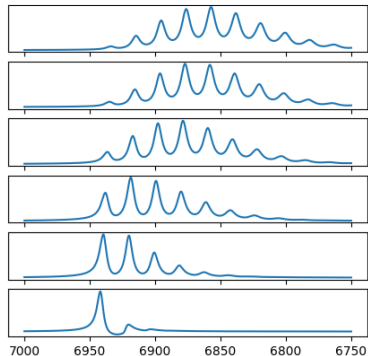
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For a fixed driving ϵ_{rf} ,
plot the reduction
 $V_0 - \langle a^\dagger + a \rangle_{ss}$ v.s. ω_s .

ϵ_{rf} is labeled by \bar{n} with
relationship:

$$\bar{n} = n_{\text{th}} + \frac{\epsilon_{\text{rf}}^2}{\delta^2 + \kappa^2/4}$$





Direct spectroscopic observation of quantized cavity photon number: compare

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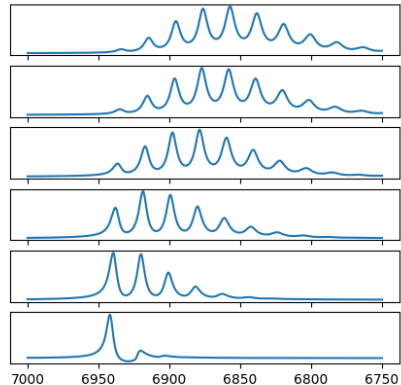
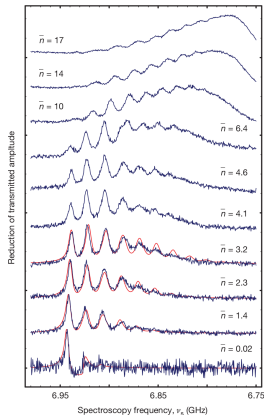
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- Fits well with small \bar{n} , but other noise becomes significant for larger \bar{n}



Strengthen?

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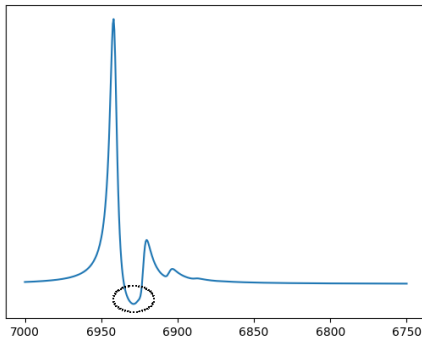
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For small signal, there's a range where the transmitted amplitude is increased. We'll explain it later.



Thermal Drive

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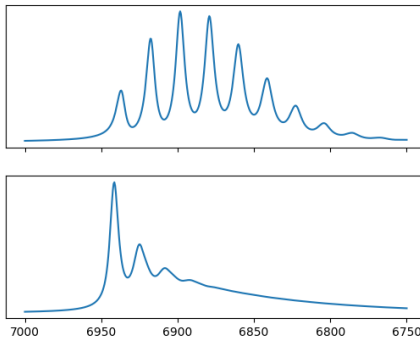
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- Thermal Drive is equivalent to setting n_{th} in collapse operator to the driving average, with small ϵ_{rf} to show the phase lock-in at the given frequency.





Thermal Drive: compare

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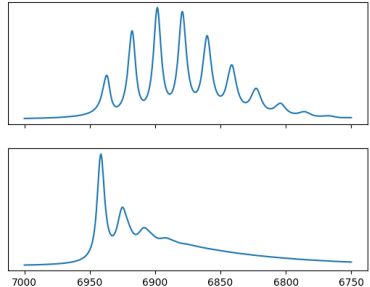
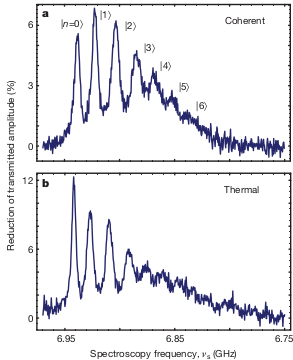
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- Note that there's no thermal drive theory fitting. Our results tracks fewer peaks, but this depends on how they do the measurement, which is not mentioned in the paper.



Discussion: The picture of what happens

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- The peaks shows discreteness in the photon state in the cavity.
- Exciting the qubit making the cavity off-resonance, which results in the reduction?



Discussion: The picture of what happens

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- The peaks shows discreteness in the photon state in the cavity.
- Exciting the qubit making the cavity off-resonance, which results in the reduction? **NOT TRUE**



Discussion: The picture of what happens

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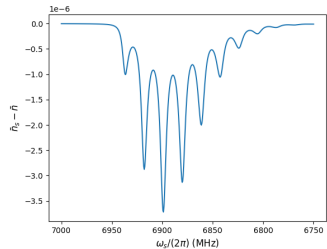
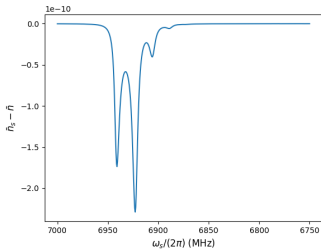
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- The peaks shows discreteness in the photon state in the cavity.
- Exciting the qubit making the cavity off-resonance, which results in the reduction? **NOT TRUE**
- Expected photon number increases at the peaks!





What happens

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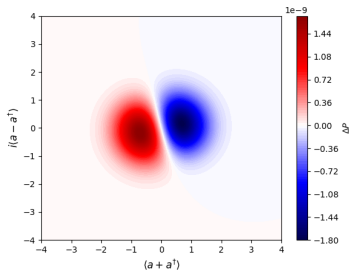
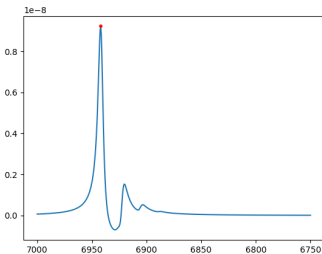
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Conclusion

- Excitation of the qubit is not the dominant effect, but the polarization of the qubit, which twists the cavity photon state.
- This can be shown from the difference of the Wigner function (quasiprobability distribution on phase diagram) with/without the signal field.





What happens

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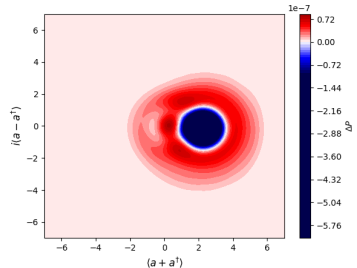
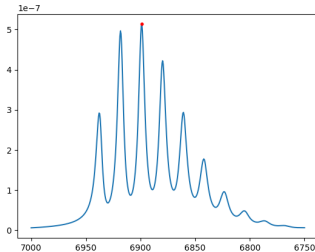
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Conclusion

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TBD

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- ▶ Alexandre Blais, Ren-Shou Huang, Andreas Wallraff, Steven M Girvin, and R Jun Schoelkopf.

Cavity quantum electrodynamics for superconducting electrical circuits: An architecture for quantum computation.

Physical Review A, 69(6):062320, 2004.

- ▶ David Isaac Schuster.

Circuit quantum electrodynamics.

Yale University, 2007.

- ▶ DI Schuster, AA Houck, JA Schreier, A Wallraff, JM Gambetta, A Blais, L Frunzio, J Majer, B Johnson, MH Devoret, et al.

Resolving photon number states in a superconducting circuit.

Nature, 445(7127):515–518, 2007.



Conclusion





Circuit Cavity QED

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Cavity

- 1D transmission line resonator
- Full-wave section of superconducting coplanar waveguides

Qubit

- Cooper pair box
- Superconducting mesoscopic island connected via a Josephson Junction to a reservoir



Circuit Cavity QED

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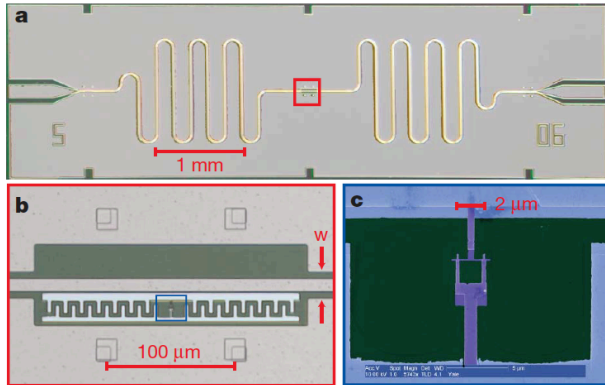


Figure: Cooper pair box inside a cavity, and spectral features of the circuit QED system.

Image from Schuster, D. I., et al. "Resolving photon number states in a superconducting circuit." Nature 445.7127 (2007): 515-518.

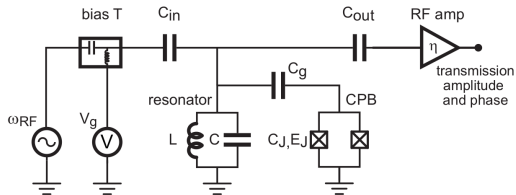


Measurement

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In the experiment, the transmitted amplitude at frequency ω_{rf} is the main observable. The exact way to measure can be found in Schuster's thesis [2]:



- What we really measure is the expectation of the voltage, or electrical field $E \propto \langle a + a^\dagger \rangle$