## Fitting procedure for InAs

Below is the detailed fitting procedure used to fit the data for Faraday rotation from InAs using single Drude function. Faraday rotation ( $\theta_F$ ) in terms of right and left circularlarly polarized THz light transmission is given by,

$$\theta_F = -\arctan\left[i\frac{T_r - T_l}{T_r + T_l}\right],$$

These transmissions in thin film limit can be written as,

$$T_r(\omega) = \frac{1+n}{1+n+Z_0 d\sigma_r(\omega)} e^{\frac{i\omega(n-1)\Delta L}{C}}, \quad T_l(\omega) = \frac{1+n}{1+n+Z_0 d\sigma_l(\omega)} e^{\frac{i\omega(n-1)\Delta L}{C}}$$

Where n is the refractive index of the substrate,  $\sigma_{r,l}$  are the conductivities for right and left circularly polarized light, d is the thickness of the film and  $\Delta L$  is the thickness difference of substrates used for reference and the film. Thus Faraday rotation can be written in terms of conductance as,

$$\theta_F = -\arctan\left[\frac{i(G_l - G_r)}{(G_l + G_r) + \frac{2(n+1)}{Z_0}}\right]$$

Conductance in magnetic field can be described by Drude model,

$$G_{\pm} = -i\epsilon_0 \omega d \left[ \frac{\omega_{pD}^2}{-\omega^2 - i\Gamma_D \omega \mp \omega_c \omega} + (\epsilon_{\infty} - 1) \right], \quad Let \ (\epsilon_{\infty} - 1) = \alpha_{\infty}$$

Where  $G_{r,l}=G_{\pm}$ ,  $\omega_{pD}^2$  is the plasma frequency,  $\Gamma_D$  is the scattering rate and  $\omega_c$  is the cyclotron resonance. We can write

$$G_R = \epsilon_0 d \left[ \frac{\omega_{pD}^2}{\Gamma_D - i(\omega + \omega_c)} - i\omega \right]_{\infty}$$

$$G_L = \epsilon_0 d \left[ \frac{\omega_{pD}^2}{\Gamma_D - i(\omega - \omega_c)} - i\omega \alpha_{\infty} \right]$$

Thus for small angles,

$$\theta_F \sim \text{Tan}(\theta_F) = \frac{i(G_R - G_L)}{(G_R + G_L) + 2(n+1)/Z_O}$$

$$Tan(\theta_F) = -\frac{\epsilon_0 d\omega_{pD}^2 \omega_c}{\epsilon_0 d\omega_{pD}^2 (\Gamma_D - i\omega) + \left[\frac{n+1}{Z_0} - i\epsilon_0 d\alpha_\infty \omega\right] [\Gamma_D^2 + \omega_c^2 - \omega^2 - 2i\Gamma_D \omega]}$$

To solve in terms of THz just put  $\epsilon_0=8.85$ . For InP (@4 K) refractive index n= 3.51, Z<sub>0</sub>= 377. For the fits I use,

$$Tan(\theta_F) = -\frac{\epsilon_0 ab}{\epsilon_0 a(c - ix) + [0.1196 - iex][c^2 + b^2 - x^2 - 2icx]}$$

Where,  $a=\omega_{pD}^2d$ ,  $b=\omega_c$ ,  $c=\Gamma_D$  and  $e=d\alpha_\infty$ . I keep a and e constant as they don't vary even when they are not constant. Thus only free parameters are cyclotron resonance  $(\omega_c/2\pi)$  and scattering rate  $(\Gamma_D/2\pi)$ . Effective mass can be extracted from the field dependence of the cyclotron resonance frequency  $(\omega_c=eB/m^*)$ . The spectral weight  $(\omega_{pD}^2d)$  is proportional to the integrated area of each feature in the real part of the conductance. It gives the ratio of carrier density to an effective transport mass:

$$\frac{2}{\pi\epsilon_0} \int G_{D1} d\omega = \omega_{pD}^2 d = \frac{n_{2D} e^2}{m^* \epsilon_0}$$

Mobility is given by,  $\mu = e\Gamma_D/m^*$ .