

# Physical Hamiltonians, the Zeros of L-Functions, and a Hybrid Quantum–Classical Search Pipeline

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## Abstract

**Short answer:** Yes—in principle. If one finds (or reliably engineers) a self-adjoint physical Hamiltonian whose spectrum equals the nontrivial zeros of an  $L$ -function, then measuring that system’s spectrum provides direct physical access to the zeros. Because the zeros determine prime-counting via explicit formulas, the primes or their distribution can be recovered from the measured spectrum instead of by evaluating the  $L$ -function numerically.

This document explains the potential, the obstacles, and a concrete hybrid pipeline (KAN  $\rightarrow$  VQE  $\rightarrow$  Krylov/SBQD  $\rightarrow$  FFT  $\rightarrow$  QPE) to search for and validate candidate Hamiltonians.

## 1 Why this would be transformative

1. **Direct physical measurement of zeros.** Quantum phase estimation (QPE) or high-resolution spectroscopy on a system whose eigenvalues coincide with the zeros will output those zeros (or highly accurate approximations) without evaluating  $\zeta(s)$  numerically.
2. **Primes follow from zeros.** The explicit trace/Weil formulas relate prime-counting functions to sums over zeros. Substituting measured zeros into those formulas reconstructs prime-distribution data (e.g.  $\pi(x)$ , Chebyshev functions) — effectively a “physical prime oracle.”
3. **New viewpoint / potential proofs.** A concrete self-adjoint operator whose spectrum equals the zeros would amount to a Hilbert–Pólya style demonstration and provide deep conceptual payoff beyond computation.

## 2 Why it is not automatic or practical yet

- **No known explicit Hamiltonian.** Several proposals exist (Berry–Keating, scattering/trace-formula models, quantum graphs) but no accepted self-adjoint Hamiltonian reproducing the nontrivial zeros exactly is known.
- **Statistics vs. identity.** Many physical systems reproduce local statistics of zeros (GUE spacings), but statistical similarity does not imply the exact spectral identity required for recovering primes.
- **Precision and resolution.** High-accuracy prime estimates require zeros to many digits. Physical measurement to that precision demands either many ancilla qubits in QPE or extremely fine spectroscopic control.
- **Encoding arithmetic structure.** The Hamiltonian must encode arithmetic features (e.g. Euler product structure) — how to embed that into realistic physical models remains open.

## 3 The Berry–Keating Model and Hermiticity

### 3.1 The bare operator

The original Berry–Keating proposal considers the classical Hamiltonian  $H = xp$ . Upon quantization in position representation, this becomes

$$\hat{H} = -i\hbar x \frac{d}{dx}.$$

However, this operator is not Hermitian (self-adjoint) on  $L^2(\mathbb{R})$ . Integration by parts shows an extra term appears in the inner product, violating  $\langle \phi | \hat{H} \psi \rangle = \langle \hat{H} \phi | \psi \rangle$ . Thus,  $\hat{H} = xp$  is not a valid physical Hamiltonian.

### 3.2 Symmetrized version

To fix Hermiticity, one defines the symmetrized operator

$$\hat{H}_{BK} = \frac{1}{2}(\hat{x}\hat{p} + \hat{p}\hat{x}) = -i\hbar \left( x \frac{d}{dx} + \frac{1}{2} \right),$$

which is formally Hermitian for wavefunctions vanishing at infinity. The spectrum of this operator is continuous (real-valued), so it is physically consistent but not yet discrete.

### 3.3 Self-adjoint extension and boundary regularization

The operator can be made self-adjoint by restricting its domain or imposing boundary conditions:

$$\psi(x_{\min}) = e^{i\theta} \psi(x_{\max}),$$

with finite cutoffs  $x_{\min}, x_{\max}$ . This compactifies the domain, giving a discrete spectrum and turning the operator into a legitimate quantum Hamiltonian. Such regularized or discretized Berry–Keating models are self-adjoint and yield real, discrete eigenvalues.

### 3.4 Discretized matrix realization

For numerical work, the discretized version uses a finite grid  $x_j \in [x_{\min}, x_{\max}]$  and represents  $\hat{p}$  by finite differences. The Hamiltonian matrix is then

$$H_{jk} = \frac{1}{2} (x_j p_{jk} + p_{jk} x_k),$$

where  $p_{jk}$  is anti-Hermitian (imaginary and skew-symmetric), making  $H_{jk}$  Hermitian. This ensures real eigenvalues and enables numerical diagonalization. In the hybrid pipeline, this discretized Hermitian form serves as the base physical model for parameter optimization.

## 4 How the hybrid pipeline helps discover such a Hamiltonian

We propose a search-and-validate workflow that uses a hybrid quantum–classical pipeline to (i) propose plausible parametric Hamiltonians, (ii) measure/estimate their spectra efficiently, and (iii) optimize parameters to match known zeros.

## 4.1 Parametric Hamiltonian family

Choose a flexible, physically plausible family  $H(\alpha)$  with tunable parameters  $\alpha$ . Examples:

- Regularized/discretized Berry–Keating model with boundary or regularization parameters.
- One-dimensional scattering or quantum-graph families with tunable potentials or edge lengths.
- Small many-body lattice Hamiltonians with adjustable couplings (e.g. BEC-inspired lattices).

## 4.2 Fast spectral probe (pipeline)

For a given  $\alpha$ :

- (a) Use **KAN** to select ansatz parameters  $\theta$  that prepare seed states with good overlap cheaply.
- (b) Run short **VQE** refinements to improve eigenstate overlaps where useful.
- (c) Build a **Krylov subspace** / **SBQD** around seed states to extract multiple approximate eigenvalues in blocks.
- (d) Use **FFT** on time-series  $C(t) = \langle \psi | e^{-iHt} | \psi \rangle$  or short-time evolutions to identify spectral peaks (coarse localization).
- (e) Run **QPE** on promising spectral windows to refine eigenvalues to target precision.

This yields a list of extracted eigenvalues  $\{\tilde{\lambda}_i(\alpha)\}$ .

## 4.3 Objective and metrics

Compare  $\{\tilde{\lambda}_i(\alpha)\}$  to true nontrivial zeros  $\{\gamma_i\}$  (computed classically for the range targeted). Suggested metrics:

- Absolute error:  $\max_i |\tilde{\lambda}_i - \gamma_i|$  for the first  $N$  zeros.
- Pair-correlation / spacing statistics: nearest-neighbor spacing distribution, two-point correlation.
- Trace-formula residual: mismatch between semiclassical spectral density  $\rho(E)$  and the zero density.
- Primes-derived residual: insert  $\tilde{\lambda}$  into an explicit formula for the Chebyshev or prime-counting functions and compare to known values.

Combine into a weighted loss

$$L(\alpha) = w_1 \text{abs\_err} + w_2 \text{stat\_err} + w_3 \text{trace\_res}.$$

Use KAN for surrogate modeling of  $L(\alpha)$  to reduce expensive evaluations.

## 4.4 Optimize over $\alpha$

Use classical optimization (derivative-free, Bayesian, or surrogate-based) to minimize  $L(\alpha)$ . Iterate until a candidate  $\alpha^*$  yields an acceptably small loss.

#### 4.5 Validate and test generalization

- Test  $\tilde{\lambda}(\alpha^*)$  against zeros outside the fitted range (out-of-sample verification).
- Test robustness to noise and parameter drift.

### 5 Practical outcomes and limitations

- If a Hamiltonian matches zeros to arbitrary high accuracy and generalizes, measuring its spectrum provides a direct route to primes — a revolutionary outcome.
- More realistically, one may find Hamiltonians that match zeros approximately over some window; these models are still valuable for exploring physics–number-theory connections.
- Resource demands for optimization and high-precision QPE will be significant; careful error mitigation is required.

### 6 Concrete next steps (implementable)

Possible tasks (user-selectable):

1. **Prototype search code:** Implement the full loop for a chosen parametric family (KAN sampling, VQE, Krylov+FFT spectral extraction, objective evaluation vs. first  $N$  classical zeros, and parameter optimization). Example target: first 20 nontrivial zeros.
2. **FFT-enhanced spectral pipeline:** Build the FFT-based spectral estimation from short-time evolutions plus Krylov/SBQD and hooks to QPE.
3. **Validation scripts:** Functions to compute explicit-formula-based prime estimates from measured zeros and compare to  $\pi(x)$  and Chebyshev functions.
4. **Scoring and visualization:** Produce plots tracking spectral match during optimization and prime-distribution residuals.

### 7 Acknowledgements and references

This document summarizes an idea and an actionable hybrid quantum–classical plan. Relevant literature includes Berry–Keating models, quantum graph spectral theory, semiclassical trace-formula methods, and modern variational/Krylov quantum eigensolvers. The reader is encouraged to consult standard references in analytic number theory and quantum computation for background.