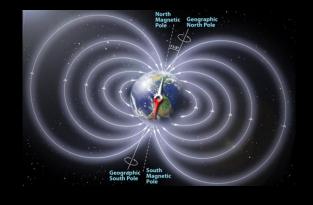
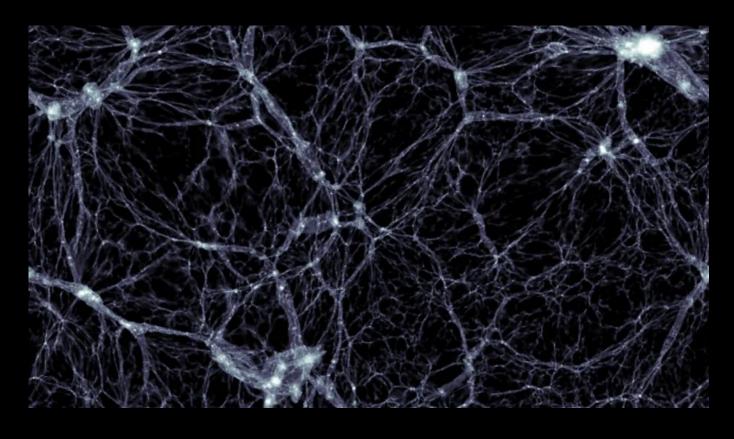


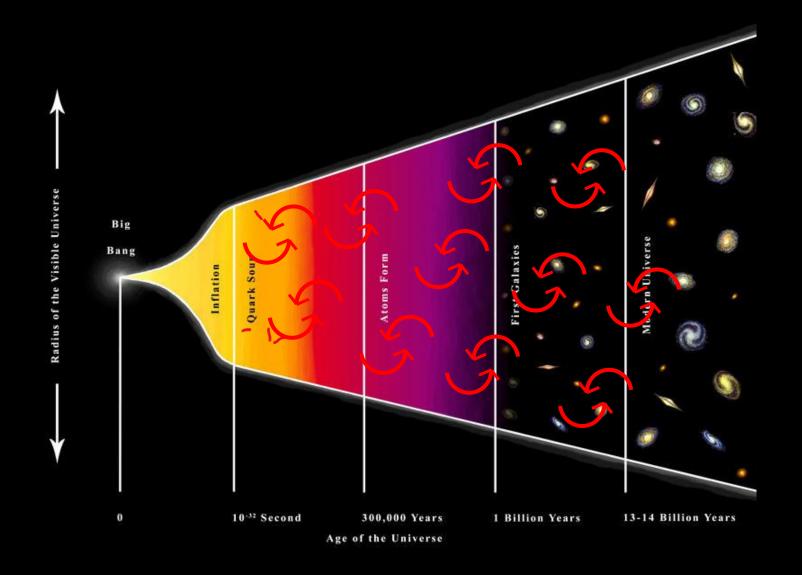
UBIQUITOUS MAGNETIC FIELDS



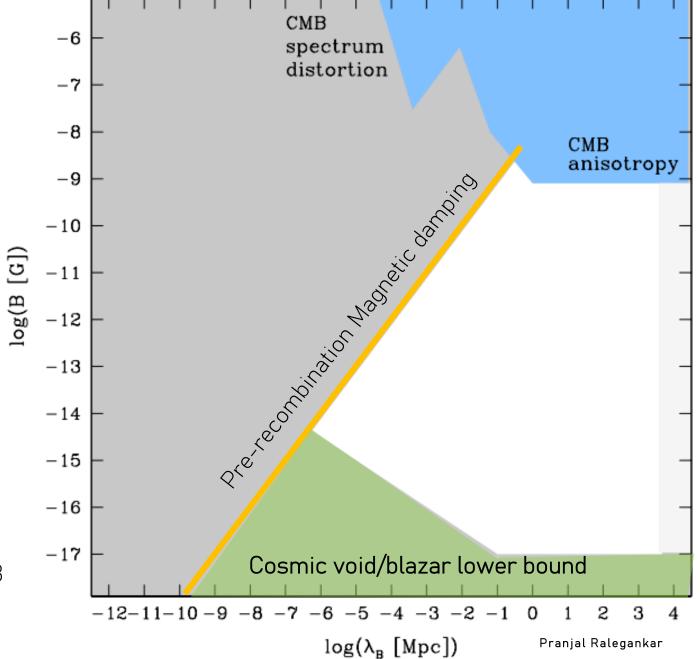




PRIMORDIAL: PRODUCED BY BIG BANG PLASMA

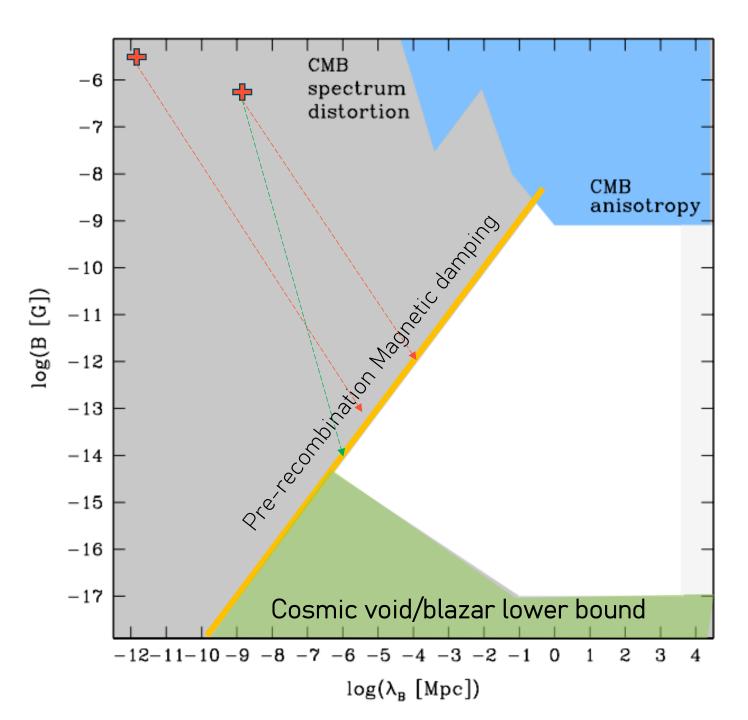


ALLOWED PMF PARAMETER SPACE

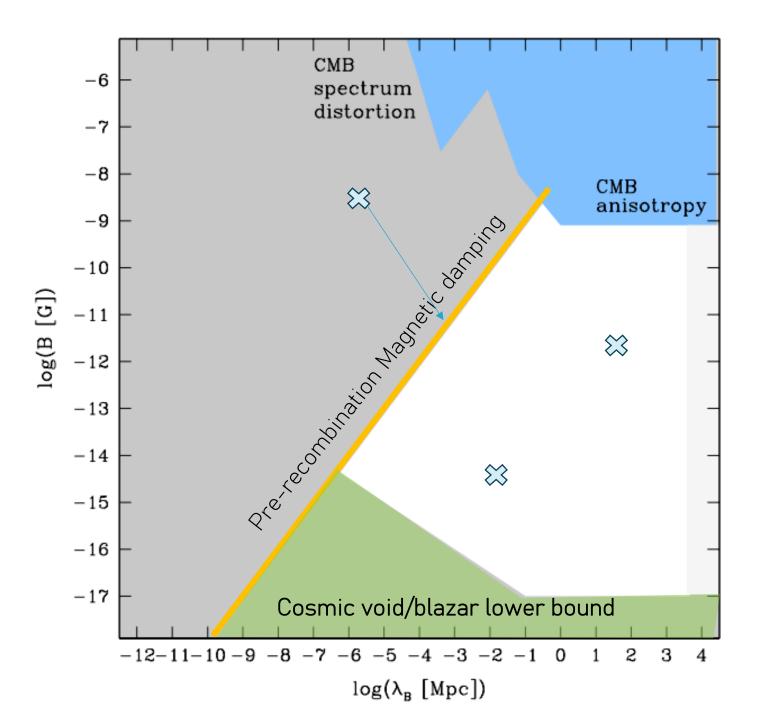


PMFS
GENERATED
POST
INFLATION LIE
ON THE
DAMPING LINE

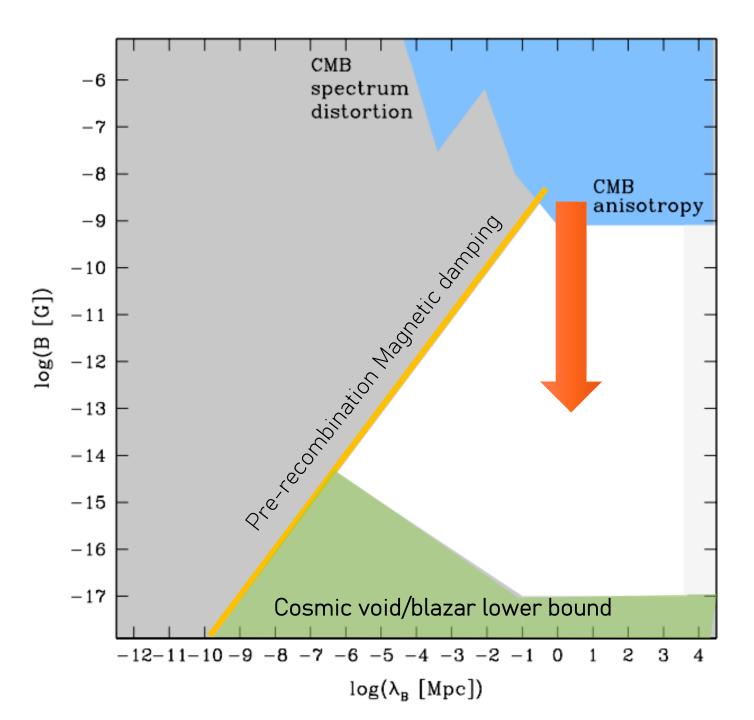
Banerjee and Jedamzik 2004



INFLATION
GENERATED
PMFS CAN BE
ANYWHERE ON
THE RIGHT OF
DAMPING LINE

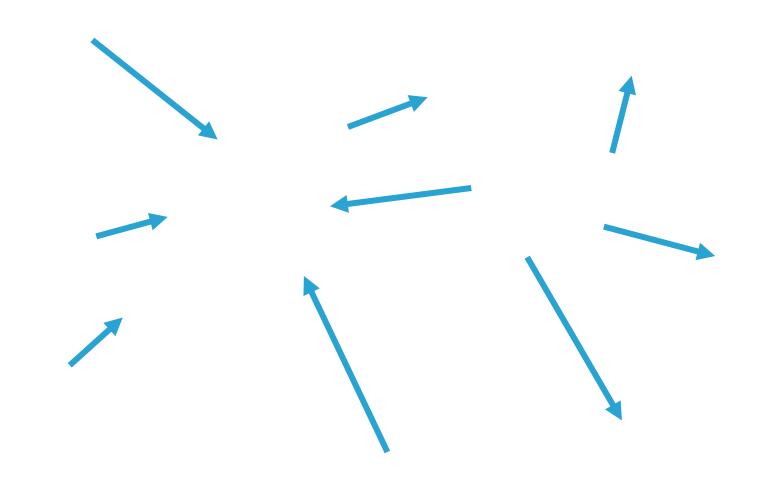


GOAL: TEST THE PRIMORDIAL HYPOTHESIS OF MAGNETIC FIELDS

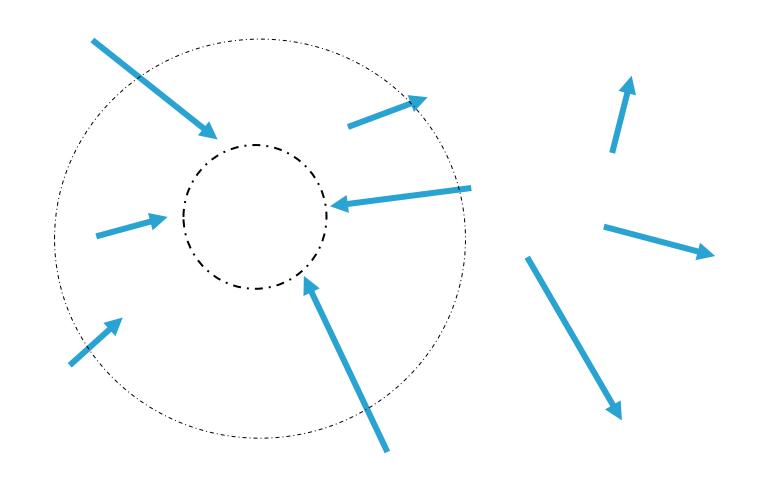


PRIMORDIAL MAGNETIC FIELDS ENHANCE DENSITY PERTURBATIONS

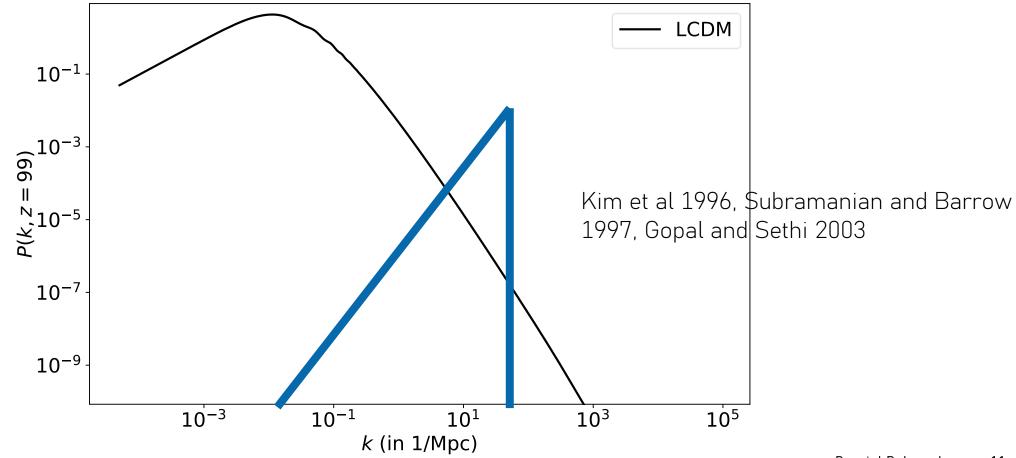
PRIMORDIAL MAGNETIC FIELDS ENHANCE DENSITY PERTURBATIONS



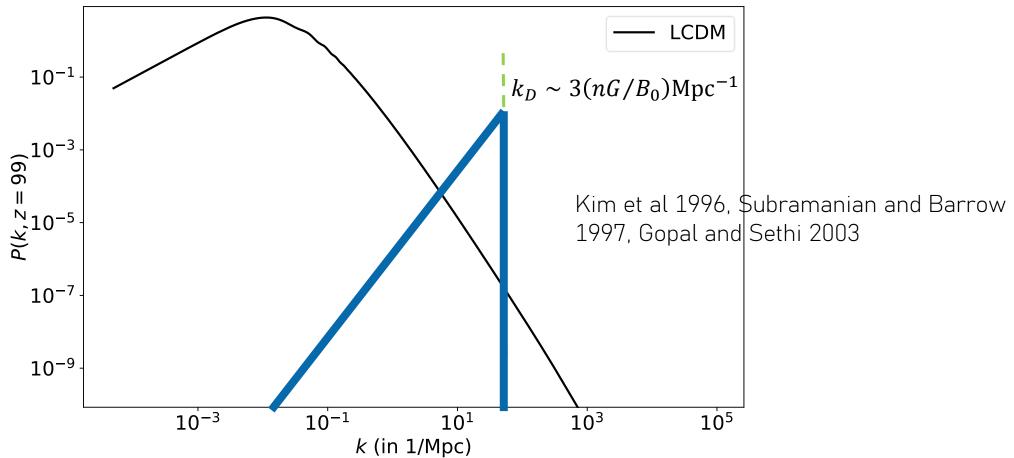
PRIMORDIAL MAGNETIC FIELDS ENHANCE DENSITY PERTURBATIONS



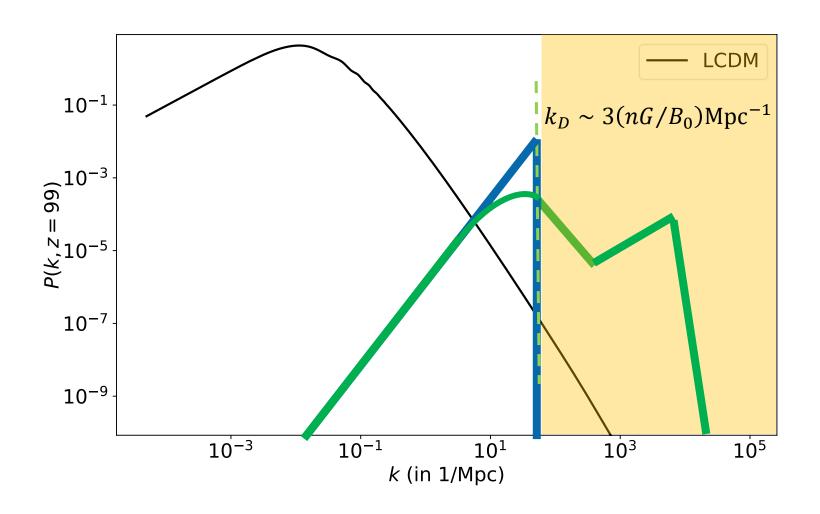
PRIMORDIAL MAGNETIC FIELDS ENHANCE POWER SPECTRUM ON SMALL SCALES



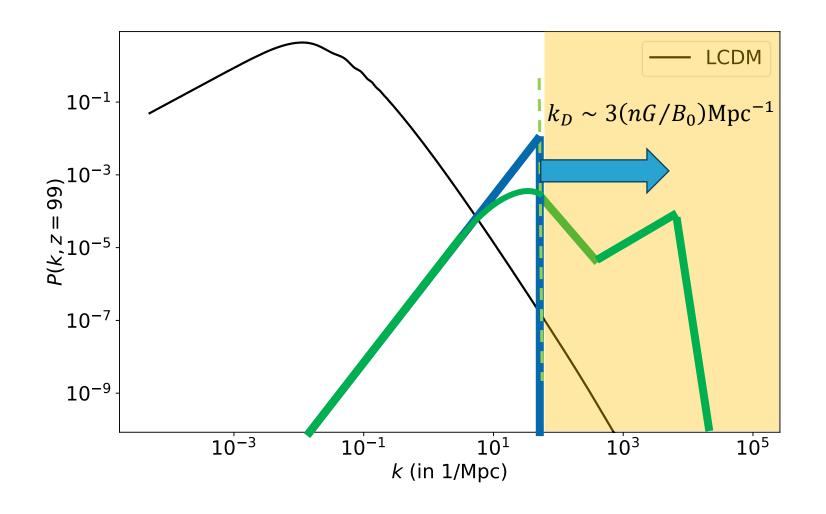
EARLY WORKS: BARYON DENSITY PERTURBATIONS SUPPRESSED BELOW MAGNETIC DAMPING (JEANS) SCALE



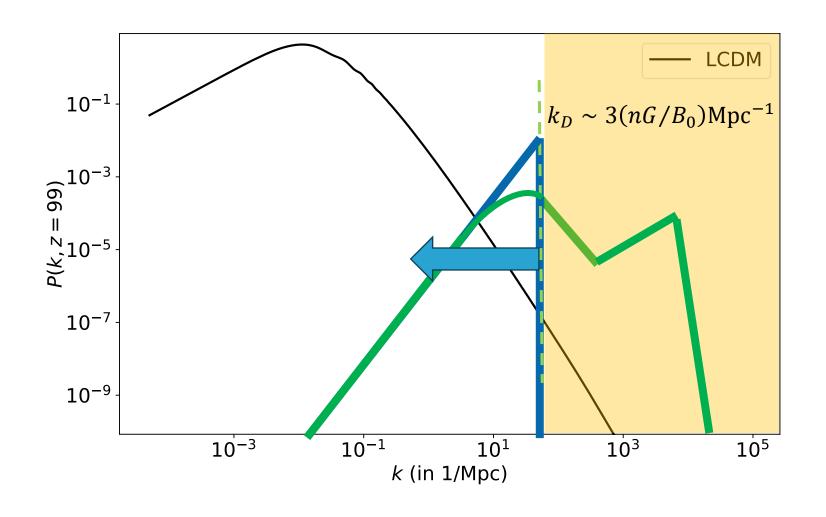
FINDING: HIGHLY ENHANCED POWER SPECTRUM BELOW JEANS SCALE



PART 1: DARK MATTER MINIHALOS BELOW JEANS SCALE



PART 2: LARGE SCALES RELEVANT FOR JWST

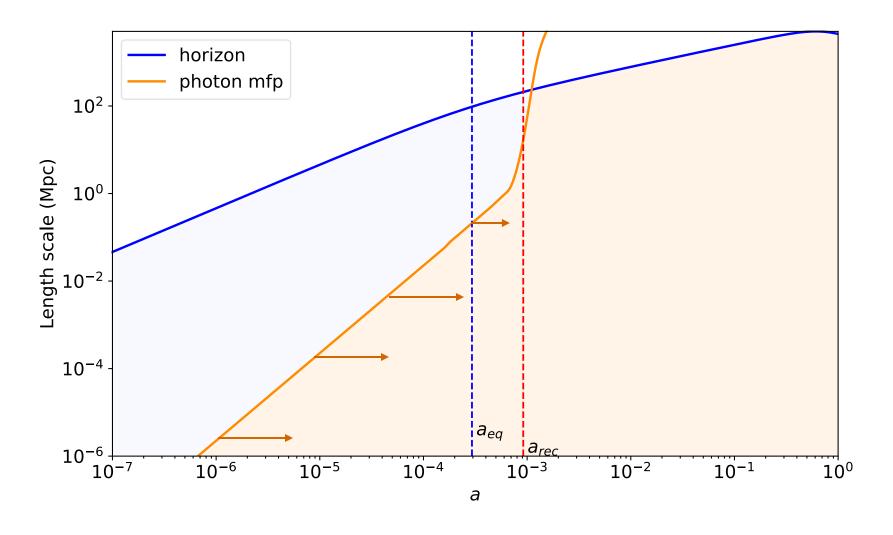


PART 1

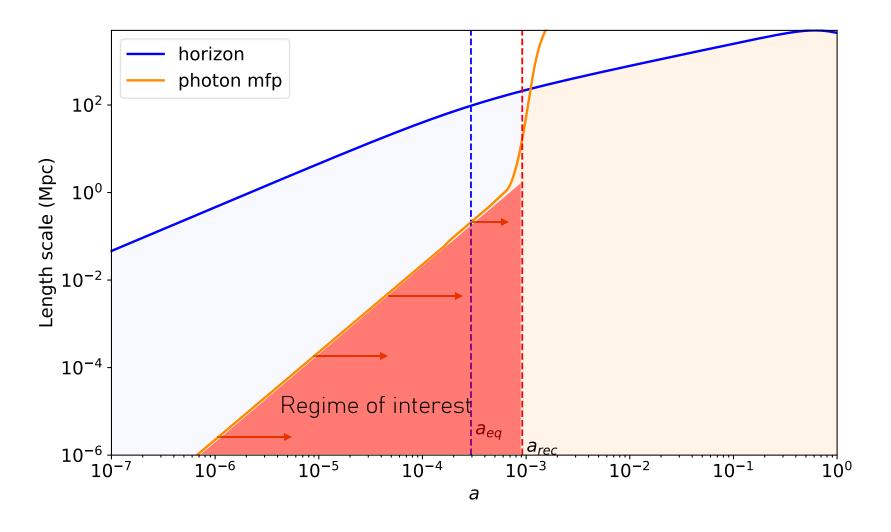
Probing Primordial magnetic fields through dark matter minihalos

ARXIV: 2303.11861

SCALES OF INTEREST: PRE-RECOMBINATION AND SCALES SMALLER THAN PHOTON MFP



SCALES OF INTEREST: PRE-RECOMBINATION AND SCALES SMALLER THAN PHOTON MFP



IDEAL MHD IN PHOTON DRAG REGIME

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b.\nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

IDEAL MHD IN PHOTON DRAG REGIME

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b, \nabla)\vec{v}_b}{\alpha} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi\alpha\rho_b} - \frac{c_b^2}{4\pi\alpha\rho_b}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a} =$$



PRE-RECOMBINATION IDEAL MHD: MAGNETIC FIELDS INFLUENCE BY BARYON FLOW

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b.\nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

PRE-RECOMBINATION IDEAL MHD: BARYONS PUSHED BY LORENTZ FORCE

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b \cdot \nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

PRE-RECOMBINATION IDEAL MHD: REMAINING EQUATIONS SAME!

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b.\nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

PRE-RECOMBINATION IDEAL MHD: LARGE PHOTON DRAG

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b \cdot \nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

PRE-RECOMBINATION IDEAL MHD: LARGE PHOTON DRAG MAKES FLOW LAMINAR

$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b + \frac{(\vec{v}_b \cdot \nabla)\vec{v}_b}{a} = \frac{(\nabla \times \vec{B}) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{\partial a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

Jedamzik and Abel 2013

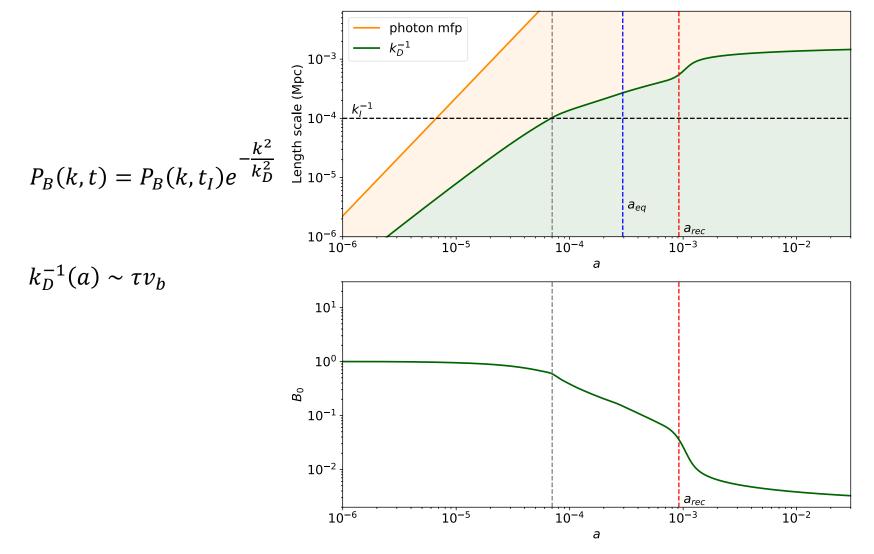
CAN ANALYTICALLY SOLVE MHD EQS: VISCOUS DAMPING

$$P_B(k,t) = P_B(k,t_I)e^{-\frac{k^2}{k_D^2}}$$

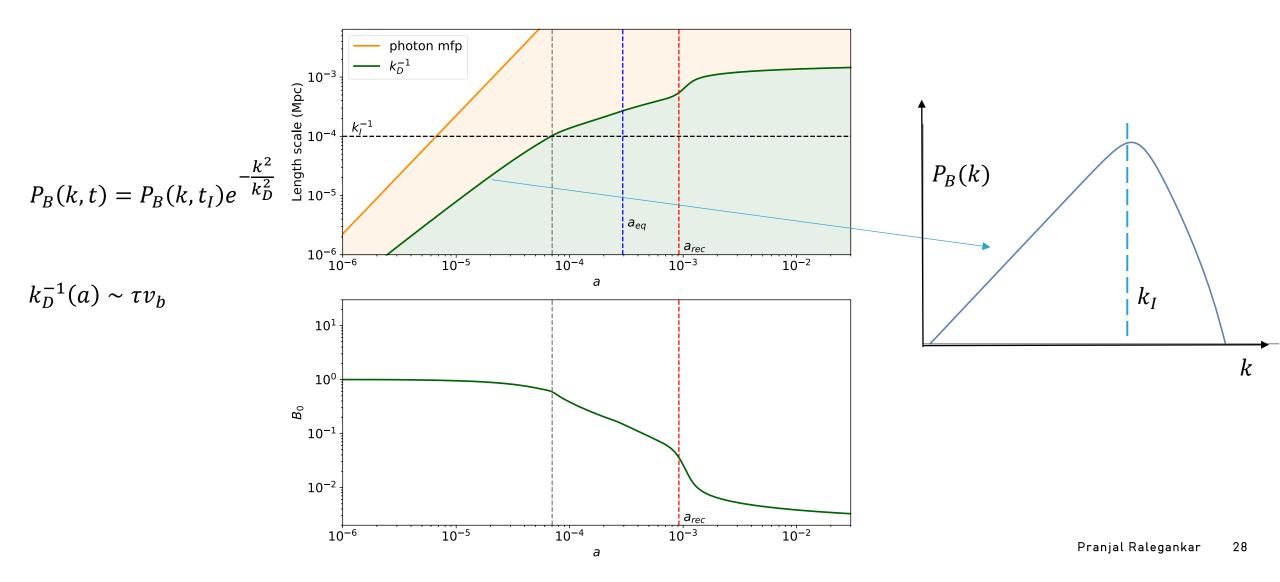
$$k_D^{-1}(a) \sim \tau v_b$$

Assumed B is always Gaussian!

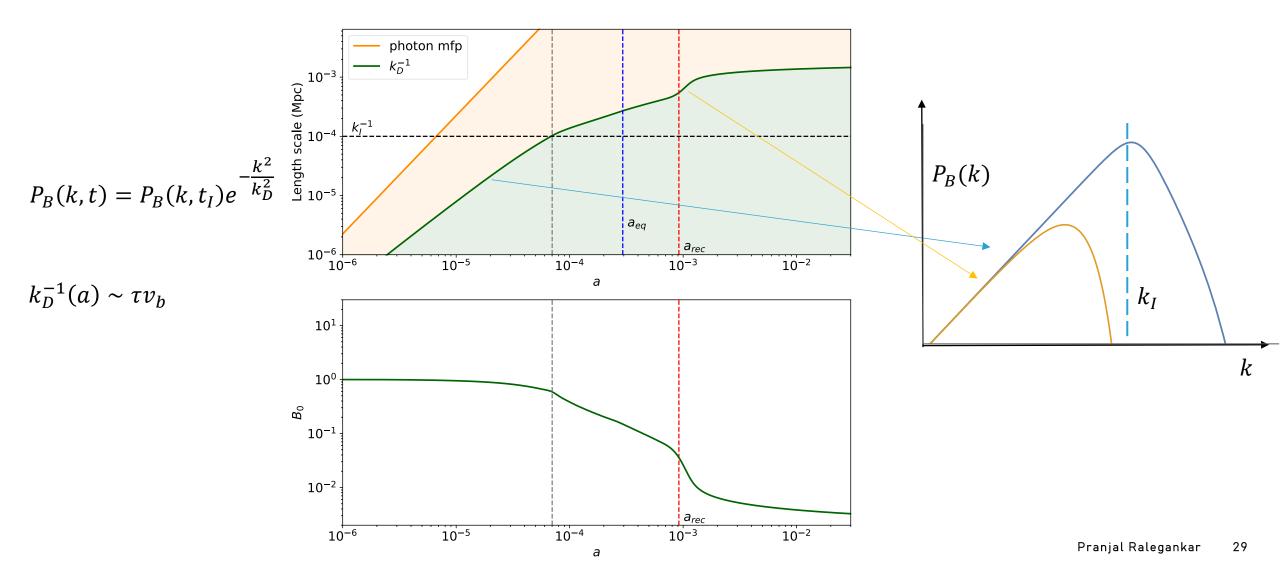
EVOLUTION OF MAGNETIC DAMPING SCALE



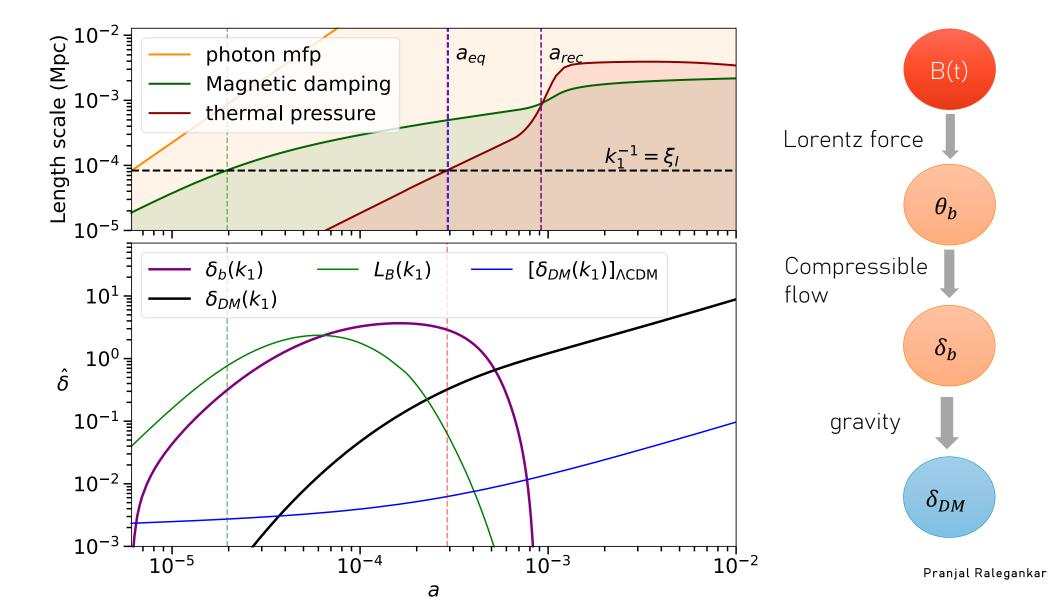
EVOLUTION OF MAGNETIC DAMPING SCALE



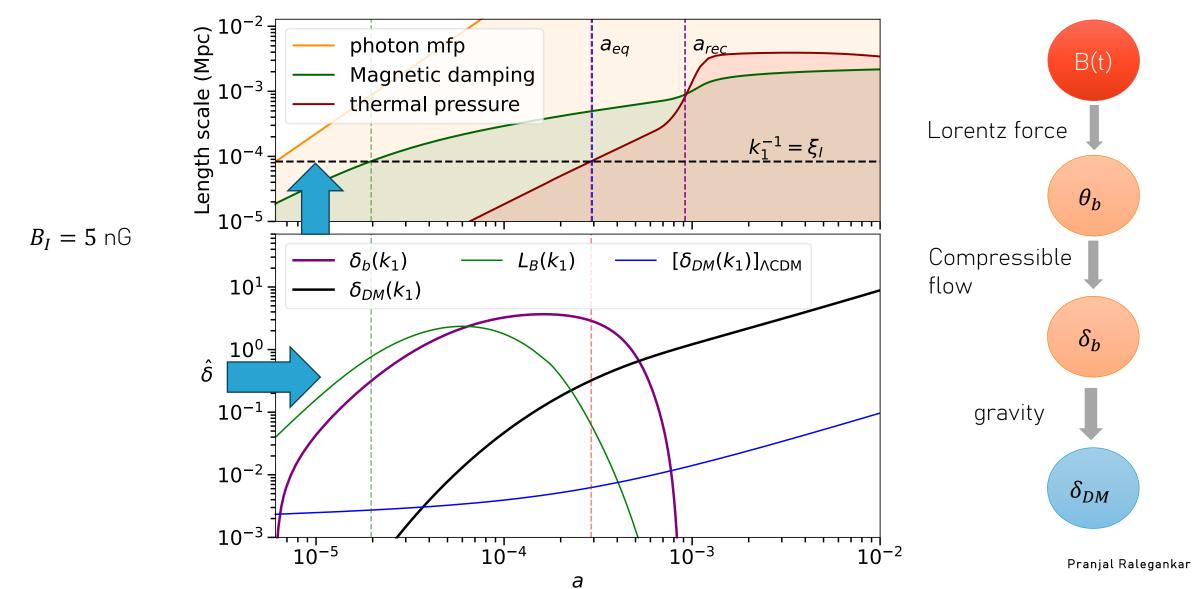
EVOLUTION OF MAGNETIC DAMPING SCALE



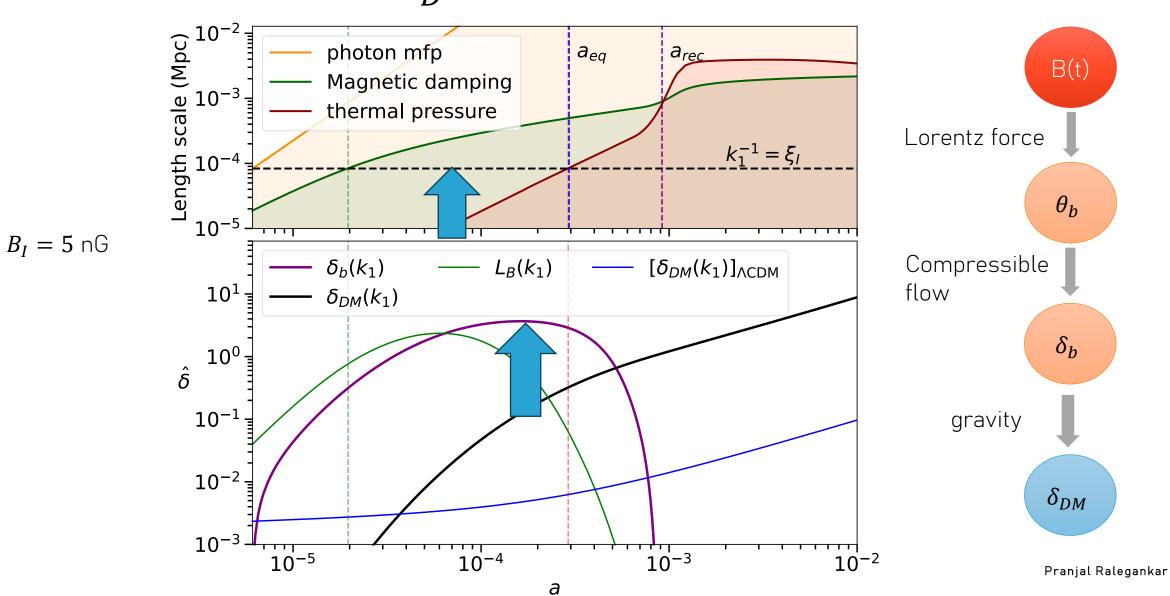
PERTURBATION EVOLUTION PLOT



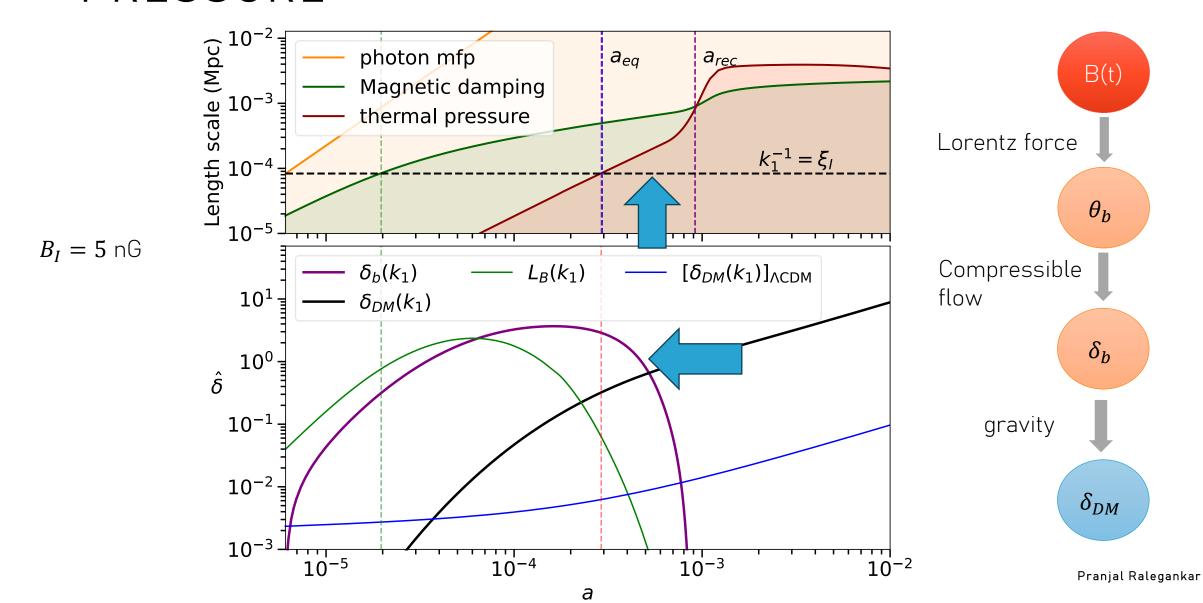
LORENTZ FORCE ENHANCES BARYON PERTURBATIONS FOR MODES OUTSIDE k_D^{-1}



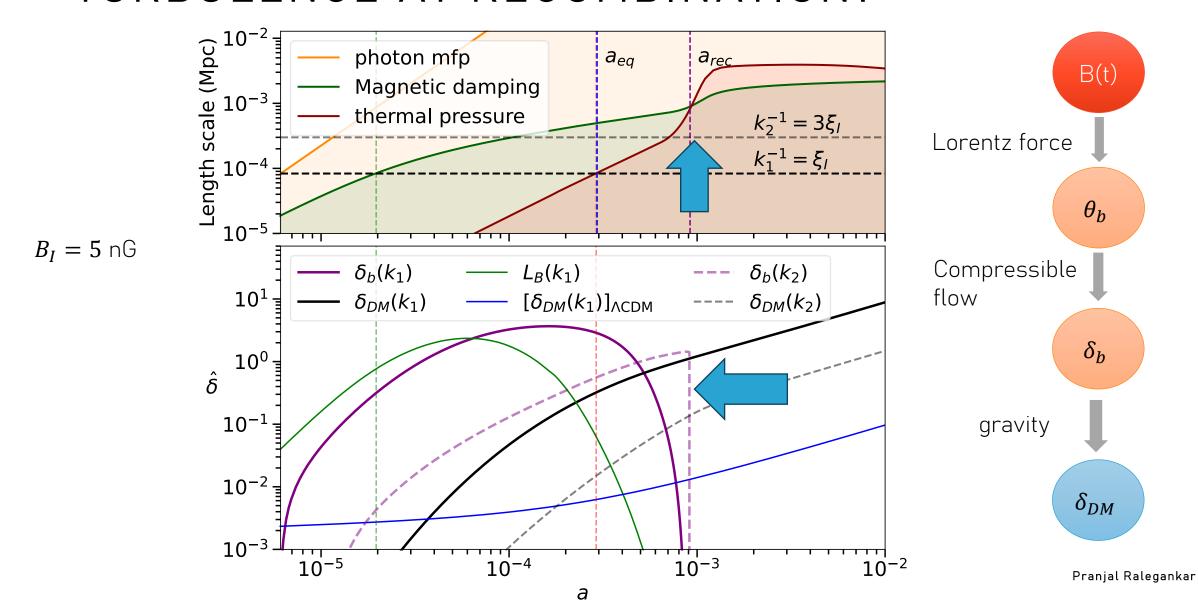
BARYON PERTURBATIONS ASYMPTOTE ONCE MODE ENTERS k_D^{-1}



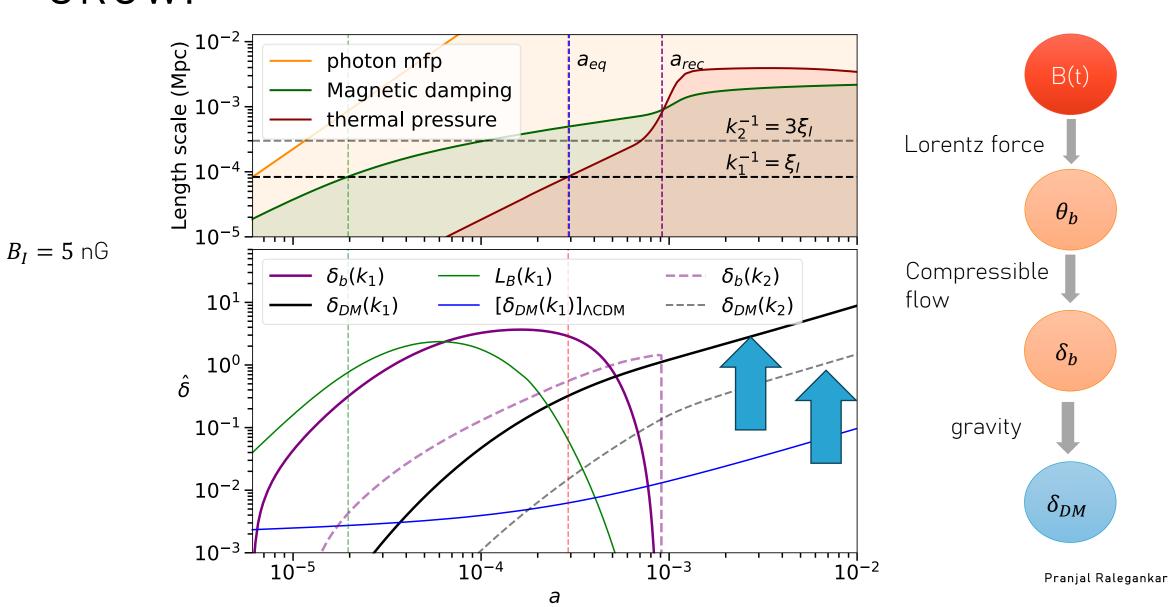
BARYON PERTURBATIONS DAMPED BY THERMAL PRESSURE



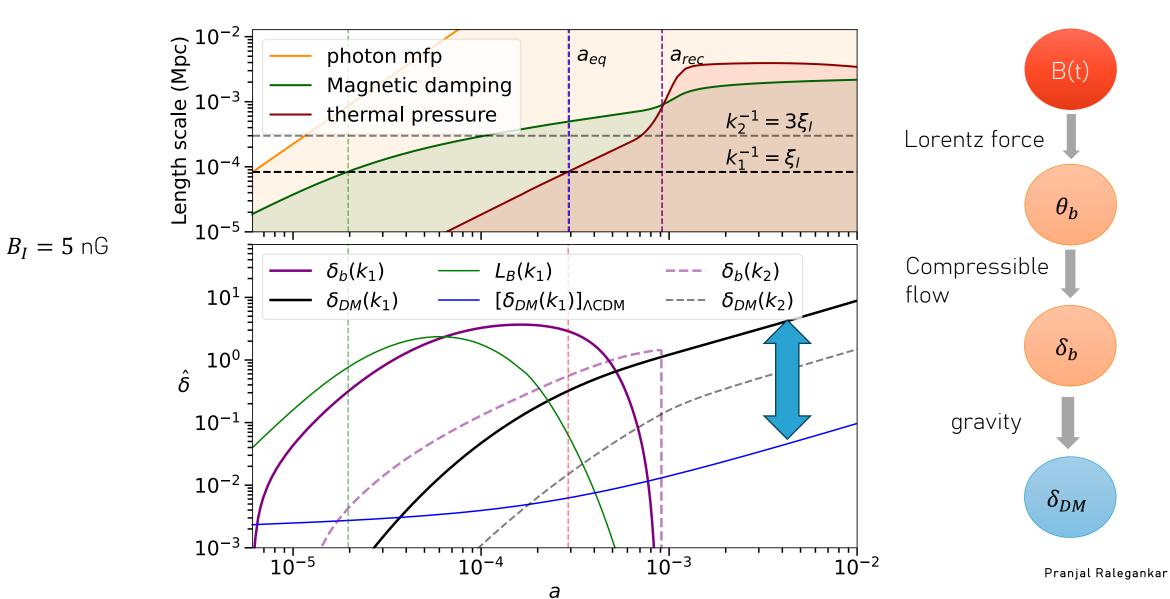
BARYON PERTURBATIONS DAMPED BY TURBULENCE AT RECOMBINATION?



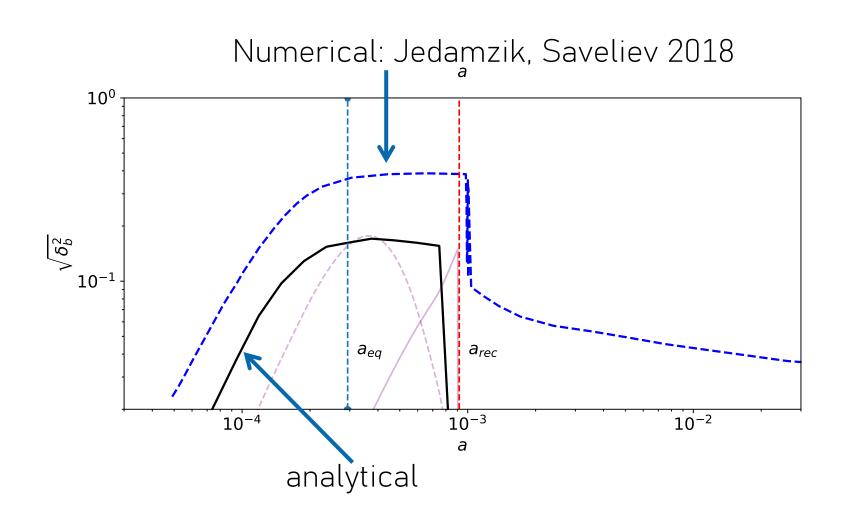
DARK MATTER PERTURBATIONS CONTINUES TO GROW!



DARK MATTER PERTURBATIONS ENHANCED BY ORDERS OF MAGNITUDE COMPARED TO ACDM

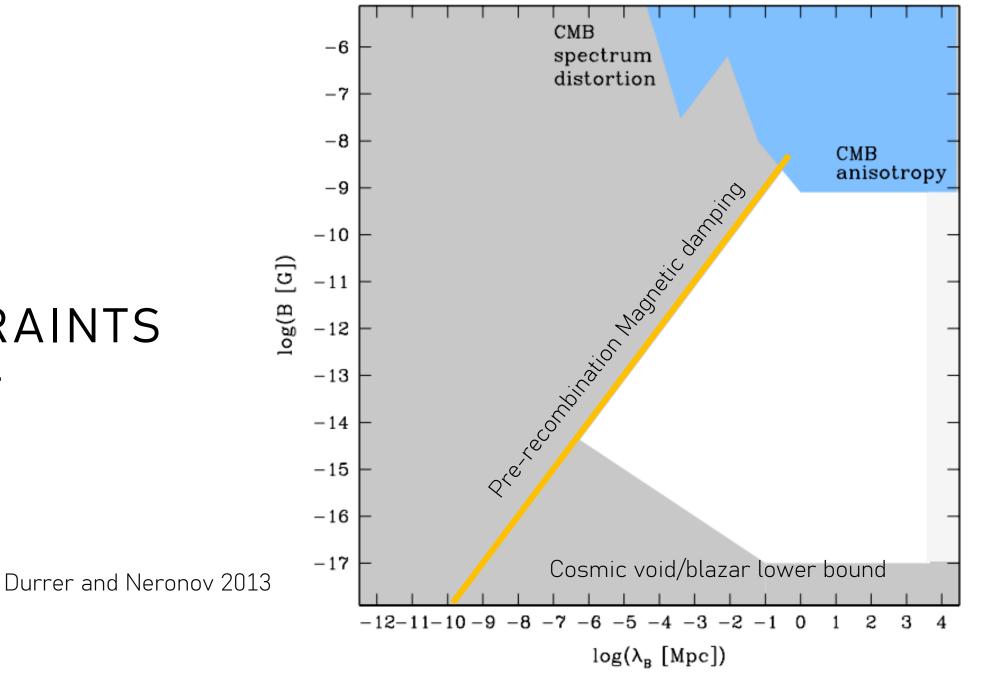


COMPARING WITH SIMULATIONS: ANALYTICAL NOT THAT BAD

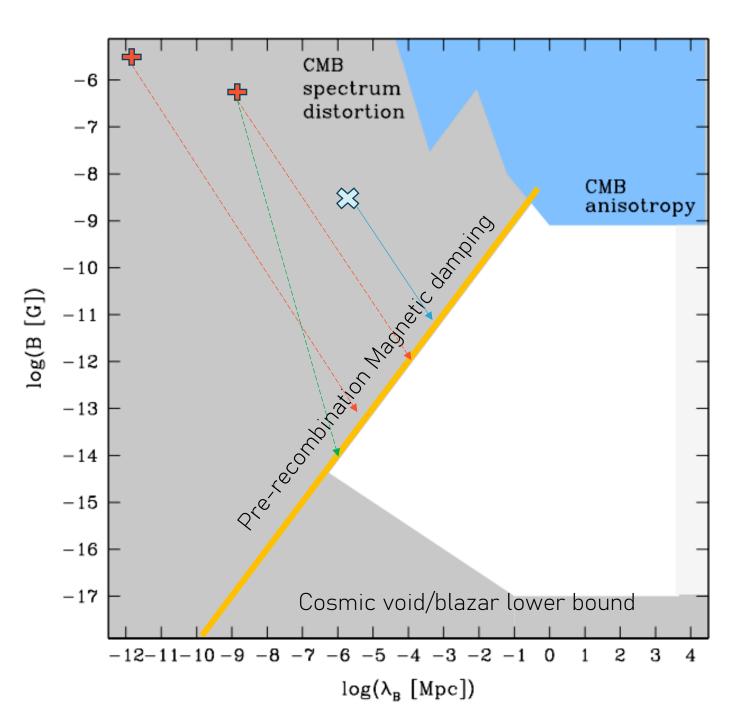


 $B_{0I} = 0.525$ nG

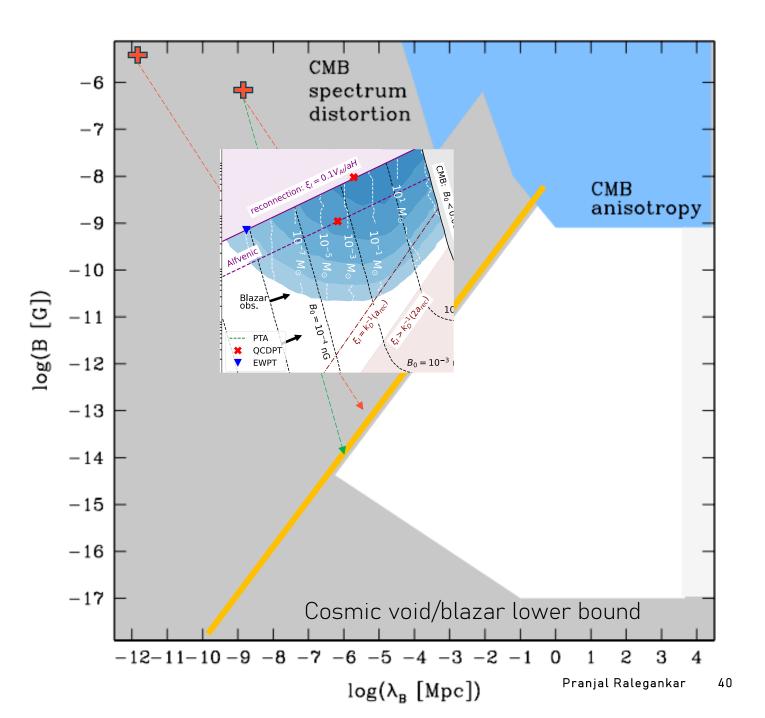
CONSTRAINTS ON PMF



EVOLUTION OF
EARLY
UNIVERSE
PMFS

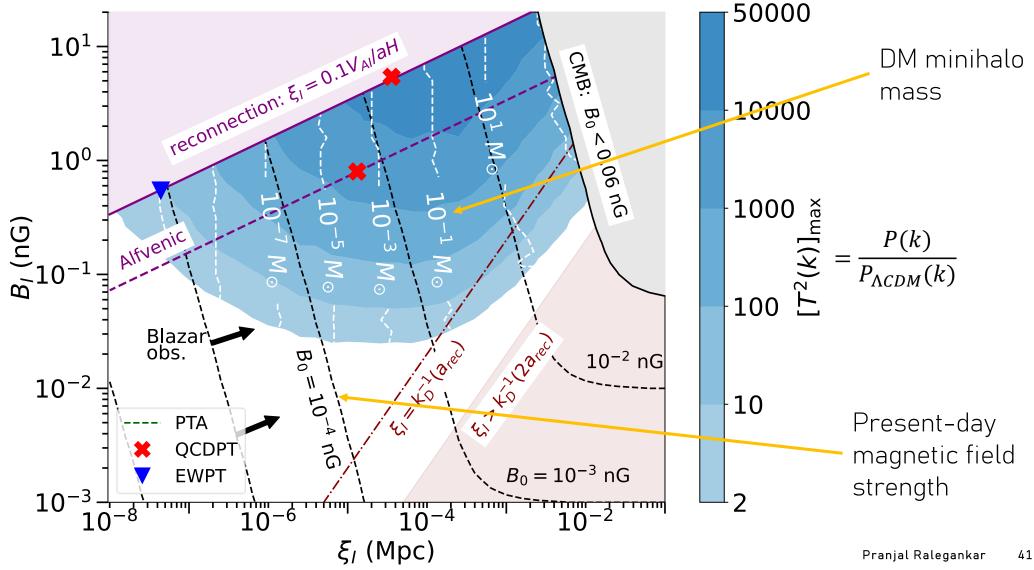


RELEVANCE
OF DARK
MATTER
MINIHALO
GENERATION



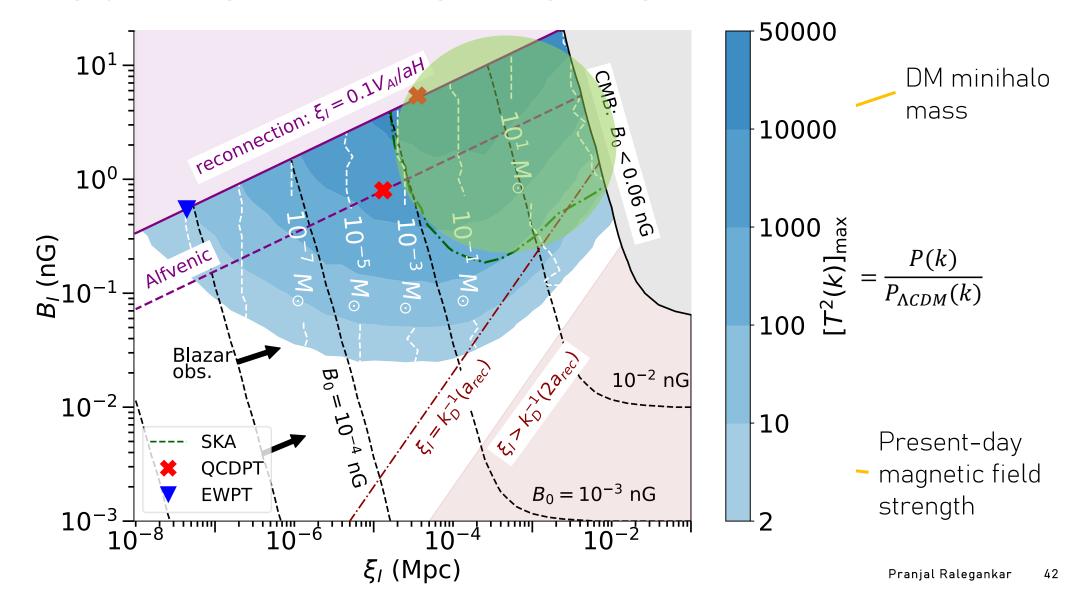
PARAMETER SPACE WITH ENHANCED POWER ON SMALL SCALES

Subscript I refers to the time at the beginning of photon drag regime



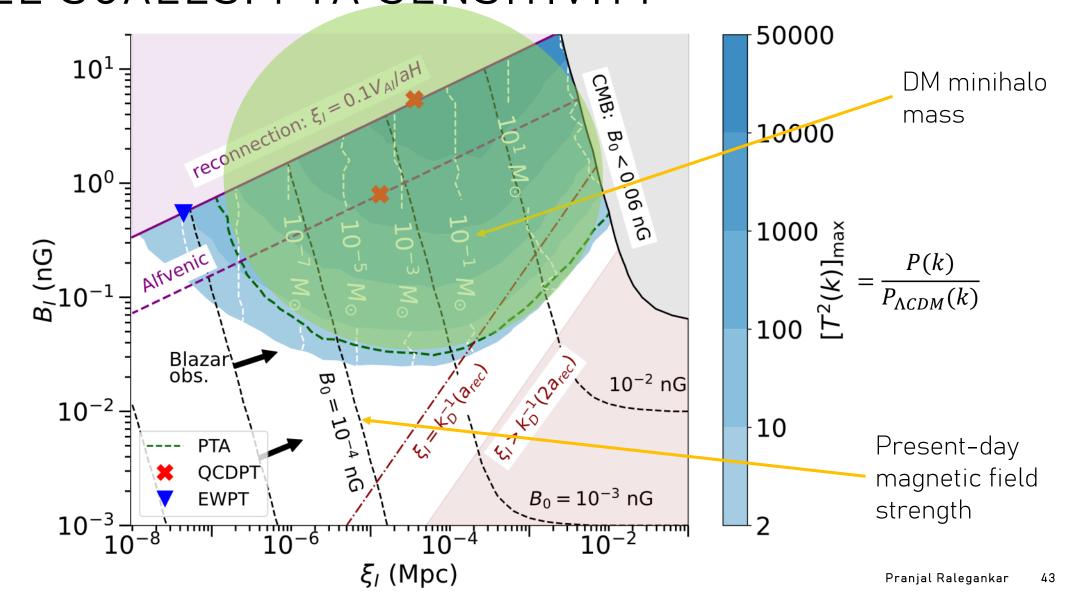
PARAMETER SPACE WITH ENHANCED POWER ON SMALL SCALES: THEIA SKA SENSITIVITY

Subscript *I* refers to the time at the beginning of laminar flow regime

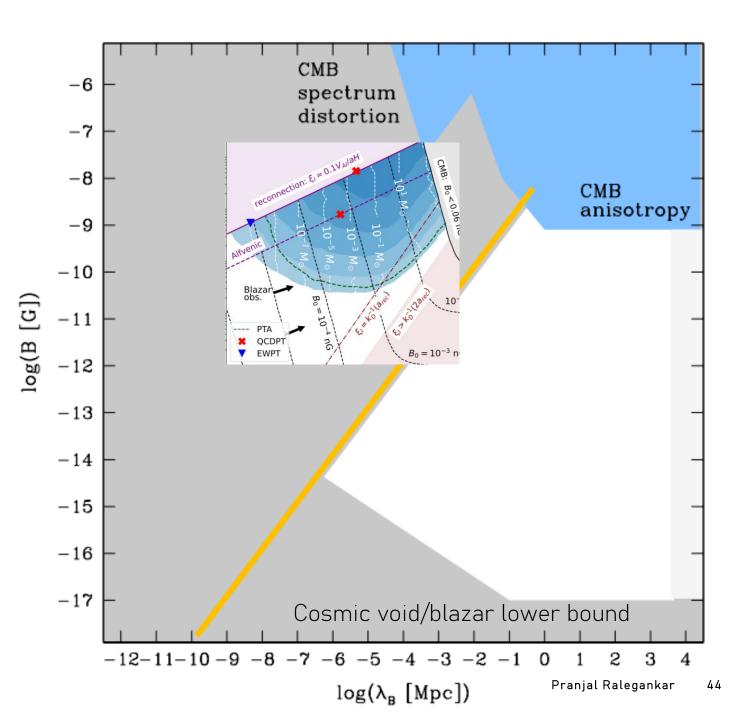


PARAMETER SPACE WITH ENHANCED POWER ON SMALL SCALES: PTA SENSITIVITY

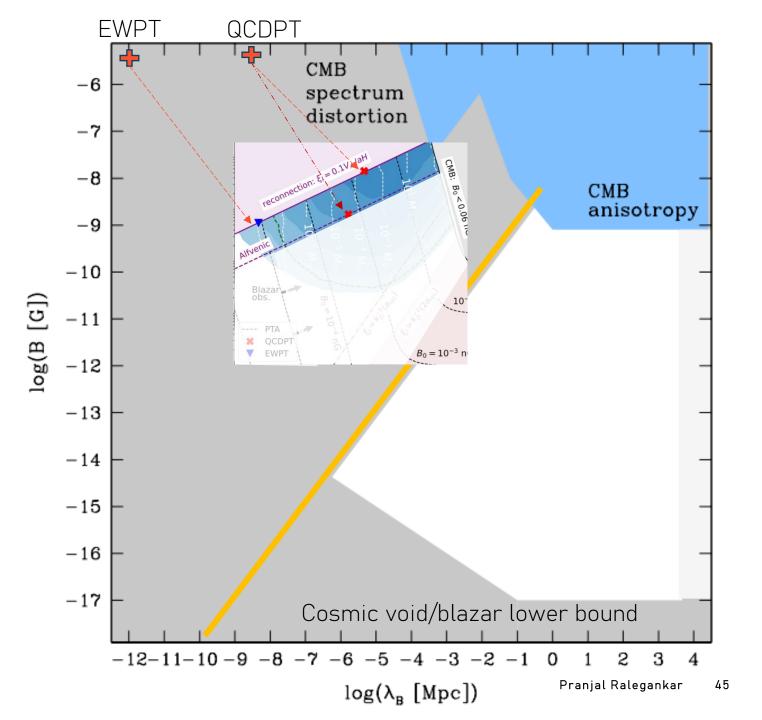
Subscript *I* refers to the time at the beginning of laminar flow regime



MINIHALOS
FROM
CAUSALLY
GENERATED
PMFS

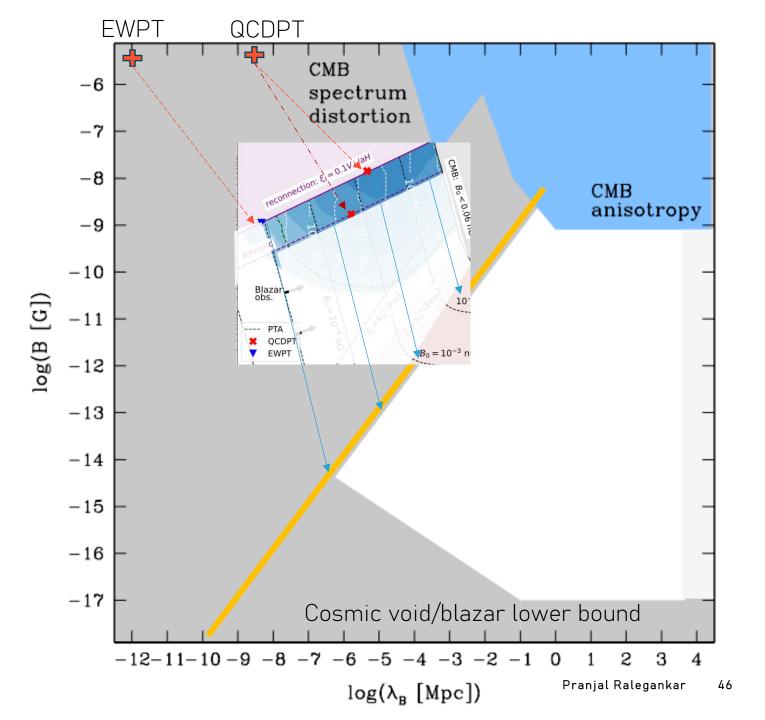


MINIHALOS
FROM
CAUSALLY
GENERATED
PMFS



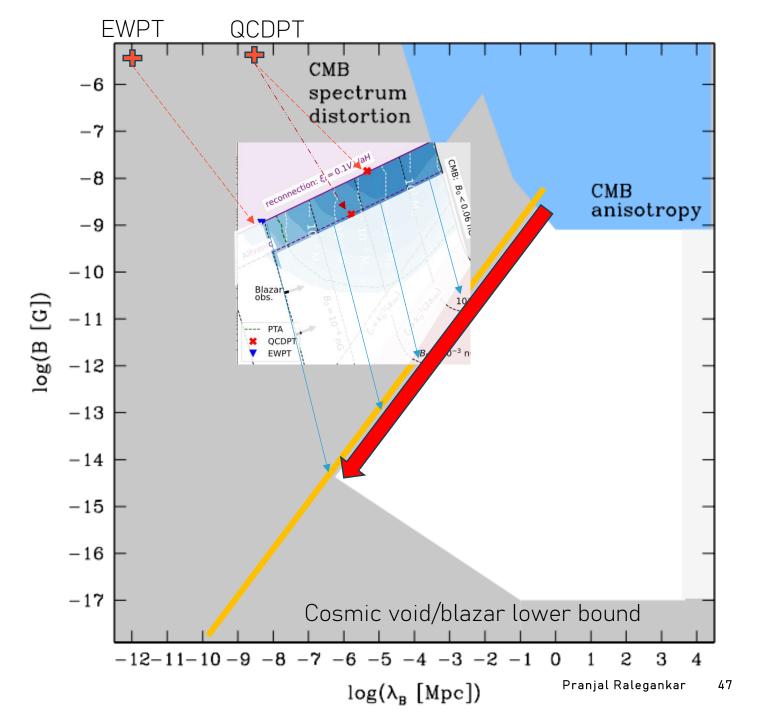
PMFS TO
EXPLAIN
COSMIC VOID
OBSERVATIONS

Assuming Batchelor spectrum!



UNIVERSE
MAYBE FILLED
WITH DARK
MATTER
MINIHALOS!!

Assuming Batchelor spectrum!



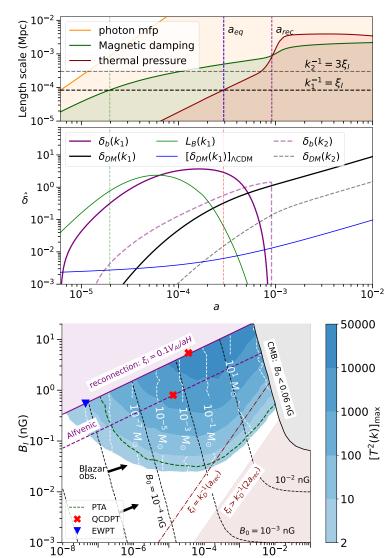
SUMMARY AND CONCLUDING REMARKS

• Magnetic fields can enhance power dark matter power spectrum below magnetic Jeans scale.

 PTA/GAIA detection of DM minihalos can provide best probe of primordial magnetic fields

 Results are qualitative: Need MHD simulations to get accurate quantitative answers.

• Ironic: how invisible dark matter can help look for visible entity: magnetic fields



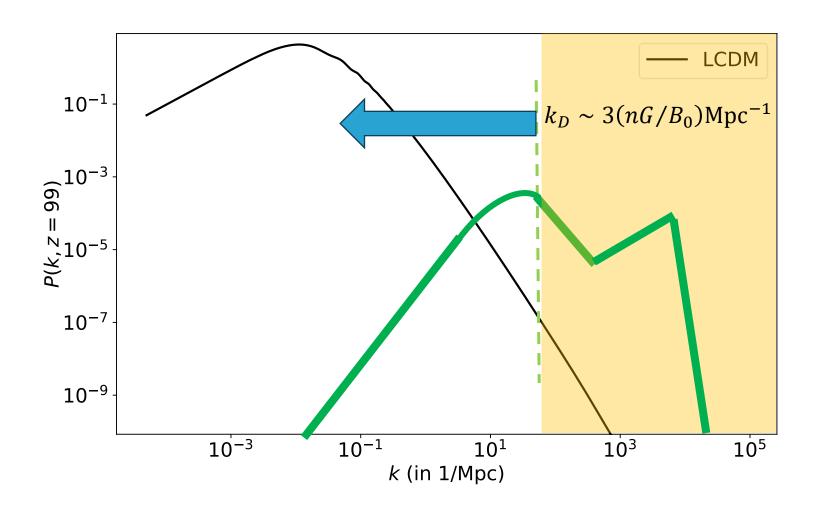


PART 2

Baryon fraction enhanced on Large scales

Arxiv: 2402.14079

PART 2: LARGE SCALES RELEVANT FOR JWST



POST-RECOMBINATION IDEAL MHD

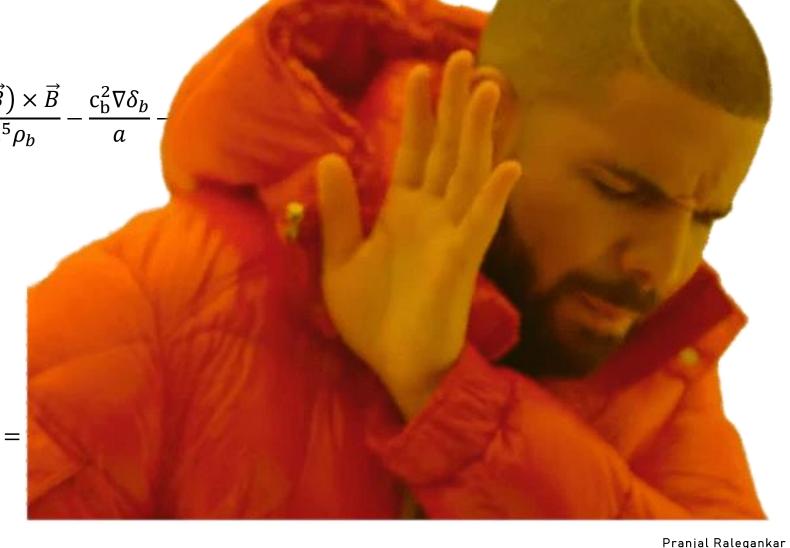
$$\frac{\partial (\vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times \vec{B})}{a}$$

$$\frac{\partial \vec{v}_b}{\partial t} + H\vec{v}_b + \frac{(\vec{v}_b \cdot \nabla)\vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a^5 \rho_b} - \frac{c_b^2 \nabla \delta_b}{a}$$

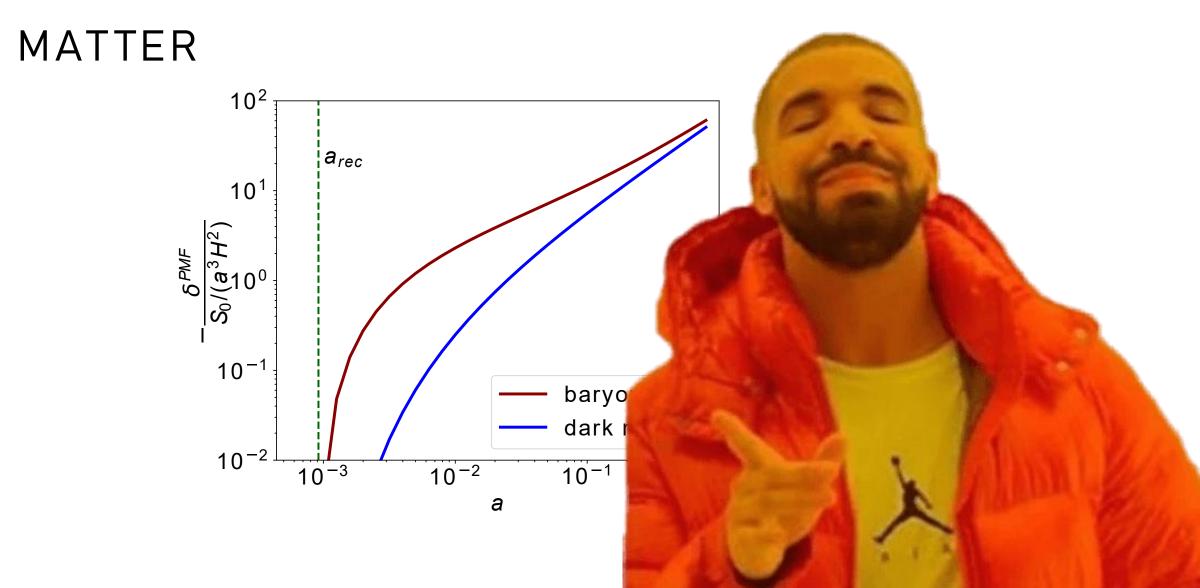
$$\frac{\partial \delta_b}{\partial t} = -\frac{\nabla \cdot \vec{v}_b}{a} - \frac{\nabla \cdot (\delta_b \vec{v}_b)}{a}$$

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

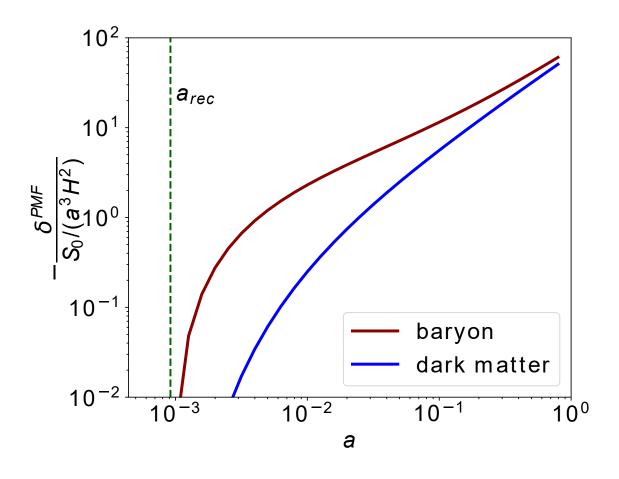
$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \left[\frac{\partial \ln(a^2 H)}{\partial \ln a} + 1 \right] \frac{\partial \delta_{DM}}{a \partial a} =$$



POST RECOMBINATION: BARYON PERTURBATIONS MORE ENHANCED THAN DARK

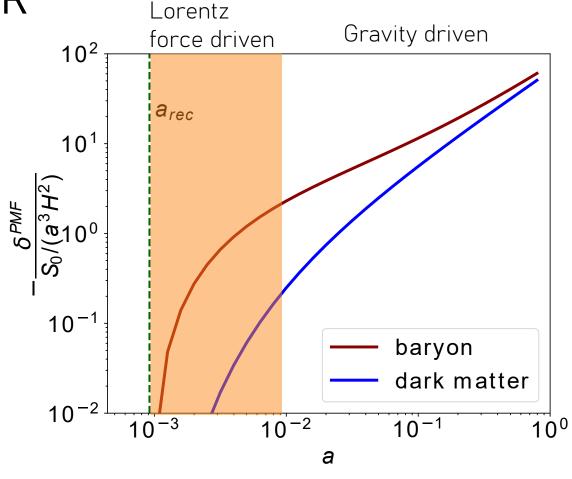


POST RECOMBINATION: BARYON PERTURBATIONS MORE ENHANCED THAN DARK MATTER

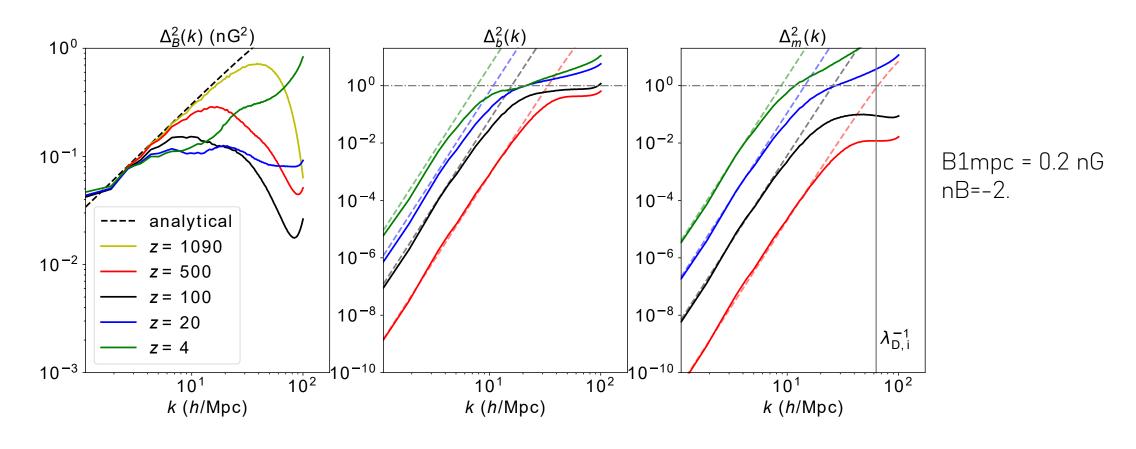


POST RECOMBINATION: BARYON PERTURBATIONS MORE ENHANCED THAN DARK

MATTER

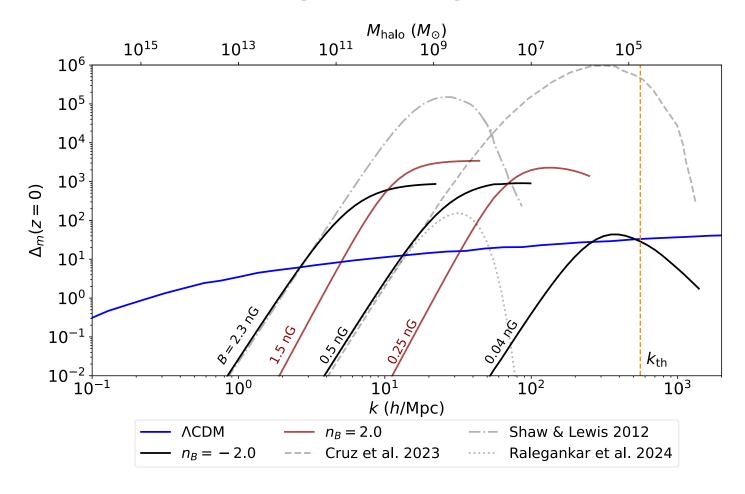


MHD SIMULATIONS: MATCHES ANALYTICAL

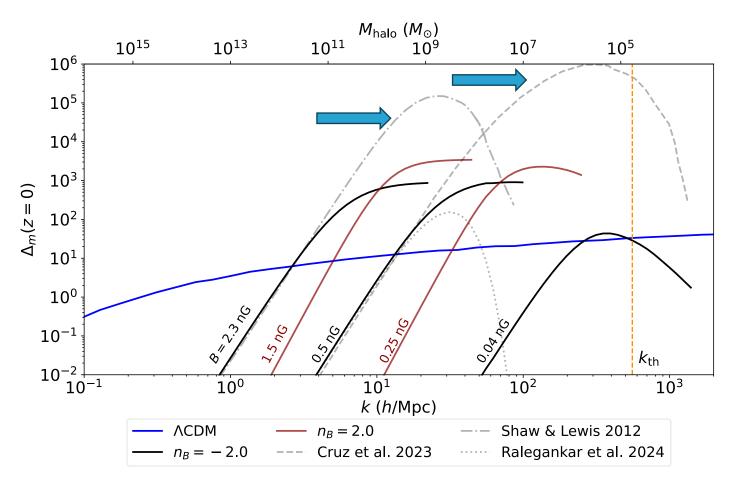


Ralegankar, Garaldi, Viel 2024

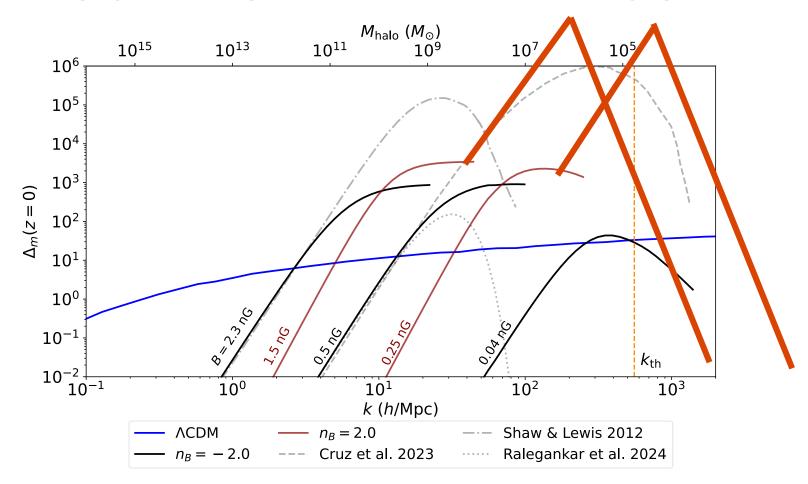
ENHANCEMENT MOVES TO SMALLER SCALES WITH SMALLER PMF STRENGTH



EARLIER ANALYTICAL STUDIES OVER-ESTIMATED MAGNETIC JEANS SCALE

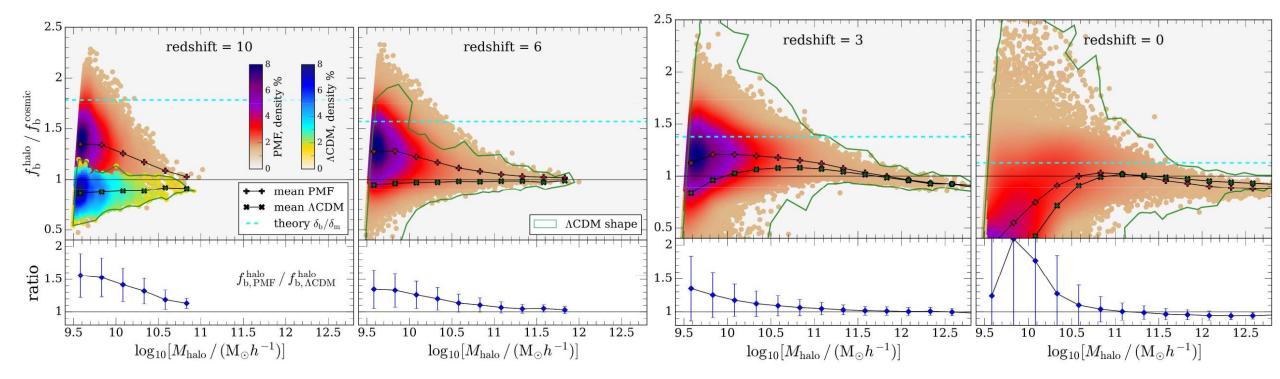


SMALLER SCALES: DM MINIHALOS



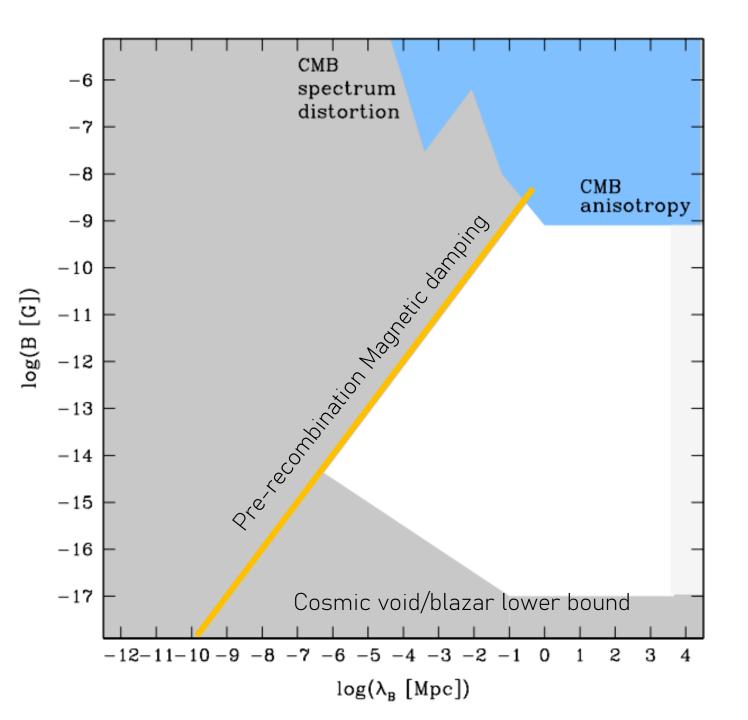
IMPLICATIONS FOR LARGE SCALES: ENHANCED BARYON FRACTION IN HALOS

IMPLICATIONS FOR LARGE SCALES: ENHANCED BARYON FRACTION IN HALOS

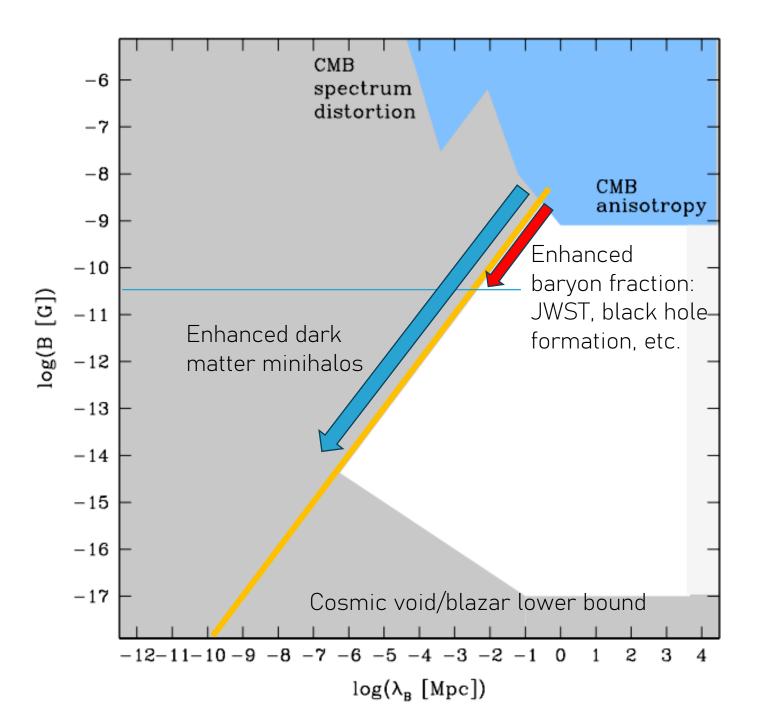


Scale invariant 1 nG PMFs

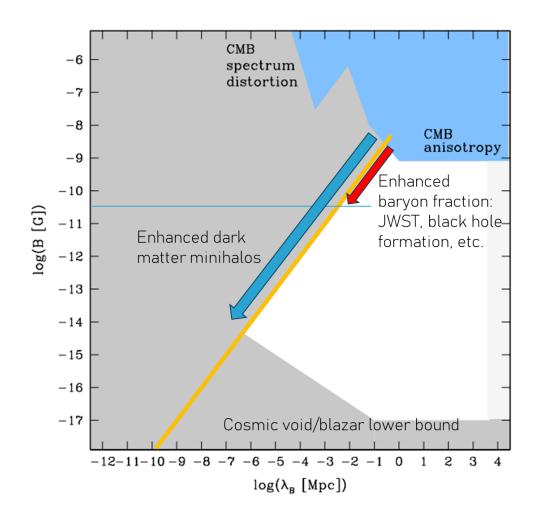
IMPLICATIONS FOR PMFS

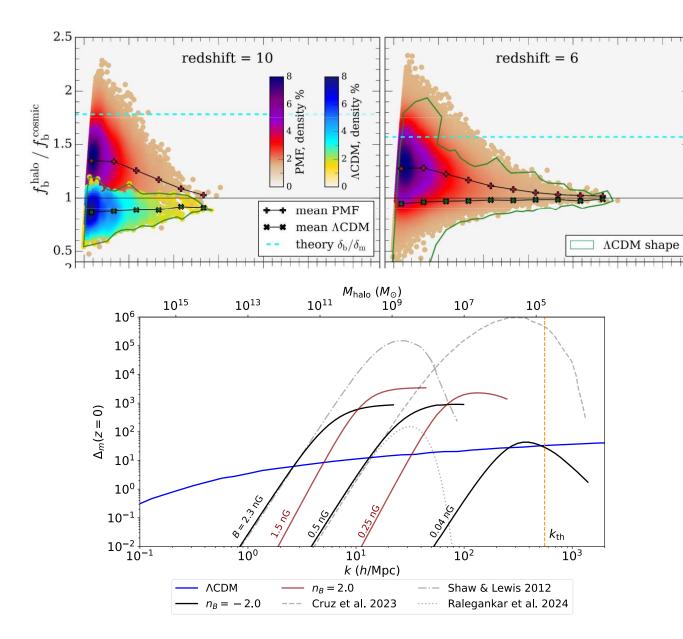


POWER **SPECTRUM ABOVE** MAGNETIC JEANS SCALE IS SENSITIVE **UPTO 0.05 NG PMFS**



SUMMARY





BACKUP SLIDES

BACKUP: ANALYTIC DERIVATION IN PRE-RECOMBINATION FLUID

IDEAL MHD IN PHOTON DRAG REGIME:

IDEAL MHD IN PHOTON DRAG REGIME: LAMINAR FLOW IN BARYONS

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha) \vec{v}_b + \frac{(\vec{v}_b, \nabla) \vec{v}_b}{a} = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

IDEAL MHD IN PHOTON DRAG REGIME: LAMINAR FLOW IN BARYONS

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

Abel and Jedamzik 2010, Campanelli 2013, Jedamzik and Saveliev 2018

IDEAL MHD IN PHOTON DRAG REGIME: KEY FORCES

$$\frac{\partial \vec{v}_b}{\partial t} + (H + \alpha)\vec{v}_b = \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b} - \frac{c_b^2 \nabla \delta_b}{a} - \frac{\nabla \phi}{a}$$

$$\text{Cravity}$$
Lorentz force
Thermal pressure

IDEAL MHD IN PHOTON DRAG REGIME: LARGE LORENTZ FORCE LIMIT

$$(H + \alpha)\vec{v}_b \approx \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b}$$

IDEAL MHD IN PHOTON DRAG REGIME: LARGE LORENTZ FORCE LIMIT

$$(H + \alpha)\vec{v}_b \approx \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b}$$

$$\frac{\partial \ (a^2 \vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times a^2 \vec{B})}{a}$$

IDEAL MHD IN PHOTON DRAG REGIME: MAGNETIC DAMPING SCALE

$$(H + \alpha)\vec{v}_b \approx \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b}$$

$$\frac{\partial (a^2 \vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times a^2 \vec{B})}{a}$$

$$P_B(k,t) = P_B(k,t_I)e^{-\frac{k^2}{k_D^2}}$$

$$k_D^{-1}(a) \sim \tau v_b$$

Campanelli 2013

IDEAL MHD IN PHOTON DRAG REGIME: MAGNETIC DAMPING SCALE

$$(H + \alpha)\vec{v}_b \approx \frac{\left(\nabla \times \vec{B}\right) \times \vec{B}}{4\pi a \rho_b}$$
$$\frac{\partial (a^2 \vec{B})}{\partial t} = \frac{\nabla \times (\vec{v}_b \times a^2 \vec{B})}{a}$$

$$P_B(k,t) = P_B(k,t_I)e^{-\frac{k^2}{k_D^2}}$$

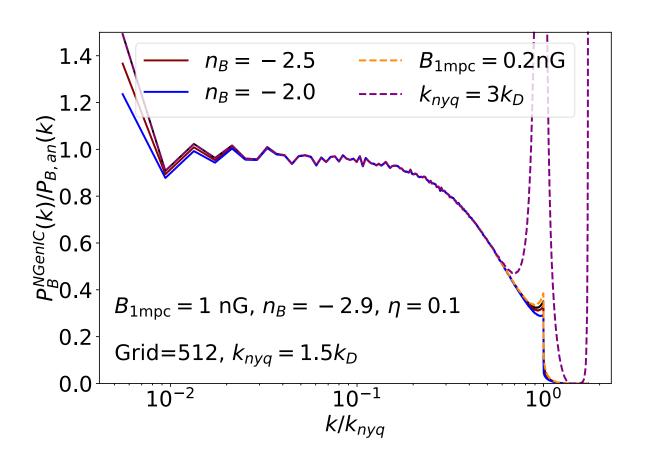
$$k_D^{-1}(a) \sim \tau v_b$$

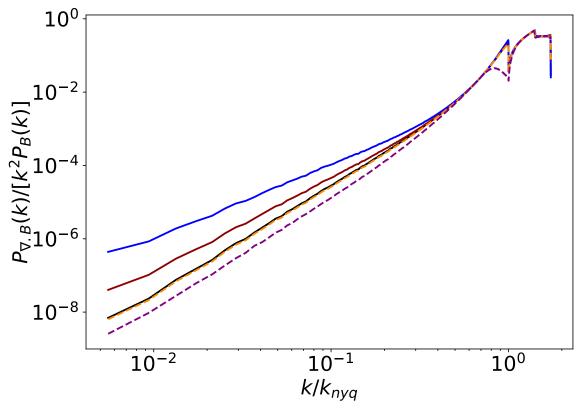
Campanelli 2013

ASSUMED B_0 Gaussian

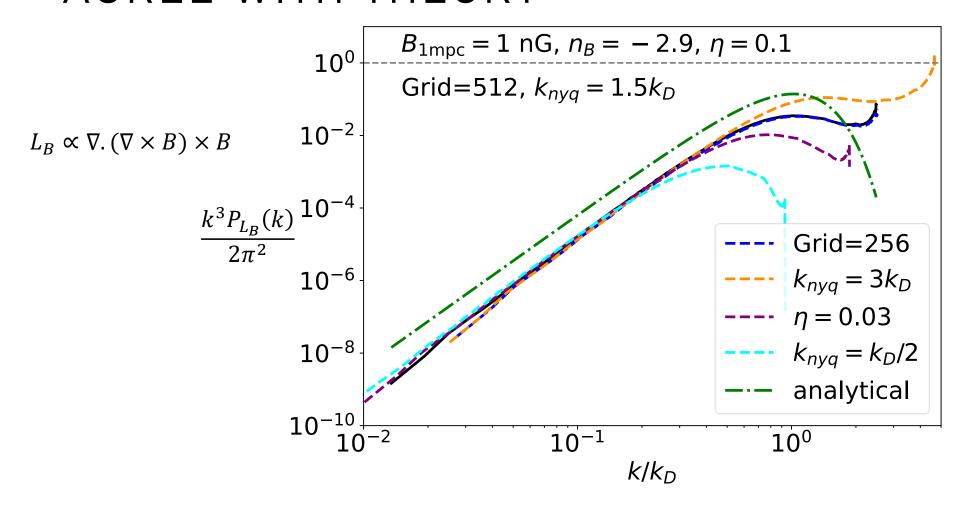
PROBLEM WITH LORENTZ FORCE IN MY LATTICE

INITIALIZING STOCHASTIC PMFS ON LATTICE

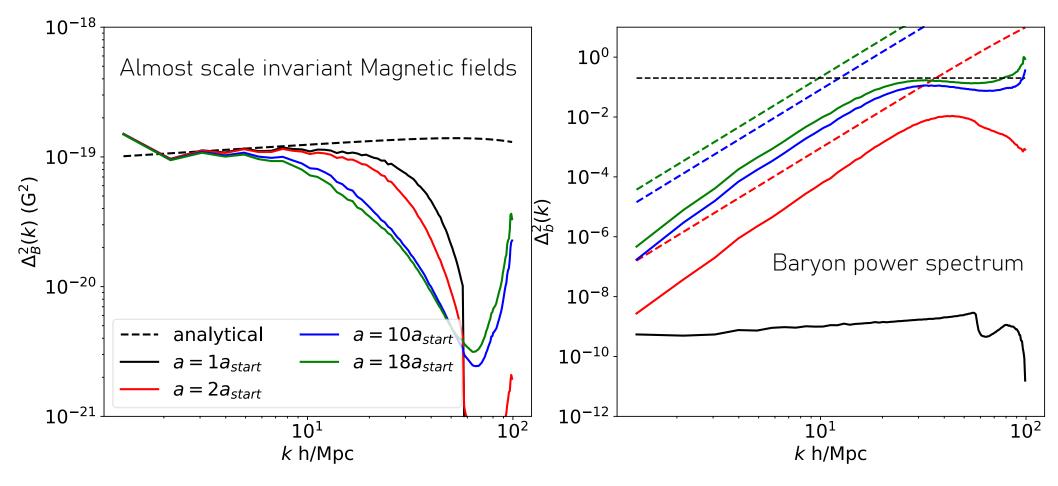




LORENTZ FORCE POWER SPECTRUM DOESN'T AGREE WITH THEORY

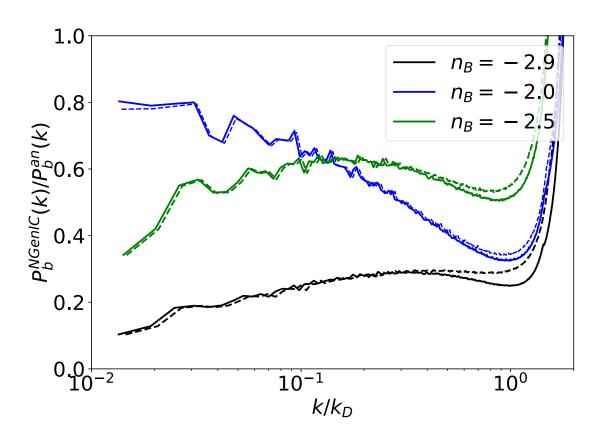


THE SUPPRESSION OF POWER IS ALSO SEEN IN AREPO (PRELIMINARY!!)



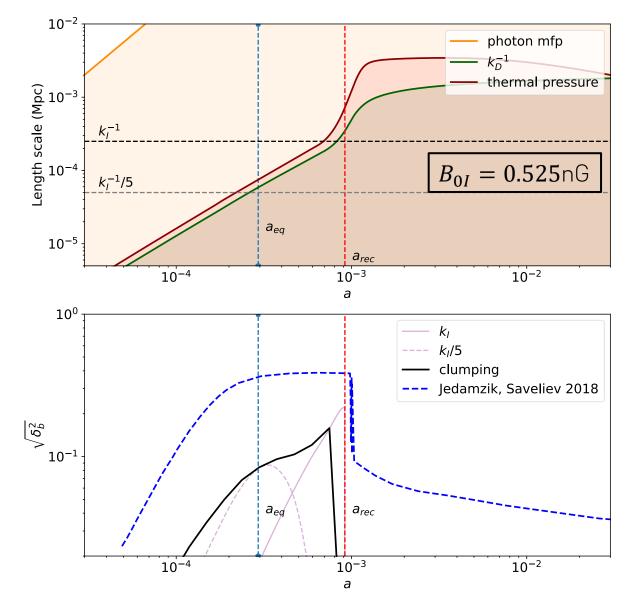
BACKUP SLIDES

SPECTRUM SHAPE DEPENDENCE OF SHIFT

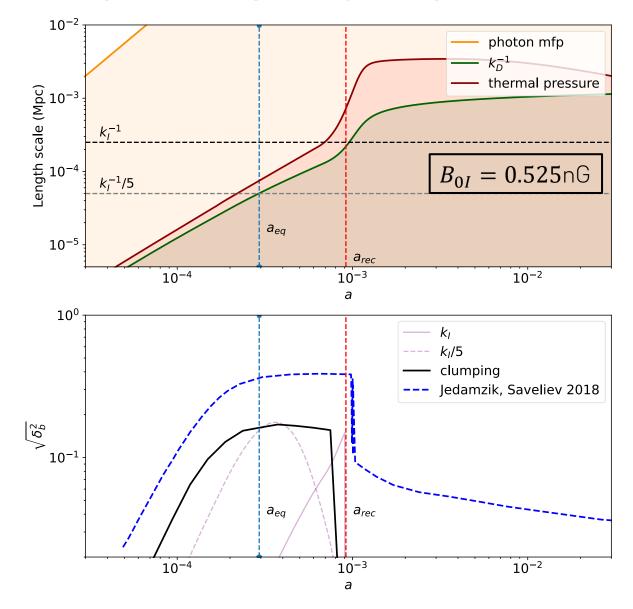


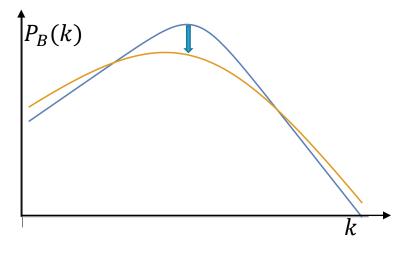
COMPARING WITH FULL MHD SIMULATIONS

COMPARING WITH SIMULATIONS: SENSITIVE TO INITIAL POWER SPECTRUM

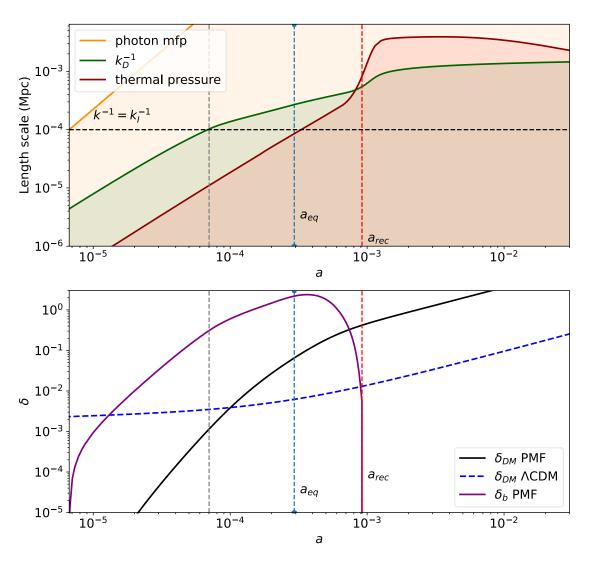


COMPARING WITH SIMULATIONS: SENSITIVE TO INITIAL POWER SPECTRUM

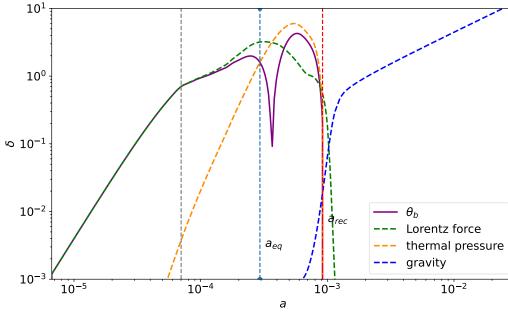




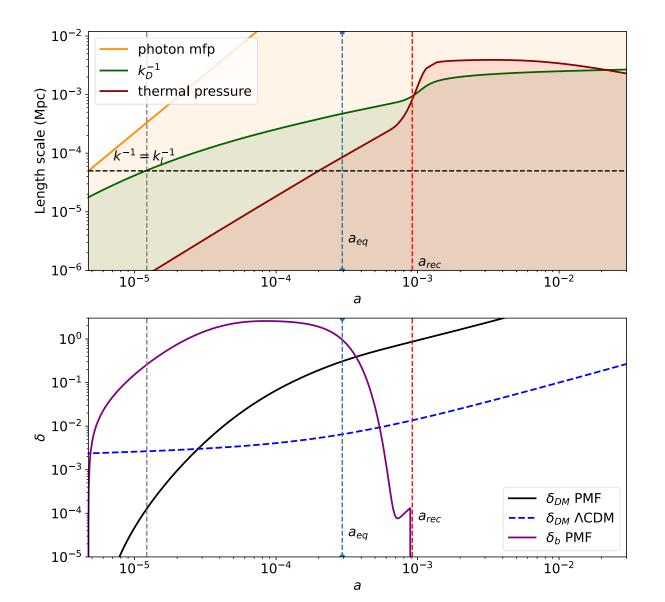
MORE PERTURBATION PLOTS



$$B_0 = 1$$
nG $k_I = 10^4 \ Mpc^{-1}$



MORE PERTURBATION PLOTS



$$B_0 = 8$$
nG $k_I = 10^4 \ Mpc^{-1}$

